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Binding energies of ground and isomeric states in neutron-rich ruthenium isotopes: Measurements at JYFLTRAP and comparison to theory

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We report on precision mass measurements of $^{113,115,117}\text{Ru}$ performed with the JYFLTRAP double Penning trap mass spectrometer at the Accelerator Laboratory of University of Jyväskylä. The phase-imaging ion-cyclotron-resonance technique was used to resolve the ground and isomeric states in $^{113,115}\text{Ru}$ and enabled for the first time a measurement of the isomer excitation energies, $E_x(^{113}\text{Ru}^m) = 100.5(8)$ keV and $E_x(^{115}\text{Ru}^m) = 129(5)$ keV. The ground state of ^{117}Ru was measured using the time-of-flight ion-cyclotron-resonance technique. The new mass-excess value for ^{117}Ru is around 36 keV lower and seven times more precise than the previous literature value. With the more precise ground-state mass values, the evolution of the two-neutron separation energies is further constrained and a similar trend as predicted by the BSkG1 model is obtained up to the neutron number $N = 71$.

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I. INTRODUCTION

Neutron-rich nuclei between zirconium ($Z = 40$) and tin ($Z = 50$) exhibit a variety of shapes; several of them even exhibit shape coexistence, where excited states are linked to shapes which differ from that of the nuclear ground state. The diverse manifestations of collectivity in general and nuclear shapes in particular in this region of the nuclear chart have been studied widely, both theoretically and experimentally; see, e.g., Ref. [1] and references therein. The relevant nuclear configurations are not limited to shapes with a comparatively high degree of symmetry such as spheres or axially symmetric ellipsoids with prolate or oblate deformation, but also includes shapes with no remaining rotational symmetry axis: triaxial shapes. There is evidence that the ground states of neutron-rich ruthenium isotopes ($Z = 44$) fall in the latter category [1,2], an interpretation that is further supported by different

models [3,4]. These models typically agree that the effect of triaxial deformation is largest at the midshell and that the effect tapers off when even more neutrons are added to the nucleus, i.e., that sufficiently neutron-rich nuclei revert to an axially symmetric or spherical shape towards the shell closure at $N = 82$.

Structural changes can be studied via a wide range of experimental methods, including laser and decay spectroscopy as well as Coulomb excitation. At the same time, Penning-trap mass spectrometry can be used to explore differences in binding energy which can reveal possible shape transitions [5–7]. With the development of the phase-imaging ion-cyclotron-resonance (PI-ICR) technique [8,9], not only the ground-state binding energies but also the isomer excitation energies down to a few tens of keV [10,11] can be extracted, allowing one to obtain new insight into the nuclear structure.

Masses of neutron-rich ruthenium isotopes up to $A = 116$ [12,13] have been measured before with the JYFLTRAP double Penning trap mass spectrometer [14]. However, for the cases where long-lived isomers are present, namely $^{113,115}\text{Ru}$, the time-of-flight ion-cyclotron-resonance (TOF-ICR) [15] technique used at that time did not provide enough resolving power to separate the ground and isomeric states in ^{113}Ru or to detect the isomer in ^{115}Ru unknown at that time. Therefore these results might have suffered from a systematic shift for

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the reported ground-state mass-excess values [11]. More exotic ruthenium isotopes were studied using the Experimental Storage Ring at GSI [16]. However, ^{117}Ru had the uncertainty increased 2.4 times by the Atomic Mass Evaluation 2020 (AME20) evaluators while the mass-excess value of ^{118}Ru was rejected due to a significant 700-keV deviation from the mass trends [17].

In this work, we report on the direct mass measurement of the ground states of $^{113,115,117}\text{Ru}$ and the isomeric states in ^{113}Ru and ^{115}Ru , the latter being the shortest-lived state ($T_{1/2} = 76(6)$ ms [18]) ever measured at JYFLTRAP so far. The role of deformation for the systematics of masses in this region and the nature of the isomeric state in ^{115}Ru are analyzed within the context of the recent global microscopic models BSkG1 [4] and BSkG2 [19,20] that are based on self-consistent Hartree-Fock-Bogoliubov (HFB) calculations using a Skyrme energy density functional (EDF).

II. EXPERIMENTAL METHOD

The masses of neutron-rich ruthenium isotopes were studied at the Ion Guide Isotope Separator On-Line (IGISOL) facility [21] using the JYFLTRAP double Penning trap mass spectrometer [14] during two experiments. The isotopes of interest were produced in proton-induced fission by impinging a 25 MeV proton beam onto a thin target, ^{232}Th for ^{113}Ru and $^{\text{nat}}\text{U}$ for $^{115,117}\text{Ru}$. First, the fission fragments were stopped in a helium gas cell operating close to 300 mbar from which they were extracted and guided using a sextupole ion guide [22]. Then, the produced ions were accelerated to $30q$ keV and mass-separated based on their mass-to-charge ratio using a 55° dipole magnet. The continuous mass-separated beam was cooled and bunched using the helium buffer gas-filled radio-frequency quadrupole cooler-buncher [23]. Finally, the ion bunches were injected into the JYFLTRAP double Penning trap.

In the first trap of JYFLTRAP, known as the purification trap, the ion bunch was cooled and centered, and the ions of interest were selected utilizing the mass-selective buffer gas cooling technique [24]. After that, the purified ion sample was sent into the second trap, called the precision trap, where the mass measurements took place.

In addition, $^{113}\text{Ru}^{2+}$ ions were produced via the in-trap decay of ^{113}Tc ($T_{1/2} = 152(8)$ ms [18]). The $^{113}\text{Tc}^{+}$ ions, produced via fission, were captured in the first trap, after which the ion motion was cooled for 102 ms. Then a dipolar excitation on the magnetron frequency was applied for 10 ms. During the trapping time a fraction of the $^{113}\text{Tc}^{+}$ ion sample β decayed to $^{113}\text{Ru}^{2+}$. Quadrupolar excitation of 100 ms was used to select the ions of interest by matching the excitation frequency of $^{113}\text{Ru}^{2+}$ ions. After, the $^{113}\text{Ru}^{2+}$ ions were sent to the second trap for the precision mass measurement.

In the presence of a magnetic field of strength B , the mass m of an ion is related to its cyclotron frequency ν_c :

$$\nu_c = \frac{1}{2\pi} \frac{q}{m} B, \quad (1)$$

where q/m is the charge-to-mass ratio of the measured ion. To determine the magnetic field strength precisely, $^{133}\text{Cs}^{+}$ ions

from the IGISOL offline surface ion source station [25] were used as a reference for the mass measurement of $^{113,115}\text{Ru}^{+}$ ground states and $^{117}\text{Ru}^{+}$. For the mass measurement of isomeric states in $^{113,115}\text{Ru}$, the ground-state masses were used as a reference. To account for the temporal magnetic field fluctuations, ruthenium ions and their references were measured alternately. The atomic mass m is determined from the frequency ratio $r = \nu_{c,\text{ref}}/\nu_c$ between the reference ions and the ions of interest:

$$M = \frac{z_i}{z_{\text{ref}}} (M_{\text{ref}} - z_{\text{ref}} m_e) r + z_i m_e, \quad (2)$$

where M_{ref} is an atomic mass of the reference, m_e is an electron mass, and z_i and z_{ref} are charge states of the ion of interest and the reference ion, respectively. The isomer excitation energies were extracted as follows:

$$E_x = (r - 1)[M_{\text{gs}} - z m_e] c^2, \quad (3)$$

where M_{gs} is the ground-state atomic mass, z is the charge state of the reference ion and the ion of interest (both either singly charged, $z = 1$, or doubly charged, $z = 2$, in this experiment), and c is the speed of light in vacuum. Contributions from electron binding energies are on the order of eV and have thus been neglected.

To measure the masses of the ground and isomeric states in $^{113,115}\text{Ru}$, the PI-ICR technique [9,26] was utilized in the precision trap. With the PI-ICR technique, the ion's cyclotron frequency is determined by measuring the sum of the accumulated residual phases of the magnetron (ϕ_-) and cyclotron (ϕ_+) motions that are projected onto a position-sensitive microchannel plate (2D MCP) detector after a phase accumulation time t_{acc} . Using the polar angles of the cyclotron (α_+) and magnetron (α_-) phase images on the detector, the angle between the motion phases with respect to the center spot is $\alpha_c = \alpha_+ - \alpha_-$. This can be used for the determination of the cyclotron frequency:

$$\nu_c = \frac{\alpha_c + 2\pi n}{2\pi t_{\text{acc}}}, \quad (4)$$

with n being the sum of full revolutions performed at the magnetron and modified cyclotron frequencies during the phase accumulation time in the precision trap. We used the following accumulation times for the PI-ICR mass measurements: 557 ms for the $^{113}\text{Ru}^{+}$ ground and isomeric state, 220 ms for the $q = 2+$ ions of ^{113}Ru isomeric state, 200 ms for the $^{115}\text{Ru}^{+}$ ground state, and 100 ms for the $^{115}\text{Ru}^{+}$ isomer (see Fig. 1). The measurement pattern utilized at JYFLTRAP is described in more detail in Refs. [26,27] and the PI-ICR measurement technique in Ref. [9].

For $^{117}\text{Ru}^{+}$, the TOF-ICR technique [15,28] was applied. The ion's cyclotron frequency ν_c in TOF-ICR technique is determined from a time-of-flight resonance measured with the 2D MCP detector, located outside the strong magnetic field of the trap. To enhance the resolving power, the Ramsey method of time-separated oscillatory fields [29,30] was utilized. A short 10-30-10 ms (on-off-on) pattern was used in order to minimize the decay losses (see Fig. 2).

In the mass measurement of ^{113}Ru and ^{115}Ru , the ground state and the isomer were in the precision trap at the same

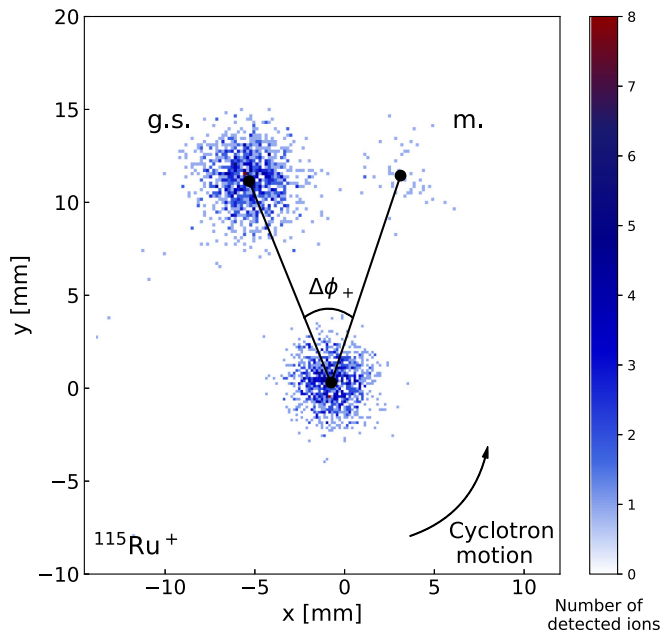


FIG. 1. A PI-ICR measurement of ^{115}Ru ground state versus the isomeric state with a 100 ms accumulation time. Only the projection of cyclotron motion on the 2D MCP is shown. The angle difference $\Delta\phi_+$ leads to an excitation energy of 129(5) keV. The center spot, i.e., without any excitation, is also shown.

time. It is known that when two or more ions of different masses are present in the trap simultaneously, the ion-ion interaction can cause a frequency shift [27]. To account for the ion-ion interaction, a count-rate class analysis [27,31], i.e., analysis of the variation of the frequency ratio with the number of ions stored simultaneously in the trap, was performed for the ground state ion of ^{115}Ru , while for other cases it was not statistically feasible. At JYFLTRAP the systematic uncertainty related to temporal magnetic field fluctuation has been determined to be $\delta B/B = 2.01(25) \times 10^{-12} \text{ min}^{-1} \times \delta t$ [27], where δt is the time between the measurements. In all of the measurements the maximum systematic uncertainty related to the temporal magnetic field fluctuations was calculated but was found to be negligible compared to the

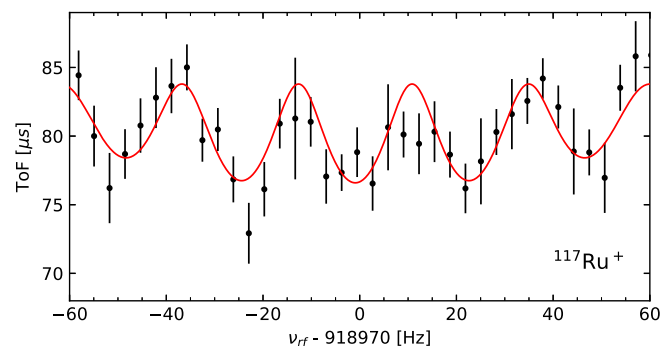


FIG. 2. A TOF-ICR measurement of $^{117}\text{Ru}^+$ using a 10-30-10 ms (on-off-on) Ramsey excitation pattern. The mean data points are shown in black, the fit of the theoretical curve [29] in red.

statistical uncertainty. We added a further mass-dependent uncertainty of $\delta_m r/r = -2.35(81) \times 10^{-10}/u \times (M_{\text{ref}} - M)$ and a residual systematic uncertainty of $\delta_{\text{res}} r/r = 9 \times 10^{-9}$ for measurements where the A/q of the reference and ion of interest were not the same, i.e., when using the ^{133}Cs ions as reference [27]. A systematic uncertainty related to the magnetron phase advancement and systematic angle error were also accounted for in the PI-ICR measurements. A more detailed description on the systematic uncertainties and their determination at JYFLTRAP can be found in Ref. [27].

III. RESULTS

The ground- and isomeric-state masses of $^{113,115}\text{Ru}$ and the ground-state mass of ^{117}Ru are reported in detail below. The measured frequency ratios (r), mass-excess (ME) values, and excitation energies (E_x) are summarized in Table I.

A. ^{113}Ru

The ground-state mass excess of ^{113}Ru , $-71\,874.6(15)$ keV, was determined using $^{133}\text{Cs}^+$ ions as a reference. The isomer excitation energy, $E_x = 100.5(8)$ keV, was determined against the ground state, both as singly charged ions produced directly in fission as well as doubly charged ions produced via in-trap decay of $^{113}\text{Tc}^+$ (for details see Sec. II). This yields a mass excess of $-71\,774.2(17)$ keV for the isomer.

The mass of ^{113}Ru was previously measured at JYFLTRAP by Hager *et al.* [12], using the TOF-ICR technique with a 400 ms quadrupolar excitation time and $^{105}\text{Ru}^+$ ions as a reference. With the AME20 [32] mass value for ^{105}Ru , this results in a mass-excess value of $-71\,826(12)$ keV. The revised value is in between the ground- and isomeric-state mass-excess values reported in this work [see Fig. 3(a)], suggesting that

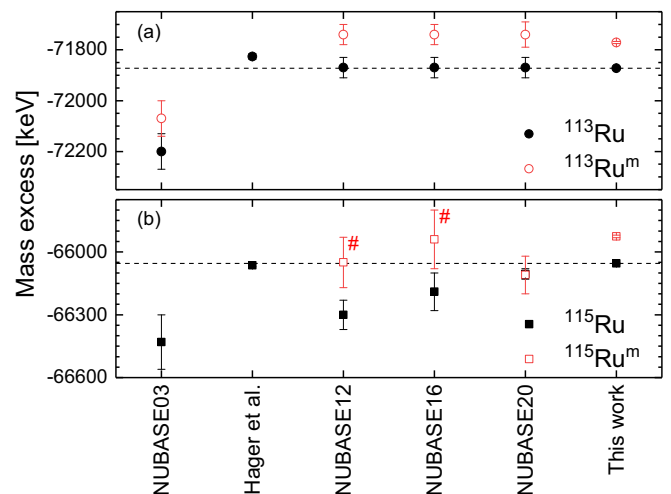


FIG. 3. The mass-excess values determined in this work for the ground states (solid black symbols) and isomers (open red symbols) in (a) ^{113}Ru and (b) ^{115}Ru , in comparison with the revised JYFLTRAP value reported by Hager *et al.* [12,32] and different NUBASE compilations [18,34–36]. The dashed black lines show the ground-state mass-excess values determined in this work. # denotes mass-excess values based on systematics.

TABLE I. The measured frequency ratios ($r = \nu_{c,ref}/\nu_c$) and corresponding mass-excess (ME) values determined in this work using the listed reference ions (Ref.). The charge state z used both for the reference ions and the ion of interest is also listed. The reported uncertainties are total uncertainties. The mass-excess values from AME20 [32] and NUBASE20 [18] ($ME_{lit.}$) and the differences $Diff. = ME - ME_{lit.}$ are given for comparison. All the half-lives $T_{1/2}$ and spin-parity assignments J^π of $^{113,113m,117}\text{Ru}$ are taken from the NUBASE2020 evaluation [18] while spin-parity assignments for $^{115}\text{Ru}^{gs,m}$ are taken from Ref. [33] and this work. # denotes that the spin is based on systematics while parentheses indicate a tentative assignment.

Nuclide	$T_{1/2}$ (ms)	J^π	Ref.	z	r	ME (keV)	$ME_{lit.}$ (keV)	E_x (keV)	$E_{x,lit.}$ (keV)	Diff. (keV)
^{113}Ru	800(50)	(1/2 ⁺)	^{133}Cs	1	0.849 647 289(12)	-71 874.6(15)	-71 870(40)			-5(40)
$^{113}\text{Ru}^m$	510(30)	(7/2 ⁻)	^{113}Ru	1	1.000 000 951(10) ^a	-71 774.6(18)		100.0(11)		
			^{113}Ru	2	1.000 000 963(14) ^b	-71 773.3(21)		101.3(15)		
					Final value:	-71 774.2(17)	-71 740(50)	100.5(8) ^c	131(33)	-34(50)
^{115}Ru	318(19)	(3/2 ⁺)	^{133}Cs	1	0.864 742 653(23)	-66 054.7(28)	-66 105(25)			50(26)
$^{115}\text{Ru}^m$	76(6)	(9/2 ⁻)	^{115}Ru	1	1.000 001 206(47)	-65 925.6(58)	-66 110(90)	129(5)	82(6)	184(91)
^{117}Ru	151(3)	3/2 ⁺ #	^{133}Cs	1	0.879 843 74(52)	-59 526(64)	-59 490(430)			-36(435)

^aMeasured with 1⁺ ions produced directly in fission.

^bMeasured with 2⁺ ions produced in in-trap-decay of ^{113}Tc .

^cWeighted average of the two measurements.

a mixture of states was measured in Ref. [12]. A similar effect was observed in Rh isotopes, as reported in Ref. [11].

The reported mass-excess values are in agreement with the NUBASE20 evaluation [18], where it was correctly assumed that the value measured in Ref. [12] was a mixture of the ground state and an isomer at 131(33) keV. To date, the isomeric-state excitation energy was not based on direct experimental observations but on the suggestion that it has to lie in between the states at 98 and 164 keV in ^{113}Ru [37,38]. In this work, we have confirmed this hypothesis by determining the excitation energy for the first time and by placing the isomer just above the 98 keV state (see Fig. 4). The production of both long-lived states in ^{113}Ru in the β decay of ^{113}Tc is also in agreement with the work by Kurpeta *et al.* [37].

B. ^{115}Ru

The ground state mass excess, $-66\,054.7(28)$ keV, was measured against a $^{133}\text{Cs}^+$ reference. The isomer excitation energy, $E_x = 129(5)$ keV, was determined against the ground

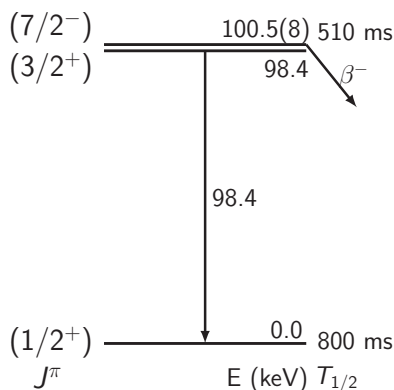


FIG. 4. Proposed partial level scheme of ^{113}Ru based on this work and Ref. [39].

state, resulting in a mass excess of $-65\,925.6(58)$ keV for the isomer.

Our ground-state mass excess value is in agreement with the previous TOF-ICR-based JYFLTRAP measurement ($ME = -66\,064.0(69)$ keV [12,32]) after adjusting for the updated mass of the reference ^{120}Sn ion. In our previous work we have observed that for nuclei with low-lying isomeric states the masses obtained with the TOF-ICR method are a weighted average of the ground state and the isomer masses [11]. In the case of ^{115}Ru , an apparent absence of the isomer influence on the measured mass can be explained by a relatively short half-life of the isomeric state ($T_{1/2} = 76$ ms [18]) compared to the 300 ms excitation time used in Ref. [12].

Figure 3(b) shows a comparison of our measurement with the values reported in NUBASE evaluations on nuclear and decay properties from 2003 [34], 2012 [35], 2016 [36], and 2020 [18] as well as the revised JYFLTRAP value of Ref. [12]. Changes between different editions of NUBASE can be explained as due to varying input data. In NUBASE03 [34], the only entry for ^{115}Ru was from a β -decay endpoint energy study [40]. After the JYFLTRAP measurement by Hager *et al.* [12], a long-lived isomeric state in ^{115}Ru was discovered [41], and the evaluators of NUBASE12 [35] applied a special procedure for mixtures of isomeric states assuming the excitation energy to be 250(100) keV. In NUBASE16 [36], the β -decay endpoint energy study was excluded from the global fit and the only remaining information was from Ref. [12]. Finally, in NUBASE20 [18], the energy of the isomeric state was adjusted to 82(6) keV based on the value originally proposed in Ref. [41]. However, the isomeric-state excitation energy seems not to be taken into account for the mass-excess value of the isomer but only for its uncertainty.

C. ^{117}Ru

The value determined in this work, $-59\,526(64)$ keV, is in agreement with AME20 [32] and it is almost seven times more precise. The mass-excess value adopted in

AME20, $-59\,490(430)$ keV [32], is based on storage-ring measurements [16,42] but with the uncertainty artificially increased by evaluators [17]. The only known isomeric state has a half-life of $2.49(6)\mu\text{s}$ [18] which is much shorter than the measurement cycle used in this work.

IV. DISCUSSION

In this section, we discuss the experimental results and compare them to the BSkG family of models of nuclear structure [4,19,20]. This section is organized as follows: we first establish the theoretical framework in Sec. IV A and then proceed to study first the trends of the ground state (g.s.) binding energies of neutron-rich Ru isotopes in Sec. IV B. Section IV C discusses the isomeric state in ^{115}Ru as well as the implication of our measurement of its excitation energy.

A. Theoretical framework

The BSkG family of models responds to the need for reliable data on the structural properties of exotic nuclei in different fields of research and in astrophysics in particular. These models are based on an empirical energy density functional (EDF) of Skyrme type that models the effective in-medium nucleon-nucleon interaction. The concept of an EDF allows for a global yet microscopic description of all relevant quantities at a reasonable computational cost. The coupling constants of the EDF are the main element of phenomenology in this type of model and have to be adjusted to experimental data. Since binding energies are crucial ingredients for the modeling of nuclear reactions, the ensemble of known nuclear masses is a key ingredient of the parameter adjustment of the BSkG models. Because of this, these models reach root-mean-square (rms) deviations better than 800 keV on the thousands of masses included in AME20 [32]. This performance is not at all competitive with the uncertainties of the measurements we report on here, but it nevertheless reflects the state-of-the-art in global mass modeling: it is only matched by some of the older BSk models that were adjusted in the same spirit [43], microscopic-macroscopic approaches [44], and empirical models [45]. The latter two types of model become particularly accurate when refined with machine learning techniques [46], but either do not extend their predictions to other observables or struggle to describe them with the same parameter values deduced from the masses.

The BSkG-family so far counts two entries: BSkG1 [4] and BSkG2 [19,20]. Both models combine a description of many hundreds of measured charge radii and realistic predictions for the properties of infinite nuclear matter with a description of the AME20 masses with similar accuracy (rms deviations of 741 and 678 keV, respectively). Although some of the BSk models reach an rms deviation below 600 keV [43], BSkG1 and BSkG2 are better adapted to study the neutron-rich Ru isotopes as they rely on a three-dimensional representation of the nucleus, thereby accommodating naturally the triaxial deformation that is known to be particularly relevant for this region of the nuclear chart. BSkG2 incorporates a full treatment of the so-called ‘time-odd’ terms in an EDF [19] and improves systematically on the description of fission properties com-

pared to its predecessor [20]. Since (i) the inclusion of the time-odd terms did not result in a meaningful improvement of our global description of binding energies and (ii) fission properties are not directly related to the masses, *a priori* we expect BSkG1 and BSkG2 to be of roughly equal quality for the task at hand and therefore we will compare experiment to both models in what follows.

Large-scale EDF-based models of nuclear structure such as the BSk- and BSkG-models describe the nucleus in terms of one single product wavefunction, typically of the Bogoliubov type. The simplicity of such an *ansatz*, as compared to the complexity of the many-body problem, is compensated for by allowing for spontaneous symmetry breaking in the mean fields. By considering such deformed configurations EDF-based models can account for a large part of the effects of nuclear collectivity on bulk properties such as masses while remaining at the mean-field level and thus keeping calculations tractable. Nevertheless, symmetry breaking comes at considerable computational cost. For all calculations that we report on, we employed the MOCCa code [47] to represent the single-nucleon wavefunctions on a three-dimensional coordinate mesh. All numerical parameters such as the mesh point spacing are identical to those employed in the adjustment of both BSkG models [4,19].

In a three-dimensional calculation, the quadrupole deformation of a nucleus of mass A can be described by way of the (dimensionless) deformation β_2 and the triaxiality angle γ , defined as

$$\beta_2 = \frac{4\pi}{3R^2A} \sqrt{Q_{20}^2 + 2Q_{22}^2}, \quad (5)$$

$$\gamma = \text{atan}(\sqrt{2}Q_{22}/Q_{20}), \quad (6)$$

where $R = 1.2A^{1/3}$ fm. The quadrupole moments Q_{20} and Q_{22} are defined in terms of integrals of the total nuclear density and spherical harmonics; see for instance Ref. [4]. Axially symmetric prolate and oblate shapes correspond to $\gamma = 0^\circ$ and 60° , respectively, while intermediate values of the triaxiality angle in between those two extremes indicate triaxial shapes.

We show in Fig. 5 the potential energy surface (PES) of ^{115}Ru in the β - γ plane as obtained with BSkG2, calculations with BSkG1 leading to a similar PES. Since ^{115}Ru has an odd number of nucleons, Fig. 5 shows the result of so-called false-vacuum calculations, where we constrained the expected number of neutrons to $\langle N \rangle = 71$, but otherwise treated the nucleus as if it were even-even. We emphasize that all the calculations for which we report masses do not rely on this approximation: for both BSkG1 and BSkG2 our treatment of the odd-mass Ru isotopes includes self-consistent blocking of a neutron quasiparticle. For BSkG2, we also include the energy contribution of the finite spin and current densities induced by the presence of the odd neutrons. For more details on our treatment of odd-mass and odd-odd nuclei, see the discussion in Ref. [19]. A complete calculation for ^{115}Ru that includes blocking leads to the deformation shown as a black star on Fig. 5; its offset with respect to the minimum of the false-vacuum calculations is due to the polarization induced by the odd neutron.

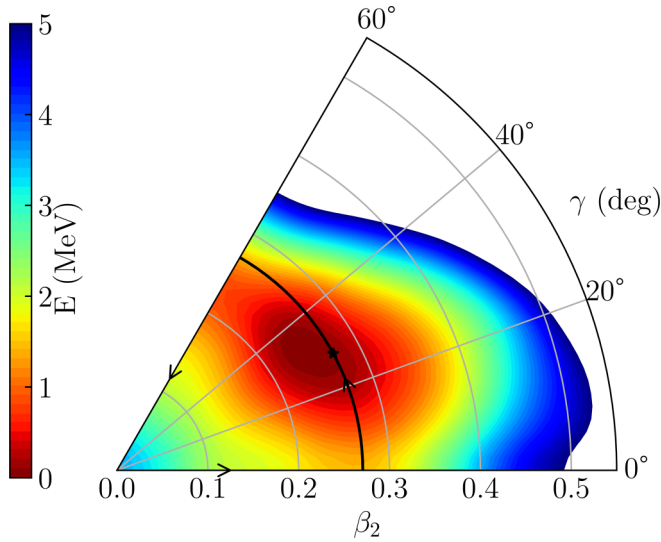


FIG. 5. Potential energy surface in the (β, γ) plane for false-vacuum calculations (see text) of ^{115}Ru with BSkG2. The trajectory followed by the Nilsson diagram in Fig. 9 is indicated by black arrows. The location of the minimum obtained in a complete calculation of ^{115}Ru is indicated by a black star.

Qualitatively, the false-vacuum PES of ^{115}Ru looks similar to the PES of ^{112}Rh that we discussed in Ref. [11]: we observe a somewhat broad triaxial minimum near $\gamma = 30^\circ$ of significant quadrupole deformation. Close inspection reveals some quantitative differences: $\beta_2 \approx 0.27$ is here somewhat smaller than the value 0.3 obtained for ^{112}Rh for instance. Another difference is the energy gain due to triaxiality: the difference between the oblate saddle point and the minimum on Fig. 5 is about 800 keV, while it exceeds 1 MeV for ^{112}Rh . This can be linked to the four additional neutrons in ^{115}Ru compared to ^{112}Rh : as we approach the $N = 82$ shell closure, the neutrons have less freedom to exploit quadrupole correlations and the importance of (static) quadrupole deformation in general and triaxial deformation in particular diminishes.

B. The g.s. masses of Ru isotopes and their trends

For the chain of Ru isotopes between $N = 65$ and $N = 73$, BSkG1 reproduces the absolute g.s. binding energies best: the deviation with respect to experiment for the absolute mass excesses averages to 360 keV and never exceeds 640 keV. The performance of BSkG2 is not as good: an average deviation of 650 keV with a deviation of up to 1.175 MeV for ^{115}Ru . Interestingly, the sign of the deviation is consistent: both models overbind these Ru isotopes and hence produce mass excesses that are too large in absolute size. As discussed before, the experimental uncertainties are several orders of magnitude beyond the accuracy of global models like BSkG1 and BSkG2: instead of comparing absolute masses in more detail, we will focus in what follows primarily on the trends of mass differences.

We start with the two-neutron separation energy S_{2n} , defined as

$$S_{2n}(Z, N) = \text{ME}(Z, N - 2) - \text{ME}(Z, N) + 2 \text{ME}(0, 1), \quad (7)$$

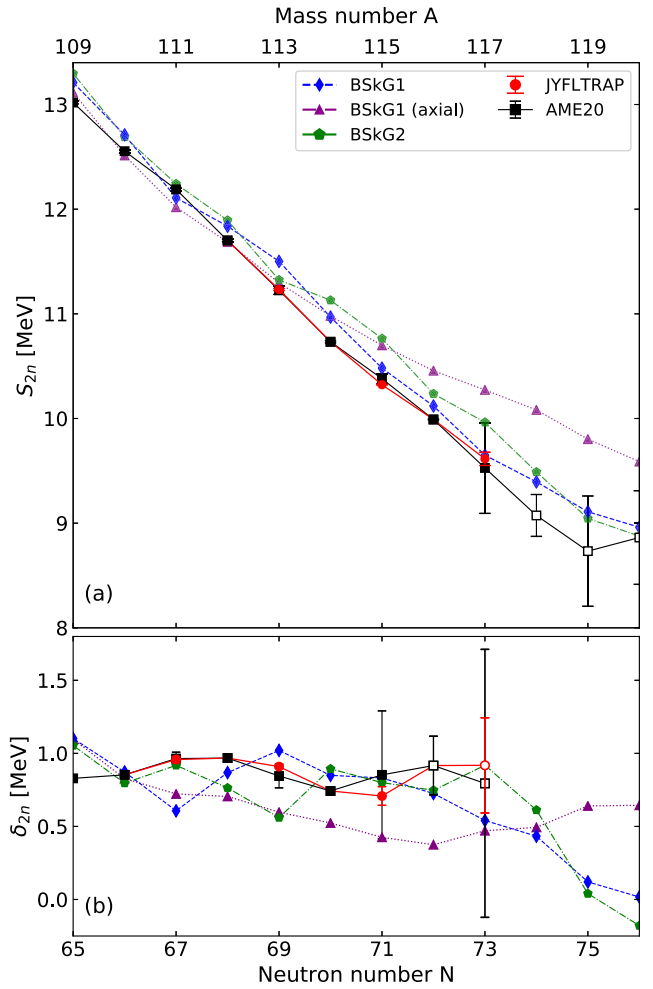


FIG. 6. Comparison of mass differences along the Ru isotopic chain: experimental values either as tabulated in AME20 (black squares) or updated with the results of our new measurements (red circles) versus calculated values obtained with BSkG1 (blue diamonds), BSkG1 axial (purple triangles, see text), and BSkG2 (green pentagons). Open markers for the experimental results represent values at least partially based on extrapolated mass values from AME20 [32]. Top panel (a): two-neutron separation energies S_{2n} . Bottom panel (b): two-neutron shell gaps δ_{2n} .

where $\text{ME}(Z, N)$ is the mass excess of a nucleus with Z protons and N neutrons and $\text{ME}(0, 1)$ is the mass excess of the neutron. The top panel of Fig. 6 compares the S_{2n} values derived from the newly measured masses to the values reported in the AME20 [32] evaluation and the two mass models. We also show the results of the less general calculations with BSkG1 reported on in Ref. [4,11], which restrict the nucleus to axially symmetric configurations.

For the less exotic $^{109,111,113}\text{Ru}$, all three calculations with BSkG-models reproduce the general trend of the experimental S_{2n} rather well, although deviations on the order of several hundred keV are clearly visible. For the BSkG1 model, the description of the more neutron-rich isotopes follows the trend of the more stable ones, systematically overestimating the S_{2n} values by a small value. BSkG2 also overestimates

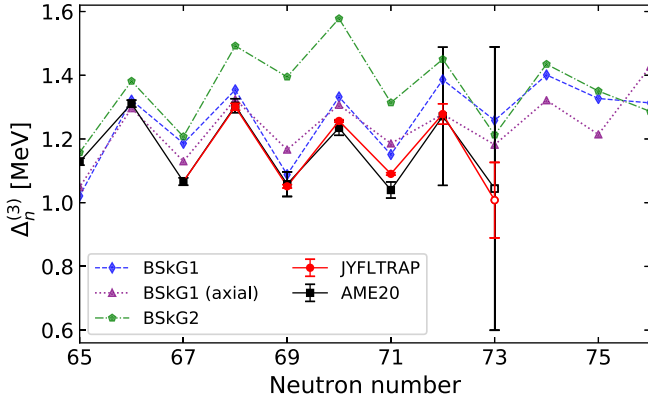


FIG. 7. Same as for Fig. 6, but for the three-point neutron gap $\Delta_n^{(3)}$.

the separation energies and describes their overall trend, but with deviations that are somewhat larger than those of its predecessor. Calculations with BSkG1 that are restricted to axial shapes, however, entirely miss the experimental trend.

We can furthermore discuss the slope of the S_{2n} curve by introducing the empirical two-neutron shell gaps δ_{2n} :

$$\delta_{2n}(Z, N) = S_{2n}(Z, N) - S_{2n}(Z, N + 2), \quad (8)$$

which we show in the bottom panel of Fig. 6. The new JYFLTRAP measurement for ^{115}Ru clearly establishes that the slope of the S_{2n} in this isotopic chain evolves smoothly at least until $N = 71$. Although the corresponding curves are less regular, the BSkG1 and BSkG2 results produce δ_{2n} values that remain close to experiment up to $N = 71$. For the heavier $N = 72, 73$, and 74 isotopes, whose experimental δ_{2n} values are at least partially based on extrapolated AME20 values, the two models predict no major change in slope either. It is only for $N = 75-76$ that BSkG1 and BSkG2 predict a change in slope that is correlated with the disappearance of triaxial deformation for $N \geq 76$. For ^{120}Ru and even more neutron-rich isotopes, the models predict axially symmetric prolate shapes with deformation that gradually diminishes towards $N = 82$.

Finally, we discuss the three-point neutron gaps $\Delta_n^{(3)}(Z, N)$:

$$\Delta_n^{(3)}(Z, N) = \frac{(-1)^N}{2} [\text{ME}(Z, N + 1) + \text{ME}(Z, N - 1) - 2 \text{ME}(Z, N)]. \quad (9)$$

This quantity estimates the average distance between the curves that interpolate the masses of the even- N and odd- N isotopes, respectively, as a function of neutron number. It is particularly sensitive to the neutron pairing, but it can also be affected by variations in the structure of these isotopes with N . The new experimental results confirm the continuation of the trend of less exotic isotopes: the three-point gaps for the even- N isotopes at $N = 66, 68, 70$, and 72 are all equal within error bars, see Fig. 7. For $N = 70$, our new result actually brings the $\Delta_n^{(3)}$ value more in line with this trend. The updated value of $\Delta_n^{(3)}$ for $N = 71$ falls significantly

out of the uncertainty range of AME20, which reflects the lack of accuracy of the AME20 estimate for the excitation energy of the isomeric state of ^{115}Ru . Nevertheless, it is not dramatically larger than the gap values for $N = 69$ and $N = 71$.

The BSkG2 model generally overestimates $\Delta_n^{(3)}$, and its curve exhibits features at $N = 68, 69$, and 70 that are not seen in the experimental data. BSkG1 on the other hand, provides a fair description of the experimental results, whether including or not triaxial deformation. Yet even this model is clearly not without flaws: the deviation of the full calculation with respect to experiment grows with N from $N = 69$ onwards. In this respect, the deviation between the calculated BSkG1 value and the updated point at $N = 73$ (which incorporates the recommended AME20 binding energy for ^{118}Ru) seems ominous. We note in passing that both BSkG models systematically overestimate $\Delta_n^{(3)}$ along odd- Z isotopic chains, which we discovered for the first time during the study of neighboring Rh isotopes in Ref. [11]. Similarly, both models overestimate the calculated three-point proton gaps in odd- N isotopic chains. The common origin of these issues is the failure of both models to account for a small amount of binding energy in odd-odd nuclei that is usually ascribed to the residual interaction between the two odd nucleons; see Ref. [19]. This issue does not affect our discussion here, but it explains why both models describe much better the three-point neutron gaps in even- Z Ru isotopes than in odd- Z Rh isotopes.

We have established that the performance of BSkG2 for the $N = 65-71$ Ru isotopes is worse than that of BSkG1 for absolute masses as well as all mass differences discussed. Since these models are the result of a complicated parameter adjustment which is global in scope, it is hard to pinpoint a particular source of this (local) deficiency. As we remarked in the previous section, we did not *a priori* expect that BSkG2 would offer an improved description of the measured masses. Although the difference we observe between models indicates BSkG1 as the tool of choice for future studies of this region, this does not imply that BSkG2 is a step backwards compared to its predecessor. The newer model presents a different compromise on the very large number of observables included in the parameter adjustment, leading to a worse description of the nuclei we study here but also to an improved description of other observables [19].

To close this section, we note again that our new measurement indicates a rather uneventful continuation to $N = 71$ of the trends of binding energies and mass differences as established for less exotic isotopes. This can be interpreted as experimental confirmation that the structural evolution of nuclei in this isotopic chain is smooth rather than dramatic. From the point of view of the BSkG models this was expected: from $N = 55$ onwards, the Ru isotopes exhibit triaxial deformation that smoothly evolves with neutron number until $N = 76$. The authors of Ref. [33] relied on the Woods-Saxon single-particle spectrum of Ref. [48] to interpret the change in (tentative) ground state spin assignment in $^{113-115}\text{Ru}$ [(1/2⁺) and (3/2⁺), respectively] as a sign of a shape transition from prolate to oblate deformation. The trend of masses and mass differences does not seem to support such a scenario.

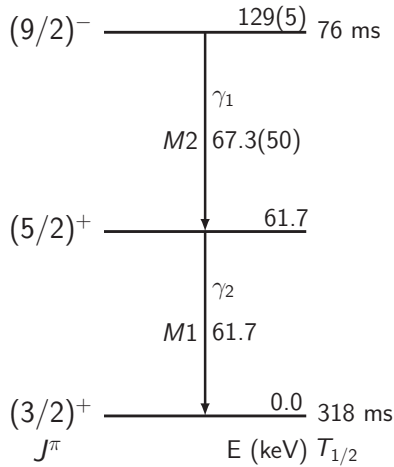


FIG. 8. Proposed level scheme of ^{115}Ru based on this work and Refs. [33,41].

C. The isomer in ^{115}Ru

The isomeric state in ^{115}Ru was reported for the first time in Ref. [41], discussing the analysis of a β decay experiment. The authors observed that the 61.7-keV γ ray is not in coincidence with a β particle or any other γ ray. In addition, the half-life extracted from this transition, $T_{1/2} = 76(6)$ ms, differed from the half-life obtained for the ^{115}Ru ground state [$T_{1/2} = 318(19)$ ms]. Consequently, it was assumed that the isomeric state deexcites via an unobserved γ ray having energy below Ru K x rays ($E \approx 20$ keV), which we label γ_1 , followed by an emission of the 61.7 keV γ ray, labeled as γ_2 .

With the assumption of the energy of γ_1 being below 20 keV, the observed ruthenium K x rays were associated solely with the emission of K internal conversion electrons from the γ_2 transition. This observation enabled a determination of the γ_2 K internal conversion coefficient ($\alpha_K = 2.7(6)$ [41]) by calculating the ratio of the ruthenium K x rays and the γ_2 transitions. The new isomer excitation energy reported in this work renders previous calculations incorrect. However, if one assumes that (i) the total intensity (γ rays and internal conversion electrons emission) of γ_1 and γ_2 is identical, (ii) γ_1 has a pure $M2$ character, and (iii) γ_2 has a pure $M1$ character, the observed ratio of the ruthenium K x rays to γ_2 would be equal to 2.8(8). Any other assumptions regarding the multipolarity of both transitions would lead to a ratio that differs significantly from the experimental value of 2.7(6) [41]. Therefore, we propose $M2$ and $M1$ multiplicities for γ_1 and γ_2 , respectively. By assigning $(3/2)^+$ as the ground-state spin-parity as proposed in [33] from a detailed β -decay spectroscopy experiment of ^{115}Ru , a tentative $(9/2)^-$ isomer assignment can be adopted; see Fig. 8.

A precise description of the level scheme of ^{115}Ru is beyond the capabilities of current large-scale models such as BSkG1 and BSkG2, but we can use them to gain a qualitative understanding of the existence of the isomeric state. To this end, we show in Fig. 9, the Fermi energy and the single-particle energies for both neutrons and protons obtained in

false-vacuum calculations for ^{115}Ru with BSkG2 along the trajectory in the β - γ plane indicated by the arrows in Fig. 5. Although symmetry-breaking allows models such as BSkG1 and BSkG2 to grasp a significant part of the effect of collectivity on nuclear structure, here is where we pay the price: we can no longer use the quantum numbers of an operator associated with a broken symmetry to label single-particle states. At the spherical point, on the utmost left and right of Fig. 9, no symmetry is broken and all single-particle levels are simultaneous eigenstates of three operators with three associated quantum numbers: the angular momentum squared \hat{J}^2 with quantum number J , parity \hat{P} with quantum number π , and the z component of the angular momentum \hat{J}_z with quantum number K . The quantum numbers of the orbitals at the spherical point are indicated in the traditional spectroscopic notation on the right of Fig. 9. Along the first segment of the path on Fig. 5, we break rotational symmetry but conserve axial symmetry: the levels in the leftmost column are no longer eigenstates of \hat{J}^2 but retain the K quantum number,¹ which is indicated by colors in Fig. 5. When exploring finite values of γ along the second segment of the path in Fig. 5, axial symmetry is broken and K can no longer be used to label the single-particle states, hence the absence of colors in the middle column of Fig. 9. The final segment of the path explores oblate shapes which are axially symmetric, such that levels in the right column of Fig. 9 can again be color coded. For all our calculations we conserve parity, such that π is a good single-particle quantum number along the entire path that we can use to distinguish between levels of positive (full lines) and negative parity (dashed lines) in all columns of Fig. 9.

This loss of single-particle quantum numbers also translates to the many-body state: the BSkG models cannot currently offer definite angular momentum assignments for calculated ground states for odd-mass and odd-odd nuclei. Doing so would require symmetry-restoration techniques [49] whose application is presently still out of the scope of global models for reasons of their numerical cost and because of formal issues with the type of EDF assumed for the BSkG models. We are, however, not entirely without options: we can calculate expectation values $\langle i | \hat{J}_z | i \rangle$, which will not be half-integer multiples of \hbar but which nevertheless tell us something about the angular momentum of the single-particle state $|i\rangle$. In the limit of a noninteracting particle-core model of the ground states of odd-mass Ru isotopes, the angular momentum expectation value of the odd neutron will also be the expectation value of angular momentum of the many-body state.

We discussed a qualitatively similar Nilsson diagram obtained for ^{112}Rh in Ref. [11] and repeat here a few observations that are common to both nuclei before discussing the isomer. Local minima in the PES correspond to deformations for which the single-particle level density near the Fermi energy is low: for nuclei with $Z = 43, 44$, and 45 , the protons drive the appearance of triaxial deformation since their single-particle spectrum at $\beta_2 \approx 0.28$ – 0.3 , $\gamma \approx 30^\circ$ is very sparse.

¹For axially symmetric configurations, we always align the symmetry axis with the z axis in the simulation volume.

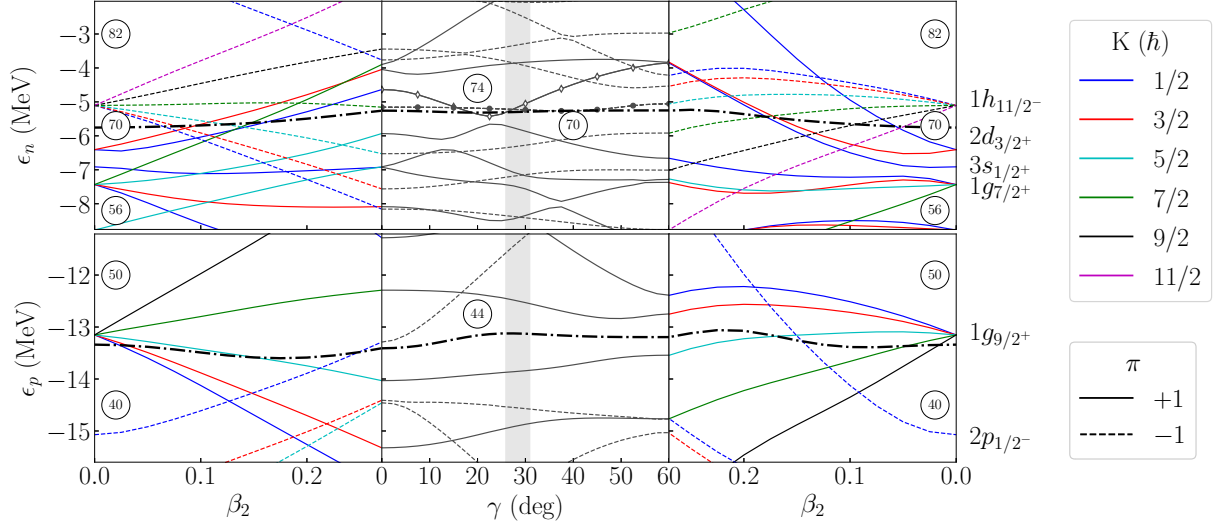


FIG. 9. Eigenvalues of the single-particle Hamiltonian for neutrons (top row) and protons (bottom row) along the path in the β - γ plane indicated by arrows in Fig. 5 for ^{115}Ru (see text for details). The Fermi energy is drawn as a dash-dotted line, while full (dashed) lines indicate single-particle levels of positive (negative) parity. The three indicated regions correspond to axially symmetric prolate shape with (left column), fixed total quadrupole deformation $\beta_2 = 0.27$ with varying γ (center column), and axially oblate shape with $\gamma = 60^\circ$ (right column). The vertical gray band in the center panels is centered at $\gamma = 28.4^\circ$, the value obtained in a complete, i.e., blocked, calculation of ^{115}Ru . The quantum numbers of the shells at sphericity are indicated on the right-hand side. Two-neutron levels near the Fermi energy are highlighted by markers in the middle column: these are the positive and negative parity levels referred to in the text, respectively, as $|\diamond\rangle$ and $|\bullet\rangle$.

In this region of the PES only positive parity states orbital are near the Fermi energy, matching the parity assignments of all even- N Tc and Rh isotopes. The single-particle level density of the neutrons on the other hand is much higher, resulting in a closely spaced set of levels with different parities near the Fermi energy.

We interpret the close interleaving of positive and negative parity neutron states with different angular momentum content as the origin of the isomeric state in ^{115}Ru . Two-neutron states are nearly degenerate near the Fermi energy at the location of the minimum of the PES: these are highlighted in the middle column of Fig. 9 and we will refer to them by their markers: $|\diamond\rangle$ and $|\bullet\rangle$. These levels differ in their parity, but also in their angular momentum content: near $\gamma = 30^\circ$ the positive parity state has an average $\langle \diamond | \hat{J}_z | \diamond \rangle \approx 0.73\hbar$, while that of the negative parity state is significantly larger, $\langle \bullet | \hat{J}_z | \bullet \rangle \approx 4.13\hbar$. Since the odd neutron can be assigned to each of these levels, we expect the appearance of two low-lying levels with opposite parity in the spectrum of ^{115}Ru that are close in energy yet differ substantially in their angular momentum, hence one of them being an isomer. Finally, the $(5/2)^+$ state in between the g.s. and the isomer on Fig. 8 could be rotational in character: taking the calculated moments of inertia of ^{115}Ru and under the assumption of a rigid triaxial rotor, a $1\hbar$ change in total angular momentum corresponds to about 88 keV of excitation energy.

Moving beyond simple arguments based on a noninteracting particle-core picture and the Nilsson diagram, we explicitly calculated the lowest-lying configuration of each parity in ^{115}Ru with both BSkG1 and BSkG2. One of these is the calculated g.s., whose binding energy figured in the previous section: for BSkG1 this is the state with positive

parity and for BSkG2 this is the one with negative parity. In both cases, we find an excited state of opposite parity at low excitation energy; 33 and 90 keV for BSkG1 and BSkG2 respectively. For BSkG2, we have direct access to the average many-body angular momentum along the z axis: a small value $\langle J_z \rangle \approx 0.7\hbar$ for the positive parity state and a large one $\langle J_z \rangle \approx 3.1\hbar$ for the negative parity state. These calculations support our conclusions drawn from the Nilsson diagram and the calculated excitation energy are very roughly comparable to the experimental isomer excitation energy. These results should not be overinterpreted: all relevant energy differences are very small and the neutron spectrum in Fig. 9 is very complicated. Small changes to any aspect of the model will affect the precise location of level crossings and therefore the ordering of levels. Our calculated excitation energies should thus not be taken as a precise prediction, but rather as a confirmation that two states of opposite parity that differ little in energy can be constructed with different angular momentum content. Predicting their ordering and energy difference with accuracy is beyond BSkG1 and BSkG2, or for that matter, any large-scale model that we are aware of.

The same mechanism can be used to interpret the isomerism in nearby $N = 71$ isotones: isomeric states with half-lives on the order of seconds or longer have been observed in ^{116}Rh , ^{118}Ag , and ^{119}Cd whereas shorter-lived isomeric states are known in ^{114}Tc and ^{117}Pd [18]. For $Z = 42$ – 46 , one can expect from Fig. 9 triaxial deformation with a sparse proton single-particle spectrum and two low-lying states arising from neutron orbitals of different parities. The experimental systematics extend much further: in the entire range of $Z = 43$ – 57 , low-lying isomers have been

observed [18]. A more in-depth study of isomerism in the $N = 71$ isotones would certainly require more diagrams like Fig. 9 for larger proton numbers and is outside of the scope of this study. Nevertheless, we remark that both BSkG1 and BSkG2 predict triaxial deformation for almost all $N = 71$ isotones in the range $Z = 40$ – 60 .²

V. SUMMARY

The masses of $^{113,115,117}\text{Ru}$ have been measured using the Penning-trap mass spectrometry at the JYFLTRAP double Penning trap. The ground- and isomeric states in $^{113,115}\text{Ru}$ have been separated and masses measured using the PI-ICR technique. The isomer excitation energies were determined directly for the first time. The high-precision measurements reported in this work place the $(7/2)^-$ isomeric state in ^{113}Ru at 100.5(8) keV, just above the $(3/2^+)$ level at 98.4(3) keV [39], but still in agreement with the previous prediction of 133(33) keV [18]. For $^{115}\text{Ru}^m$, the excitation energy was found to be 129(5) keV, which is significantly larger than proposed in Ref. [41] or the value listed in the most recent NUBASE evaluation, 82(6) keV [18].

The determined ground-state masses of $^{113,117}\text{Ru}$ are in excellent agreement with the atomic mass evaluation [32]. For ^{115}Ru , we report a mass-excess value which is 50(26) keV larger than reported in AME20 [32]. However, it is in agreement with the previous JYFLTRAP mass measurement by Hager *et al.* [12]. With the mass values determined in this work, the trend in the two-neutron separation energies continues smoothly.

The experimental results have been compared with the global BSkG1 [4] and BSkG2 [19,20] models, which allow for triaxially deformed shapes. Detailed calculations were performed for the structure of ^{115}Ru . In the predicted triaxial deformation, the proton single-particle spectrum was found to

²The only exceptions occur for BSkG1 near the $Z = 50$ shell closure: ^{118}Ag , ^{119}Cd , ^{120}In , and ^{121}Sn remain axially symmetric.

be sparse and the predicted low-lying states arise from neutron orbitals with different parities. More systematic studies on the isomeric states in this triaxially deformed region would be needed to shed more light on the reasons for the isomerism in these nuclei.

Note added in proof. Recently, we have been made aware of an independent mass measurement of ^{113}Ru performed at RIKEN using the ZD-MRTOF-MS device [50]. The reported result, $\text{ME}_{\text{lit.}} = -71\,866(7)$ keV, differs by $-8.6(72)$ keV from the ground state mass-excess value reported in this work, $\text{ME} = -71\,874.6(15)$ keV. This small deviation might be related to a limited resolving power of the ZD-MRTOF-MS which did not enable separation of the low-lying isomeric state, $E_x = 100.5(8)$ keV. A similar conclusion was made by the authors of Ref. [50] in the case of ^{112}Rh .

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