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PESTOV IDENTITIES AND X-RAY TOMOGRAPHY ON MANIFOLDS OF LOW REGULARITY

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ABSTRACT. We prove that the geodesic X-ray transform is injective on scalar functions and (solenoidally) on one-forms on simple Riemannian manifolds (M, g) with $g \in C^{1,1}$. In addition to a proof, we produce a redefinition of simplicity that is compatible with rough geometry. This $C^{1,1}$ -regularity is optimal on the Hölder scale. The bulk of the article is devoted to setting up a calculus of differential and curvature operators on the unit sphere bundle atop this non-smooth structure.

1. INTRODUCTION

How regular does a Riemannian metric have to be for the geodesic X-ray transform to be injective? It is well known (see e.g. [Muk77, Muk81, Rom86, AR97]) that on a smooth simple Riemannian manifold this injectivity property holds. If the regularity is too low, the question itself falls apart: If the Riemannian metric is $C^{1,\alpha}$ for $\alpha < 1$, then the geodesic equation can fail to have unique solutions [Har50, SS18]. Therefore it is indeed in a sense optimal on the Hölder scale when we prove that on a $C^{1,1}$ -smooth simple Riemannian manifold the geodesic X-ray transform is injective on scalars and one-forms, the latter one up to natural gauge.

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The geodesic X-ray transform is ubiquitous in the theory of geometric inverse problems. It appears either directly or through linearization in many imaging problems of anisotropic and inhomogeneous media. Most inverse problems have been studied in smooth geometry but the nature is not smooth. The irregularities of the structure of the Earth range from individual rocks (zero-dimensional, small) to interfaces like the core–mantle boundary (two-dimensional, global scale). Irregularity across various scales and dimensions are most conveniently captured in a single geometric structure of minimal regularity assumptions. Specific kinds of irregularities can well be analyzed further, but we restrict our attention to a uniform and global but low regularity.

We prove this injectivity result by using a Pestov identity, an approach that can well be called classical (cf. [Muk77, Muk81, Rom86, AR97, PSU14b, IM19, Uhl14, PSU23]). What requires care is keeping track of regularity. The manifold does not have natural structure beyond $C^{1,1}$, so regularity beyond is both useless and inaccessible. The natural differential operators on the manifold and its unit sphere bundle are not smooth, and only a couple of derivatives of any kind can be taken at all. The various commutators that appear in the calculations have to be interpreted in a suitable way, so that $[A, B]$ exists reasonably even when the products AB and BA do not. We employ two methods around these obstacles: approximation by smooth structures and careful analysis in the non-smooth geometry.

We say that a function is in the class $C^{1,1}$ if it is continuously differentiable and the derivative is Lipschitz, and we define in definition 5 what a $C^{1,1}$ simple Riemannian metric is. Throughout the article our manifolds are assumed to be connected and to have dimension $n \geq 2$.

Theorem 1. *Let (M, g) be a simple $C^{1,1}$ manifold in the sense of definition 5.*

- (1) *If f is a Lipschitz function on M that integrates to zero over all maximal geodesics of M , then $f = 0$.*
- (2) *Let h be a Lipschitz 1-form on M that vanishes on the boundary ∂M . Then h integrates to zero over all maximal geodesics of M if and only if there is a scalar function $p \in C^{1,1}(M)$ vanishing on the boundary ∂M so that $h = dp$.*

We have to redefine simplicity to be tractable in our rough setup, and we regard this new definition as one of our main results. To verify that our redefinition is a valid one, we prove that it agrees with the classical definition when the metric is smooth. The classical definition of a smooth simple manifold implies the existence of global coordinates, but in the $C^{1,1}$ case we assume the coordinates in the definition — in light of the following theorem the coordinate assumption is not superfluous.

Theorem 2. *In smooth geometry definitions 4 and 5 are equivalent in the following sense:*

- (1) *If M is a simple C^∞ Riemannian manifold (see definition 4), then it is diffeomorphic to a closed ball in \mathbb{R}^n and it is a simple $C^{1,1}$ Riemannian manifold (see definition 5).*
- (2) *If M is a simple $C^{1,1}$ Riemannian manifold (see definition 5) and its metric tensor is C^∞ -smooth, then M is a smooth simple Riemannian manifold (see definition 4).*

Remark 3. The assumption $h|_{\partial M} = 0$ in claim (2) of theorem 1 is probably not necessary. Not assuming this is fine in smooth geometry but leads to technical

difficulties in our rough setup. This added assumption is the only way in which our results fail to correspond to the classical smooth results.

1.1. Related results. Geodesic X-ray transforms have been studied a lot on smooth manifolds equipped with C^∞ -smooth Riemannian metrics. Injectivity of the transform is reasonably well understood both on manifolds with a boundary and on closed manifolds. On manifolds with boundary one integrates over maximal geodesics between two boundary points, whereas on closed manifolds one integrates over periodic geodesics.

After Mukhometov's introduction of the Pestov identity for scalar tomography [Muk75, Muk77, Muk81], the method has been applied to 1-forms and higher order tensor fields [AR97, PS88, PSU13, PSU15] on many simple manifolds. When one passes from simple manifolds with boundary to closed Anosov manifolds, the Pestov identity remains the same but the other tools around it change somewhat [CS98, DS03, PSU14a, PSU15, US00]. Cartan–Hadamard manifolds are a non-compact analogue of simple manifolds, and the familiar Pestov identity works well [Leh16, LRS04]. Other variations of the problem change the Pestov identity, but a variant remains true and useful: In the presence of reflecting rays a boundary term on the reflector is added [IS16, IP18], an attenuation or a Higgs field [SU11, PSU12, GPSU16] and magnetic flows [DPSU07, Ain13, MP11] add a term to the geodesic vector field, non-abelian versions of the problem remove the concept of a line integral entirely [FU01, PS20, MNP21], and on Finsler surfaces a number of new terms are needed to account for non-Riemannian geometry [AD18]. On pseudo-Riemannian manifolds a Pestov identity useful for the light ray transform only seems to exist in product geometry of at least $2 + 2$ dimensions [Ilm16].

Pestov identities are not the only tool in the box for studying ray transforms on manifolds. For the variety of other methods we refer the reader to the review [IM19].

Inverse problems in integral geometry have been mostly studied on manifolds whose Riemannian metric is smooth or otherwise substantially above our $C^{1,1}$ in regularity. Injectivity of the scalar X-ray transform is known on spherically symmetric manifolds of regularity $C^{1,1}$ satisfying the so-called Herglotz condition when the conformal factor of the metric is in $C^{1,1}$ [dHIK22].

Some geometric inverse problems outside integral geometry have been solved in low regularity. A manifold with a metric tensor in a suitable Zygmund class is determined by its boundary spectral data [AKK⁺04], interior spectral data [BKL22] or by its boundary distance function [KKL07].

1.2. Preliminaries. In this subsection we will set up enough language to be able to state our definitions and give our proofs on a higher level. For a similar framework in the traditional smooth setting, see e.g. [PSU15]. We will cover the foundations in more detail in section 4 before embarking on the detailed proofs of our key lemmas.

The Riemannian manifold (M, g) , where g is $C^{1,1}$ regular, comes equipped with the unit sphere bundle $\pi: SM \rightarrow M$. The geodesic flow is a dynamical system on SM and its generator X is called the geodesic vector field. Properties and coordinate representations of X will be given later.

We will make frequent use of the bundle N over SM defined next. If π^*TM is the pullback of TM over SM , then N is the subbundle of π^*TM with fibers $N_{(x,v)} = \{v\}^\perp \subseteq T_xM$. It is well known (see [Pat99]) that the tangent bundle TSM

of SM has an orthogonal splitting

$$TSM = \mathbb{R}X \oplus \mathcal{H} \oplus \mathcal{V} \quad (1)$$

with respect to the so-called Sasaki metric, where \mathcal{H} and \mathcal{V} are called horizontal and vertical subbundles respectively. Roughly speaking, $\mathcal{H}_{(x,v)}$ corresponds to derivatives on SM in the base without components in the direction of v and $\mathcal{V}_{(x,v)}$ corresponds to derivatives on a fiber $S_x M$. It is natural to identify $\mathcal{H}_{(x,v)} = N_{(x,v)}$ and $\mathcal{V}_{(x,v)} = N_{(x,v)}$.

Given $z \in SM$, let γ_z be the unique geodesic corresponding to the initial condition z . We define the geodesic flow to be the collection of (partially defined) maps $\phi_t: SM \rightarrow SM$, $\phi_t(z) = (\gamma_z(t), \dot{\gamma}_z(t))$, where t goes through the values for which the right side is defined on SM . For any $z \in SM$ the geodesic γ_z is defined on a maximal interval $[\tau_-(z), \tau_+(z)]$. The travel time function $\tau: SM \rightarrow \mathbb{R}$ describes the first time a geodesic exists the manifold and it is defined by $\tau(z) = \tau_+(z)$ for $z \in SM$. Clearly $\gamma_z(\tau(z)) \in \partial M$ for any $z \in SM$.

A function f on M can be identified with the function $\pi^* f$ on SM . If h is a 1-form on M , then it can be considered as a function $\tilde{h}: SM \rightarrow \mathbb{R}$ through the identification $\tilde{h}(x, v) = h_x(v)$ for $(x, v) \in SM$. Since $h_x: T_x M \rightarrow \mathbb{R}$ is linear, \tilde{h} uniquely corresponds to h . The integral function $u^f: SM \rightarrow \mathbb{R}$ of $f \in \text{Lip}(SM)$ is defined by

$$u^f(x, v) := \int_0^{\tau(x, v)} f(\phi_t(x, v)) dt \quad (2)$$

for all $(x, v) \in SM$.

The lift of a unit speed curve $\gamma: I \rightarrow M$ is $\tilde{\gamma}: I \rightarrow SM$ given by $\tilde{\gamma}(t) = (\gamma(t), \dot{\gamma}(t))$. The curve γ is a geodesic if and only if the lift satisfies $\dot{\tilde{\gamma}}(t) = X(\tilde{\gamma}(t))$. The geodesic vector field X acts naturally on scalar fields by differentiation, and on sections V of N it acts by

$$XV(z) = D_t V(\phi_t(z))|_{t=0}, \quad (3)$$

where D_t is the covariant derivative along the curve $t \mapsto \gamma_z(t)$. This operator maps indeed sections of N to sections of N .

According to (1) the gradient of a C^1 function u on SM we can be written as

$$\nabla_{SM} u = (Xu)X + \overset{\text{h}}{\nabla} u + \overset{\text{v}}{\nabla} u. \quad (4)$$

This gives rise to two new differential operators $\overset{\text{v}}{\nabla}$ and $\overset{\text{h}}{\nabla}$, called, respectively, the vertical and the horizontal gradient. Both $\overset{\text{v}}{\nabla} u$ and $\overset{\text{h}}{\nabla} u$ are naturally interpreted as sections of N ; see [PSU23] for details. There are natural L^2 spaces for functions on the sphere bundle as well as for the sections of the bundle N . These will be denoted $L^2(SM)$ and $L^2(N)$ and we will often label the corresponding inner products explicitly. Formal adjoints of $\overset{\text{v}}{\nabla}$ and $\overset{\text{h}}{\nabla}$ with respect to appropriate L^2 inner products are the vertical and horizontal divergences $-\overset{\text{v}}{\text{div}}$ and $-\overset{\text{h}}{\text{div}}$ respectively. The mapping properties of the operators in $C^{1,1}$ regular metric setting are

$$X: C^1(SM) \rightarrow C(SM) \quad (5)$$

$$X: C^1(N) \rightarrow C(N), \quad (6)$$

$$\overset{\text{v}}{\nabla}, \overset{\text{h}}{\nabla}: C^1(SM) \rightarrow C(N), \quad \text{and} \quad (7)$$

$$\overset{v}{\text{div}}, \overset{h}{\text{div}}: C^1(N) \rightarrow C(SM). \quad (8)$$

These mapping properties are easily verified by inspecting the explicit formulas in local coordinates; see section 4.

We will deal with Sobolev spaces $H_{(0)}^1(SM)$ and $H_{(0)}^1(N)$ defined as completions of $C_{(0)}^1$ regular functions or sections in the relevant norms (see section 4), where the optional subscript 0 indicates zero boundary values. Similarly, we denote by $\text{Lip}_0(M)$ and $\text{Lip}_0(SM)$ the spaces of Lipschitz functions zero boundary values. As the last function space we introduce a Sobolev space $H_{(0)}^1(N, X)$, which only gives control over the operator X operating on sections of N . From definitions of various Sobolev norms it will be clear that all differential operators are bounded $H^1 \rightarrow L^2$ and thus extend to operators between Sobolev spaces.

Finally, there is a special quadratic form Q appearing in the Pestov identity. To define it, we use the Riemannian curvature tensor $R: L^\infty(N) \rightarrow L^\infty(N)$ acting on sections of N by

$$R(x, v)V(x, v) = R(V(x, v), v)v. \quad (9)$$

In order to verify the mapping property of R , observe that the second partial derivatives of $g \in C^{1,1} = W^{2,\infty}$ are in L^∞ . We define Q by letting

$$Q(W) = \|XW\|_{L^2(N)}^2 - (RW, W)_{L^2(N)}. \quad (10)$$

for all $W \in H^1(N, X)$.

To conclude the preliminaries we recall in definition 4 the traditional definition of a simple Riemannian manifold (cf. [PSU15]). In what follows a manifold satisfying conditions A1 and A2 is called *simple C^∞ manifold*. In definition 5 we redefine the notion of simplicity on manifolds equipped with non-smooth Riemannian metrics.

Definition 4 (Simple C^∞ manifold). Let (M, g) be a compact smooth Riemannian manifold with a smooth boundary. The manifold (M, g) is called *simple C^∞ Riemannian manifold*, if the following hold:

- A1: The boundary ∂M is strictly convex in the sense of the second fundamental form.
- A2: Any two points on M can be joined by a unique geodesic in the interior of M , and its length depends smoothly on its end points.

Definition 5 (Simple $C^{1,1}$ manifold). Let $M \subseteq \mathbb{R}^n$ be the closed unit ball and g a $C^{1,1}$ regular Riemannian metric on M . We say that (M, g) is a *simple $C^{1,1}$ Riemannian manifold* if the following hold:

- B1: There is $\varepsilon > 0$ so that $Q(W) \geq \varepsilon \|W\|_{L^2(N)}^2$ for all $W \in H_0^1(N, X)$.
- B2: Any two points of M can be joined by a unique geodesic in the interior of M , whose length depends continuously on its end points.
- B3: The function τ^2 is Lipschitz on SM .

Remark 6. In definition 5 the assumption that M is the closed unit ball is not restrictive — any simple C^∞ Riemannian manifold is diffeomorphic to a closed ball in a Euclidean space. In the absence of conjugate points the exponential map \exp_x , related to an interior point $x \in \text{int}(M)$, maps its maximal domain D_x diffeomorphically to M and D_x is itself diffeomorphic to the closed unit ball in \mathbb{R}^n (see [PSU23]). We use global coordinates on a simple $C^{1,1}$ Riemannian manifold and we have decided to include their existence in the definition.

Remark 7. If one is to define a rough simple manifold as the limit of smooth simple manifolds, the simplicity needs to be quantified. The example of a hemisphere as the limit of expanding polar caps shows that the smooth limit of smooth simple manifolds can be a smooth but non-simple manifold. The limit procedure can introduce conjugate points and failure of strict convexity on the boundary. An example of quantified simplicity can be found in [dHILS21], but we do not take this limit route in our definition here.

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2. PROOF OF THEOREM 1

This section contains the proof of theorem 1. The proofs of the necessary lemmas are postponed to section 5. More detailed definitions of function spaces and operators can be found from section 4.

We will freely identify a scalar function f and a one-form h on M with scalar functions on SM as described above. Interpreting f and h as functions on SM we can apply formula (2) to both.

Lemma 8 (Regularity of integral functions). *Let (M, g) be a simple $C^{1,1}$ manifold.*

- (1) *Let f be a Lipschitz function on M that integrates to zero over all maximal geodesics of M and let u^f be the integral function of f defined by (2). Then $u^f \in \text{Lip}_0(SM)$, $Xu^f \in H^1(SM)$ and $\overset{\vee}{\nabla} u^f \in H_0^1(N, X)$.*
- (2) *Let h be a Lipschitz 1-form on M that integrates to zero over all maximal geodesics of M and vanishes on the boundary ∂M . If u^h is the integral function of h defined by (2), then $u^h \in \text{Lip}_0(SM)$, $Xu^h \in H^1(SM)$ and $\overset{\vee}{\nabla} u^h \in H_0^1(N, X)$.*

Lemma 9 (Pestov identity). *Let (M, g) be a simple $C^{1,1}$ manifold and let $u \in \text{Lip}_0(SM)$ be such that $Xu \in H^1(SM)$ and $\overset{\vee}{\nabla} u \in H^1(N, X)$. Then*

$$\left\| \overset{\vee}{\nabla} Xu \right\|_{L^2(N)}^2 = Q \left(\overset{\vee}{\nabla} u \right) + (n-1) \|Xu\|_{L^2(SM)}^2. \quad (11)$$

Lemma 8 provides enough regularity to apply the Pestov identity (11) to the integral functions u^f and u^h because we will see in remark 15 that $\text{Lip}(SM) \subseteq H^1(SM)$ even if the metric tensor is only in $C^{1,1}$. The following lemma shows that certain norms of the integral function u^h of a 1-form cancel in the identity.

Lemma 10. *Let (M, g) be a simple $C^{1,1}$ manifold and let h be a Lipschitz 1-form on M . Then*

$$\left\| \overset{\vee}{\nabla} h \right\|_{L^2(N)}^2 = (n-1) \|h\|_{L^2(SM)}^2. \quad (12)$$

We are ready to prove theorem 1.

Proof of theorem 1. (1) The integral function u^f of $f \in \text{Lip}(M)$ satisfies $Xu^f \in H^1(SM)$ and $\overset{\vee}{\nabla} u^f \in H^1(N, X)$ by lemma 8. Thus we can apply the Pestov identity of lemma 9 to u^f . By the fundamental theorem of calculus $Xu^f = -f$ and

thus $\overset{v}{\nabla}Xu^f = 0$, since f does not depend on the direction $v \in S_xM$. By $C^{1,1}$ simplicity (definition 5) of (M, g) , the quadratic form Q is non-negative. Thus the Pestov identity reduces to

$$0 \geq (n-1) \|Xu^f\|_{L^2(SM)}^2. \quad (13)$$

Hence $f = -Xu^f = 0$ in $L^2(SM)$ as claimed.

(2) If $h = dp$ for some scalar function $p \in C^{1,1}(M)$ with $p|_{\partial M} = 0$, then by the fundamental theorem of calculus h integrates to zero over all maximal geodesics of M .

Let h be a Lipschitz 1-form on M that integrates to zero over all maximal geodesic of M and vanishes on the boundary ∂M . We will show that $h = dp$ for some function $p \in C^{1,1}(M)$ vanishing on ∂M . Lemma 8 allows us to apply the Pestov identity to the integral function u^h of h . Due to lemma 10, the identity reduces to

$$Q\left(\overset{v}{\nabla}u^h\right) = 0. \quad (14)$$

Since the manifold is simple $C^{1,1}$, this can only happen if $\overset{v}{\nabla}u^h = 0$. The function u^h is Lipschitz and independent of the direction $v \in S_xM$ on each fiber and therefore there is a Lipschitz scalar function p on M so that $u^h = -\pi^*p$ on SM . Additionally, $p|_{\partial M} = u^h|_{\partial(SM)} = 0$, since h integrates to zero over all maximal geodesics of M . Since $Xu^h = -h$, we have shown that $dp = h$ in the weak sense. Because h is Lipschitz-continuous by assumption, we have that dp is Lipschitz and thus $p \in C^{1,1}$ and the proof is complete. \square

3. PROOF OF THEOREM 2

In this section we prove that in the smooth setting definition 5 of $C^{1,1}$ simplicity is equivalent to definition 4 of C^∞ simplicity. Proofs of lemmas 11 and 12 are given in section 6. Theorem 2 readily follows from lemmas 11 and 12.

Lemma 11. *Let (M, g) be a simple $C^{1,1}$ manifold with C^∞ -smooth Riemannian metric g . Then there are no conjugate points in M , not even on the boundary.*

Lemma 12. *Let M be a compact Riemannian manifold with smooth boundary and a C^∞ -smooth Riemannian metric g . Suppose that (M, g) is non-trapping. Then ∂M is strictly convex in the sense of the second fundamental form if and only if $\tau^2 \in \text{Lip}(SM)$.*

Proof of theorem 2. By remark 6 each simple C^∞ Riemannian manifold is diffeomorphic to the closed unit ball \bar{B} in \mathbb{R}^n . Thus we may assume that $M = \bar{B}$ and let g be a C^∞ -smooth Riemannian metric on M . It suffices to show that (M, g) satisfies conditions A1–A2 in definition 4 if and only if it satisfies conditions B1–B3 in definition 5. We have illustrated these implications in figure 1.

By lemma 12 conditions A1 and B3 are equivalent. By lemma 11 the condition B1 implies that there are no conjugate points on M . Thus we can promote the continuous dependence in B2 to smooth dependence A2. Therefore simple $C^{1,1}$ manifolds satisfy both conditions A1 and A2 of C^∞ simplicity. Conversely, simple C^∞ manifolds satisfy B1 (see [PSU15, Lemma 11.2]) and clearly B2 is strictly weaker than A2. \square

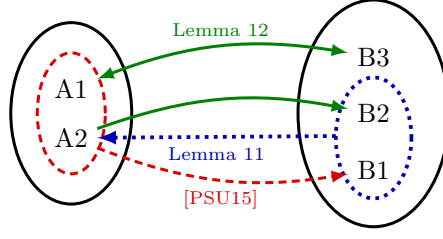


FIGURE 1. Illustration of the proof of theorem 2. The arrows represent implications except the one double headed arrow, which represents equivalence. The green (solid) arrows connect one condition to another. The red (dashed) and the blue (dotted) arrows indicate that one condition follows from the two conditions circled with the same color (style).

4. BUNDLES, FUNCTION SPACES AND OPERATORS

This section complements the preliminaries in subsection 1.2. The main focus is on a detailed description of structures, functions spaces and operators build on a compact Riemannian manifold (M, g) with a $C^{1,1}$ regular Riemannian metric.

4.1. Function spaces on smooth manifolds. Let M be a compact smooth manifold with a smooth boundary. The space of smooth functions on M is denoted $C^\infty(M)$ and the space of differentiable functions with Lipschitz derivatives is denoted $C^{1,1}(M)$. We let $C^{1,1}(T^2M)$ denote the space of 2-tensor fields on M , whose component functions are in $C^{1,1}(M)$.

If h is a smooth Riemannian metric on M , then $L_h^2(M)$ and $L_h^\infty(M)$ will respectively denote spaces of square integrable and essentially bounded functions on M , where the Riemannian volume form of h is used as the measure. Similarly, $W_h^{1,p}(M)$ and $W_h^{2,p}(M)$ will respectively denote Sobolev spaces with p -integrable covariant derivatives of the first order and of the second order. Norms of the covariant derivatives on the tangent spaces are always defined by the metric h .

4.2. Structures in low regularity. Let (M, g) be a compact Riemannian manifold with a smooth boundary. We assume that $g \in C^{1,1}(T^2M)$. The unit sphere bundle $SM = \{v \in TM : |v| = 1\}$ is a submanifold of TM , but not in general a smooth one. Despite the non-smoothness of $SM \subseteq TM$ as a submanifold, it can be equipped with an induced smooth structure: SM is naturally homeomorphic to the quotient space $(TM \setminus 0)/\sim$, where $v \sim \lambda v$ for all $\lambda > 0$ and $v \in T_x M$. Metric structures like the Sasaki metric are still non-smooth, so this smooth structure is of little use. We will only see SM as a submanifold of TM .

For $k \in \{0, 1\}$ a function $u: SM \rightarrow \mathbb{R}$ is said to be in $C^k(SM)$ if u is k times continuously differentiable — for $k \geq 2$ this concept is undefined in our setting. As a C^1 submanifold of TM the sphere bundle has enough regularity to define both $C(SM)$ and $C^1(SM)$. The subset $C_0^k(SM)$ of $C^k(SM)$ consists of functions vanishing on

$$\partial(SM) = \{(x, v) \in SM : x \in \partial M\}. \quad (15)$$

The set of Lipschitz functions on SM is denoted by $\text{Lip}(SM)$. We denote the inward unit normal vector field to the boundary ∂M by ν . The boundary $\partial(SM)$

is divided into parts pointing inwards and outwards, respectively denoted by

$$\partial_{\text{in}}(SM) := \{ (x, v) \in \partial(SM) : \langle v, \nu(x) \rangle \geq 0 \} \quad (16)$$

and

$$\partial_{\text{out}}(SM) := \{ (x, v) \in \partial(SM) : \langle v, \nu(x) \rangle \leq 0 \}. \quad (17)$$

Their intersection consists of tangential directions

$$\partial_0(SM) := \partial_{\text{in}}(SM) \cap \partial_{\text{out}}(SM). \quad (18)$$

Many differential operators considered in this article operate on sections of the bundle N . To describe C^k spaces of sections of N , recall that N is the subbundle of π^*TM with fibers $N_{(x,v)} = \{v\}^\perp \subseteq T_xM$. A section V of the bundle N is a section of the bundle π^*TM with the property that $\langle V(x, v), v \rangle_{g(x)} = 0$ for all $(x, v) \in SM$. We say that such a section is in $C^k(N)$ for $k \in \{0, 1\}$ if the corresponding section of π^*TM is k times continuously differentiable. Differentiability of a section W of π^*TM is well defined since W is a certain function between two differentiable manifolds SM and TM . The subspace $C_0^k(N) \subseteq C^k(N)$ consists of sections V of N that vanish on $\partial(SM)$.

Let (x, v) be a local coordinate system on TM and let ∂_{x^j} and ∂_{v^k} be corresponding coordinate vector fields. We introduce new vector fields $\delta_{x^j} = \partial_{x^j} - \Gamma_{jk}^l v^k \partial_{v^l}$ on TM , where Γ_{jk}^l are the Christoffel symbols of the metric g . As the metric tensor in our results is of regularity $C^{1,1}$, it follows that the Christoffel symbols and thus the vector fields δ_{x^j} are only Lipschitz.

4.3. Differential operators. Next we define differential operators on SM and N . The basic coordinate derivatives of a function $u \in C^1(SM)$ are defined by

$$\delta_j u := \delta_{x^j}(u \circ r)|_{SM} \quad \text{and} \quad \partial_k u := \partial_{v^k}(u \circ r)|_{SM}, \quad (19)$$

where $r: TM \setminus 0 \rightarrow SM$ is the radial function $r(x, v) = (x, v|v|_{g(x)}^{-1})$. We denote $\delta^j := g^{jk} \delta_k$ and $\partial^j := g^{jk} \partial_k$. We use the basic derivatives to define operators in local coordinates.

The geodesic vector field X is a differential operator that acts both on functions on SM and on sections of the bundle N . The actions on a scalar function u and on a section V are defined by

$$Xu = v^j \delta_j u \quad \text{and} \quad XV = (XV^j) \partial_{x^j} + \Gamma_{jk}^l v^j V^k \partial_{x^l}. \quad (20)$$

Vertical and horizontal gradients are differential operators defined respectively by

$$\overset{v}{\nabla} u = (\partial^j u) \partial_{x^j} \quad \text{and} \quad \overset{h}{\nabla} u = (\delta^j u - (Xu) v^j) \partial_{x^j}. \quad (21)$$

Coordinate formulas indicate that $\overset{v}{\nabla}$ is the gradient in v and $\overset{h}{\nabla}$ is the gradient in x with the direction of v being projected out. The adjoint operators of $\overset{v}{\nabla}$ and $\overset{h}{\nabla}$ are the vertical and the horizontal divergences

$$\overset{v}{\text{div}} V = \partial_j V^j \quad \text{and} \quad \overset{h}{\text{div}} V = (\delta_j + \Gamma_{ji}^i) V^j. \quad (22)$$

The Riemannian curvature tensor R of the metric g has an action on sections of N defined by

$$RV = R^l_{ijk} V^i v^j v^k \partial_{x^l}. \quad (23)$$

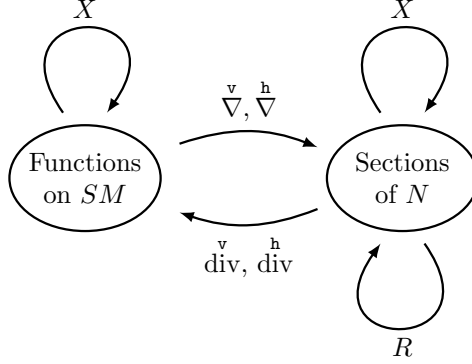


FIGURE 2. Interplay of the operators defined in subsection 4.3. The gradients map functions on SM to sections of N . The divergences map sections of N back to function on SM . The geodesic vector field maps functions to functions and sections to sections. The curvature operator acts only on sections and produces sections.

4.4. Integration and Sobolev spaces. A simple $C^{1,1}$ manifold M is orientable, so the Riemannian volume form on it can be defined in local coordinates as

$$dV_g(x) := |\det(g(x))|^{1/2} dx^1 \wedge \cdots \wedge dx^n. \quad (24)$$

For any $x \in M$ the pair $(S_x M, g(x))$ is a Riemannian manifold. Let dS_x be the associated Riemannian volume form on $S_x M$. We use dV_g and dS_x to define the volume form $d\Sigma_g$ on SM , given in local coordinates by

$$d\Sigma_g(x, v) = dS_x(v) \wedge dV_g(x). \quad (25)$$

The form $d\Sigma_g$ is natural as it coincides with the Riemannian volume form of the Sasaki metric on SM . Since dV_g has as much regularity as g , so does $d\Sigma_g$.

The L^2 -norm of a scalar function u on SM is denoted by $\|u\|_{L^2(SM)}$ and the L^2 -norm of a section V of N is denoted by $\|V\|_{L^2(N)}$. The L^2 -norms are induced by the inner products

$$(u, w)_{L^2(SM)} := \int_{SM} uw \, d\Sigma_g \quad (26)$$

and

$$(V, W)_{L^2(N)} := \int_{SM} g_{ij} V^i W^j \, d\Sigma_g. \quad (27)$$

We define the $\|\cdot\|_{H^1(SM)}$ -norm of a function $u \in C^1(SM)$ by

$$\|u\|_{H^1(SM)}^2 = \|u\|_{L^2(SM)}^2 + \|Xu\|_{L^2(SM)}^2 + \left\| \overset{v}{\nabla} u \right\|_{L^2(N)}^2 + \left\| \overset{h}{\nabla} u \right\|_{L^2(N)}^2. \quad (28)$$

The Sobolev space $H^1(SM)$ is defined to be the completion of the subset of $C^1(SM)$ that consists of functions with finite $H^1(SM)$ -norm. We denote by $H_0^1(SM)$ the closure of $C_0^1(SM)$ in $H^1(SM)$.

Sobolev spaces for sections of N are defined in an analogous fashion. For a section $V \in C^1(N)$ we define the two Sobolev norms

$$\|V\|_{H^1(N)}^2 = \|V\|_{L^2(N)}^2 + \|XV\|_{L^2(N)}^2 + \left\| \overset{v}{\operatorname{div}} V \right\|_{L^2(SM)}^2 + \left\| \overset{h}{\operatorname{div}} V \right\|_{L^2(SM)}^2 \quad (29)$$

and

$$\|V\|_{H^1(N,X)}^2 = \|V\|_{L^2(N)}^2 + \|XV\|_{L^2(N)}^2. \quad (30)$$

The corresponding Sobolev spaces (the completions of $C^1(N)$ under these norms) are denoted by $H^1(N)$ and $H^1(N, X)$, and the Sobolev spaces of sections vanishing on the boundary $\partial(SM)$ are denoted by $H_0^1(N)$ and $H_0^1(N, X)$.

Remark 13. Contrary to what one might expect, the norm on $H^1(N)$ defined above does not contain derivatives in all possible directions, as it only includes divergences in the vertical and horizontal directions. We will use these norms only to estimate from above, so this omission of derivatives makes no difference.

In the case where g is a C^∞ -smooth Riemannian metric, we introduce one more Sobolev space, $K^2(SM)$. The defining norm on the dense subspace $C^2(N)$ is

$$\|u\|_{K^2(SM)}^2 = \|u\|_{H^1(SM)}^2 + \|Xu\|_{H^1(SM)}^2 + \left\| \overset{v}{\nabla} u \right\|_{H^1(N,X)}^2. \quad (31)$$

Remark 14. It is important to realize that we cannot define Sobolev spaces using smooth test functions as in the smooth case. The reason is two-fold. First, the natural structure of SM as a submanifold of TM is not regular enough to define the function class $C^\infty(SM)$. Second, the differential operators themselves are not smooth. Applying any of the differential operators immediately drops regularity to that of the coefficients, and they involve the metric tensor.

4.5. Differential operators on Sobolev spaces. It is clear from the definitions that all of our differential operators are bounded $H^1 \rightarrow L^2$. Thus all classically defined operators extend to operators between Sobolev spaces. We therefore have the continuous operators

$$X: H^1(SM) \rightarrow L^2(SM), \quad (32)$$

$$X: H^1(N) \rightarrow L^2(N), \quad (33)$$

$$\overset{v}{\nabla}, \overset{h}{\nabla}: H^1(SM) \rightarrow L^2(N), \quad \text{and} \quad (34)$$

$$\overset{v}{\operatorname{div}}, \overset{h}{\operatorname{div}}: H^1(N) \rightarrow L^2(SM). \quad (35)$$

Basic integration by parts holds for the extended operators: If $u, w \in H^1(SM)$ and $V, W \in H^1(N)$ and w and W vanish on the boundary, then

$$(Xu, w)_{L^2(SM)} = -(u, Xw)_{L^2(SM)}, \quad (36)$$

$$(XV, W)_{L^2(N)} = -(V, XW)_{L^2(N)}, \quad (37)$$

$$\left(\overset{v}{\nabla} u, W \right)_{L^2(N)} = - \left(u, \overset{v}{\operatorname{div}} W \right)_{L^2(SM)}, \quad \text{and} \quad (38)$$

$$\left(\overset{h}{\nabla} u, W \right)_{L^2(N)} = - \left(u, \overset{h}{\operatorname{div}} W \right)_{L^2(SM)}. \quad (39)$$

We can use the space $C_0^1(SM)$ as test functions and $C_0^1(N)$ as test sections.

4.6. Switching between different unit sphere bundles. Suppose we have two Riemannian metrics $g, h \in C^{1,1}(T^2M)$ on the manifold M . Let S_gM and S_hM denote the corresponding unit sphere bundles. There is a natural radial $C^{1,1}$ -diffeomorphism

$$s: S_gM \rightarrow S_hM, \quad s(x, v) = (x, v|v|_h^{-1}). \quad (40)$$

In section 5 we will have three Riemannian metrics $g \in C^{1,1}(T^2M)$ and $\tilde{g}, h \in C^\infty(T^2M)$ with certain roles. In this case we will denote the corresponding radial $C^{1,1}$ -diffeomorphisms by

$$\tilde{s}: S_{\tilde{g}}M \rightarrow S_hM, \quad \tilde{r}: S_gM \rightarrow S_{\tilde{g}}M \quad \text{and} \quad s: S_gM \rightarrow S_hM. \quad (41)$$

In section 5 we frequently use the convention that the bundles related to \tilde{g} are denoted $\tilde{S}M := S_{\tilde{g}}M$ and $\tilde{N} := N_{\tilde{g}}$, the operators related to \tilde{g} are decorated with α on top or as a subscript, the bundles and the operators related to h are decorated with subscripts h , and the bundles and the operators related to the metric g are written without decorations.

Remark 15. We can switch between sphere bundles and corresponding Sobolev spaces using pullbacks along radial functions. If u is a scalar function on SM , then $(s^{-1})^*u$ is a scalar function on S_hM . To see that pullback behaves well on the Sobolev scale, note that the $H^1(SM)$ -norm controls all possible derivatives on SM since $TSM = \mathbb{R}X \oplus \mathcal{V} \oplus \mathcal{H}$. Thus the $H^1(SM)$ -norm is equivalent to

$$\|u\| = \|u\|_{L^2(SM)} + \|d_{SM}u\|_{L^2(T^*SM)} \quad (42)$$

with the norm of the differential interpreted with respect to any Riemannian metric on SM . With the norm (42) we see that regularity on Sobolev scale is preserved, since $(s^{-1})^*(d_{SM}u) = d_{S_hM}(u \circ s^{-1})$ by standard properties of the pullback.

Remark 15 allows us to prove continuous Sobolev embeddings between Sobolev spaces of low regularity metrics. We present one example that will be useful to us later. Let $g \in C^{1,1}(T^2M)$ and $h \in C^\infty(T^2M)$ be two Riemannian metrics on M . If $u \in \text{Lip}(SM)$, then $(s^{-1})^*u \in \text{Lip}(S_hM)$. Since h is C^∞ -smooth, we have $(s^{-1})^*u \in H^1(S_hM)$. Then since $\|u\|_{H^1(SM)} \lesssim \|(s^{-1})^*u\|_{H^1(S_hM)}$ by remark 15, we see that $u \in H^1(SM)$. We have shown that the inclusion $\text{Lip}(SM) \subseteq H^1(SM)$ holds even when the metric tensor is only $C^{1,1}$.

5. LEMMAS IN LOW REGULARITY

5.1. The Pestov identity. In this subsection (M, g) is a simple $C^{1,1}$ Riemannian manifold. We prove a variant of the commutator formula $[X, \tilde{\nabla}] = -\tilde{\nabla}$ and the Pestov identity on (M, g) . First, we show that both results are valid for Sobolev functions on a manifold equipped with a C^∞ -smooth Riemannian metric. Then we show that only $C^{1,1}$ regularity of the Riemannian metric is needed. The main focus of the subsection is on proving the Pestov identity of lemma 9.

Lemma 16. *Let (M, h) be a compact smooth manifold with a smooth boundary, where h is a C^∞ -smooth Riemannian metric. The commutator formula $[X, \tilde{\nabla}] = -\tilde{\nabla}$ holds in the H^1 sense on (M, h) : For $u \in H_0^1(S_hM)$ and $V \in C^1(N_h)$, we have*

$$\left(\tilde{\nabla}_h u, V \right)_{L^2(N_h)} = \left(\tilde{\nabla}_h u, X_h V \right)_{L^2(N_h)} - \left(X_h u, \text{div}_h V \right)_{L^2(S_hM)}. \quad (43)$$

Proof. Let $u \in H_0^1(S_h M)$ and $V \in C^\infty(N_h)$. Since V is smooth, by [PSU15, Lemma 2.1.] we have

$$X_h \operatorname{div}_h^v V - \operatorname{div}_h^v X_h V = -\operatorname{div}_h^h V. \quad (44)$$

Thus

$$\begin{aligned} \left(\operatorname{div}_h^h u, V \right)_{L^2(N_h)} &= - \left(u, \operatorname{div}_h^v X_h V \right)_{L^2(S_h M)} + \left(u, X_h \operatorname{div}_h^v V \right)_{L^2(S_h M)} \\ &= \left(\operatorname{div}_h^v u, X_h V \right)_{L^2(N_h)} - \left(X_h u, \operatorname{div}_h^v V \right)_{L^2(S_h M)}, \end{aligned} \quad (45)$$

where the last equality holds since $u \in H_0^1(S_h M)$, and since $X_h V \in C^\infty(N_h)$ and $\operatorname{div}_h^v V \in C^\infty(S_h M)$. The same identity holds for $V \in C^1(N_h)$ by approximation, since only first order derivatives appear in the statement. \square

Lemma 17. *Let (M, h) be a compact smooth manifold with a smooth boundary, where h is a C^∞ -smooth Riemannian metric. Suppose that $u \in K^2(S_h M)$ vanishes on the boundary $\partial(S_h M)$. Then*

$$\left\| \operatorname{div}_h^v X_h u \right\|_{L^2(N_h)}^2 = Q_h \left(\operatorname{div}_h^v u \right) + (n-1) \|X_h u\|_{L^2(S_h M)}^2, \quad (46)$$

where Q_h is the quadratic form defined for $W \in H_0^1(N_h, X_h)$ by

$$Q_h(W) = \|X_h W\|_{L^2(N_h)}^2 - (R_h W, W)_{L^2(N_h)}. \quad (47)$$

Proof. Since $u \in K^2(S_h M)$ and u vanishes on the boundary $\partial(S_h M)$, there is a sequence $(u^\beta)_{\beta \in \mathbb{N}}$ of smooth functions on $S_h M$ vanishing on $\partial(S_h M)$ so that $u^\beta \rightarrow u$ in $K^2(S_h M)$. We see that

$$\left\| \operatorname{div}_h^v X_h u^\beta - \operatorname{div}_h^v X_h u \right\|_{L^2(N_h)}^2 \leq \|X_h u^\beta - X_h u\|_{H^1(S_h M)}^2 \leq \|u^\beta - u\|_{K^2(S_h M)}^2 \quad (48)$$

and

$$\|X_h u^\beta - X_h u\|_{L^2(N_h)}^2 \leq \|u^\beta - u\|_{H^1(S_h M)}^2 \leq \|u^\beta - u\|_{K^2(S_h M)}^2. \quad (49)$$

Therefore $\operatorname{div}_h^v X_h u^\beta \rightarrow \operatorname{div}_h^v X_h u$ in $L^2(N_h)$ and $X_h u^\beta \rightarrow X_h u$ in $L^2(S_h M)$ as $\beta \rightarrow \infty$. Additionally, since the curvature operator R of the metric h continuously maps $L^\infty(N_h) \rightarrow L^\infty(N_h)$, we have

$$\left\| \operatorname{div}_h^v u^\beta - \operatorname{div}_h^v u \right\|_{L^2(N_h)}^2 \leq \|u^\beta - u\|_{H^1(S_h M)}^2 \leq \|u^\beta - u\|_{K^2(S_h M)}^2 \quad (50)$$

and

$$\left\| R_h \operatorname{div}_h^v u^\beta - R_h \operatorname{div}_h^v u \right\|_{L^2(N_h)}^2 \lesssim \|u^\beta - u\|_{H^1(S_h M)}^2 \leq \|u^\beta - u\|_{K^2(S_h M)}^2. \quad (51)$$

Thus $Q_h(\operatorname{div}_h^v u^\beta) \rightarrow Q_h(\operatorname{div}_h^v u)$ as $\beta \rightarrow \infty$. By the Pestov identity for smooth functions and metrics (see [PSU15, Remark 2.3.]) we have

$$\left\| \operatorname{div}_h^v X_h u^\beta \right\|_{L^2(N_h)}^2 = Q_h \left(\operatorname{div}_h^v u^\beta \right) + (n-1) \|X_h u^\beta\|_{L^2(S_h M)}^2. \quad (52)$$

We now let $\beta \rightarrow \infty$ in (52). By our estimates (48), (49), (50), and (51) we end up with the claimed identity (46). \square

The rest of this section focuses on showing that we can replace the C^∞ -smooth Riemannian metric h in lemmas 16 and 17 by a $C^{1,1}$ regular Riemannian metric. Let (M, g) be a $C^{1,1}$ simple Riemannian manifold. Next, we construct approximations of g by C^∞ -smooth Riemannian metrics \tilde{g} .

Let (x^1, \dots, x^n) be the usual Cartesian coordinates on the Euclidean closed ball $M \subset \mathbb{R}^n$ and extend all components $g_{ij} \in C^{1,1}(M)$ of g to functions $\bar{g}_{ij} \in C^{1,1}(\mathbb{R}^n)$. Such extensions exist since $C^{1,1} = W^{2,\infty}$ and the boundary of M is smooth (see [Ste70, Chapter 6, Theorem 5]). Let us then choose a non-negative compactly supported smooth function $\varphi: \mathbb{R}^n \rightarrow \mathbb{R}$ with unit integral and define a sequence of standard mollifiers $\tilde{\varphi}(x) = \alpha^n \varphi(\alpha x)$ for $\alpha \in \mathbb{N}$. Then we define

$$\tilde{g}_{ij} := (\tilde{\varphi} * \bar{g}_{ij})|_M \in C^\infty(M). \quad (53)$$

Lemma 18. *Let (M, g) be a simple $C^{1,1}$ manifold. Let h be a smooth reference metric on M . There exists a sequence $(\tilde{g})_{\alpha \in \mathbb{N}}$ of C^∞ -smooth metrics on M such that*

- (1) $\tilde{g}_{ij} \rightarrow g_{ij}$ in $W_h^{2,2}(M)$ and in $W_h^{1,\infty}(M)$,
- (2) $\tilde{g}^{ij} \rightarrow g^{ij}$ in $W_h^{1,2}(M)$ and in $L_h^\infty(M)$,
- (3) $\tilde{\Gamma}_{jk}^i \rightarrow \Gamma_{jk}^i$ in $W_h^{1,1}(M)$ and in $L_h^\infty(M)$,
- (4) $\tilde{R}_{jkl}^i \rightarrow R_{jkl}^i$ in $L_h^1(M)$.

Proof. For each $\alpha \in \mathbb{N}$ let $\tilde{g} := \tilde{g}_{ij} dx^i \otimes dx^j \in C^\infty(T^2M)$, where \tilde{g}_{ij} are as in (53). We will show that a subsequence of the sequence $(\tilde{g})_{\alpha \in \mathbb{N}}$ consists of smooth Riemannian metrics and satisfies conditions (1)–(4).

We see that for large α each \tilde{g} is a C^∞ -smooth Riemannian metric. Smoothness follows standard properties of the mollifiers $\tilde{\varphi}$. Each \tilde{g} is symmetric by construction. For large α each \tilde{g} is positive definite since this is an open condition and pointwise convergence $\tilde{g}_{ij} \rightarrow g_{ij}$ follows from continuity and item (1).

(1) Because \bar{g}_{ij} is compactly supported and in both spaces $W^{2,2}(\mathbb{R}^n)$ and $W^{2,\infty}(\mathbb{R}^n)$, the convolution converges $\tilde{\varphi} * \bar{g}_{ij} \rightarrow \bar{g}_{ij}$ in both spaces $W^{2,2}(\mathbb{R}^n)$ and $W^{1,\infty}(\mathbb{R}^n)$. This implies convergence in the corresponding function spaces over the subdomain $M \subset \mathbb{R}^n$.

(2) Let us denote the adjugate of a matrix A by $\text{adj}(A)$; we interpret rank two tensor fields as matrix-valued functions on $M \subset \mathbb{R}^n$. By item (1) we have

$$\det(\tilde{g}) \rightarrow \det(g) \quad \text{and} \quad \text{adj}(\tilde{g})^{ij} \rightarrow \text{adj}(g)^{ij} \quad (54)$$

in $L_h^\infty(M)$. Thus for sufficiently large α the matrices \tilde{g} are uniformly invertible in the sense that

$$\|\det(\tilde{g})^{-1}\|_{L_h^\infty(M)} \leq C. \quad (55)$$

Since

$$\tilde{g}^{ij}(x) = \det(\tilde{g}(x))^{-1} \text{adj}(\tilde{g}(x))^{ij}, \quad (56)$$

we have that $\tilde{g}^{ij} \rightarrow g^{ij}$ in $L_h^\infty(M)$. Derivatives of the inverse satisfy $\partial_k \tilde{g}^{ij} = -\tilde{g}^{il} (\partial_k \tilde{g}_{lm}) \tilde{g}^{mj}$, which implies convergence of the derivatives in $L_h^2(M)$.

(3) This follows from

$$\tilde{\Gamma}_{jk}^i = \frac{1}{2} \tilde{g}^{il} (\partial_j \tilde{g}_{kl} + \partial_k \tilde{g}_{jl} - \partial_l \tilde{g}_{jk}) \quad (57)$$

and items (1) and (2).

(4) This follows from

$$\mathring{R}_{ijk}{}^l = \partial_i \mathring{\Gamma}^l{}_{jk} - \partial_j \mathring{\Gamma}^m{}_{jk} + \mathring{\Gamma}^m{}_{jk} \mathring{\Gamma}^l{}_{im} - \mathring{\Gamma}^m{}_{ik} \mathring{\Gamma}^l{}_{jm} \quad (58)$$

and item (3). \square

Next we prove the Pestov identity for $C^{1,1}$ regular metrics. In the context of lemma 9 the manifold (M, g) is simple $C^{1,1}$, the Riemannian metric g is $C^{1,1}$ regular, the function u is in $\text{Lip}_0(SM)$ and satisfies $Xu \in H^1(SM)$ and $\mathring{\nabla} u \in H^1(N, X)$.

Proof of lemma 9. Choose a smooth reference Riemannian metric h on M and let $(\mathring{g})_{\alpha \in \mathbb{N}}$ be a sequence of smooth metrics approximating g as in lemma 18. For each $\alpha \in \mathbb{N}$ denote $\mathring{u} := u \circ \mathring{r}$. Then by remark 15 we have $\mathring{u} \in H_0^1(\mathring{S}M)$, $\mathring{X}\mathring{u} \in H^1(\mathring{S}M)$ and $\mathring{\nabla}_\alpha \mathring{u} \in H^1(\mathring{N}, \mathring{X})$, which implies that $\mathring{u} \in K^2(\mathring{S}M)$ and $\mathring{u}|_{\partial(\mathring{S}M)} = 0$. For each α an application of lemma 17 gives

$$\left\| \mathring{\nabla}_\alpha \mathring{X}\mathring{u} \right\|_{L^2(\mathring{N})}^2 = Q_\alpha \left(\mathring{\nabla}_\alpha \mathring{u} \right) + (n-1) \left\| \mathring{X}\mathring{u} \right\|_{L^2(\mathring{S}M)}^2. \quad (59)$$

We will show that

$$\lim_{\alpha \rightarrow \infty} \left\| \mathring{\nabla}_\alpha \mathring{X}\mathring{u} \right\|_{L^2(\mathring{N})}^2 = \left\| \mathring{\nabla} Xu \right\|_{L^2(N)}^2. \quad (60)$$

Similar arguments can be used to deduce that

$$\lim_{\alpha \rightarrow \infty} Q_\alpha \left(\mathring{\nabla}_\alpha \mathring{u} \right) = Q \left(\mathring{\nabla} u \right) \quad (61)$$

and

$$\lim_{\alpha \rightarrow \infty} \left\| \mathring{X}\mathring{u} \right\|_{L^2(\mathring{S}M)}^2 = \|Xu\|_{L^2(SM)}^2. \quad (62)$$

Then letting $\alpha \rightarrow \infty$ in equation (59) proves the claim of the theorem. Since the arguments showing equations (61) and (62) are analogous to what is presented below, we omit them. (The fact that components of the curvature tensor only converge in L^1 and not in L^∞ is where the assumption $u \in \text{Lip}(SM)$ is useful.) Coordinate formulas required to show equations (61) and (62) are given in appendix A.

For any L^p convergence to make sense, we fix $S_h M$ to be our common reference bundle for objects to be integrated on. First we study how the $L^2(\mathring{N})$ norm on the left-hand side of (60) transforms under \mathring{s} . Let $\tilde{u} := \mathring{u} \circ \mathring{s}^{-1}$ and fix $\hat{z} \in S_h M$ and $z \in \mathring{S}M$ such that $\mathring{s}(z) = \hat{z}$. By basic properties of pushforwards we have

$$\left(\left((\mathring{s}_* \mathring{X}) \tilde{u} \right) \circ \mathring{s} \right) (z) = (\mathring{s}_* \mathring{X})_{\hat{z}} \tilde{u} = \mathring{X}_z (\tilde{u} \circ \mathring{s}) = \mathring{X}_z \mathring{u}. \quad (63)$$

Thus

$$(\mathring{s}_* \mathring{\partial}^j)_z \left((\mathring{s}_* \mathring{X}) \tilde{u} \right) = \mathring{\partial}_z^j \left((\mathring{s}_* \mathring{X}) \tilde{u} \circ \mathring{s} \right) = \mathring{\partial}_z^j \left(\mathring{X} \mathring{u} \right). \quad (64)$$

Since $\pi(z) = \pi(\hat{z}) \in M$ we get

$$\begin{aligned}
\left\| \overset{v}{\nabla}_\alpha \overset{\alpha}{X} \tilde{u} \right\|_{L^2(\overset{\alpha}{N})}^2 &= \int_{z \in \overset{\alpha}{S}M} \overset{\alpha}{g}_{ij}(\pi(z)) \left(\overset{\alpha}{\partial}_z^i (\overset{\alpha}{X} \tilde{u}) \right) \left(\overset{\alpha}{\partial}_z^j (\overset{\alpha}{X} \tilde{u}) \right) d\overset{\alpha}{\Sigma}(z) \\
&= \int_{\hat{z} \in S_h M} \overset{\alpha}{g}_{ij}(\pi(\hat{z})) \left(\overset{\alpha}{\partial}_{\overset{\alpha}{s}^{-1}(\hat{z})}^i (\overset{\alpha}{X} \tilde{u}) \right) \\
&\quad \times \left(\overset{\alpha}{\partial}_{\overset{\alpha}{s}^{-1}(\hat{z})}^j (\overset{\alpha}{X} \tilde{u}) \right) |\det(d\overset{\alpha}{s}_{\hat{z}}^{-1})| d\Sigma_h(\hat{z}) \\
&= \int_{\hat{z} \in S_h M} \overset{\alpha}{g}_{ij}(\pi(\hat{z})) \left((\overset{\alpha}{s}_* \overset{\alpha}{\partial}^i)_{\hat{z}} \left((\overset{\alpha}{s}_* \overset{\alpha}{X}) \tilde{u} \right) \right) \\
&\quad \times \left((\overset{\alpha}{s}_* \overset{\alpha}{\partial}^j)_{\hat{z}} \left((\overset{\alpha}{s}_* \overset{\alpha}{X}) \tilde{u} \right) \right) |\det(d\overset{\alpha}{s}_{\hat{z}}^{-1})| d\Sigma_h(\hat{z}).
\end{aligned} \tag{65}$$

An analogous formula holds for the right-hand side of equation (60). Note that $\tilde{u} = u \circ s^{-1}$. Thus we see that to prove equation (60) we need to prove the following two items.

- (i) $(\overset{\alpha}{s}_* \overset{\alpha}{\partial}^j) \left((\overset{\alpha}{s}_* \overset{\alpha}{X}) \tilde{u} \right) \rightarrow (s_* \partial^j) ((s_* X) \tilde{u})$ in $L^\infty(S_h M)$.
- (ii) $\det(d\overset{\alpha}{s}^{-1}) \rightarrow \det(ds^{-1})$ in $L^\infty(S_h M)$.

The push-forward $\overset{\alpha}{s}_*$ has a useful block matrix representation in the coordinates of the unit sphere bundles. Let $(x, \tilde{w}) \in \overset{\alpha}{S}M$ and $(x, v) \in S_h M$ correspond to each other through $\overset{\alpha}{s}(x, \tilde{w}) = (x, v)$. To (x, \tilde{w}) and (x, v) we associate the coordinate vector fields $\partial_{x^1}, \dots, \partial_{x^n}, \partial_{\tilde{w}^1}, \dots, \partial_{\tilde{w}^n}$ and $\partial_{x^1}, \dots, \partial_{x^n}, \partial_{v^1}, \dots, \partial_{v^n}$ respectively. The matrix representation of $\overset{\alpha}{s}_*$ in a block form with respect to the bases $\partial_{x^1}, \dots, \partial_{x^n}, \partial_{\tilde{w}^1}, \dots, \partial_{\tilde{w}^n}$ and $\partial_{x^1}, \dots, \partial_{x^n}, \partial_{v^1}, \dots, \partial_{v^n}$ is

$$\overset{\alpha}{s}_* = \begin{pmatrix} I & 0 \\ \partial_x v & \partial_{\tilde{w}} v \end{pmatrix}. \tag{66}$$

To find coordinate expressions for $\overset{\alpha}{s}_* \overset{\alpha}{\partial}^j$ and $\overset{\alpha}{s}_* \overset{\alpha}{X}$, the vector fields $\overset{\alpha}{\partial}^j$ and $\overset{\alpha}{X}$ need to be expressed in the basis $\partial_{x^1}, \dots, \partial_{x^n}, \partial_{\tilde{w}^1}, \dots, \partial_{\tilde{w}^n}$. As long as everything is only evaluated on $\overset{\alpha}{S}M$, we have $\overset{\alpha}{X} \tilde{u} = \tilde{w}^j \delta_j^{\tilde{u}} \tilde{u}$ for any $\tilde{u}: TM \setminus 0 \rightarrow \mathbb{R}$ such that $\tilde{u}|_{\overset{\alpha}{S}M} = \tilde{u}$. Therefore, in local coordinates and as long as we are careful to only evaluate only $\overset{\alpha}{S}M$, we have

$$\overset{\alpha}{X} = \tilde{w}^j \partial_{x^j} - \tilde{\Gamma}_{jl}^k \tilde{w}^j \tilde{w}^l \partial_{\tilde{w}^k}. \tag{67}$$

Similarly we get $\overset{\alpha}{\partial}^j = \tilde{g}^{jl} \partial_{\tilde{w}^l}$.

Coordinate formulas for the push-forwards vector fields can be found by multiplying with $\overset{\alpha}{s}_*$. This time only evaluating on $S_h M$, we have

$$\overset{\alpha}{s}_* \overset{\alpha}{X} = \tilde{w}^j \partial_{x^j} + \left(\tilde{w}^k \partial_{x^k} v^j - \tilde{\Gamma}_{lm}^k \tilde{w}^l \tilde{w}^m (\partial_{\tilde{w}^k} v^j) \right) \partial_{v^j} \tag{68}$$

and

$$\overset{\alpha}{s}_* \overset{\alpha}{\partial}^j = \tilde{g}^{jl} (\partial_{\tilde{w}^l} v^k) \partial_{v^k}. \tag{69}$$

From these expressions it is clear that convergence in item (i) comes down to three matters. In the base there are derivatives of components of \tilde{g} up to the second order and derivatives of components of \tilde{g}^{-1} up to the first order. Again, components of the metric \tilde{g} appear in the coefficients \tilde{w}^j . The behaviour on the limit of all of these matters is controlled by lemma 18. We have concluded item (i).

To prove item (ii), we write the matrix $d\tilde{s}^{\alpha-1}$ in the block form

$$d\tilde{s}^{\alpha-1} = \begin{pmatrix} I & 0 \\ \partial_x \tilde{w} & \partial_v \tilde{w} \end{pmatrix}. \quad (70)$$

Clearly $\det(d\tilde{s}^{\alpha-1}) = \det(\partial_v \tilde{w})$. Therefore the behaviour as $\alpha \rightarrow \infty$ depends on sums and products of derivatives $\partial_{v^k}(v|v|_{\alpha}^{-1})$, which comes down to the metric \tilde{g} . By lemma 18 we have $\det(d\tilde{s}^{\alpha-1}) \rightarrow \det(ds^{-1})$ in $L^{\infty}(S_h M)$.

We have shown both items (i) and (ii) and thus we have proved equation (60). This finishes the proof of the lemma. \square

Lemma 19. *Let (M, g) be a simple $C^{1,1}$ manifold. The commutator formula $[X, \tilde{\nabla}] = -\tilde{\nabla}$ holds on (M, g) in the H^1 sense: For all $u \in H_0^1(SM)$ and $V \in C^1(N)$ we have*

$$\left(\tilde{\nabla} u, V \right)_{L^2(N)} = \left(\tilde{\nabla} u, XV \right)_{L^2(N)} - \left(Xu, \operatorname{div} V \right)_{L^2(SM)}. \quad (71)$$

Proof. Let h be a smooth reference metric on M and choose a sequence $(\tilde{g})_{\alpha \in \mathbb{N}}$ of smooth metrics approximating g as in lemma 18. For each $\alpha \in \mathbb{N}$ denote $\tilde{u} := u \circ \tilde{r}$ and $\tilde{V} := V \circ \tilde{r}$. Then by remark 15 we have $\tilde{u} \in H_0^1(\tilde{S}M)$ and $\tilde{V} \in C^1(\tilde{N})$. We apply lemma 16 to \tilde{u} and \tilde{V} to get

$$\left(\tilde{\nabla}_{\alpha} \tilde{u}, \tilde{V} \right)_{L^2(\tilde{N})} = \left(\tilde{\nabla}_{\alpha} \tilde{u}, \tilde{X} \tilde{V} \right)_{L^2(\tilde{N})} - \left(\tilde{X} \tilde{u}, \operatorname{div}_{\alpha} \tilde{V} \right)_{L^2(\tilde{S}M)}. \quad (72)$$

Letting $\alpha \rightarrow \infty$ in equation (72) proves the claimed identity (71), after we have shown that

$$\lim_{\alpha \rightarrow \infty} \left(\tilde{\nabla}_{\alpha} \tilde{u}, \tilde{V} \right)_{L^2(\tilde{N})} = \left(\tilde{\nabla} u, V \right)_{L^2(N)}, \quad (73)$$

$$\lim_{\alpha \rightarrow \infty} \left(\tilde{\nabla}_{\alpha} \tilde{u}, \tilde{X} \tilde{V} \right)_{L^2(\tilde{N})} = \left(\tilde{\nabla} u, XV \right)_{L^2(N)} \quad (74)$$

and

$$\lim_{\alpha \rightarrow \infty} \left(\tilde{X} \tilde{u}, \operatorname{div}_{\alpha} \tilde{V} \right)_{L^2(\tilde{S}M)} = \left(Xu, \operatorname{div} V \right)_{L^2(SM)}. \quad (75)$$

All formulas (73), (74) and (75) can be shown by arguments analogous to those used in proving the formula (60) in the proof of lemma 9. Thus we omit the details. Coordinate formulas needed to complete the proofs of the formulas are given in appendix A. \square

5.2. Regularity of the integral function. Let (M, g) be a simple $C^{1,1}$ manifold. In this section we will prove lemma 8 concerning regularity properties of the integral functions of Lipschitz functions and one-forms. We prove a Lipschitz property for the geodesic flow in lemma 20. The Lipschitz property lets us prove that the integral functions are Lipschitz in lemma 21. To prove $H^1(N, X)$ regularity for the vertical gradients of the integral functions in lemma 22, we use the commutator formula from lemma 19.

On a compact manifold M and its unit sphere bundle SM all reasonable notions of distance are bi-Lipschitz equivalent. Since in this section M will be the Euclidean closed ball, we choose Euclidean distances.

Lemma 20. *Let (M, g) be a simple $C^{1,1}$ manifold. For $z \in SM$ let $[\tau_-(z), \tau_+(z)]$ be the maximal interval of existence of the geodesic γ_z . The geodesic flow ϕ_t is Lipschitz continuous in $z \in SM$: There is a uniform $L > 0$ so that for all $z, \hat{z} \in SM$ and $t \in [\tau_-(z), \tau_+(z)] \cap [\tau_-(\hat{z}), \tau_+(\hat{z})]$ we have*

$$d_{SM}(\phi_t(z), \phi_t(\hat{z})) \leq L d_{SM}(z, \hat{z}). \quad (76)$$

Proof. Let $z, \hat{z} \in SM$. Note that both lifted geodesics $t \mapsto \phi_t(z)$ and $t \mapsto \phi_t(\hat{z})$ satisfy the equation $X(\psi(t)) = \psi(t)$ on the interval $[\tau_-(z), \tau_+(z)] \cap [\tau_-(\hat{z}), \tau_+(\hat{z})]$. Since the distance d_{SM} can be taken to be Euclidean and the Christoffel symbols of the metric $g \in C^{1,1}(T^2M)$ are Lipschitz continuous, we get by Grönwall's inequality that

$$d_{SM}(\phi_t(z), \phi_t(\hat{z})) \leq e^{K|t-0|} d_{SM}(\phi_0(z), \phi_0(\hat{z})) = e^{K|t|} d_{SM}(z, \hat{z}), \quad (77)$$

for some $K > 0$ independent of t, z and \hat{z} . Since M is simple $C^{1,1}$, the function τ is uniformly bounded on SM . Thus we find a constant $L > 0$ independent of z and \hat{z} such that $e^{K|t|} \leq L$ uniformly for $t \in [\tau_-(z), \tau_+(z)] \cap [\tau_-(\hat{z}), \tau_+(\hat{z})]$, which finishes the proof. \square

Lemma 21. *Let (M, g) be a simple $C^{1,1}$ manifold. Let $f \in \text{Lip}_0(SM)$ and let u^f be the integral function of f defined by (2). Then $u^f \in \text{Lip}(SM)$.*

Proof. Let $z, \hat{z} \in SM$ be so that $\tau(\hat{z}) \leq \tau(z)$. Then by a simple calculation

$$\begin{aligned} |u^f(z) - u^f(\hat{z})| &\leq (\tau(z) - \tau(\hat{z})) \sup_{t \in [\tau(\hat{z}), \tau(z)]} |f(\phi_t(z))| \\ &\quad + \int_0^{\tau(\hat{z})} |f(\phi_t(z)) - f(\phi_t(\hat{z}))| dt. \end{aligned} \quad (78)$$

We will show that both summands on the right-hand side of equation (78) are bounded by $C d_{SM}(z, \hat{z})$ for some constant $C > 0$ independent of z and \hat{z} .

First, we treat the second term on the right-hand side of (78). Since by lemma 20 the geodesic flow ϕ_t and f both are Lipschitz, there is a constant $K > 0$ independent of t, z and \hat{z} so that

$$|f(\phi_t(z)) - f(\phi_t(\hat{z}))| \leq K d_{SM}(z, \hat{z}). \quad (79)$$

Since the manifold M is simple $C^{1,1}$, there is a constant $L > 0$ independent of \hat{z} so that $\tau(\hat{z}) \leq L$. It follows that

$$\int_0^{\tau(\hat{z})} |f(\phi_t(z)) - f(\phi_t(\hat{z}))| dt \leq K L d_{SM}(z, \hat{z}), \quad (80)$$

which proves the desired bound for the second term.

Then we turn to the first term on the right-hand side of (78). Since f is Lipschitz and vanishes on the boundary $\partial(SM)$, for all $t \in [\tau(\hat{z}), \tau(z)]$ we have

$$\begin{aligned} |f(\phi_t(z))| &= |f(\phi_t(z)) - f(\phi_{\tau(z)}(z))| \\ &\leq \text{Lip}(f) d_{SM}(\phi_t(z), \phi_{\tau(z)}(z)) \\ &\leq \text{Lip}(f)(\tau(z) - t) \\ &\leq \text{Lip}(f)(\tau(z) - \tau(\hat{z})) \\ &\leq \text{Lip}(f)(\tau(z) + \tau(\hat{z})). \end{aligned} \quad (81)$$

The function τ^2 is Lipschitz since the manifold is simple $C^{1,1}$, and so

$$\begin{aligned} (\tau(z) - \tau(\hat{z})) \sup_{t \in [\tau(\hat{z}), \tau(z)]} |f(\phi_t(z))| &\leq \text{Lip}(f)(\tau^2(z) - \tau^2(\hat{z})) \\ &\leq \text{Lip}(f) \text{Lip}(\tau^2) d_{SM}(z, \hat{z}), \end{aligned} \quad (82)$$

as desired.

Combining estimates (78), (80) and (82) yields a Lipschitz estimate for the integral function u^f . \square

Lemma 22. *Let (M, g) be a simple $C^{1,1}$ manifold. Assume that $f \in \text{Lip}_0(SM)$ integrates to zero over all maximal geodesics in M . Then $\overset{\vee}{\nabla} u^f \in H^1(N, X)$, where u^f is the integral function of f defined by equation (2).*

Proof. The integral function u^f is in $\text{Lip}(SM)$ by lemma 21 and $u^f|_{\partial(SM)} = 0$ since f integrates to zero over all maximal geodesics of M . Thus by remark 15 we have $u^f \in H_0^1(SM)$. Then an application of lemma 19 gives

$$\left(\overset{\vee}{\nabla} u^f, XV \right)_{L^2(N)} = \left(\overset{\text{h}}{\nabla} u^f, V \right)_{L^2(N)} - \left(Xu^f, \overset{\vee}{\text{div}} V \right)_{L^2(SM)} \quad (83)$$

for any $V \in C^1(N)$. Here $Xu^f \in H^1(SM)$, since $Xu^f = -f \in \text{Lip}(SM)$. As $Xu^f = -f = 0$ at ∂SM , for any $V \in C^1(N)$ we can integrate by parts in (83) to get

$$\left(\overset{\vee}{\nabla} u^f, XV \right)_{L^2(N)} = \left((\overset{\text{h}}{\nabla} - \overset{\vee}{\nabla} X) u^f, V \right)_{L^2(N)}. \quad (84)$$

Therefore $X \overset{\vee}{\nabla} u^f = (\overset{\vee}{\nabla} X - \overset{\text{h}}{\nabla}) u^f \in L^2(N)$, which shows that $\overset{\vee}{\nabla} u^f \in H^1(N, X)$. \square

Lemma 23. *Let (M, g) be a simple $C^{1,1}$ manifold. Then for any $x \in \partial M$ and $v \in S_x(\partial M)$, there is a sequence of vectors $v_k \in S_x M$ so that $\tau(x, v_k) > 0$, $v_k \rightarrow v$ and $\tau(x, v_k) \rightarrow 0$ as $k \rightarrow \infty$.*

Proof. Let $x \in \partial M$ and $v \in S_x(\partial M)$. Choose a C^1 boundary curve σ defined on an interval I so that $\sigma(0) = x$ and $\dot{\sigma}(0) = v$. Choose a sequence (x_k) of boundary points on $\sigma(I)$ so that $x_k \rightarrow x$. For each k let $v_k \in S_x M$ be the initial velocity of the unique geodesic γ_k joining x to x_k in the interior of M — the geodesic γ_k exists by simplicity. Then $\tau(x, v_k) > 0$ for each k . Since the lengths of the geodesics depend continuously on their end points, we get $\tau(x, v_k) = l(\gamma_k) \rightarrow 0$ as $k \rightarrow \infty$.

It remains to verify that $v_k \rightarrow v$. The geodesic equation gives

$$|\ddot{\gamma}_k^i(t)| = \left| \sum_{j,l} \Gamma_{jl}^i(\gamma_k(t)) \dot{\gamma}_k^j(t) \dot{\gamma}_k^l(t) \right| \leq n^2 \sup_x |\Gamma_{jl}^i(x)| \sup_t |\dot{\gamma}_k(t)|^2, \quad (85)$$

where all norms are the Euclidean ones of the global coordinates and the suprema range over all coordinates. Therefore $|\ddot{\gamma}_k(t)|$ is uniformly bounded for all t and k , and so by Taylor approximation in the coordinates

$$x_k = \gamma_k(\tau_k) = x + \tau_k v_k + \mathcal{O}(\tau_k^2), \quad (86)$$

where the error term is uniform over k . Therefore (in local coordinates)

$$\begin{aligned}
v &= \dot{\sigma}(0) \\
&= \lim_{k \rightarrow \infty} \frac{x_k - x}{\tau_k} \\
&= \lim_{k \rightarrow \infty} \frac{\tau_k v_k + \mathcal{O}(\tau_k^2)}{\tau_k} \\
&= \lim_{k \rightarrow \infty} v_k
\end{aligned} \tag{87}$$

as claimed. \square

Lemma 24. *Let (M, g) be a simple $C^{1,1}$ manifold. Suppose that $f \in \text{Lip}(M)$ integrates to zero over all maximal geodesics of M . Then f vanishes on the boundary ∂M .*

Proof. Let $x \in \partial M$ be a boundary point. Suppose that $v \in S_x(\partial M)$. By lemma 23 there is a sequence of tangent vectors $v_k \in S_x M$ so that $\tau(x, v_k) > 0$, $\tau(x, v_k) \rightarrow 0$ and $v_k \rightarrow v$ when $k \rightarrow \infty$. Since integrals of f over all maximal geodesics vanish, the integral function u^f of f vanishes on the boundary $\partial(SM)$. As the lengths of the geodesics approach zero we get

$$f(x) = \lim_{k \rightarrow \infty} \frac{1}{\tau(x, v_k)} \int_0^{\tau(x, v_k)} f(\gamma_{x, v_k}(t)) dt = \lim_{k \rightarrow \infty} \frac{1}{\tau(x, v_k)} u^f(x, v_k) = 0 \tag{88}$$

as claimed. \square

If $f \in \text{Lip}(SM)$, then the proof above only gives $f|_{\partial_0(SM)} = 0$, not $f|_{\partial(SM)} = 0$. This is true also in the smooth case, and this conclusion is optimal for general functions on SM . If a function on SM arises from a tensor field, then the natural boundary determination is more involved in low regularity and we shall not discuss it here; cf. remark 3.

Proof of lemma 8. (1) Let f be a Lipschitz function on M that integrates to zero over all maximal geodesics of M . Define the integral function u^f of f as in (2). We have $f \in \text{Lip}_0(M)$ by lemma 24. Thus $u^f \in \text{Lip}_0(SM)$ by lemma 21. We have $\overset{\vee}{\nabla} u^f \in H_0^1(N, X)$ by lemma 22 since u^f vanishes on $\partial(SM)$. The last claim $Xu^f = -\pi^* f \in \text{Lip}(SM) \subseteq H^1(SM)$ follows from the fundamental theorem of calculus.

(2) Let h be a Lipschitz 1-form on M that integrates to zero over all maximal geodesics of M and vanishes on the boundary ∂M . Let u^h be the integral function of h defined by (2). Then $u^h \in \text{Lip}_0(SM)$ by lemma 21. We see that $Xu^h \in H^1(SM)$ and $\overset{\vee}{\nabla} u^h \in H_0^1(N, X)$ as in item (1). \square

5.3. The integral function in the Pestov identity. This subsection concludes the proofs of the lemmas required to prove theorem 1. We verify that the integral function of a Lipschitz 1-form h on M behaves in the same way in the Pestov identity as it does in the smooth case. Recall that (M, g) is a simple $C^{1,1}$ Riemannian manifold and particularly g is a $C^{1,1}$ regular Riemannian metric on M .

Proof of lemma 10. Let h be a Lipschitz 1-form on M and denote by \tilde{h} the associated function on SM . We will show that

$$\left\| \overset{v}{\nabla} \tilde{h} \right\|_{L^2(N)}^2 = (n-1) \left\| \tilde{h} \right\|_{L^2(SM)}^2. \quad (89)$$

The Lipschitz assumption guarantees that the left-hand side of (89) is well defined. Let ω stand for the $(n-1)$ -dimensional measure of the unit sphere in \mathbb{R}^n . By [Ilm16, Lemma 4] we have

$$\int_{S_x M} \left| \tilde{h}(x, v) \right|^2 dS_x = |h(x)|^2 \frac{\omega}{n} \quad (90)$$

and

$$\int_{S_x M} \left| \overset{v}{\nabla} \tilde{h}(x, v) \right|^2 dS_x = |h(x)|^2 \frac{\omega(n-1)}{n} \quad (91)$$

on every fiber $S_x M$ of the unit sphere bundle. We may integrate over x just as in [Ilm16, Lemma 4] despite having less regularity, and we find

$$\left\| \overset{v}{\nabla} \tilde{h} \right\|_{L^2(N)}^2 = (n-1) \int_M |h(x)|^2 \frac{\omega}{n} dV_g = (n-1) \left\| \tilde{h} \right\|_{L^2(SM)}^2 \quad (92)$$

as claimed. \square

6. LEMMAS IN SMOOTH GEOMETRY

This final section contains the proofs of the lemmas used to verify that the two definitions of simplicity (definitions 4 and 5) agree when the geometry is C^∞ -smooth. We assume that $M \subseteq \mathbb{R}^n$ is the closed unit ball and we let g be a C^∞ -smooth Riemannian metric on M .

We denote by I_γ the index form along a geodesic γ of M . Recall that if there are interior conjugate points along γ , then I_γ is indefinite and if the end points of γ are conjugate to each other along γ , then there is a normal vector field $V \neq 0$ along γ so that $I_\gamma(V) = 0$. If V is a normal vector field along γ vanishing at the end points of γ , we abbreviate $I_\gamma(V) := I_\gamma(V, V)$.

Proof of lemma 11. Let (M, g) be a simple $C^{1,1}$ manifold and assume that the Riemannian metric g is C^∞ -smooth. Let $\gamma_0: [a, b] \rightarrow M$ be a maximal geodesic in M and let $V \neq 0$ be a normal vector field along γ_0 vanishing at the end points of γ_0 . We will show that $I_{\gamma_0}(V) > 0$, proving that there cannot be conjugate points along γ even at its end points.

Let $(\gamma_0(0), \dot{\gamma}_0(0)) =: z_0 \in \partial_{\text{in}}(SM)$ be the initial data of a geodesic γ_0 and let $\tilde{\gamma}_0$ be the lift to the sphere bundle. The pullback bundle $\tilde{\gamma}_0^* N$ consists precisely of all normal vector fields along γ_0 . Particularly, V is a section of $\tilde{\gamma}_0^* N$ vanishing at the end points, so by lemma 25 (in appendix B) there is a smooth section \tilde{V} of N vanishing on the boundary and satisfying $\tilde{V}|_{\tilde{\gamma}_0} = V$. We may assume that \tilde{V} is supported in a small neighborhood of $\tilde{\gamma}_0$.

Choose for each $k \in \mathbb{N}$ a smooth function $a_k: \partial_{\text{in}}(SM) \rightarrow \mathbb{R}$ so that $a_k^2 \rightarrow \delta_{z_0}$ in the weak sense and $\int_{\partial_{\text{in}}(SM)} a_k^2 d\mu = 1$, where $d\mu(x, v) = \langle \nu(x), v \rangle d\Sigma_g(x, v)$. Since we are working locally around z_0 , it is enough to find such a sequence of functions in Euclidean space and we see that a sequence of square roots of positive standard mollifiers will suffice.

For each $k \in \mathbb{N}$ let $W_k \in C^\infty(N)$ be a section such that $W_k(\phi_t(z)) = a_k(z)\tilde{V}(\phi_t(z))$ for all $z \in \partial_{\text{in}}(SM)$ and $t \in [0, \tau(z)]$. By Santaló's formula (see [Sha99, Lemma 3.3.2]) it follows that as $k \rightarrow \infty$ we have

$$\begin{aligned} Q(W_k) &= \int_{z \in \partial_{\text{in}}(SM)} I_{\gamma_z}(W_k|_{\tilde{\gamma}_z}) d\mu(z) \\ &= \int_{z \in \partial_{\text{in}}(SM)} a_k^2(z) I_{\gamma_z}(\tilde{V}|_{\tilde{\gamma}_z}) d\mu(z) \\ &\rightarrow \int_{z \in \partial_{\text{in}}(SM)} \delta_{z_0}(z) I_{\gamma_z}(\tilde{V}|_{\tilde{\gamma}_z}) d\mu(z) \\ &= I_{\gamma_0}(\tilde{V}|_{\tilde{\gamma}_0}) = I_{\gamma_0}(V). \end{aligned} \tag{93}$$

Here we have written the distribution δ_{z_0} as a function on SM to simplify notation. Similarly as $k \rightarrow \infty$ we get

$$\begin{aligned} \|W_k\|_{L^2(N)}^2 &= \int_{z \in \partial_{\text{in}}(SM)} \int_0^{\tau(z)} |W_k|_{\tilde{\gamma}_z}|^2 dt d\mu(z) \\ &= \int_{z \in \partial_{\text{in}}(SM)} a_k^2(z) \left(\int_0^{\tau(z)} |\tilde{V}|_{\tilde{\gamma}_z}|^2 dt \right) d\mu(z) \\ &\rightarrow \int_{z \in \partial_{\text{in}}(SM)} \delta_{z_0}(z) \left(\int_0^{\tau(z)} |\tilde{V}|_{\tilde{\gamma}_z}|^2 dt \right) d\mu(z) \\ &= \int_0^{\tau(z_0)} |\tilde{V}|_{\tilde{\gamma}_0}|^2 dt = \int_0^{\tau(z_0)} |V|^2 dt. \end{aligned} \tag{94}$$

By $C^{1,1}$ simplicity of (M, g) and zero boundary values of W_k there is $\varepsilon > 0$ so that $Q(W_k) \geq \varepsilon \|W_k\|_{L^2(N)}^2$ for all k . We conclude that

$$I_{\gamma_0}(V) \geq \varepsilon \int_0^{\tau(z_0)} |V|^2 dt > 0, \tag{95}$$

which proves that there cannot be conjugate points along γ_0 even at its end points. \square

Proof of lemma 12. Let (M, g) be a compact smooth Riemannian manifold with a smooth boundary. We assume that the Riemannian metric g is C^∞ -smooth.

First, we will prove that strict convexity implies Lipschitz continuity of τ^2 . As the boundary is strictly convex, all geodesics starting in the interior $\text{int}(SM)$ meet the boundary transversally. The implicit function theorem implies that τ is smooth in $\text{int}(SM)$. As $\tau: SM \rightarrow \mathbb{R}$ is continuous on all of SM , it suffices to show that the gradient of τ^2 (in Sasaki or any other Riemannian metric on SM) is uniformly bounded in the interior.

Let $z \in SM$ be an interior point and let $s \mapsto z_s$ be a smooth curve of interior points, where $s \in (-\varepsilon, \varepsilon)$ and $z_0 = z$. Choose $s \mapsto z_s$ to have unit speed with respect to the Sasaki metric related to the C^∞ -smooth metric g . The implicit function theorem gives an explicit formula for the differential $d\tau$ of τ . Applying the implicit function theorem to $\rho(\gamma_{z_s}(t))$ yields

$$\frac{d}{ds} \tau(z_s) = - \frac{\left\langle \frac{d}{ds} \gamma_{z_s}(t), \nu(\gamma_{z_s}(t)) \right\rangle}{\left\langle \dot{\gamma}_{z_s}(t), \nu(\gamma_{z_s}(t)) \right\rangle} \Big|_{t=\tau(z_s)}, \tag{96}$$

where ρ is a boundary defining function. To prove that $d(\tau^2) = 2\tau d\tau$ is uniformly bounded in the interior, we will show that

$$\tau(z_s) \frac{d}{ds} \tau(z_s) \quad (97)$$

is bounded by some absolute constant near $s = 0$. Boundedness of (97) will follow after we have shown that¹

$$\tau(z) \lesssim |\langle \dot{\gamma}_z(\tau(z)), \nu(\gamma_z(\tau(z))) \rangle| \quad (98)$$

for all $z \in \text{int}(SM)$, since by growth estimates for Jacobi fields and $|\dot{z}_0| = 1$ we have

$$\left| \left\langle \frac{d}{ds} \gamma_{z_s}(\tau(z_s)), \nu(\gamma_{z_s}(\tau(z_s))) \right\rangle \right| \leq C, \quad (99)$$

where C is a constant depending only on curvature bounds and diameter. Since the right-hand side of (98) is constant along the geodesic γ_z , it is enough to prove boundedness for $z \in \partial_{\text{in}}(SM)$.

Outside any neighbourhood of the compact set $\partial_0(SM)$, the right-hand side of (98) is uniformly bounded from below by a positive constant and $\tau(z)$ is also uniformly bounded from above. Thus if we can prove that there is a neighbourhood of the set $\partial_0(SM)$ where (98) holds, it will hold everywhere on $\partial(SM)$.

Take any $x \in \partial M$ and an inward pointing vector $v \in S_x M$. Let

$$\hat{x} := \gamma_{x,v}(\tau(x, v)) \quad \text{and} \quad \hat{v} := -\dot{\gamma}_{x,v}(\tau(x, v)). \quad (100)$$

Let ν be the inward unit normal vector at the boundary. We decompose the vector \hat{v} as $\hat{v}^\perp \nu + \hat{v}^\parallel$, where $\hat{v}^\perp > 0$ and \hat{v}^\parallel is parallel to the boundary. It follows from [Ilm14, Lemma 12] that as $\hat{v}^\perp \rightarrow 0$, we have

$$\tau(\hat{x}, \hat{v}) = 2v^\perp S(\hat{v}^\parallel, \hat{v}^\parallel)^{-1} + \mathcal{O}((\hat{v}^\perp)^2), \quad (101)$$

where S is the second fundamental form of ∂M and the error term is locally uniform. As the boundary is strictly convex, the second fundamental form is bounded uniformly from below by $c > 0$. Thus as $\hat{v}^\perp \rightarrow 0$ we get

$$\tau(\hat{x}, \hat{v}) \leq 3c^{-1} \hat{v}^\perp = 3c^{-1} \langle \nu(\hat{x}), \hat{v} \rangle. \quad (102)$$

Therefore, since $\tau(x, v) = \tau(\hat{x}, \hat{v})$, as $v^\perp \rightarrow 0$ we get

$$\tau(x, v) = \tau(\hat{x}, \hat{v}) \lesssim |\langle \nu(\hat{x}), \hat{v} \rangle| = |\langle \nu(\gamma_{x,v}(\tau(x, v))), \dot{\gamma}_{x,v}(\tau(x, v)) \rangle|. \quad (103)$$

This shows that (98) holds in a neighbourhood of the tangential point (x, v^\parallel) . Thus estimate (98) holds in a neighbourhood of $\partial_0(SM)$.

Next we turn to the opposite statement. We assume that τ^2 is Lipschitz. If the boundary were not to be strictly convex everywhere, there would be a $v \in S_x(\partial M)$ so that $S_x(v, v) \leq 0$.

As τ^2 is Lipschitz, the function τ itself is Hölder-continuous. Because the continuous function τ vanishes on $\partial_{\text{out}} SM \setminus \partial_0(SM)$ (the geodesics stop immediately), we have

$$\tau|_{\partial_0(SM)} = 0 \quad (104)$$

as well.

We use boundary normal coordinates near the base point $x \in \partial M$. We construct a family $(\gamma_h)_{h \in [0,1]}$ of geodesics as follows. Parallel translate the vector v for time h

¹This estimate follows from [Sha94, Lemma 4.1.2], but we reprove it here. Our method of proof is different and may be of interest to some readers.

along the geodesic starting normally inwards from x . Call this vector $v_h \in T_{x_h} M$. Let γ_h be the geodesic with the initial data $\dot{\gamma}_h(0) = v_h$. The geodesic γ_0 (with initial direction $v_0 = v$ at $x_0 = x$) starts at the boundary and may, depending on the convexity of the boundary, be only defined at $t = 0$.

As in [Ilm14, Eq. (2)] we extend the second fundamental form in the boundary normal coordinates near x . Denote $S_h(t) := S_{\gamma_h(t)}(\dot{\gamma}_h(t), \dot{\gamma}_h(t))$. Since $S_x(v, v) \leq 0$ we have

$$S_h(t) = S_0(0) + \mathcal{O}(h) + \mathcal{O}(|t|) \leq C(h + |t|), \quad (105)$$

for some $C > 0$ when h and $|t|$ are small. If $z_h(t)$ is the distance from $\gamma_h(t)$ to the boundary, we have $z_h(0) = h$ and $\dot{z}_h(0) = 0$. By writing the geodesic equation in boundary normal coordinates (as in [Ilm14, Eq. (8)]) we find that

$$\ddot{z}_h(t) = -S_h(t) \geq -C(h + |t|). \quad (106)$$

The total length τ_h of the geodesic γ_h can be divided into forward and backward parts, denoted respectively by τ_h^+ and τ_h^- . We want to find estimates for τ_h^+ and τ_h^- from below.

Let us first consider the case of positive time, $t > 0$. Integrating the estimate (106) leads to

$$z_h(t) = h + \int_0^t \int_0^s \ddot{z}_h(r) dr ds \geq h - \frac{C}{2} h t^2 - \frac{C}{6} t^3 =: \hat{z}_h(t) \quad (107)$$

for all $t > 0$. If we choose $A := \min\left(\sqrt{\frac{2}{3C}}, \sqrt[3]{\frac{2}{C}}\right)$ and $\hat{\tau}_h^+ := Ah^{1/3}$, then for all $t \in [0, \hat{\tau}_h^+]$ we have

$$\hat{z}_h(t) \geq h \left[1 - \frac{1}{2} C h^{2/3} A^2 - \frac{1}{6} C A^3 \right] \geq \frac{h}{3}. \quad (108)$$

Therefore $z_h(t) \geq \hat{z}_h(t) > 0$ for $t \in [0, \hat{\tau}_h^+]$. This shows that $\tau_h^+ \geq \hat{\tau}_h^+$.

The case of negative time can be reduced to previous case by substituting $t = -s$, $s > 0$ and similarly we get $\tau_h^- \geq Ah^{1/3}$. Equation (104) implies $\tau_0 = 0$, and this together with the Lipschitz continuity of τ^2 implies that there is $B > 0$ so that $\tau_h^2 \leq Bh$. As $0 < h \ll 1$, this gives us

$$Bh \geq \tau_h^2 = (\tau_h^+ + \tau_h^-)^2 \geq 4A^2 h^{2/3}, \quad (109)$$

which is impossible for small h . This is a contradiction so the boundary has to be strictly convex. \square

APPENDIX A. COORDINATE FORMULAS AND NORMS

We have collected here the remaining formulas from proofs of lemmas 9 and 19. In the context of the proof of lemma 9 following formulas hold. The Q_α -term in identity (59) is

$$Q_\alpha \left(\overset{\vee}{\nabla}_\alpha \overset{\vee}{u} \right) = \left\| \overset{\vee}{X} \overset{\vee}{\nabla}_\alpha \overset{\vee}{u} \right\|_{L^2(\overset{\vee}{N})} - \left(\overset{\vee}{R} \overset{\vee}{\nabla}_\alpha \overset{\vee}{u}, \overset{\vee}{\nabla}_\alpha \overset{\vee}{u} \right)_{L^2(\overset{\vee}{N})}. \quad (110)$$

For L^2 quantities in identity (59) we have

$$\begin{aligned} \left\| \overset{\alpha}{X} \overset{\vee}{\nabla}_{\alpha} \tilde{u} \right\|_{L^2(\overset{\alpha}{N})}^2 &= \int_{S_h M} \tilde{g}_{ij} \left(\tilde{w}^k \left(\overset{\alpha}{s}_* \overset{\alpha}{\delta}_k \right) \left((\overset{\alpha}{s}_* \overset{\alpha}{\partial}^i) \tilde{u} \right) + \overset{\alpha}{\Gamma}^i_{lk} \tilde{w}^l (\overset{\alpha}{s}_* \overset{\alpha}{\partial}^k) \tilde{u} \right) \\ &\quad \times \left(\tilde{w}^k \left(\overset{\alpha}{s}_* \overset{\alpha}{\delta}_k \right) \left((\overset{\alpha}{s}_* \overset{\alpha}{\partial}^j) \tilde{u} \right) + \overset{\alpha}{\Gamma}^j_{lk} \tilde{w}^l (\overset{\alpha}{s}_* \overset{\alpha}{\partial}^k) \tilde{u} \right) \\ &\quad \times |\det(\mathrm{d}\overset{\alpha}{s}^{-1})| \, \mathrm{d}\Sigma_h \end{aligned} \quad (111)$$

and

$$\left\| \overset{\alpha}{X} \tilde{u} \right\|_{L^2(\overset{\alpha}{S}M)}^2 = \int_{S_h M} \left| \left(\overset{\alpha}{s}_* \overset{\alpha}{X} \right) \tilde{u} \right|^2 |\det(\mathrm{d}\overset{\alpha}{s}^{-1})| \, \mathrm{d}\Sigma_h \quad (112)$$

and

$$\begin{aligned} \left(\overset{\alpha}{R} \overset{\vee}{\nabla}_{\alpha} \tilde{u}, \overset{\vee}{\nabla}_{\alpha} \tilde{u} \right)_{L^2(\overset{\alpha}{N})} &= \int_{S_h M} \tilde{g}_{ij} \left(\overset{\alpha}{R}^i_{jkl} \left((\overset{\alpha}{s}_* \overset{\alpha}{\partial}^j) \tilde{u} \right) \tilde{w}^k \tilde{w}^l \right) \\ &\quad \times \left((\overset{\alpha}{s}_* \overset{\alpha}{\partial}^j) \tilde{u} \right) |\det(\mathrm{d}\overset{\alpha}{s}^{-1})| \, \mathrm{d}\Sigma_h. \end{aligned} \quad (113)$$

For the vector fields $\overset{\alpha}{s}_* \overset{\alpha}{\delta}_k$, $\overset{\alpha}{s}_* \overset{\alpha}{X}$ and $\overset{\alpha}{s}_* \overset{\alpha}{\partial}^j$ appearing in formulas (111), (112) and (113) we have coordinate formulas

$$\overset{\alpha}{s}_* \overset{\alpha}{\delta}_k = \partial_{x^k} + (\partial_{x^k} v^j) \partial_{v^j} - \overset{\alpha}{\Gamma}^i_{kj} \tilde{w}^j (\partial_{\tilde{w}^i} v^l) \partial_{v^l}, \quad (114)$$

$$\overset{\alpha}{s}_* \overset{\alpha}{X} = \tilde{w}^j \partial_{x^j} + \left(\tilde{w}^k \partial_{x^k} v^j - \overset{\alpha}{\Gamma}^k_{lm} \tilde{w}^l \tilde{w}^m (\partial_{\tilde{w}^k} v^j) \right) \partial_{v^j} \quad \text{and} \quad (115)$$

$$\overset{\alpha}{s}_* \overset{\alpha}{\partial}^j = \tilde{g}^{jl} (\partial_{\tilde{w}^l} v^k) \partial_{v^k}. \quad (116)$$

In the context of the proof of lemma 19 the following formulas hold. For the L^2 inner products in equation (72) we have

$$\left(\overset{\vee}{\nabla}_{\alpha} \tilde{u}, \overset{\alpha}{V} \right)_{L^2(\overset{\alpha}{N})} = \int_{S_h M} \tilde{g}_{ij} \left((\overset{\alpha}{s}_* \overset{\alpha}{\partial}^i) \tilde{u} + \tilde{w}^i (\overset{\alpha}{s}_* \overset{\alpha}{X}) \tilde{u} \right) \tilde{V}^j |\det(\mathrm{d}\overset{\alpha}{s}^{-1})| \, \mathrm{d}\Sigma_h \quad (117)$$

and

$$\begin{aligned} \left(\overset{\vee}{\nabla}_{\alpha} \tilde{u}, \overset{\alpha}{X} \overset{\alpha}{V} \right)_{L^2(\overset{\alpha}{N})} &= \int_{S_h M} \tilde{g}_{ij} \left((\overset{\alpha}{s}_* \overset{\alpha}{\partial}^i) \tilde{u} \right) \left((\overset{\alpha}{s}_* \overset{\alpha}{X}) \tilde{V}^j + \overset{\alpha}{\Gamma}^j_{lk} \tilde{w}^l \tilde{V}^k \right) \\ &\quad \times |\det(\mathrm{d}\overset{\alpha}{s}^{-1})| \, \mathrm{d}\Sigma_h \end{aligned} \quad (118)$$

and

$$\left(\overset{\alpha}{X} \tilde{u}, \overset{\vee}{\mathrm{div}}_{\alpha} \tilde{V} \right)_{L^2(\overset{\alpha}{S}M)} = \int_{S_h M} \left((\overset{\alpha}{s}_* \overset{\alpha}{X}) \tilde{u} \right) \left((\overset{\alpha}{s}_* \overset{\alpha}{\partial}^j) \tilde{V}^j \right) |\det(\mathrm{d}\overset{\alpha}{s}^{-1})| \, \mathrm{d}\Sigma_h. \quad (119)$$

New vector fields $\overset{\alpha}{s}_* \overset{\alpha}{\partial}^j$ and $\overset{\alpha}{s}_* \overset{\alpha}{\delta}^k$ appear in equations (117), (118) and (119). For them we have the coordinate formulas

$$\overset{\alpha}{s}_* \overset{\alpha}{\partial}^j = (\partial_{\tilde{w}^j} v^k) \partial_{v^k}, \quad (120)$$

and

$$\overset{\alpha}{s}_* \overset{\alpha}{\delta}^k = \tilde{g}^{kl} \partial_{x^l} + (\tilde{g}^{kl} (\partial_{x^l} v^j)) \partial_{v^j} - \tilde{g}^{kl} \overset{\alpha}{\Gamma}^i_{lm} \tilde{w}^m (\partial_{\tilde{w}^i} v^j) \partial_{v^j}. \quad (121)$$

APPENDIX B. SMOOTH EXTENSION FROM A CURVE

This appendix is devoted to the proof of the following lemma. We will comment on some of the definitions and give examples after the statement. In this appendix everything is smooth and all manifolds and bundles have finite dimension.

Lemma 25. *Let M be a smooth manifold with boundary and $\pi: B \rightarrow M$ a bundle over it whose fiber is a closed manifold. Let $\Pi: E \rightarrow B$ be a vector bundle over B .*

Let $\sigma: [a, b] \rightarrow B$ be a smooth curve without self-intersections so that the end points $\pi(\sigma(a))$ and $\pi(\sigma(b))$ are on ∂M and $\pi(\sigma(t)) \in \text{int}(M)$ for all $t \in (a, b)$. Suppose the exit directions $\dot{\sigma}(a)$ and $\dot{\sigma}(b)$ are not tangent to the boundary $\partial B := \pi^{-1}(\partial M)$.

*Let V be a smooth section of the pullback bundle σ^*E so that $V(a) = V(b) = 0$. Then there is a smooth section W of E so that $W|_{\partial B} = 0$ and $W(\sigma(t)) = V(t)$ for all $t \in [a, b]$.*

The fiber of the bundle B is a smooth and compact manifold of any finite dimension, including zero. The result is valid in the trivial case where the fiber is a singleton and $B = M$. If E is the trivial line bundle $B \times \mathbb{R}$, then sections of it are merely scalar functions $B \rightarrow \mathbb{R}$. Therefore the lemma covers extensions of scalar functions from smooth curves γ on M but also much more. The result will only be applied in the case $B = SM$ and $E = N$, but we record it in more generality as it adds no cost.

As $\sigma: [a, b] \rightarrow B$ is an injective smooth map, a section of the pullback bundle σ^*E is simply a smooth map $W: [a, b] \rightarrow E$ so that $\Pi(W(t)) = \sigma(t)$ for all $t \in [a, b]$.

Proof of lemma 25. Denote the projected curve by $\gamma := \pi \circ \sigma: [a, b] \rightarrow M$. The assumption that $\dot{\sigma}(a)$ and $\dot{\sigma}(b)$ are not tangential to ∂B implies that the end directions $\dot{\gamma}(a)$ and $\dot{\gamma}(b)$ on the base are not tangential to ∂M .

The point $x = \gamma(a)$ has a neighborhood $\omega_1 \subset M$ where we may choose local coordinates $\phi: \omega_1 \rightarrow \mathbb{R}^n$ so that $\phi(\partial M \cap \omega_1) = \{x_n = 0\}$ and for all interior points $y \in M \setminus \partial M$ we have $\phi_n(y) > 0$. In these coordinates the initial direction satisfies $\dot{\gamma}_n(a) > 0$, and so the map

$$\theta: [a, a + \varepsilon) \ni t \mapsto \gamma_n(t) \in [0, h) \quad (122)$$

is a diffeomorphism for some choice of $\varepsilon, h > 0$.

We may shrink ω_1 so that $\phi(\omega_1) \subset \mathbb{R}^{n-1} \times [0, h)$ and the bundle B is locally trivial: $B \supset \pi^{-1}(\omega_1) \approx \omega_1 \times F$, where F is a closed manifold (the typical fiber of B). Denote $y = \sigma(a) \in B_x = F$. There is a neighborhood $U \ni y$ in F so that the bundle E is trivial over $\omega_1 \times U =: \Omega_1 \subset \pi^{-1}(\omega_1) \subset B$ (with the product in the sense of the local trivialization of B) in the sense that $\Pi^{-1}(\Omega_1) \approx \Omega_1 \times \mathbb{R}^K$, where $K \in \mathbb{N}$ is the dimension of the fiber of E . In these coordinates the section V of σ^*E may be written as a smooth function $[a, b] \rightarrow \mathbb{R}^K$, and we denote the component functions as $V_k: [a, b] \rightarrow \mathbb{R}$. By the non-intersecting property of σ we may assume the neighborhoods $\omega_1 \subset M$ and $\Omega_1 \subset B$ to be so small that the curve σ does not return to Ω_1 after leaving it.

We define a function $W_1: \Omega_1 \rightarrow \mathbb{R}^K$ by letting its components be

$$W_1(z)_k = V_k(\theta^{-1}(\pi(z)_n)). \quad (123)$$

This defines a section W_1 of the bundle E in a neighborhood of the point $(x, y) \in B$. By construction $W_1(z) = 0$ when $z \in \partial B$, as that corresponds to the set where

$\pi(z)_n = 0$ and we have $V(a) = 0$. This section W_1 satisfies the required restriction property where it is defined: Whenever $t \in [a, b]$ satisfies $\sigma(t) \in \Omega_1$, we have $W_1(\sigma(t)) = V(t)$.

Similarly, there is a neighborhood Ω_2 of $(\gamma(b), \sigma(b)) \in B$ and a local section $W_2: \Omega_2 \rightarrow E$ with the same property: Whenever $t \in [a, b]$ satisfies $\sigma(t) \in \Omega_2$, we have $W_2(\sigma(t)) = V(t)$.

In addition to satisfying the restriction property, both of these local sections W_1 and W_2 of E vanish on the boundary ∂B when defined there. The point of the construction in (123) is to ensure that the local extension vanishes on the boundary.

For any $t \in (a, b)$ it is easy to provide local extensions as σ has no self-intersections and there are no boundary conditions to worry about. Using compactness of $\sigma([a, b])$ to pass to a finite subcover, we find sets $\Omega_3, \dots, \Omega_J \subset B \setminus \partial B$ and local sections $W_j: \Omega_j \rightarrow E$ of E so that $W_j(\sigma(t)) = V(t)$ whenever $\sigma(t) \in \Omega_j$ and $\sigma([a, b]) \subset \bigcup_{j=1}^J \Omega_j$.

We also let $\Omega_0 = B \setminus \sigma([a, b])$ and let $W_0: \Omega_0 \rightarrow E$ be the zero section. The vector field W_0 has the same boundary conditions and restriction properties as the other W_j s but for trivial reasons.

The sets $\Omega_0, \dots, \Omega_J$ are an open cover of the smooth manifold B with boundary ∂B . Let the functions $\psi_0, \dots, \psi_J \in C_c^\infty(B)$ be a partition of unity subordinate to this cover in the sense that each ψ_j is supported in Ω_j and $\sum_{j=0}^J \psi_j(z) = 1$ for all $z \in B$. The functions $B \rightarrow E$ defined by $\psi_j(z)W_j(z)$ are smooth (interpreted to be zero outside Ω_j where W_j is defined) and the global smooth section $W: B \rightarrow E$ given by

$$W(z) = \sum_{j=0}^J \psi_j(z)W_j(z) \quad (124)$$

is quickly verified to have all the required properties. \square

REFERENCES

- [AD18] Yernat Assylbekov and Nurlan Dairbekov. The X-ray transform on a general family of curves on Finsler surfaces. *J. Geom. Anal.*, 28:1428–1455, 2018.
- [Ain13] Gareth Ainsworth. The attenuated magnetic ray transform on surfaces. *Inverse Probl. Imaging.*, 7 (1):27–46, 2013.
- [AKK⁺04] Michael Anderson, Atsushi Katsuda, Yaroslav Kurylev, Matti Lassas, and Michael Taylor. Boundary regularity for the Ricci equation, geometric convergence, and Gel’fand’s inverse boundary problem. *Invent. Math.*, 158:261–321, 2004.
- [AR97] Yu. E. Anikonov and V. G. Romanov. On uniqueness of determination of a form of first degree by its integrals along geodesics. *J. Inverse Ill-Posed Probl.*, 5(6):487–490, 1997.
- [BKL22] Roberta Bosi, Yaroslav Kurylev, and Matti Lassas. Reconstruction and stability in Gelfand’s inverse interior spectral problem. *Anal. PDE*, 15(2):273–326, 2022.
- [CS98] Christopher Croke and Vladimir Sharafutdinov. Spectral rigidity of a compact negatively curved manifold. *Topology*, 37:1265–1273, 1998.
- [dHIK22] Maarten V. de Hoop, Joonas Ilmavirta, and Vitaly Katsnelson. Spectral rigidity for spherically symmetric manifolds with boundary. *J. Math. Pures Appl. (9)*, 160:54–98, 2022.
- [dHILS21] Maarten V. de Hoop, Joonas Ilmavirta, Matti Lassas, and Teemu Saksala. Stable reconstruction of simple Riemannian manifolds from unknown interior sources, 2021. arXiv:2102.11799 [math.DG].
- [DPSU07] Nurlan Dairbekov, Gabriel P. Paternain, Plamen Stefanov, and Gunther Uhlmann. The boundary rigidity problem in the presence of a magnetic field. *Adv. Math.*, 216:535–609, 2007.

- [DS03] Nurlan Dairbekov and Vladimir Sharafutdinov. Some problems of integral geometry on Anosov manifolds. *Ergod. Theory Dyn. Syst.*, 23, 2003.
- [FU01] David Finch and Gunther Uhlmann. The X-ray transform for a non-Abelian connection in two dimensions. *Inverse Probl.*, 17:695–701, 2001.
- [GPSU16] Colin Guillarmou, Gabriel P. Paternain, Mikko Salo, and Gunther Uhlmann. The X-ray transform for connections in negative curvature. *Comm. Math. Phys.*, 343(1):83–127, 2016.
- [Har50] Philip Hartman. On the local uniqueness of geodesics. *Am. J. Math.*, 72(4):723–730, 1950.
- [Ilm14] Joonas Ilmavirta. Boundary reconstruction for the broken ray transform. *Ann. Fenn. Math.*, 39:485–502, 2014.
- [Ilm16] Joonas Ilmavirta. X-ray transforms in pseudo-Riemannian geometry. *J. Geom. Anal.*, 28:606–626, 2016.
- [IM19] Joonas Ilmavirta and François Monard. Integral geometry on manifolds with boundary and applications. In Ronny Ramlau and Otmar Scherzer, editors, *The Radon Transform: The First 100 Years and Beyond*. de Gruyter, April 2019.
- [IP18] Joonas Ilmavirta and Gabriel P. Paternain. Broken ray tensor tomography with one reflecting obstacle. *Commun. Anal. Geom.*, May 2018. To appear.
- [IS16] Joonas Ilmavirta and Mikko Salo. Broken ray transform on a Riemann surface with a convex obstacle. *Commun. Anal. Geom.*, 24:379–408, 2016.
- [KKL07] Atsushi Katsuda, Yaroslav Kurylev, and Matti Lassas. Stability of boundary distance representation and reconstruction of Riemannian manifolds. *Inverse Probl. Imaging.*, 1:135–157, 2007.
- [Leh16] Jere Lehtonen. The geodesic ray transform on two-dimensional Cartan-Hadamard manifolds, 2016. arXiv:1612.04800 [math.DG].
- [LRS04] Jere Lehtonen, Jesse Railo, and Mikko Salo. Tensor tomography on Cartan-Hadamard manifolds. *Inverse Probl.*, 34:044004, 2004.
- [MNP21] François Monard, Richard Nickl, and Gabriel P. Paternain. Consistent inversion of noisy non-Abelian X-ray transforms. *Commun. Pure Appl. Math.*, 74(5):1045–1099, 2021.
- [MP11] Will Merry and Gabriel P. Paternain. Lecture notes: Inverse problems in geometry and dynamics, 2011.
- [Muk75] R.G. Mukhometov. Inverse kinematic problem of seismic on the plane. *Math. Problems of Geophysics, Akad. Nauk. SSSR, Sibirsk. Otdel., Vychisl. Tsentr, Novosibirsk*, 6:243–252, 1975.
- [Muk77] R.G. Mukhometov. The reconstruction problem of a two-dimensional Riemannian metric (Russian). *Dokl. Akad. Nauk SSSR*, 232(1):32–35, 1977.
- [Muk81] R.G. Mukhometov. On one problem of reconstruction of Riemannian metric (Russian). *Siberian Math. J.*, 22(3):119–135, 1981.
- [Pat99] Gabriel P. Paternain. *Geodesic Flows*. Birkhäuser Basel, Boston, 1999.
- [PS88] Leonid Pestov and Vladimir Sharafutdinov. Integral geometry of tensor fields on manifold of negative curvature. *Sib. Math. J.*, 29:427–441, 1988.
- [PS20] Gabriel P. Paternain and Mikko Salo. The non-Abelian X-ray transform on surfaces, 2020. arXiv:2006.02257 [math.DG].
- [PSU12] Gabriel P. Paternain, Mikko Salo, and Gunther Uhlmann. The attenuated ray transform for connections and Higgs fields. *Geom. Funct. Anal.*, 22:1460–1489, 2012.
- [PSU13] Gabriel P. Paternain, Mikko Salo, and Gunther Uhlmann. Tensor tomography on surfaces. *Inventiones Mathematicae*, 193:229–247, 2013.
- [PSU14a] Gabriel P. Paternain, Mikko Salo, and Gunther Uhlmann. Spectral rigidity and invariant distributions on Anosov surfaces. *J. Differ. Geom.*, 98(1):147–181, 2014.
- [PSU14b] Gabriel P. Paternain, Mikko Salo, and Gunther Uhlmann. Tensor tomography: Progress and challenges. *Chin. Ann. Math. Ser. B*, 35:399–428, 2014.
- [PSU15] Gabriel P. Paternain, Mikko Salo, and Gunther Uhlmann. Invariant distributions, Beurling transforms and tensor tomography in higher dimensions. *Math. Ann.*, 363:305–362, 2015.
- [PSU23] Gabriel P. Paternain, Mikko Salo, and Gunther Uhlmann. *Geometric inverse problems—with emphasis on two dimensions*, volume 204 of *Cambridge Studies in*

- Advanced Mathematics*. Cambridge University Press, Cambridge, 2023. With a foreword by András Vasy.
- [Rom86] Vladimir Romanov. *Inverse Problems of Mathematical Physics*. De Gruyter, Berlin, Boston, 1986.
- [Sha94] V. A. Sharafutdinov. *Integral Geometry of Tensor Fields*. De Gruyter, Berlin, New York, 1994.
- [Sha99] Vladimir Sharafutdinov. *Ray Transform on Riemannian Manifolds: Eight lectures on integral geometry*. University of Oulu, 1999. Lecture notes.
- [SS18] Clemens Sämann and Roland Steinbauer. On geodesics in low regularity. *J. Phys. Conf. Ser.*, 968:012010, 14, 2018.
- [Ste70] Elias Stein. *Singular Integrals and Differentiability Properties of Functions*. Princeton University Press, Princeton, 1970.
- [SU11] Mikko Salo and Gunther Uhlmann. The attenuated ray transform on simple surfaces. *J. Differ. Geom.*, 88(1):161–187, 2011.
- [Uhl14] Gunther Uhlmann. Inverse problems: seeing the unseen. *Bulletin of Mathematical Sciences*, 4:209–279, 2014.
- [US00] Gunther Uhlmann and Vladimir Sharafutdinov. On deformation boundary rigidity and spectral rigidity of Riemannian surfaces with no focal points. *J. Differ. Geom.*, 56(1):93–110, 2000.