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Author(s): Mardor, I.; Ayet San Andrés, S.; Dickel, T.; Amanbayev, D.; Beck, S.; Bergmann, J.; Geissel, H.; Gröf, L.; Haettner, E.; Hornung, C.; Kalantar-Nayestanaki, N.; Kripko-Koncz, G.; Miskun, I.; Mollaebrahimi, A.; Plaß, W. R.; Scheidenberger, C.; Weick, H.; Bagchi, Soumya; Balabanski, D. L.; Bezbakh, A. A.; Brencic, Z.; Charviakova, O.; Chudoba, V.; Constantin, Paul; Dehghan, M.; Fomichev, A. S.; Grigorenko, L. V.;

Title: Mass measurements of As, Se, and Br nuclei, and their implication on the proton-neutron interaction strength toward the N=Z line

Year: 2021

Version: Published version

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Please cite the original version:

Mardor, I., Ayet San Andrés, S., Dickel, T., Amanbayev, D., Beck, S., Bergmann, J., Geissel, H., Gröf, L., Haettner, E., Hornung, C., Kalantar-Nayestanaki, N., Kripko-Koncz, G., Miskun, I., Mollaebrahimi, A., Plaß, W. R., Scheidenberger, C., Weick, H., Bagchi, S., Balabanski, D. L., . . . Zhao, J. (2021). Mass measurements of As, Se, and Br nuclei, and their implication on the proton-neutron interaction strength toward the N=Z line. *Physical Review C*, 103(3), Article 034319. <https://doi.org/10.1103/PhysRevC.103.034319>

Mass measurements of As, Se, and Br nuclei, and their implication on the proton-neutron interaction strength toward the $N = Z$ line

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(Received 4 December 2020; accepted 11 March 2021; published 23 March 2021)

Mass measurements of the nuclides ⁶⁹As, ^{70,71}Se, and ⁷¹Br, produced via fragmentation of a ¹²⁴Xe primary beam at the Fragment Separator (FRS) at GSI, have been performed with the multiple-reflection time-of-flight mass spectrometer (MR-TOF-MS) of the FRS Ion Catcher with an unprecedented mass resolving power of almost 1 000 000. Such high resolving power is the only way to achieve accurate results and resolve overlapping peaks of short-lived exotic nuclei, whose total number of accumulated events is always limited. For the nuclide ⁶⁹As, this is the first direct mass measurement. A mass uncertainty of 22 keV was achieved with only ten events. For the nuclide ⁷⁰Se, a mass uncertainty of 2.6 keV was obtained, corresponding to a relative accuracy of $\delta m/m = 4.0 \times 10^{-8}$, with less than 500 events. The masses of the nuclides ⁷¹Se and ⁷¹Br have been measured with an uncertainty of 23 and 16 keV, respectively. Our results for the nuclides ^{70,71}Se and ⁷¹Br are in good agreement with the 2016 Atomic Mass Evaluation, and our result for the nuclide ⁶⁹As resolves the discrepancy between the previous indirect measurements. We measured also the mass of the molecule ¹⁴N ¹⁵N ⁴⁰Ar ($A = 69$) with a relative accuracy of $\delta m/m = 1.7 \times 10^{-8}$, the highest yet achieved with an MR-TOF-MS. Our results show that the measured restrengthening of the proton-neutron interaction (δV_{pn}) for odd-odd nuclei along the $N = Z$

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line above $Z = 29$ (recently extended to $Z = 37$) is hardly evident at the $N - Z = 2$ line, and not evident at the $N - Z = 4$ line. Nevertheless, detailed structure of δV_{pn} along the $N - Z = 2$ and $N - Z = 4$ lines, confirmed by our mass measurements, may provide a hint regarding the ongoing ≈ 500 keV discrepancy in the mass value of the nuclide ^{70}Br , which prevents including it in the world average of Ft value for superallowed $0^+ \rightarrow 0^+$ β decays. The reported work sets the stage for mass measurements with the FRS Ion Catcher of nuclei at and beyond the $N = Z$ line in the same region of the nuclear chart, including the nuclide ^{70}Br .

DOI: [10.1103/PhysRevC.103.034319](https://doi.org/10.1103/PhysRevC.103.034319)

I. INTRODUCTION

Nuclear masses are key properties of atomic nuclei, as they are a measure of the nuclear binding energy, which reflects the details of nuclear structure and forces between the nucleons. The extension of nuclear mass measurements towards rare, short-lived nuclei far from the valley of stability is important for widening nuclear structure models to all regions of the nuclear chart [1,2] and for astrophysical nucleosynthesis calculations [3].

Nuclear masses of nuclei around the $N = Z$ line provide crucial input for determining the path of the rapid proton-capture (rp) process that takes place in type-I x-ray bursts and steady-state nuclear burning in rapidly accreting neutron stars [4,5].

Experimental masses of numerous odd-odd $N = Z$ nuclei, which decay via superallowed $0^+ \rightarrow 0^+ \beta^+$ transitions, determine Q_{EC} values for these transitions. These values, together with the half-lives and decay branching ratios, formulate the Ft values that in turn test the unitarity of the CKM matrix, a fundamental principle of the standard model [6,7].

However, the Ft value of ^{70}Br that is derived from its precise half-life and decay branching ratio [8] and only direct mass measurement (Penning trap [9]) deviates by approximately 12σ from the world average. Since the measurements of Ref. [8] are accurate enough and not under dispute, the source of this discrepancy must lie in its mass value. The mass measurement of Ref. [9] differs by approximately 500 keV from that obtained via positron end-point energy [10], which gives a consistent Ft value, but with a large uncertainty.

Actually, the Penning trap result was derived from measuring the mass of the ^{70m}Br isomer ($t_{1/2} = 2.2$ s) and subtracting its known excitation energy, because the cycle time employed in the measurement [9] was much longer compared to the half-life of the nuclide ^{70}Br ((78.42 ± 0.51) ms [8]). Therefore, the Penning trap measurement of Ref. [9] is not included in the current Ft value world average, and a new direct mass measurement of the nuclide ^{70}Br is called for in Refs. [6,7].

In this article, we focus on insights that can be derived from accurate mass measurements on the interaction strength between the last proton and the last neutron of intermediate-mass nuclei towards the $N = Z$ line. By ‘last’ we refer to the nucleons that occupy the highest energy orbitals in the nucleus. This interaction strength, denoted as δV_{pn} , can be deduced from the double difference of binding energies (obtained directly from mass measurements) of the nucleus of interest and its neighbors. For odd-odd nuclei with N neutrons and Z protons, whose binding energy is $B(N, Z)$, the expression

is [11]

$$\delta V_{pn}(N, Z) = [B(N, Z) - B(N - 1, Z)] - [B(N, Z - 1) - B(N - 1, Z - 1)]. \quad (1)$$

Assuming that the nuclear core remains essentially unchanged in the four nuclei of Eq. (1), such a double difference largely cancels out the p - p and n - n pairing interactions and the mean field component of the binding energy, isolating the empirical p - n interaction strength [12–15]. Moreover, δV_{pn} can be used as a sensitive filter to study anomalies on the mass surface [16].

It has been predicted that δV_{pn} depends on the spatial overlap of wave functions of the last neutron and proton [17]. This prediction seems to be corroborated by experimental data around doubly magic nuclei [18,19], and especially by the dramatic increase of δV_{pn} at $N = Z$ nuclei, which is attributed to the manifestation of Wigner’s SU(4) isospin and spin symmetry [11].

Due to the relation to the overlap of the wave functions, it is expected that for $N = Z$ nuclei (and in general), δV_{pn} will decrease with increasing Z . This is because of three cumulative reasons. First, as Z increases, the nuclear radii increase, the average distance between the last proton and neutron increases, and their mutual interaction strength decreases even for nucleons occupying the same orbitals. Second, as Z increases, the Coulomb force causes the proton and neutron single-particle energies to differ even when they are in the same orbital [17]. Third, the spin-orbit term in the nuclear mean-field potential increases with mass, gradually breaking the spin part of Wigner’s SU(4) symmetry [11]. The decrease with Z has been observed experimentally for all measured even-even nuclei [17] and odd-odd nuclei up to $Z = 29$ [20].

The authors of Ref. [20] note that from $Z = 29$ onwards, ‘there is an indication of a trend toward the restrengthening of δV_{pn} ’. We submit that based on a newer direct mass measurement of ^{65}As [21] and a very recent indirect mass measurement of ^{73}Rb [22], this restrengthening trend may be confirmed, as shown in Fig. 1, which depicts the experimental data for $N = Z$ odd-odd nuclei from $Z = 29$ to $Z = 37$.

For the nuclide ^{74}Rb , we included in Fig. 1 both the δV_{pn} that is extracted from the AME16 [23] mass value of ^{73}Rb and that from Ref. [22], to show how the 2020 measurement emphasizes the restrengthening trend. For the nuclide ^{70}Br , we included in Fig. 1 both the AME16 value that is based on the Penning trap measurement of [9], and the older indirect measurement based on positron endpoint energy [10], due to the approximate 500 keV discrepancy described above.

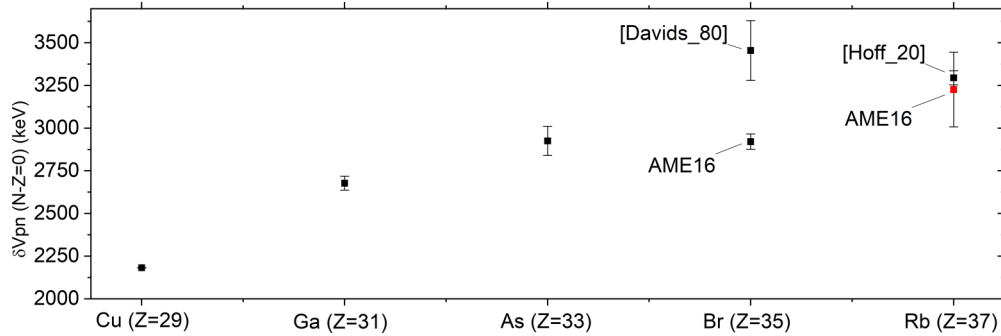


FIG. 1. Proton-neutron interaction strength as a function of Z for odd-odd $N = Z$ nuclei from the nuclide ^{58}Cu to the nuclide ^{74}Rb . Black points are experimental data or AME16 [23] evaluations based on experimental data. The red point for the nuclide ^{74}Rb is the AME16 evaluation based on systematics. Hoff_20 refers to [22] and Davids_80 refers to [10].

The different values for the nuclide ^{70}Br give systematically different local trends of δV_{pn} , either peaking at this nuclide [10] or flattening out [9] before continuing to increase towards the nuclide ^{74}Rb . A new mass measurement of the nuclide ^{70}Br and heavier nuclides in the $N \leq Z$ region should resolve the question which is the correct trend.

For the mass measurement of such nuclei, which are mostly very short-lived (10s of ms) and with low production cross sections, an optimal device is the multiple-reflection time-of-flight mass spectrometer (MR-TOF-MS) [24], which is sensitive even to nuclei that are produced at rates as low as a few events per hour or day [25,26]. The MR-TOF-MS has a unique combination of performance parameters: fast (cycle times of a few ms), accurate (relative mass-measurement uncertainty down to the 10^{-8} level), sensitive (only a few detected ions per nucleus are required for accurate mass determination), non-scanning (simultaneous measurement of many different nuclei), and can be used to spatially resolve isobars and even isomers [26–28].

In this work, we present accurate MR-TOF-MS mass measurements of the nuclides ^{69}As , $^{70,71}\text{Se}$, and ^{71}Br , all near the $N = Z$ line. These results help to confirm trends of δV_{pn} as a function of Z for odd-odd $N - Z = 2$ and $N - Z = 4$ nuclides in the region $Z = 29\text{--}37$, which may provide hints on the expected $N = Z$ δV_{pn} behavior around the nuclide ^{70}Br .

II. EXPERIMENTAL SETUP

The nuclei presented in this work were produced and spatially separated at relativistic energies via projectile fragmentation at the Fragment Separator (FRS) at GSI [29], and delivered to the FRS Ion Catcher [30] where the ions were slowed down and thermalized in a gas-filled cryogenic stopping cell (CSC) [31–33]. Subsequently, the ions were extracted and transported via a radiofrequency quadrupole (RFQ) beam-line [34,35] to an MR-TOF-MS [36,37], to perform mass measurements and ion counting [28] with mass resolving powers at full width at half-maximum (FWHM) of almost 1 000 000. We previously reached a resolving power above 1 000 000 for stable ^{133}Cs ions with our MR-TOF-MS [38,39], and this is the first time we approach this value with unstable ions.

The production mechanism was fragmentation of a ^{124}Xe primary beam impinging on a Be target of 4009 mg/cm^2 thickness at 982 MeV/u , with intensities up to 3×10^9 ions per spill and a typical spill length of 2.5 s. The FRS was set up in achromatic mode from the target to the mid focal plane by using a wedge shaped aluminum degrader (141 mrad , 4005 mg/cm^2) at the focal plane F1. At the midfocal plane a secondary reaction target was installed. From the mid to the final focus the FRS was set to a dispersive mode. A new degrader system, including a variable angle degrader, was used at the final focal plane to range-bunch the ions. The measured range distribution after range bunching was about 50 mg/cm^2 [40]. All nuclides in this work were measured with the same FRS setting, and only the degrader thickness in the final focal plane was changed. In addition to their accurate mass measurement in the MR-TOF-MS, the ions were also identified by the FRS particle ID system.

This enabled the parallel operation of two independent experiments with the same secondary beam—proton decaying nuclei were investigated with the EXPERT setup (see, e.g., [41]) installed after the secondary reaction target at the mid-focal plane, and the longer-living nuclei that are described in this work were studied at the FRS Ion Catcher at the final focus of the FRS.

The CSC temperature during the experiment was 130 K , higher than the standard working temperature ($\sim 85 \text{ K}$), since the CSC had to be warmed up for repair just prior to the experiment, and the available cooling time was not sufficient for reaching the standard temperature. The potential impact of a higher CSC temperature is an increase of the amount and variety of contaminants in the buffer gas, most of which are frozen at the standard working temperature. Nevertheless, the total amount of contaminants in our acquired mass spectra was small, and their overall effect on our results was minute and under control. The CSC was operated with an areal density of 3.3 mg/cm^2 and a mean extraction time of about 80 ms .

III. DATA ANALYSIS

We analyzed the data according to the procedure presented in [28]. Here, we emphasize the parts that were particularly important for this work. The drifts of the time-of-flight data during the experiment were corrected by performing a

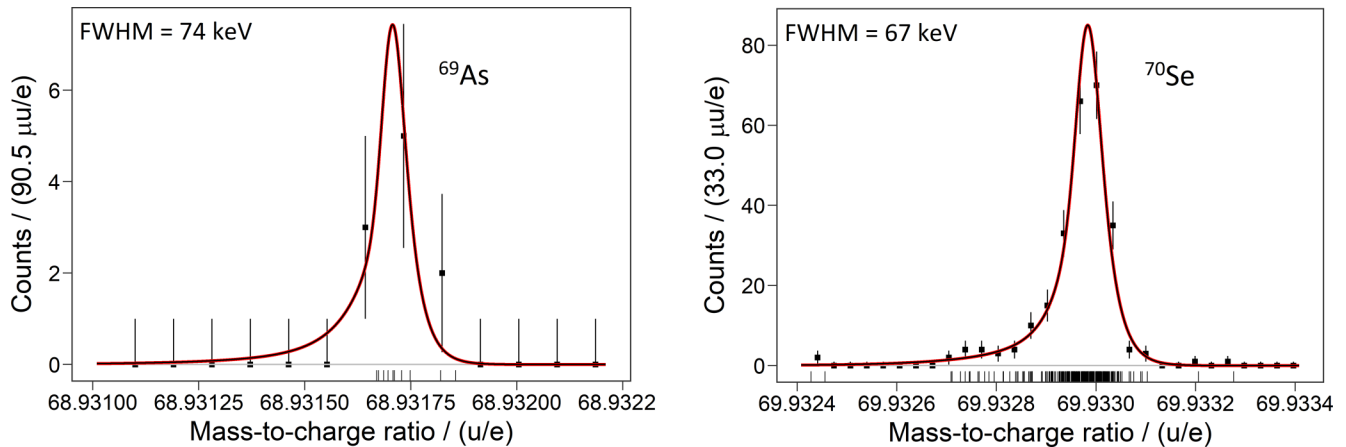


FIG. 2. Left panel: Measured mass-to-charge ratio spectrum of singly charged ^{69}As nuclides. The square points represent the histogram of the unbinned data, which is shown as well in the lower part of the plot. A hyper-EMG function with two exponentials on the left side and one exponential on the right side [hyper-EMG(2,1)] with a FWHM of 74 keV (mass resolving power of 870 000 FWHM) was used for fitting the un-binned data ('rug' graph below the histogram), the shape parameters of which were obtained from the calibrant data (red line). Right panel: Same as on the left, for the nuclide ^{70}Se , showing the results after the wMLE fit of a hyper-EMG(2,1) with a FWHM of 67 keV (mass resolving power of 970 000 FWHM) to a set of data containing 256 counts of the nuclide ^{70}Se .

time-resolved calibration using a well-known mass. The peak shape was obtained from a high-count reference and used for fitting the ion of interest. The analytical function describing the peaks is the hyper- Exponentially Modified Gaussian (EMG) [42], and a weighted maximum likelihood estimate (wMLE) was used to fit this function to the unbinned data. All measured nuclei in this work were measured as singly charged positive ions.

The relationship between time-of-flight and mass-to-charge is defined in Eq. (2),

$$\frac{m}{q} = \frac{c(t_{\text{exp}} - t_0)^2}{(1 + N_{\text{it}}b)^2}, \quad (2)$$

where $b = l_{\text{it}}/l_{\text{tfs}}$ (l_{it} is the path length for one turn in the analyzer and l_{tfs} is the path length from the injection trap to the detector), $c = 2U_{\text{eff}}/l_{\text{tfs}}^2$ (U_{eff} is the effective voltage), t_{exp} is the total measured time-of-flight that takes into account a time delay between the start signal and the real start of the ions (t_0), and N_{it} is the number of turns in the analyzer (see [28] for a more detailed description). The parameter t_0 was common for all the data analyzed in the context of this work and was determined to be 274 ± 2 ns.

Mass values were obtained using calibrant ions that made the same and different turn numbers with respect to the ion of interest. We used four stable molecules that were present in the spectrum and underwent three different turn numbers in the analyzer, and whose relative mass accuracy $\delta m/m$ is in the 10^{-11} level (much better than the accuracy that we expect for the unstable nuclei of interest): $^{12}\text{C}_5^1\text{H}_{10}$ ($A = 70$) performing 847 turns, $^{12}\text{C}_5^1\text{H}_9$ ($A = 69$) performing 853 turns, and $^{14}\text{N}^{15}\text{N}^{40}\text{Ar}$ ($A = 69$) and $^{12}\text{C}^{19}\text{F}_3$ ($A = 69$) each performing 854 turns. We always used $^{12}\text{C}^{19}\text{F}_3$ (the most abundant peak in the spectrum) together with two other molecules to perform the calibration for the mass measurement of the fourth molecule.

In practice, we used $^{12}\text{C}^{19}\text{F}_3$ ($A = 69$) together with a) $^{12}\text{C}_5^1\text{H}_{10}$ ($A = 70$) and $^{12}\text{C}_5^1\text{H}_9$ to measure $^{14}\text{N}^{15}\text{N}^{40}\text{Ar}$

($A = 69$) for the same-turn accuracy test, with b) $^{12}\text{C}_5^1\text{H}_{10}$ and $^{14}\text{N}^{15}\text{N}^{40}\text{Ar}$ ($A = 69$) to measure $^{12}\text{C}_5^1\text{H}_9$, and with c) $^{12}\text{C}_5^1\text{H}_9$ and $^{14}\text{N}^{15}\text{N}^{40}\text{Ar}$ ($A = 69$) to measure $^{12}\text{C}_5^1\text{H}_{10}$.

The measurement with the same-turn calibration of $^{14}\text{N}^{15}\text{N}^{40}\text{Ar}$ ($A = 69$) showed a deviation from literature of (0.6 ± 1.1) keV, which represents a relative mass uncertainty 1.7×10^{-8} . For the multiturn calibration, the measurement of $^{12}\text{C}_5^1\text{H}_{10}$ and $^{12}\text{C}_5^1\text{H}_9$ resulted in a deviation from literature of (-1.2 ± 2.9) keV ($\delta m/m$ of 4.4×10^{-8}) and (0.9 ± 2.9) keV ($\delta m/m$ of 4.5×10^{-8}), respectively, with only about 300 events for each molecule.

The above analyses of the molecular masses show that both of our calibration procedures (same-turn and multiturn) provide accuracy in the low 10^{-8} level. We have reached in the past relative uncertainties of 6.0×10^{-8} with our system, but then it was achieved with approximately 6000 events [28]. The relative mass uncertainty of 1.7×10^{-8} achieved for $^{14}\text{N}^{15}\text{N}^{40}\text{Ar}$ ($A = 69$) in this work is the most accurate measurement yet performed with an MR-TOF-MS, due to our unprecedented mass resolving power of almost 1 000 000 FWHM. The uncertainty evaluation of the molecules and unstable nuclides of this work include all the systematic contributions that were presented in Ref. [28].

IV. RESULTS

A. Mass measurement of the nuclide ^{69}As

The mass of the nuclide ^{69}As has been previously measured only indirectly, via the endpoint energy of its β^+ decay [43,44] and the endpoint energy of its mother nucleus (the nuclide ^{69}Se) β^+ decay [44]. The values from these three measurements (with uncertainties of 50 keV) present some inconsistency with each other, but nevertheless the value at AME16 is a weighted average of them, with a minor weight from the β^+ decay of the nuclide ^{69}Se [45], and an overall uncertainty of 30 keV [23].

TABLE I. Results table of the different nuclides measured in this work and their comparison with AME16 [23]. Half-lives were extracted from [51]. The symbol † denotes a literature value based on indirect mass measurements. See Sec. IV for information about the mass references used in the measurement of each nuclide.

Nuclei	Half-life	ME _{FRS-IC} [keV]	ME _{AME16} [keV]	ME _{FRS-IC} - ME _{AME16} [keV]	Events
⁶⁹ As	15.2 ± 0.2 min	-63116 ± 22	-63110 ± 30†	-6 ± 37	10
⁷⁰ Se	41.1 ± 0.3 min	-61926.0 ± 2.6	-61929.9 ± 1.6	3.9 ± 3.0	485
⁷¹ Se	4.74 ± 0.05 min	-63147 ± 23	-63146.5 ± 2.8	0 ± 23	7
⁷¹ Br	21.4 ± 0.6 s	-56481 ± 16	-56502 ± 5	21 ± 17	19

In this work, we performed the first direct mass measurement of the nuclide ⁶⁹As. The acquired data with the MR-TOF-MS contained only ten events, highlighting the sensitivity of our system, as can be seen in the left side of Fig. 2. The total time of flight for this measurement was about 23.1 ms and the ⁶⁹As ions performed 901 turns inside the analyzer.

Calibration was obtained with the molecular isobar ¹²C¹⁹F₃ ($A = 69$), which performed the same number of turns as the ⁶⁹As ion. Some of these molecules are produced by ionization in the CSC and others from the electron impact source of the MR-TOF-MS (fragment of the molecule C₃F₈).

The ten events of the nuclide ⁶⁹As were acquired in 40 min. When setting the degrader thickness in the FRS final focal plane to stop $A = 70$ nuclides in CSC, we measured also the $A = 69$ region for about 5 h and did not observe any events in the mass region of the nuclide ⁶⁹As. This provided confidence that the acquired ten events are all ⁶⁹As nuclides. Nevertheless, presence of unknown contamination in the peak could not be completely ruled out, and was thus taken into account in the uncertainty evaluation, according to the procedure described in Ref. [28].

Due to the low number of events, the mass uncertainty of the nuclide ⁶⁹As is dominated by statistics and a minor contribution from possible unknown contamination, because of the relatively low amount of unidentified peaks in the entire mass spectrum [28]. Its overall value is 22 keV, corresponding to a relative accuracy of $\delta m/m = 2.8 \times 10^{-7}$.

The mass-excess value and uncertainty of the nuclide ⁶⁹As from this work are given in Table I and are compared to the AME16 value and previous indirect measurements in Fig. 3. It can be seen in Fig. 3 that our result resolves the discrepancy between the previous indirect mass measurements of the nuclide ⁶⁹As.

B. Mass measurement of the nuclide ⁷⁰Se

The mass of the nuclide ⁷⁰Se has been previously measured both indirectly and directly. The direct measurements have been performed with time-of-flight (GANIL [46,47]), Penning traps (ISOLTRAP at ISOLDE/CERN [48] and LEBIT at NSCL [9]) and isochronous mass measurements (ESR at GSI [49]). The AME16 [23] mass value and uncertainty (1.6 keV) are based mainly on the results from Ref. [9]. The more recent Penning trap result is consistent with [9] but with a larger uncertainty (17 keV) [48]. The isochronous measurement result [49] has a large uncertainty (70 keV) and deviates by 2σ from the measurements with the Penning traps. The

direct time-of-flight results are with even larger uncertainties, 130 keV [46] and 460 keV [47], and both are consistent with the Penning trap measurements.

In this work we performed the first mass measurement of the nuclide ⁷⁰Se with an MR-TOF-MS. The total number of accumulated events is 485 with two different turn numbers (848 and 895) with a total TOF of about 21.9 and 23.1 ms, respectively.

In the data sets where the ⁷⁰Se ion underwent 848 turns, an isobaric molecule ¹³C¹⁹F₃ ($A = 70$) undergoing also 848 turns (produced partially in the CSC and partially from the electron impact source in the MR-TOF-MS) was used as a calibrant. In the measurement with 895 turns the ion of interest performed a different number of turns than the calibrant. Therefore, an accurate multi-turn calibration was performed by the determination of the parameter c [see Eq. (2)] using other singly charged species performing different turn numbers, namely: ¹²C₅¹H₁₀ ($A = 70$) performing 847 turns, ¹²C₅¹H₉ ($A = 69$) performing 853 turns, and ¹⁴N¹⁵N⁴⁰Ar ($A = 69$) and ¹²C¹⁹F₃ ($A = 69$) performing 854 turns.

The parameter b [see Eq. (2)] was determined in the different measurements from ¹²C¹⁹F₃ ($A = 69$) or ¹⁴N¹⁵N⁴⁰Ar ($A = 69$) performing 854 turns. Following the procedure described in [28], an uncertainty of 2.6 keV was obtained, which corresponds to a relative accuracy $\delta m/m$ of 4.0×10^{-8} . A sample spectrum containing 256 counts is shown in Fig. 2.

The mass-excess value and uncertainty of the nuclide ⁷⁰Se from this work are given in Table I and are compared together with previous measurements to the value given in AME16 [23] in Fig. 3. It can be seen in Fig. 3 that our result is consistent with previous high accuracy measurements and has a precision similar to the best one.

Our relative uncertainty for the nuclide ⁷⁰Se is similar to the best relative uncertainty recorded so far for an unstable nucleus by an MR-TOF-MS, $\delta m/m = 3.5 \times 10^{-8}$ for ⁶⁵Ga in GARIS-II at RIKEN [50]. Notice that our level of uncertainty was obtained with almost a factor of 40 less events due to the much higher mass resolving power in our MR-TOF-MS.

C. Mass measurement of the nuclide ⁷¹Se

The mass of the nuclide ⁷¹Se has been previously measured both indirectly and directly. The direct measurements have been performed with time-of-flight (GANIL [46,47]), Penning trap (ISOLTRAP at ISOLDE/CERN [48]), and isochronous mass measurements (ESR at GSI [49]). The AME16 [23] mass value and uncertainty (2.8 keV) are based mainly on the results from Ref. [48]. The isochronous measurement

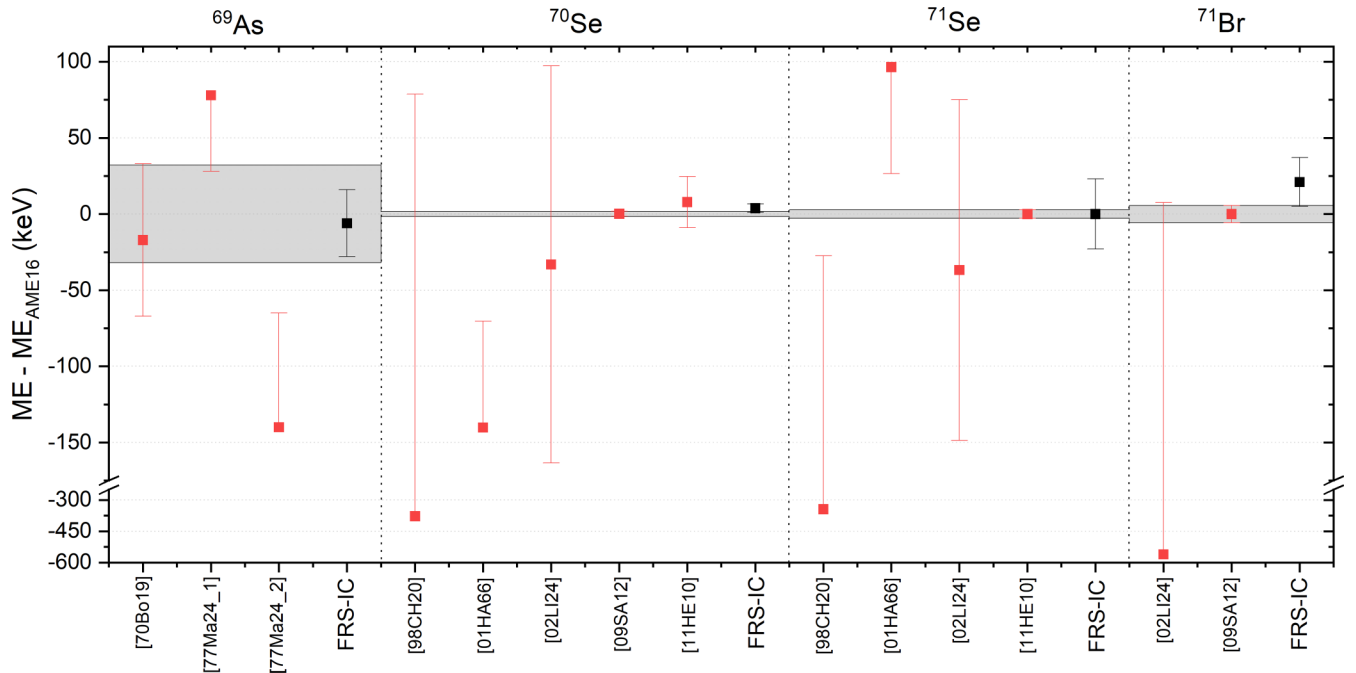


FIG. 3. Deviation of the measured mass-excess of this work (see Table I), denoted by ‘FRS-IC’, and previous ones from the values given in AME16 [23]. [70Bo19] refers to [43], [77Ma22_1] and [77Ma22_2] refer to [44], which includes the published endpoint energies of the decays $^{69}\text{As}(\beta+)^{69}\text{Ge}$ and $^{69}\text{Se}(\beta+)^{69}\text{As}$, respectively, [98CH20] refers to [47], [01HA66] refers to [49], [02LI24] refers to [46], [09SA12] refers to [9], and [11HE10] refers to [48]. The gray band around the horizontal axis represents the AME16 uncertainty. Breaks in the y axis with different scale widths are included in order to be able to see high-accuracy data as well as low-accuracy data. Note that if one of the borders of the error bar is out of the displayed region, then it is not shown.

result [49] has a large uncertainty (70 keV) and deviates by 1.3σ from the Penning trap measurements. The direct time-of-flight results have even larger uncertainties, 112 keV [46] and 317 keV [47], and both are consistent with the Penning trap measurements (see Fig. 3 for a comparison with AME16).

The mass of the nuclide ^{71}Se was measured with only seven counts. The ions underwent 895 turns in about 23.1 ms. There was no same-turn calibration possible and the measurement was performed calculating the parameter c [see Eq. (2)] from the multiturn calibration used with ^{70}Se and using $^{13}\text{C}^{19}\text{F}_3$ ($A = 70$) as a precision calibrant, which performed six turns more than ^{71}Se (901 turns). The uncertainty achieved was 23 keV or a relative mass accuracy $\delta m/m$ of 3.4×10^{-7} , where the main contributions are statistical error and a nonperfect correction of the TOF drifts.

D. Mass measurement of the nuclide ^{71}Br

Due to its closest location to the $N = Z$ line, the mass of the nuclide ^{71}Br has been previously measured only twice, by time of flight (GANIL [46]) and a Penning trap (LEBIT at NSCL [9]). The AME16 mass value and uncertainty (5.4 keV) are based solely on the results of Ref. [9]. The time-of-flight result has a much larger uncertainty, 570 keV [46], which is consistent with the value from Ref. [9]. The comparison can be seen in Fig. 3.

The measurement of the nuclide ^{71}Br included a total of 19 counts. The ions underwent 895 turns in about 23.1 ms. This nuclide was measured together with the nuclide ^{71}Se and

the same calibration procedure was performed. The mass is in good agreement with the one presented in AME16 and presents an uncertainty of 16 keV or a relative mass accuracy $\delta m/m$ of 2.5×10^{-7} .

V. DISCUSSION

The proton-neutron interaction strength for $N = Z$ nuclei, $\delta V_{pn}(N = Z)$, is expected to decrease with increasing mass (see Sec. I [11,17]) whereas recent mass measurements show a restrengthening of this parameter in the range $Z = 29-37$, as depicted in Fig. 1. This could be explained by various physical phenomena, including the amount of overlap between wave functions of valence protons and neutrons [17–19], the growth of deformation in nuclei [14,52], and the manifestation of Wigner’s SU(4) symmetry at $N = Z$ nuclei [11].

In order to discern which interpretation may be behind the surprising trend in $\delta V_{pn}(N = Z)$, we applied the measured masses of this work and corresponding data from AME16 [23] and primary literature [43,44] to calculate the quantity δV_{pn} for $N - Z = 2$ and $N - Z = 4$ nuclides in the proton number range $Z = 29-37$. The results are shown in Fig. 4.

From comparing Fig. 4 to Fig. 1, it is observed that $\delta V_{pn}(N = Z)$ is much higher (up to a factor of almost 3) than the values for $N - Z = 2$ and $N - Z = 4$, as expected [11]. Our new direct measurement of the nuclide ^{69}As confirms an essentially monotonic increase of δV_{pn} staggering as a function of Z , as one moves away from $N = Z$ towards $N - Z = 2$ and $N - Z = 4$. It is thus evident that the significant

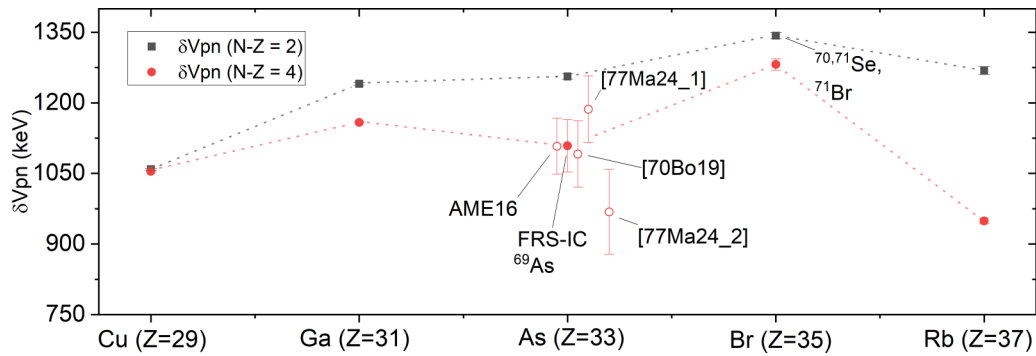


FIG. 4. Proton-neutron interaction strength as a function of Z for odd-odd $N - Z = 2$ and $N - Z = 4$ nuclei from Cu isotopes to Rb isotopes. Solid symbols are experimental data or AME16 evaluations [23] based on experimental data. The open symbols for $\delta V_{pn}(N - Z = 4)$ of the As isotope are values based on the previous indirect mass measurements of the nuclide ^{69}As . [70Bo19] refers to [43], [77Ma22_1] and [77Ma22_2] refer to [44] that includes the published endpoint energies of the decays $^{69}\text{As}(\beta^+)^{69}\text{Ge}$ and $^{69}\text{Se}(\beta^+)^{69}\text{As}$, respectively. Nuclide symbols near the points show the values which the measurement of this work had an impact on (the nuclide ^{69}As) or confirmed (the nuclides $^{70,71}\text{Se}$, ^{71}Br). Dotted lines are shown to guide the eye between δV_{pn} points of the same $N - Z$ value.

restrengthening of δV_{pn} (approximately 1200 keV, see Fig. 1) is a sole characteristic of the $N = Z$ nuclei. The $N - Z = 2$ net increase is only about 200 keV in the same range, and at $N - Z = 4$ there is staggering that ends with a net decrease of approximately 100 keV in the overall range.

This may imply that the restrengthening of $\delta V_{pn}(N = Z)$ is related to partial restoration of a symmetry that is unique to these nuclei. As explained in Sec. I, Wigner's spin-isospin $SU(4)$ symmetry is known to be broken in heavier nuclei [11, 17]. Nevertheless, it has been proposed that protons and neutrons occupying the pf shell above the magic number 28 exhibit a pseudo- $SU(4)$ symmetry [53], which follows from the combined invariance in pseudospin and isospin. Pseudospin is defined for the pf -shell orbitals as a reduction by one unit of angular momentum, so they behave as pseudo sd -shell orbitals. Therefore, their pseudospin-orbit interaction is reduced accordingly, pseudo- $SU(4)$ symmetry is restored, and one may conclude that this causes the increase of the $\delta V_{pn}(N = Z)$ values.

The pseudo- $SU(4)$ scheme has been used for a qualitative understanding of the observed Gamow-Teller β decay of pf -shell nuclei [53]. The authors of Ref. [53] stated that pseudo- $SU(4)$ could be tested via the trend of $\delta V_{pn}(N = Z)$ values in the pf shell. However, the needed mass measurements were lacking at that time. The experimental information that is presented here emphasizes the unique behavior of the $N = Z$ nuclei in the pf shell and can contribute to the test of the pseudo- $SU(4)$ concept.

It is interesting to study also the detailed structure of δV_{pn} as a function of Z for the $N - Z = 2$ and $N - Z = 4$ nuclei. The accurate measurements reveal the staggering and overall increase in δV_{pn} from the Cu ($Z = 29$) isotopes to the Br ($Z = 35$) isotopes, and show the beginning of a decrease at the Rb ($Z = 37$) isotopes. Since the δV_{pn} values beyond the closed shell 28 seem to be influenced by the shell correction and by deformation changes in the four neighboring nuclei, these results may provide input to detailed calculations of nuclear deformation [16].

Some detailed aspects of the δV_{pn} trends of $N - Z = 4$, $N - Z = 2$, and $N = Z$ nuclei may be similar. The most strik-

ing feature in the $N - Z = 4$ and $N - Z = 2$ trends may be the maximal value of δV_{pn} at the Br isotopes. If one assumes that this structure is maintained also at $N = Z$, then Fig. 1 indicates that the 'correct' value for the mass of the nuclide ^{70}Br is closer to the indirect mass measurement [10] rather than the more recent Penning trap measurement [9]. Such a suggested 'peak' at the nuclide ^{70}Br may be connected to the observed transition from spherical to deformed nuclear shapes along the $N = Z$ line at $Z = 35$ [54]. Recall that the indirect mass result of ^{70}Br [10] also leads to a Ft value that is consistent with the world average [6].

VI. SUMMARY AND CONCLUSIONS

In this work we present the mass measurement of four very-neutron-deficient nuclides with the MR-TOF-MS at the FRS Ion Catcher at GSI, with an unprecedented mass resolving power of almost 1 000 000. This is a very important milestone for mass measurements of rare short-lived nuclides, since high resolving power is the only way to achieve the required uncertainties and resolve overlapping peaks when the number of accumulated events is inherently limited. We performed the first direct mass measurement of the nuclide ^{69}As , and the first MR-TOF-MS mass measurements of the nuclides $^{70,71}\text{Se}$ and ^{71}Br .

We carried out these measurements simultaneously with another experiment that used the same secondary beam to study proton decaying nuclei at the midocal plane of the FRS. This is an important demonstration of efficient use of rare ion beams that are currently in great demand at large accelerator facilities.

For the nuclide ^{70}Se , we reached a relative mass uncertainty of $\delta m/m = 4.0 \times 10^{-8}$ with less than 500 events, similar to the best ever measurement of an unstable nucleus with an MR-TOF-MS [50]. For the other three nuclides, we reached a relative mass uncertainty of about 2×10^{-7} with less than 20 events. All our results are in good agreement with the AME16 values.

We measured the mass of the molecule $^{14}\text{N}^{15}\text{N}^{40}\text{Ar}$ ($A = 69$) with a relative accuracy of $\delta m/m = 1.7 \times 10^{-8}$, the highest yet achieved with an MR-TOF-MS.

Our first direct measurement of the nuclide ^{69}As has a better uncertainty than both previous indirect measurements and the AME16 evaluation. Therefore, our results significantly increase the confidence in all four AME16 values, and the consequential conclusions of the nuclear structure and astrophysics models that rely on them.

Our results impact (the nuclide ^{69}As) and confirm (the nuclides $^{70,71}\text{Se}$, ^{71}Br) interesting trends in the proton-neutron interaction strength δV_{pn} for $N - Z = 2$ and $N - Z = 4$ in the $Z = 29\text{--}37$ range, which singled out the $N = Z$ nuclei in this region as exhibiting the restrengthening of δV_{pn} with increasing Z , in contrast with the expected behavior. Further, the maximal δV_{pn} value at the Br isotopes with $N - Z = 2$ and $N - Z = 4$ may indicate that the indirect mass measurement [10] of the nuclide ^{70}Br may be more consistent with the regional δV_{pn} trend than the Penning trap measurement [9]. This gives another important incentive to remeasure the mass of this nuclide, in addition to the apparent deviation of its Q_{EC} from the value that is consistent with the Ft world average [6,7]. From the technical point of view, the high accuracy achieved in the measurements of this work with a small number of events ensure that the MR-TOF-MS at the FRS Ion Catcher at GSI is ready for accurate measurements of the masses of rare $N \leq Z$ nuclides in the $A \approx 70$ region. In particular, our mass measurement of the nuclide ^{71}Br shows that our system is able to extract the specially reactive element bromine, and thus our path to measure the nuclide ^{70}Br and heavier $N = Z$ nuclides is open and will be realized.

ACKNOWLEDGMENTS

The results presented here were obtained in experiment S459+, which was performed at the FRS at the GSI Helmholtzzentrum für Schwerionenforschung, Darmstadt (Germany) in the framework of FAIR Phase-0. This work was carried out within the scientific program of the Super-FRS Experiment Collaboration. This work was supported by the German Federal Ministry for Education and Research (BMBF) under Contracts No. 05P12RGFN8, No. 05P16RGFN1, and No. 05P19RGFN1, Justus-Liebig-Universität Gießen and GSI under the JLU-GSI strategic Helmholtz partnership agreement, HGS-HIRE, the Hessian Ministry for Science and Art (HMWK) through the LOEWE Center HICforFAIR, Polish National Science Centre (Grant No. 2016/21/B/ST2/01227), UKRI STFC Grant No. ST/P004008/1, Ministry of Education, Youth and Sports, Czech Republic (Projects No. LTT17003 and LM2018112), and by the Israel Ministry of Energy, Research Grant No. 217-11-023. This project has received funding from the European Unions Horizon 2020 research and innovation programme under Grant Agreements No. 654002 and No. 771036 (ERC CoG MAIDEN). The ELI-NP team acknowledges the support from the Extreme Light Infrastructure Nuclear Physics (ELI-NP) Phase II, a project co-financed by the Romanian Government and the European Union through the European Regional Development Fund - the Competitiveness Operational Programme (1/07.07.2016, COP, ID 1334). One of the authors (N.K.) would like to acknowledge the support from the ExtreMe Matter Institute EMMI at the GSI Helmholtzzentrum für Schwerionenforschung, Darmstadt (Germany).

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