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Tailoring Vacancy Defects in Isolated Atomically Precise Silver Clusters through Mercury-Doped Intermediates

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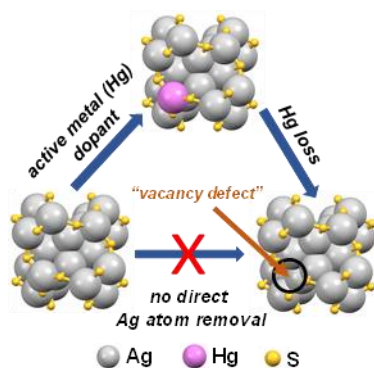
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ABSTRACT. Vacancy defects are known to have significant effects on the physical and chemical properties of nanomaterials. However, the formation and structural dynamics of vacancy defects in atomically precise coinage metal clusters has hardly been explored due to the challenges associated with isolation of such defected clusters. Herein, we isolate $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ (BDT is 1,3-benzenedithiol), a cluster with a ‘missing atom’ site compared to $[\text{Ag}_{29}(\text{BDT})_{12}]^{3-}$, whose precise structure is known from X-ray diffraction. $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ was formed in the gas-phase by collisional heating of $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$, a Hg-doped analogue of the parent cluster. The structural changes resulting from the loss of the Hg heteroatom were investigated by trapped ion mobility mass spectrometry. Density functional theory calculations were performed to provide further insights into the defect structures and molecular dynamics simulations revealed defect site dependent structural relaxation processes.

TOC GRAPHICS



KEYWORDS. nanoclusters, defects, doping, ion mobility, mass spectrometry

Exploring the chemistry and self-healing of defects in materials science is an important area of research as defects play significant roles in controlling optical, chemical, mechanical and electronic properties of the system.¹ For example, vacancy defects on the surface of nanomaterials are known to dramatically affect their catalytic properties.^{2, 3} Ligand-protected clusters of noble metals are an emerging class of nanomaterials with unique physical and chemical properties.⁴⁻⁹ These clusters are atomically precise with their structures typically known from single crystal X-ray diffraction. While the effect of introducing dopant atoms as “antisite” point defects on the structure and properties of these clusters has been investigated in detail,^{10, 11} the consequences of creating “vacancy site” defects by removing atoms are less well explored. Hollow, 12-atom icosahedral metal cores have been observed in the crystal structures of several clusters like $\text{Ag}_{44}(\text{SR})_{30}$, $\text{Au}_{144}(\text{SR})_{60}$, etc. (SR is a thiolate ligand) although filled 13-atom icosahedral cores are expected to be structurally more stable.¹²⁻¹⁴ Zhou et al. have demonstrated that the removal of two kernel Au atoms from $\text{Au}_{48}(\text{SR})_{26}$ by thermally activated solution chemistry has a significant effect on the photoluminescence properties of the resulting cluster.¹⁵ Dong et al. reported well-defined vacancy defects in the surface of ligand-stabilized $\text{Cu}_{36}\text{H}_{10}$ nanoclusters.¹⁶ Such studies based mostly on bottom up synthesis indicate the possibilities of tailoring vacancy defects in atomically precise clusters to study their structure-property correlations in condensed phase. However, removal of single atoms from preformed clusters to create vacancy sites is thermodynamically quite unfavorable due to their high cohesive energy.¹⁷ This imposes significant challenges for synthesis and isolation of such clusters.

Herein, we characterized $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ (BDT is 1,3-benzene dithiol) in the gas-phase, a cluster with a missing Ag atom site compared to $[\text{Ag}_{29}(\text{BDT})_{12}]^{3-}$ whose molecular structure is precisely known from X-ray crystallography.¹⁸ The vacancy defect was created from a heteroatom

(Hg) doped analog of $[\text{Ag}_{29}(\text{BDT})_{12}]^{3-}$ to overcome the challenges associated with single Ag atom removal from the homonuclear cluster. First, we introduced Hg, an active metal dopant, into the preformed cluster by solution reaction. This resulted in the formation of $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ which upon electrospray ionization and subsequent collisional heating led to the efficient loss of the Hg atom and formation of $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$. The isolated $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ fragment species was characterized in the gas-phase using electrospray ionization mass spectrometry (ESI MS) and structural changes upon loss of the heteroatom were studied through trapped ion mobility spectrometry (TIMS). Possible structures of $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ were investigated using density functional theory (DFT) calculations. Moreover, molecular dynamics (MD) simulations were performed to understand the dynamic changes resulting from the creation of a vacancy defect in the cluster.

The $[\text{Ag}_{29}(\text{BDT})_{12}(\text{TPP})_4]^{3-}$ (TPP is triphenyl phosphine) cluster was synthesized following a reported method¹⁸ and characterized using ESI MS as shown in Figure 1A. The mass spectrum showed a dominant peak at m/z 1603 corresponding to $[\text{Ag}_{29}(\text{BDT})_{12}]^{3-}$. The known crystal structure of the cluster is presented in the inset of Figure 1A. The cluster consists of four distinct types of silver atoms: 1) the central Ag atom of the icosahedral core, 2) the 12 equivalent positions on the icosahedral shell, 3) the four Ag atoms at tetrahedral positions on the surface to which TPP ligands are bound in the condensed phase structure but lost during electrospray ionization^{19, 20} and 4) the 12 Ag atoms comprising the four Ag_3S_3 motifs on the surface. So, it is possible in principle, that a heteroatom would be doped in any of these four possible positions in the cluster.

We reacted the cluster with $\text{Hg}(\text{OAc})_2$ and $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ was formed as the product which was characterized by ESI MS (peak at m/z 2452, see Figure 1B). The conversion of $[\text{Ag}_{29}(\text{BDT})_{12}]^{3-}$ to $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ was dependent on the amount of $\text{Hg}(\text{OAc})_2$ used in the

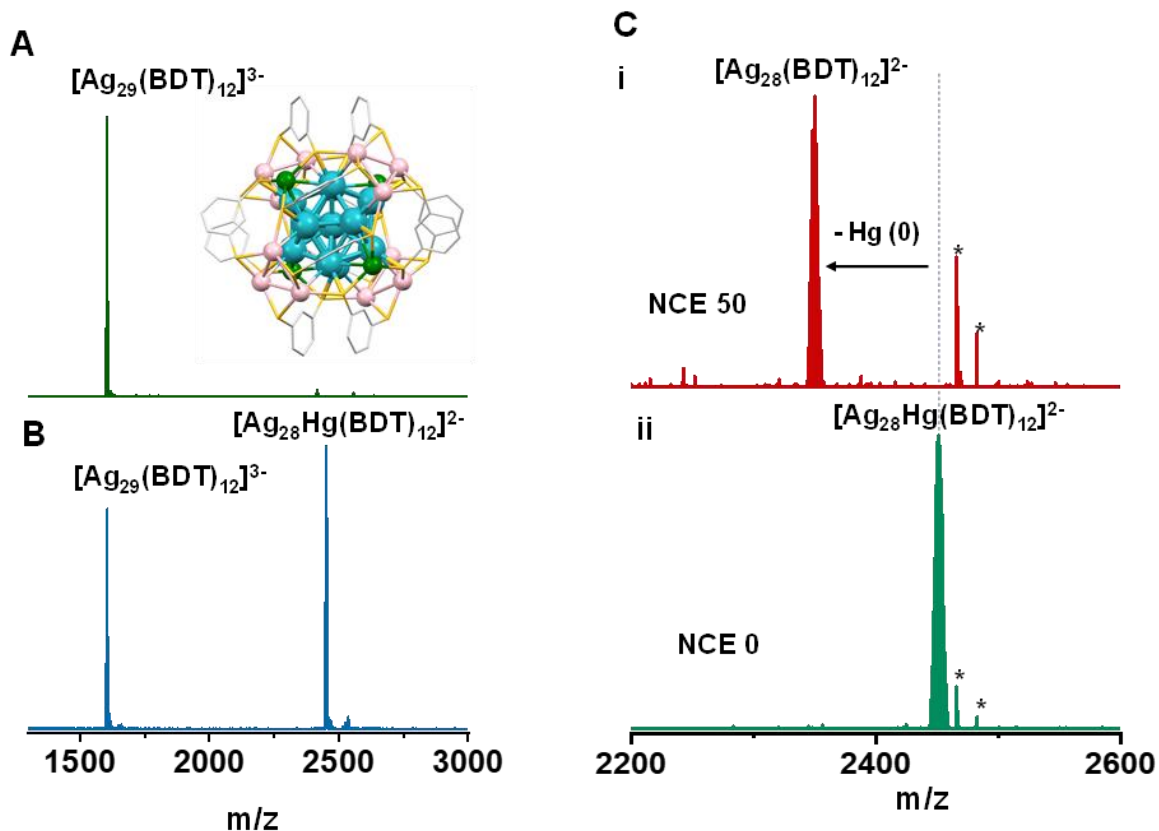
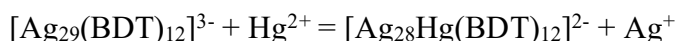


Figure 1. ESI MS of A) pure $[\text{Ag}_{29}(\text{BDT})_{12}]^{3-}$ cluster, B) $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ and $[\text{Ag}_{29}(\text{BDT})_{12}]^{3-}$ mixture and C) Collision-induced dissociation of $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ at normalized collision energy (NCE) of i) 50 and ii) 0. DFT-optimized structure of $[\text{Ag}_{29}(\text{BDT})_{12}]^{3-}$ (using the coordinates of crystal structure) is presented in the inset of A. Color codes: blue: Ag atoms in icosahedron core of the cluster, pink: Ag atoms of the Ag_3S_3 surface "staple" motifs, green: four Ag atoms at tetrahedral positions on the cluster surface where phosphine ligands are attached in the crystal structure, yellow: S, grey: C. H atoms are not shown for clarity. Peaks marked with * are electronic artefacts.

reaction (Figure S1). The relative abundance of the doped product $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ increased with increase in the concentration of $\text{Hg}(\text{OAc})_2$ (30-80 μM) added to the cluster (3 μM). However,

at further higher concentration of Hg(OAc)₂ degradation of the cluster was observed (Figure S1A). In particular, no higher amount of Hg doping was observed, which can be realized e.g. from the absence of any peak for [Ag₂₇Hg₂(BDT)₁₂]¹⁻ (expected m/z 4996) in ESI MS (Figure S1). The experimental and calculated isotopic distributions of the peaks (Figure S2) and changes in optical absorption spectra (Figure S3) are presented in the supporting information. Broadening of the peak at 447 nm was observed in optical absorption (Figure S3). Moreover, time-dependent UV-vis (Figure S4) showed no significant change in the minute time-scale and time-dependent ESI MS (Figure S5) showed that the abundance of [Ag₂₉(BDT)₁₂]³⁻ and [Ag₂₈Hg(BDT)₁₂]²⁻ remained almost unchanged with time. This revealed that at room temperature, the reaction reaches equilibrium very rapidly (compared to our instrumental transfer time of ca. 1 minute). The doping of Hg into [Ag₂₉(BDT)₁₂]³⁻ in solution presumably occurs by an anti-galvanic reduction pathway where Hg²⁺ is first reduced by the cluster and subsequently Hg is doped into the cluster. The proposed reaction pathway is as follows:



Analogous doping by Hg atoms in solution has previously been observed in case of [Au₂₅(SR)₁₈]⁻ where the X-ray structure resolved Hg in the outer shell positions of the cluster.^{21, 22} Note that doping of heteroatoms into coinage metal clusters is known to affect their properties and applications^{10, 11, 23} and similar behavior is expected also in case of [Ag₂₉(BDT)₁₂]³⁻. These are specific to the nature of the dopant atom and a brief comparison of the Hg doping vs. doping with other metal atoms (Au, Cu, Pt)^{19, 23, 24} in [Ag₂₉(BDT)₁₂]³⁻ is discussed in the supporting information.

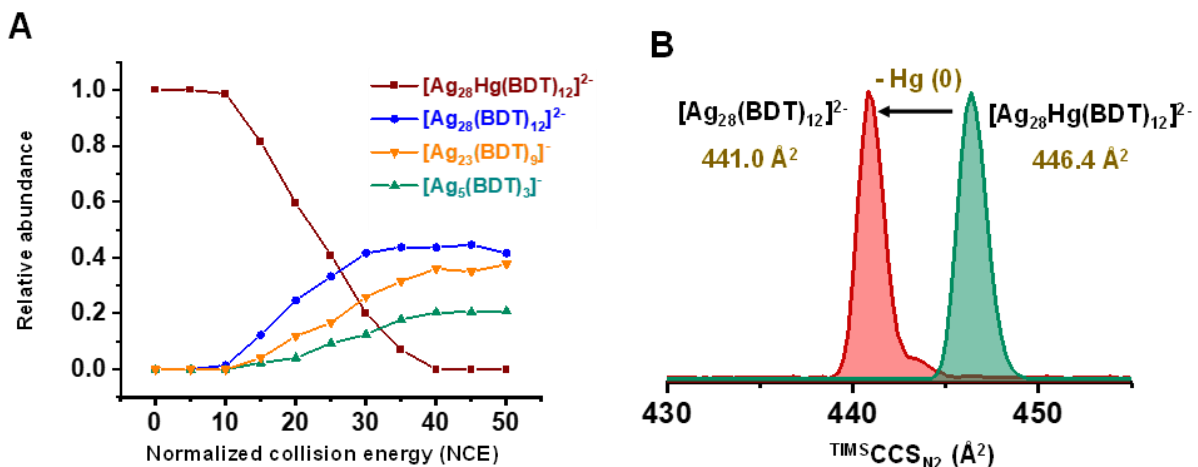


Figure 2. A) NCE-resolved fragmentation curves of $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ and B) TIMS mobilogram of $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ and $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ cluster.

Collisional activation of $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ resulted in the efficient loss of Hg atom and formation of $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$, as shown in Figure 1C. The loss of neutral Hg (0) can be inferred from the unchanged charge states (2-) of $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ and $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ clusters. Typical CID spectra at varying normalized collision energies (NCE) are presented in Figure S6. Following the removal of the Hg atom, the same fragmentation channel as observed for the reference $[\text{Ag}_{29}(\text{BDT})_{12}]^{3-}$ cluster,^{25, 26} i.e. $[\text{Ag}_5(\text{BDT})_3]^{-}$ loss was seen (Figure S6). The corresponding NCE-resolved fragmentation curves are presented in Figure 2A. Such selective loss of the heteroatom during collisional activation was not observed in analogous doped clusters like $[\text{Ag}_{28}\text{Au}(\text{BDT})_{12}]^{3-}$ (Figure S7) where Au was retained in the larger fragment as $[\text{Ag}_{23}\text{Au}(\text{BDT})_9]^{2-}$. In contrast to that, for the Cu doped analogue of the cluster, $[\text{Ag}_{28}\text{Cu}(\text{BDT})_{12}]^{3-}$, Cu was retained in the fragments as $[\text{Ag}_4\text{Cu}(\text{BDT})_3]^{-}$ and $[\text{Ag}_{23}\text{Cu}(\text{BDT})_9]^{2-}$.²⁴

The removal of a single atom from the cluster will result in a vacancy defect site which may affect its overall structural constitution. To explore the structural changes, we performed TIMS studies on isolated $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ and $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ clusters to obtain collision cross sections (CCS) of the ions in N_2 gas and hence information on the structure.²⁷ The fragment ion was produced by increasing the ion transmission voltages before the TIMS separation (Figure S8) following previously established methodologies.^{26, 28} Respective TIMS mobilograms of parent $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ and fragment $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ are in Figure 2B. The CCSs of $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ and $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ were 446.4 \AA^2 and 441.0 \AA^2 , respectively. Thus, TIMS studies revealed a slight contraction (1.2 %) in the structure upon removal of the Hg atom. The CCS resolution achieved for these ions was ~ 200 and this slight relative change of 1.2% is well above the resolution limit. Also, there was no indication of a second isomer with $> 5\%$ relative intensity within this limit of resolution.

To understand the structural changes upon fragmentation, we first computed possible structures of the $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ parent ion considering doping of Hg in four possible sites (as discussed above) of the $[\text{Ag}_{29}(\text{BDT})_{12}]^{3-}$ precursor (Figure S9A further clarifies the four possible sites). DFT optimized structures (retaining staple connectivity) are shown in Figure 3. In the energetically lowest structure of $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$, Hg occupies the icosahedral shell position ($\text{Ag}^{28}\text{HgIso2}$, Figure 3B). The nearest neighbour Hg-Ag bonds on the surface of the icosahedron in the best $\text{Ag}^{28}\text{HgIso2}$ are in range $2.86 \text{ \AA} - 3.29 \text{ \AA}$ while in $[\text{Ag}_{29}(\text{BDT})_{12}]^{3-}$ cluster, the corresponding Ag-Ag distances are in range $2.85 \text{ \AA} - 3.15 \text{ \AA}$, which gives on average an increase of 2 %. The bond distance to the central Ag atom increases only slightly from 2.84 \AA to 2.85 \AA by 0.3% and to the sulfur from 2.52 \AA to 2.59 \AA by 2.7 %. The bonds are visualized in Figure S10A. Compared to

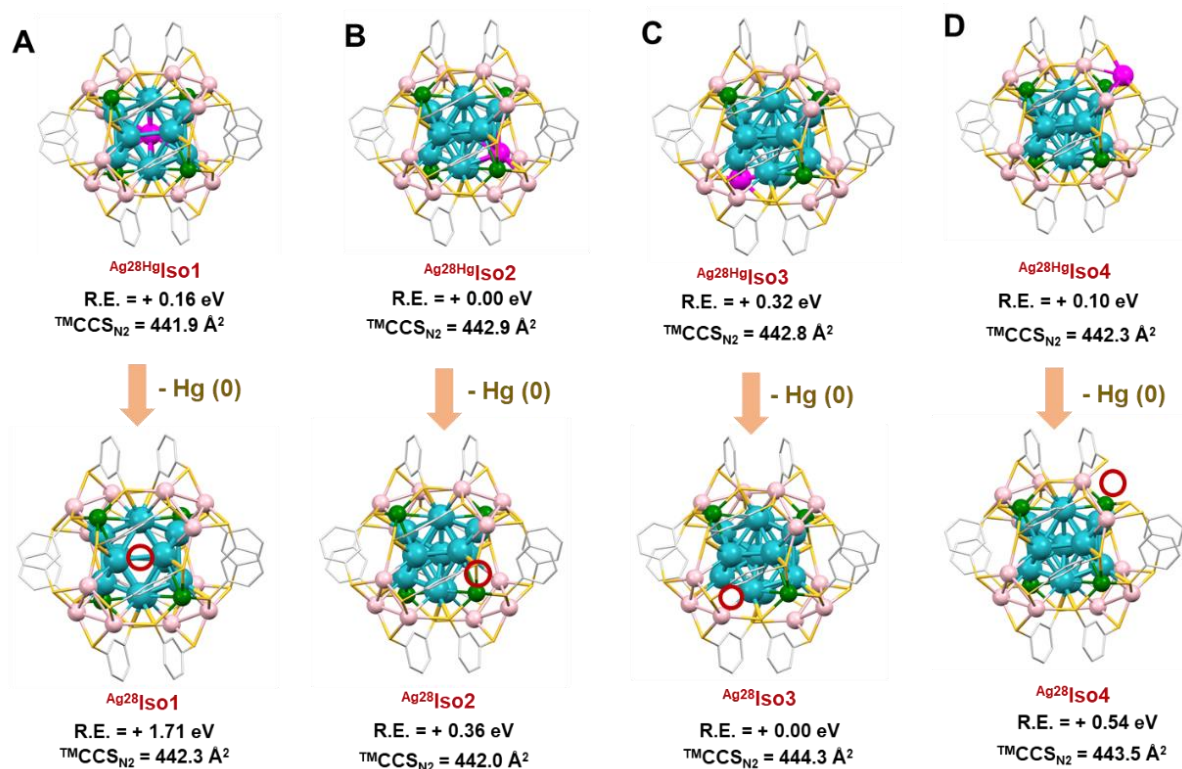


Figure 3. DFT optimized structures of $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ considering doping of the Hg atom in four different positions A) centre of the icosahedron, B) surface of icosahedron, C) at one of the tetrahedral positions on the surface and D) on the Ag_3S_3 surface “staple” motifs of the cluster. Structures of $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ formed by removal of the Hg atom in each case is shown. The vacant atom sites are indicated by red circles. Color codes: blue: Ag atoms in icosahedron core of the cluster, pink: Ag atoms of the Ag_3S_3 surface “staple” motifs, green: four Ag atoms at tetrahedral positions on the cluster surface, purple: Hg, yellow: S, grey: C. H atoms are not shown for clarity.

$[\text{Ag}_{29}(\text{BDT})_{12}]^{3-}$, a local expansion of the structure in the Hg-doped cluster is seen but overall the structure remains essentially the same. We note however, that all four of the $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ isomers considered were within 0.32 eV in energy, suggesting that Hg occupancy in the other sites cannot entirely be ruled out.

Next, starting structures for $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ were constructed by removing Hg from each of the four isomers of $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ and reoptimizing by DFT (Figure 3). Creating a vacancy at the tetrahedral position $^{\text{Ag}28}\text{Iso}3$ is the most favourable. The other isomers that have their vacancy either at the surface of the icosahedral core or at the Ag_3S_3 surface motifs are +0.36 eV and +0.54 eV higher in energy. A vacancy at the center of metal core is least probable being +1.71 eV higher in energy compared to $^{\text{Ag}28}\text{Iso}3$. Structure of the Ag_{28} core only (without the ligands) is also presented in Figure S9B to clearly represent the difference in the vacant sites of the $^{\text{Ag}28}\text{Iso}1-4$ isomers. For the lowest energy $^{\text{Ag}28}\text{Iso}3$ isomer, the effect of the vacancy on the metal ligand interface is best seen by looking at the Ag-S bonds that lie next to the vacancy site connecting the tetrahedral and Ag_3S_3 metal-ligand surface motifs. These bonds contract on average by 4.8 % (from 2.69 Å to 2.56 Å) due to the vacancy as compared to the intact $[\text{Ag}_{29}(\text{BDT})_{12}]^{3-}$ cluster. Ag-S bond distances to the metal core also get shorter by 2.5% from 2.53 Å to 2.47 Å. These bonds around the vacancy site are highlighted in Figure S10B. The effect is very local and the changes elsewhere in the structure are subtle.

If $^{\text{Ag}28}\text{Iso}3$, the product predicted to be most stable is formed, the calculated fragmentation energies for the loss of an Hg atom ($[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-} \rightarrow [\text{Ag}_{28}(\text{BDT})_{12}]^{2-} + \text{Hg}$) lie in the range of ~ 0.86-1.19 eV, depending on which $^{\text{Ag}28}\text{HgIso}1-4$ is assumed to be the parent. Note, that although the fragmentation energies of the four different isomers are within 0.33 eV, activation barriers of the respective processes are expected to be very different. For example, removing Hg from the center of the cluster would involve first transferring the dopant to the surface before detachment.

CCSs of these calculated structures were modelled using the trajectory method (TM) as implemented in IMoS 1.09.²⁹ In our previous report, we discussed the reliability of the calculation

method by comparison of the experimental CCS value (476 \AA^2) of $[\text{Ag}_{29}(\text{BDT})_{12}]^{3-}$ to the theoretical CCS based on its crystal structure.²⁶ Note that the lower CCS of $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ compared to that of parent $[\text{Ag}_{29}(\text{BDT})_{12}]^{3-}$ cluster is primarily associated with the change in its charge state and not the geometrical structure as the charge-induced dipole interactions also effect the CCS.³⁰ $^{\text{TM}}\text{CCS}_{\text{N}_2}$ of $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ isomers were $\sim 441.9 - 442.9 \text{ \AA}^2$ in agreement with the experimental value of 446.4 \AA^2 . No significant decrease in $^{\text{TM}}\text{CCS}_{\text{N}_2}$ was found for the calculated isomers of $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ compared to those calculated for $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$. In experiment, $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ is 1.2 % smaller than $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$. We regard this overall 1-2% level of agreement between experiment and our structural models as excellent – particularly given that IMoS so far uses the general Lennard-Jones parameters for Ag and Hg atoms which makes the corresponding $^{\text{TM}}\text{CCS}_{\text{N}_2}$ simulations subject to larger errors than for molecules containing only lighter elements. A vacancy defect at the previous site of the Hg atom in $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ therefore provides a good description of the $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ structure but on the basis of CCS_{N_2} alone we cannot decide which of the four possible vacancy defect(s) actually form(s).

To better understand structural stability of vacancy sites at finite temperatures, we also performed MD simulations on $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$. An MD simulation (at roughly 290 K) on $\text{Ag}_{28}\text{Iso1}$ with a hollow icosahedron starting structure is presented in Figure 4. Structural changes with time of simulation are reflected in the changes in radius of gyration (Rg) and by several snapshots at intermediate times as indicated. The data includes thermalization from 0 to 3 ps and annealing for 2 ps at the end of the run. During the simulation, it was observed that an Ag atom from the icosahedron shell rapidly migrates to the hollow site at the cluster centre and correspondingly the vacancy site migrates to the icosahedron surface, as shown in Figure 4A. The snapshots show that following this the vacancy in the icosahedral shell then hardly moves on the 12 ps timescale. Only

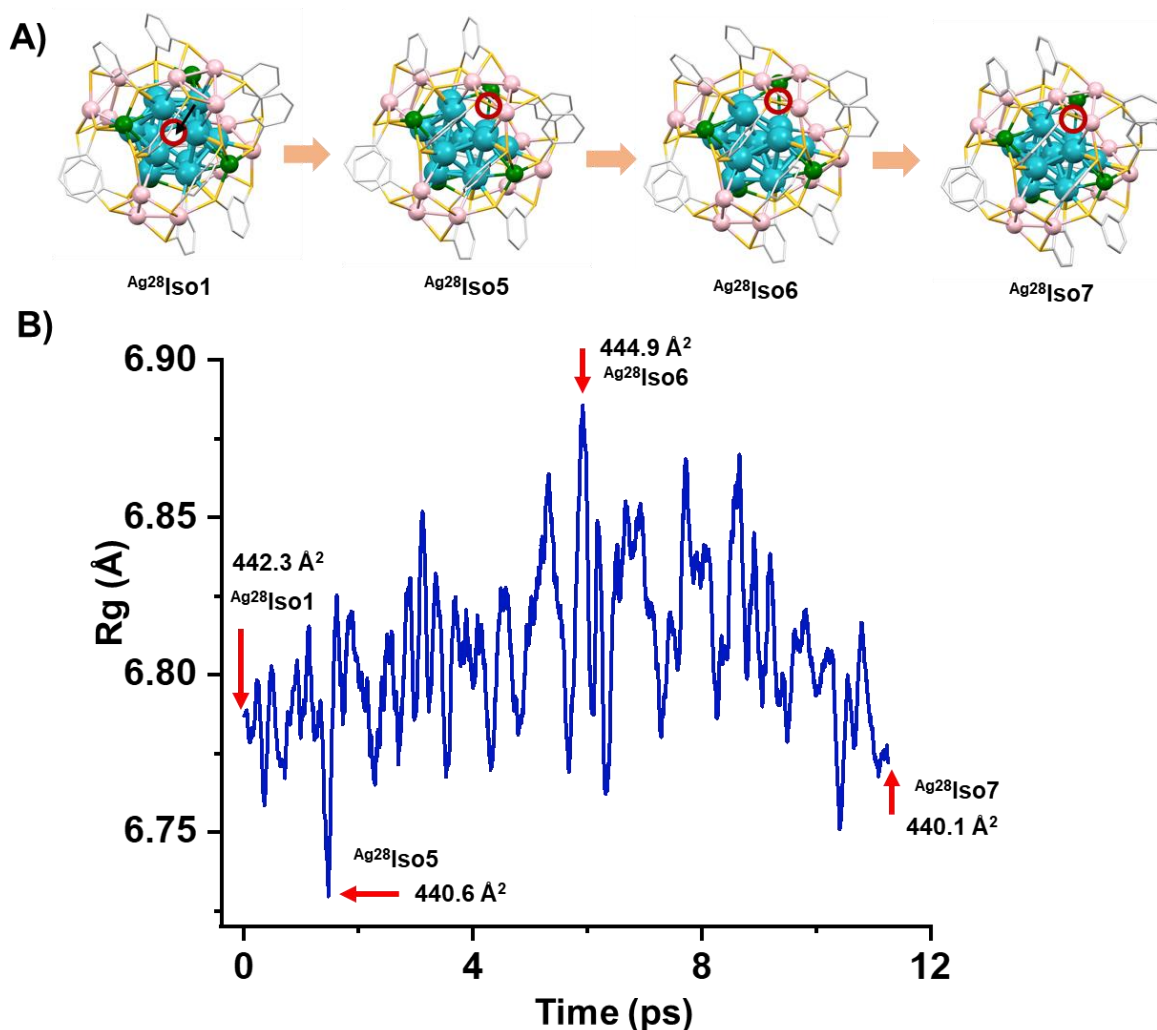


Figure 4. MD simulation on $[Ag_{28}(BDT)_{12}]^{2-}$ cluster ($Ag^{28}Iso1$), A) shows the structural changes during the MD simulation at roughly the same cluster orientation and B) shows the variation of radius of gyration (Rg) with time of simulation (ps). The vacant atom sites are indicated by red circles and movement of Ag atom is marked by an arrow. Color codes: blue: Ag atoms in icosahedron core of the cluster, pink: Ag atoms of the Ag_3S_3 surface “staple” motifs, green: four Ag atoms at tetrahedral positions on the cluster surface, yellow: S, grey: C. H atoms are not shown for clarity.

slight changes ($\pm 0.1 \text{ \AA}$) in size (R_g) of the cluster were observed with time of simulation as shown in Figure 4B. For some of the snapshot structures from the MD run we calculated ${}^{\text{TM}}\text{CCS}_{\text{N}_2}$ (Figure S11) and a slight contraction in CCS was observed in ${}^{\text{Ag}28}\text{Iso5}$ (440.6 \AA^2 which is 0.5 % smaller than ${}^{\text{Ag}28\text{Hg}}\text{Iso3}$) and ${}^{\text{Ag}28}\text{Iso7}$ (440.1 \AA^2 which is 0.6 % smaller than ${}^{\text{Ag}28\text{Hg}}\text{Iso3}$) (Figure 4). Similar fluctuations in R_g ($\pm 1.5\%$) and hence ${}^{\text{TM}}\text{CCS}_{\text{N}_2}$ were also observed during MD simulation (at 278 K) on ${}^{\text{Ag}28}\text{Iso3}$ (Figure S12). Again, the location of the (tetrahedral) vacancy site remains essentially invariant over the simulation time. The relative energies and ${}^{\text{TM}}\text{CCS}_{\text{N}_2}$ of the possible isomers of $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ are listed in Table S1. We therefore conclude from our MD simulations, that at room temperature core vacancies would rapidly anneal and are therefore unlikely to form in this system whereas surface vacancies are stable over significantly longer timescales (during which they do not interconvert). In future experimental work it will be interesting to study to what extent the chemical reactivity of $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ is influenced by such defects.

In summary, we report substitution of an Hg atom into $[\text{Ag}_{29}(\text{BDT})_{12}]^{3-}$ and demonstrate that the resulting doped $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ cluster forms $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ upon collisional activation in gas-phase. The selective loss of the Hg atom results in a vacancy site, which slightly but measurably contracts the overall cluster as measured by its CCS. DFT calculations show that the contraction is seen mainly locally around the vacancy site. We have also used MD simulations to probe associated finite temperature effects. These indicate that at room temperature migration of Ag atoms surrounding a hypothetical vacancy at the center of the cluster (to form a corresponding surface vacancy) would be almost instantaneous. By contrast, surface vacancies appear significantly longer lived against interconversion. In future, it will be of interest to synthesize and isolate such clusters with vacancy defects in the condensed phase, which may be attained through

thermal treatment on preformed Hg doped silver clusters and observing the condensed-phase structures of $[\text{Ag}_{28}\text{Hg}(\text{BDT})_{12}]^{2-}$ and $[\text{Ag}_{28}(\text{BDT})_{12}]^{2-}$ by single crystal X-ray diffraction. It will also be of interest to study the effect of such defects on the reactivity and condensed phase properties of these cluster materials.

ASSOCIATED CONTENT

Supporting Information.

The supporting information is available free of charge.

Experimental and computational methods, UV-vis studies, additional ESI MS and CID data, DFT structures, coordinates of DFT structures

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Notes

The authors declare no competing financial interests.

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