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Title: On Positivity Sets for Helmholtz Solutions

Year: 2023

Version: Published version

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Please cite the original version:

Kow, P.-Z., Salo, M., & Shahgholian, H. (2023). On Positivity Sets for Helmholtz Solutions. Vietnam Journal of Mathematics, 51(4), 985-994. <https://doi.org/10.1007/s10013-023-00646-y>



On Positivity Sets for Helmholtz Solutions

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Received: 25 October 2022 / Accepted: 2 March 2023
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Abstract

We address the question of finding global solutions of the Helmholtz equation that are positive in a given set. This question arises in inverse scattering for penetrable obstacles. In particular, we show that there are solutions that are positive on the boundary of a bounded Lipschitz domain.

Keywords Helmholtz equation · Acoustic equation · Lipschitz domain · Inverse scattering problem

Mathematics Subject Classification (2010) 35J05 · 35J15 · 35J20 · 35R30 · 35R35

1 Introduction

The objective in this short note is to consider the following problem.

Question 1.1 Let $k > 0$ and let E be a subset of \mathbb{R}^n ($n \geq 2$). Does there exist a solution of $(\Delta + k^2)u = 0$ in \mathbb{R}^n with $u|_E > 0$?

Note that any solution of the Helmholtz equation $(\Delta + k^2)u = 0$ is C^∞ , and thus the condition $u|_E > 0$ can be understood pointwise. There is a substantial literature on zero sets of solutions of elliptic equations and eigenfunctions, as discussed in the review [11]. In our setting, any real valued solution of $(\Delta + k^2)u = 0$ in \mathbb{R}^n must have a zero in any closed ball

Dedicated to Carlos E. Kenig on the occasion of his 70th birthday

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of radius $j_{\frac{n-2}{2},1}k^{-1}$ where $j_{\frac{n-2}{2},1}$ is the first zero of the Bessel function $J_{\frac{n-2}{2}}$ (see e.g. [14, Lemma 3.1]). Question 1.1 above is related to producing a global solution whose zero set avoids a given set E .

Our motivation comes from inverse scattering theory and the works [2, 9, 14]. In these works, one considers a bounded open set $D \subset \mathbb{R}^n$ (penetrable obstacle) together with a coefficient $h \in L^\infty(\mathbb{R}^n)$ with $|h| \geq c > 0$ a.e. near ∂D (contrast), and asks whether it is possible to find a solution $u_0 \not\equiv 0$ of $(\Delta + k^2)u_0 = 0$ in \mathbb{R}^n (incident wave) such that the obstacle D with contrast h does not produce any scattering response. The last condition can be precisely formulated as the existence of a function u solving

$$\begin{aligned}
 (\Delta + k^2 + h\chi_D)u &= 0 \text{ in } \mathbb{R}^n, \\
 u &= u_0 \text{ outside some ball.}
 \end{aligned}$$

If this happens for some contrast h , then the obstacle D is called a *non-scattering domain* and it will be invisible with respect to probing with the incident wave u_0 .

It was proved in [14, Theorem 2.1] that if D has real-analytic boundary and if there is an incident wave u_0 with $u_0|_{\partial D} > 0$, then D is a non-scattering domain. Similarly, the work [9] introduced the notion of quadrature domains for the Helmholtz operator $\Delta + k^2$ and proved that if D is such a domain, and if there is an incident wave u_0 with $u_0|_{\partial D} > 0$, then D is a non-scattering domain. On the other hand, the works [2, 14] show that under a nonvanishing condition for u_0 on ∂D , the boundary of a non-scattering domain can be interpreted as a free boundary in an obstacle-type problem and hence such a domain must be either regular or have thin complement near any boundary point.

It was also proved in [14] that one may be able to find incident waves that are positive on the boundary of a bounded C^1 domain (Lipschitz if $n = 2, 3$). Our first main result extends this to Lipschitz domains in any dimension.

Theorem 1.1 *Let $D \subset \mathbb{R}^n$ ($n \geq 2$) be a bounded Lipschitz domain such that $\mathbb{R}^n \setminus \overline{D}$ is connected. Suppose that $k^2 > 0$ is not a Dirichlet eigenvalue of $-\Delta$ in D . Then there exists a Herglotz wave function u_0 (see Definition 2.1) satisfying*

$$(\Delta + k^2)u_0 = 0 \text{ in } \mathbb{R}^n \text{ and } u_0|_{\partial D} > 0.$$

The proof of Theorem 1.1 is done in two steps. One first constructs a solution v of $(\Delta + k^2)v = 0$ in D with $v|_{\partial D} > 0$ by solving a Dirichlet problem. Then one approximates v in D by a suitable Herglotz wave u_0 in \mathbb{R}^n via a Runge approximation argument. This approximation needs to be done in a suitable norm to obtain the pointwise condition $u_0|_{\partial D} > 0$, but since D only has Lipschitz boundary the solution v is not very regular and this limits the choice of possible norms. We will work with fractional Sobolev spaces $H^{s,p}$ and invoke the theory of boundary value problems in Lipschitz domains.¹

We remark that the assumption in Theorem 1.1 that k^2 is not an eigenvalue is necessary, at least when D is a ball (see Example 2.5). For the first eigenvalue this was pointed out in [14, Remark 3.2].

Another instance of subsets $E \subset \mathbb{R}^n$ where one can arrange $u_0|_E > 0$ is given in the following result.

Theorem 1.2 *Let $k > 0$, and let $D \subset \mathbb{R}^n$ ($n \geq 2$) be a bounded Lipschitz domain such that $\mathbb{R}^n \setminus \overline{D}$ is connected and $|D| \leq |B_r|$ where $r = j_{\frac{n-2}{2},1}k^{-1}$. If $E \subset D$ is compact, then there*

¹ This is one of the areas where Carlos Kenig has made pioneering contributions.

exists a Herglotz wave function u_0 (see Definition 2.1) satisfying

$$(\Delta + k^2)u_0 = 0 \text{ in } \mathbb{R}^n \text{ and } u_0|_E > 0. \tag{1.1}$$

The proof is similar to that of Theorem 1.1, except that in the first step we use the Faber–Krahn inequality to produce a solution v that is positive near E .

Remark 1.3 If E is sufficiently nice and low dimensional, it may be possible to use Theorem 1.2 to find solutions that are positive on E . For example, let E be a smooth compact manifold with $\dim(E) = m \leq n - 2$ embedded in \mathbb{R}^n , which is homeomorphic to a compact submanifold E_1 of $\mathbb{R}^{n-1} \cong \mathbb{R}^{n-1} \times \{0\} \subset \mathbb{R}^n$. This holds e.g. when $m < n/2$ by the Whitney embedding theorem, or when E is homeomorphic to S^m . Since $\mathbb{R}^n \setminus E_1$ is connected, by [12, Corollary 7.9] one sees that $\mathbb{R}^n \setminus E$ is (pathwise) connected. One can construct a tubular neighborhood $D = \{x \in \mathbb{R}^n : d(x, E) < \varepsilon\}$ of E having smooth boundary ∂D and arbitrarily small measure [12, Theorem 9.23 and Remark 9.24] (see also [10, Theorem 6.24]). Since $\mathbb{R}^n \setminus E$ is connected, one can connect any two points in $\mathbb{R}^n \setminus D$ by a curve γ in $\mathbb{R}^n \setminus E$. By considering the curve $F(\gamma)$ where F is a continuous map on \mathbb{R}^n that fixes $\mathbb{R}^n \setminus D$ and collapses $D \setminus E$ to ∂D , we see that $\mathbb{R}^n \setminus D$ is connected. Since D has smooth boundary, also $\mathbb{R}^n \setminus \overline{D}$ is connected. (See [5, pp. 61–62] for a related discussion.) Thus we may apply Theorem 1.2 to find a Herglotz wave function u_0 satisfying (1.1). Note that the connectedness of $\mathbb{R}^n \setminus E$ can fail when E has dimension $n - 1$.

2 Solutions Satisfying the Positivity Condition

In this section we will prove Theorems 1.1 and 1.2. We begin with some preparations.

2.1 Fractional Sobolev Spaces

For each $s \in \mathbb{R}$ and $1 < p < \infty$, the fractional Sobolev space $H^{s,p}(\mathbb{R}^n)$ is the Banach space equipped with the norm

$$\|u\|_{H^{s,p}(\mathbb{R}^n)} := \|\langle D \rangle^s u\|_{L^p(\mathbb{R}^n)},$$

where $\langle D \rangle^s$ is the the Bessel potential of order s , i.e. the Fourier multiplier corresponding to $\langle \xi \rangle^s = (1 + |\xi|^2)^{\frac{s}{2}}$. In particular when $s = k \geq 1$ is an integer, we also have $H^{k,p}(\mathbb{R}^n) = W^{k,p}(\mathbb{R}^n)$, where

$$W^{k,p}(\mathbb{R}^n) = \{u \in L^p(\mathbb{R}^n) \mid D^\alpha u \in L^p(\mathbb{R}^n) \text{ for all multi-indices } \alpha \text{ with } |\alpha| \leq k\}.$$

From [1, Corollary 6.2.8], we have the duality statement

$$(H^{s,p}(\mathbb{R}^n))^* = H^{-s,p'}(\mathbb{R}^n) \text{ for all } s \in \mathbb{R} \text{ and } 1 < p < \infty, \tag{2.1}$$

where $(p')^{-1} + p^{-1} = 1$. We also recall the Sobolev embedding ([1, Theorem 6.5.1]):

$$H^{s,p}(\mathbb{R}^n) \subset H^{s_1,p_1}(\mathbb{R}^n)$$

whenever $1 < p \leq p_1 < \infty$, $-\infty < s_1 \leq s < \infty$, and $s - \frac{n}{p} = s_1 - \frac{n}{p_1}$.

Let D be an open set in \mathbb{R}^n . We define

$$H^{s,p}(D) := \{u|_D \mid u \in H^{s,p}(\mathbb{R}^n)\} \text{ for all } s \in \mathbb{R} \text{ and } 1 < p < \infty.$$

This is a Banach space equipped with the quotient norm

$$\|v\|_{H^{s,p}(D)} := \inf\{\|u\|_{H^{s,p}(\mathbb{R}^n)} \mid u|_D = v\}.$$

When D is a bounded Lipschitz domain, from [8, Theorem 2.3] we know that there exists a bounded linear extension operator

$$E : H^{s,p}(D) \rightarrow H^{s,p}(\mathbb{R}^n) \quad \text{with } Eu = u \text{ in } D \text{ for all } u \in H^{s,p}(D).$$

If $F \subset \mathbb{R}^n$ is closed, we define

$$H_F^{s,p}(\mathbb{R}^n) := \{u \in H^{s,p}(\mathbb{R}^n) \mid \text{supp}(u) \subset F\}.$$

If D is a bounded Lipschitz domain, the following result can be found in [8, Remark 2.7]:

$$C_c^\infty(D) \text{ is dense in } H_D^{s,p}(\mathbb{R}^n) \text{ for each } s \in \mathbb{R} \text{ and } 1 < p < \infty. \tag{2.2}$$

2.2 Runge–Herglotz Approximation

The next objective is to prove a result stating that solutions in $H^{s,p}(D)$ can be approximated in D by Herglotz waves. We first give a definition.

Definition 2.1 Let $k > 0$ and consider the operator $P_k : C^\infty(S^{n-1}) \rightarrow C^\infty(\mathbb{R}^n)$ defined by

$$(P_k f)(x) := \int_{S^{n-1}} e^{ikx \cdot \hat{z}} f(\hat{z}) d\hat{z}, \quad x \in \mathbb{R}^n.$$

The functions $u = P_k f$ with $f \in C^\infty(S^{n-1})$ are called *Herglotz waves*, and they are particular solutions of $(\Delta + k^2)u = 0$ in \mathbb{R}^n .

Proposition 2.2 Let $k > 0$, $0 < s \leq 1$, $1 < p < \infty$, and let $D \subset \mathbb{R}^n$ ($n \geq 2$) be a bounded Lipschitz domain such that $\mathbb{R}^n \setminus \overline{D}$ is connected. Given any $v \in H^{s,p}(D)$ with $(\Delta + k^2)v = 0$ in D , there exist Herglotz waves $u_j \in C^\infty(\mathbb{R}^n)$ such that

$$\|u_j - v\|_{H^{s,p}(D)} \rightarrow 0 \quad \text{as } j \rightarrow \infty.$$

If v is real-valued, then so are u_j .

The proof of Proposition 2.2 is very similar to [14, Proposition 3.4] that considered approximation in $W^{1,p}(D)$. Here we need to work with fractional Sobolev spaces instead.

Proof In view of the Hahn–Banach theorem, it is enough to prove that any bounded linear functional $\ell : H^{s,p}(D) \rightarrow \mathbb{C}$ that vanishes on $\{P_k f|_D \mid f \in C^\infty(S^{n-1})\}$ must also vanish on $\{v \in H^{s,p}(D) \mid -(\Delta + k^2)v = 0 \text{ in } D\}$. Let ℓ be such a linear functional, and define a bounded linear functional $\ell_1 : H^{s,p}(\mathbb{R}^n) \rightarrow \mathbb{C}$ by $\ell_1(u) := \ell(u|_D)$. By duality (2.1), there exists a unique $\mu \in H^{-s,p'}(\mathbb{R}^n)$ such that

$$\ell_1(u) = (u, \mu) \quad \text{for all } u \in H^{s,p}(\mathbb{R}^n),$$

where (\cdot, \cdot) is the sesquilinear distributional pairing in \mathbb{R}^n . It is easy to see that $\mu = 0$ in $\mathbb{R}^n \setminus \overline{D}$, and the condition $\ell(P_k f|_D) = 0$ for all $f \in C^\infty(S^{n-1})$ implies that

$$(P_k f, \mu) = 0 \quad \text{for all } f \in C^\infty(S^{n-1}). \tag{2.3}$$

We now define the distribution $w := \Phi_k * \mu$, where

$$\Phi_k(x) = \frac{ik^{\frac{n-2}{2}}}{4(2\pi)^{\frac{n-2}{2}}} |x|^{-\frac{n-2}{2}} H_{\frac{n-2}{2}}^{(1)}(k|x|)$$

is the outgoing fundamental solution of the Helmholtz operator $-(\Delta + k^2)$ and $H_\alpha^{(1)}$ is the Hankel function (see [15, §1.2.3]). Then w is a distributional solution of

$$-(\Delta + k^2)w = \mu \quad \text{in } \mathbb{R}^n. \tag{2.4}$$

Elliptic regularity yields $w \in H_{\text{loc}}^{2-s,p'}(\mathbb{R}^n)$, and since $\text{supp}(\mu) \subset \overline{D}$ we also have that w is C^∞ in $\mathbb{R}^n \setminus \overline{D}$.

Given any $f \in C^\infty(\mathcal{S}^{n-1})$, we write $u = P_k f \in C^\infty(\mathbb{R}^n)$. Using (2.3) and the fact that μ has compact support, we have

$$0 = (u, \mu) = \lim_{r \rightarrow \infty} (u, \mu)_{B_r}, \tag{2.5}$$

where $(\cdot, \cdot)_{B_r}$ is the sesquilinear distributional pairing in the ball B_r . We now consider a cut-off function $\chi \in C_c^\infty(\mathbb{R}^n)$ satisfying $0 \leq \chi \leq 1$ and $\chi = 1$ near \overline{D} . Using (2.4), we can write (2.5) as

$$\begin{aligned} 0 &= \lim_{r \rightarrow \infty} [(\chi u, (\Delta + k^2)w)_{B_r} + ((1 - \chi)u, (\Delta + k^2)w)_{B_r}] \\ &= \lim_{r \rightarrow \infty} [((\Delta + k^2)(\chi u), w)_{B_r} + ((\Delta + k^2)((1 - \chi)u), w)_{B_r} \\ &\quad + \int_{\partial B_r} (u \overline{\partial_{|x|} w} - (\partial_{|x|} u) \overline{w}) dS] \\ &= \lim_{r \rightarrow \infty} \int_{\partial B_r} (u \overline{\partial_{|x|} w} - (\partial_{|x|} u) \overline{w}) dS, \end{aligned} \tag{2.6}$$

where $\partial_{|x|} = \hat{x} \cdot \nabla$ denotes the radial derivative. Here we also used the fact that $(\Delta + k^2)u = 0$ in \mathbb{R}^n .

Using [13, Lemma 1.2 and equation (1.18)], we know that the Herglotz function $u = P_k f$ has the following asymptotics as $|x| \rightarrow \infty$:

$$u(x) = c'_{n,k} |x|^{-\frac{n-1}{2}} \left(e^{ik|x|} f(\hat{x}) + i^{n-1} e^{-ik|x|} f(-\hat{x}) \right) + O(|x|^{-\frac{n+1}{2}}), \tag{2.7a}$$

$$\partial_{|x|} u(x) = c'_{n,k} |x|^{-\frac{n-1}{2}} ik \left(e^{ik|x|} f(\hat{x}) - i^{n-1} e^{-ik|x|} f(-\hat{x}) \right) + O(|x|^{-\frac{n+1}{2}}), \tag{2.7b}$$

where $c'_{n,k} = k^{\frac{n-1}{2}} e^{\frac{\pi(n-1)i}{4}} (2\pi)^{-\frac{n-1}{2}}$. On the other hand, from [15, equation (2.27)], we know that w has the asymptotics

$$w(x) = c''_{n,k} |x|^{-\frac{n-1}{2}} e^{ik|x|} \hat{\mu}(k\hat{x}) + O(|x|^{-\frac{n+1}{2}}) \quad \text{as } |x| \rightarrow \infty, \tag{2.7c}$$

$$\partial_{|x|} w(x) = c''_{n,k} |x|^{-\frac{n-1}{2}} ik e^{ik|x|} \hat{\mu}(k\hat{x}) + O(|x|^{-\frac{n+1}{2}}) \quad \text{as } |x| \rightarrow \infty, \tag{2.7d}$$

where $c''_{n,k} = 2^{-1} e^{-\frac{\pi(n-3)i}{4}} (2\pi)^{-\frac{n-1}{2}} k^{\frac{n-3}{2}}$ and $\hat{\mu} \in C^\infty(\mathbb{R}^n)$ is the Fourier transform of the compactly supported distribution μ .

Combining (2.6) with (2.7a)–(2.7d), we obtain

$$\int_{\mathcal{S}^{n-1}} f(\hat{x}) \overline{\hat{\mu}(k\hat{x})} d\hat{x} = 0.$$

By the fact that $f \in C^\infty(\mathcal{S}^{n-1})$ was arbitrary, we conclude $\hat{\mu}(k\hat{x}) = 0$ for all $\hat{x} \in \mathcal{S}^{n-1}$. Consequently, (2.7c) becomes

$$w(x) = O(|x|^{-\frac{n+1}{2}}) \text{ as } |x| \rightarrow \infty.$$

In other words, the far-field pattern of w is vanishing. By the Rellich uniqueness theorem [4, 7], the unique continuation principle and the connectedness of $\mathbb{R}^n \setminus \overline{D}$, we conclude that

$$w = 0 \text{ in } \mathbb{R}^n \setminus \overline{D}.$$

Since $w \in H_{\text{loc}}^{2-s,p'}(\mathbb{R}^n)$, we also conclude that $w \in H_D^{2-s,p'}(\mathbb{R}^n)$.

Now let $v \in H^{s,p}(D)$ be any solution of $(\Delta + k^2)v = 0$ in D , and let $\tilde{v} \in H^{s,p}(\mathbb{R}^n)$ be such that $\tilde{v}|_D = v$. We see that

$$\ell(v) = \ell_1(\tilde{v}|_D) = (\tilde{v}, \mu) = (\tilde{v}, (\Delta + k^2)w).$$

From (2.2), we know that there are $w_j \in C_c^\infty(D)$ with $w_j \rightarrow w$ in $H^{2-s,p'}(\mathbb{R}^n)$. Since $(\Delta + k^2)\tilde{v} = 0$ in D , we finally conclude that

$$\ell(v) = \lim_{j \rightarrow \infty} (\tilde{v}, (\Delta + k^2)w_j) = \lim_{j \rightarrow \infty} ((\Delta + k^2)\tilde{v}, w_j) = 0,$$

which is our desired result. □

2.3 Proof of the main result

Theorem 1.1 is an immediate consequence of the following result:

Theorem 2.3 *Let D be a bounded Lipschitz domain in \mathbb{R}^n ($n \geq 2$) such that $\mathbb{R}^n \setminus \overline{D}$ is connected. Suppose that $k^2 > 0$ is not a Dirichlet eigenvalue of $-\Delta$ in D . Given any constant $c_0 \in \mathbb{R}$, there exist Herglotz wave functions $u_j \in C^\infty(\mathbb{R}^n)$ solving $(\Delta + k^2)u_j = 0$ in \mathbb{R}^n such that*

$$\lim_{j \rightarrow \infty} \|u_j - c_0\|_{L^\infty(\partial D)} = 0.$$

Before we prove Theorem 2.3 we need the following result, which is a special case of [8, Theorems 1.1 & 1.3].

Proposition 2.4 *Let D be a bounded Lipschitz domain in \mathbb{R}^n ($n \geq 2$). If $2 \leq p < \infty$ and $f \in H^{s-2,p}(D)$ where*

$$\frac{1}{p} < s < \frac{3}{p},$$

then there exists a unique $u \in H^{s,p}(D)$ satisfying $-\Delta u = f$ in D and $u = 0$ on ∂D .

Proof We first consider the case when $n \geq 3$. Let p_0 be as in [8, Theorem 1.1] (with $\Omega = D$). If $p'_0 \leq p < \infty$, the result follows from [8, Theorem 1.1(c)]. On the other hand, if $2 \leq p < p'_0$, the result follows from [8, Theorem 1.1(a)] since $s < \frac{3}{p} \leq 1 + \frac{1}{p}$. The case when $n = 2$ can be proved using identical reasoning using [8, Theorem 1.3] and the observation $\frac{3}{p} \leq \frac{2}{p} + \frac{1}{2}$. □

Proof of Theorem 2.3 Since k^2 is not a Dirichlet eigenvalue in D , there exists a unique solution $v \in H^{1,2}(D)$ such that

$$(\Delta + k^2)v = 0 \text{ in } D \text{ and } v = c_0 \text{ on } \partial D.$$

If $v \in H^{s,p}(D)$ for some $0 < s \leq 1$ and $p > n/s$, using Proposition 2.2, we know that there exist Herglotz waves $u_j \in C^\infty(\mathbb{R}^n)$ such that

$$\|u_j - c_0\|_{L^\infty(\partial D)} = \|u_j - v\|_{L^\infty(\partial D)} \leq \|u_j - v\|_{C(\bar{D})} \leq C \|u_j - v\|_{H^{s,p}(D)} \rightarrow 0,$$

where we used the Sobolev embedding.

It remains to show that $v \in H^{s,p}(D)$ for some s, p with $s > n/p$, and this follows from a standard bootstrap argument based on Proposition 2.4. We claim that

$$v \in H^{\frac{2}{p_j}, p_j}(D) \quad \text{for } 0 \leq j < \frac{n-2}{4}, \tag{2.5}$$

where

$$\frac{1}{p_j} = \frac{1}{2} - j \frac{2}{n-2}.$$

The case $j = 0$ follows since $v \in H^{1,2}(D)$. We argue by induction and assume that this holds for j . Define $w := v - c_0$ and note that w solves

$$-\Delta w = k^2 v \in H^{\frac{2}{p_j}, p_j}(D), \quad w|_{\partial D} = 0.$$

We next use the Sobolev embedding $H^{\frac{2}{p_j}, p_j}(D) \subset H^{\frac{2}{q}-2, q}(D)$ where $\frac{2}{p_j} > \frac{2}{q} - 2$ and

$$\frac{2}{p_j} - \frac{n}{p_j} = \frac{2}{q} - 2 - \frac{n}{q}.$$

It follows that $q = p_{j+1}$ and then indeed $\frac{2}{p_j} > \frac{2}{q} - 2$. In particular $-\Delta w \in H^{\frac{2}{p_{j+1}}-2, p_{j+1}}(D)$ with $w|_{\partial D} = 0$, and we may use Proposition 2.4 to conclude that $w \in H^{\frac{2}{p_{j+1}}, p_{j+1}}(D)$. This completes the induction step and proves (2.5).

We have proved that $v \in H^{\frac{2}{p_j}, p_j}(D)$ where j is the largest integer $< \frac{n-2}{4}$. Using the above notation, we have $\Delta w \in H^{\frac{2}{p_j}, p_j}(D)$ and $w|_{\partial D} = 0$. By Sobolev embedding we have $\Delta w \in H^{s-2, p}(D)$ whenever $p \geq p_j$ and

$$\frac{2}{p_j} - \frac{n}{p_j} = s - 2 - \frac{n}{p}.$$

The last condition implies that

$$s - \frac{n}{p} = 2 + \frac{2-n}{p_j} = 2 + \frac{2-n}{2} + 2j \geq 0$$

since $j \geq \frac{n-2}{4} - 1$. If $j > \frac{n-2}{4} - 1$, using Proposition 2.4 once again we obtain that w and hence v is in $H^{s,p}$ for some $s > n/p$. On the other hand, if $j = \frac{n-2}{4} - 1$ we iterate the argument once more to get $v \in H^{s,p}$ for some $s > n/p$. This concludes the proof. \square

The next simple example shows that the condition that k^2 is not an eigenvalue is necessary at least for balls.

Example 2.5 Let $v(x) := |x|^{\frac{2-n}{2}} J_{\frac{n-2}{2}}(|x|)$. We see that $v \in C^\infty(\mathbb{R}^n)$ and $(\Delta + 1)v = 0$ in \mathbb{R}^n . Suppose that u_1 is a real-valued function satisfying $(\Delta + 1)u_1 = 0$ in \mathbb{R}^n . Since

$$v(x) = 0 \quad \text{when } |x| = j_{\frac{n-2}{2}, m} \quad \text{for any } m \geq 1,$$

where $j_{\frac{n-2}{2},m}$ denotes the m th positive zero of $J_{\frac{n-2}{2}}$, we have

$$\int_{|x|=j_{\frac{n-2}{2},m}} u_1 \frac{\partial v}{\partial r} dS = \int_{|x|<j_{\frac{n-2}{2},m}} (u_1 \Delta v - v \Delta u_1) dx = 0.$$

Since

$$(-1)^m \frac{\partial v}{\partial r}(x) > 0 \quad \text{when } |x| = j_{\frac{n-2}{2},m},$$

it follows that u_1 must change sign on $|x| = j_{\frac{n-2}{2},m}$.

Similarly, if $R > 0$ and if u_0 solves $(\Delta + k_m^2)u_0 = 0$ in \mathbb{R}^n where $k_m = R^{-1}j_{\frac{n-2}{2},m}$, define u_1 via the rescaling

$$u_0(x) = u_1 \left(R^{-1} j_{\frac{n-2}{2},m} x \right) \quad \text{for } x \in \mathbb{R}^n.$$

We see that $(\Delta + 1)u_1 = 0$ in \mathbb{R}^n . The above discussion shows that u_0 must change sign on ∂B_R .

The following strong maximum principle can be found in [9, Appendix A]. However, for readers' convenience, here we exhibit the statement as well as its proof.

Lemma 2.6 (Strong maximum principle) *Let D be a bounded Lipschitz domain in \mathbb{R}^n ($n \geq 2$), and let $k^2 < \lambda_1(D)$, where $\lambda_1(D) > 0$ denotes the smallest $H_0^1(D)$ -eigenvalue of $-\Delta$. If the solution $u \in H^1(D)$ satisfies*

$$(\Delta + k^2)u = 0 \text{ in } D, \quad u \geq 0 \text{ on } \partial D,$$

then for each open component G of D we have either $u \equiv 0$ in G or $u > 0$ in G (note that $u \in C^\infty(G)$ by elliptic regularity).

Proof It is easy to see that for each component G of D we have $k^2 < \lambda_1(G)$ and

$$(\Delta + k^2)u = 0 \text{ in } G, \quad u \geq 0 \text{ on } \partial G.$$

Testing the equation above by $u_- \in H_0^1(G)$ and using Poincaré inequality, we have

$$\int_G |u_-|^2 dx \leq \frac{1}{\lambda_1(G)} \int_G |\nabla u_-|^2 dx = \frac{k^2}{\lambda_1(G)} \int_G |u_-|^2 dx.$$

Since $\frac{k^2}{\lambda_1(G)} < 1$, then $u_- \equiv 0$ in G , that is,

$$u \geq 0 \text{ in } G. \tag{2.6}$$

Let $x_0 \in G$ such that $u(x_0) = 0$. The mean value theorem for Helmholtz equation (see e.g. [9, Appendix A]) gives that

$$\int_{B_\varepsilon(x_0)} u(x) dx = 0 \tag{2.7}$$

for all sufficiently small $\varepsilon > 0$ so that $\overline{B_\varepsilon(x_0)} \subset G$. Since u is continuous in G , combining (2.6) and (2.7) we know that $u = 0$ in $B_\varepsilon(x_0)$, and this shows that $\{x \in G \mid u(x) = 0\}$ is both open and closed in G . Since G is connected, then we have either

$$\{x \in G \mid u(x) = 0\} = G \quad \text{or} \quad \{x \in G \mid u(x) = 0\} = \emptyset,$$

which concludes our lemma. □

Finally, we give the proof of Theorem 1.2.

Proof of Theorem 1.2 Since $|D| \leq |B_r|$ where $r = j_{\frac{n-2}{2},1} k^{-1}$, the Faber–Krahn inequality (see e.g. [3, Theorem III.3.1]) implies that each connected component G of D satisfies

$$\lambda_1(G) \geq \lambda_1(B_r) = k^2.$$

Case 1. If $\lambda_1(G) = k^2$, we choose v to be an eigenfunction corresponding to the first eigenvalue with $v > 0$ in G , i.e. v solves $(\Delta + k^2)v = 0$ in G with $v \in H_0^1(G)$, see e.g. [6, Theorem 2(ii) in Section 6.5.1].

Case 2. If $\lambda_1(G) > k^2$, then there exists a unique solution $v \in H^1(G)$ such that

$$(\Delta + k^2)v = 0 \text{ in } G, \quad v = 1 \text{ on } \partial G.$$

Using the strong maximum principle in Lemma 2.6, we know that $v > 0$ in G .

Next we choose a bounded Lipschitz domain D_1 that satisfies $E \subset D_1$, $\overline{D_1} \subset D$, and $\mathbb{R}^n \setminus \overline{D_1}$ is connected. The function $v|_{D_1}$ is in $H^{1,p}(D_1)$ for any $p > n$ and satisfies $v|_{\overline{D_1}} > 0$. The approximation result in Proposition 2.2 yields a sequence of Herglotz waves u_j satisfying

$$\|u_j|_{D_1} - v\|_{H^{1,p}(D_1)} \rightarrow 0 \quad \text{as } j \rightarrow \infty.$$

If j is sufficiently large, the Sobolev embedding ensures that $u_j|_E > 0$. □

Acknowledgements This project was finalized while the authors stayed at Institute Mittag Leffler (Sweden), during the program Geometric aspects of nonlinear PDE. Kow and Salo were partly supported by the Academy of Finland (Centre of Excellence in Inverse Modelling and Imaging, 312121) and by the European Research Council under Horizon 2020 (ERC CoG 770924). Shahgholian was supported by Swedish Research Council.

Funding Open access funding provided by Royal Institute of Technology.

Data Availability All data needed are contained in the manuscript.

Declaration

Conflicts of interest The authors declare that there are no competing or conflict of interests.

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References

1. Bergh, J., Löfström, J.: Interpolation Spaces. An Introduction. Grundlehren der Mathematischen Wissenschaften, vol. 223. Springer-Verlag, Berlin, Heidelberg (1976)
2. Cakoni, F., Vogelius, M.S.: Singularities almost always scatter: Regularity results for non-scattering inhomogeneities. [arXiv:2104.05058](https://arxiv.org/abs/2104.05058) (2021)
3. Chavel, I.: Isoperimetric Inequalities. Differential Geometric and Analytic Perspectives. Cambridge Tracts in Mathematics, vol. 145. Cambridge University Press, Cambridge (2001)
4. Colton, D., Kress, R.: Inverse Acoustic and Electromagnetic Scattering Theory, 4th edn. Applied Mathematical Sciences, vol. 93. Springer, Cham (2019)

5. Cronwell, R.H., Fox, R.H.: Introduction to Knot Theory. Graduate Text in Mathematics, vol. 57. Springer-Verlag, New York, Heidelberg (1977)
6. Evans, L.C.: Partial Differential Equations, 2nd edn. Graduate Studies in Mathematics, vol. 19. American Mathematical Society, Providence, RI (2010)
7. Hörmander, L.: Lower bounds at infinity for solutions of differential equations with constant coefficients. *Isr. J. Math.* **16**, 103–116 (1973)
8. Jerison, D., Kenig, C.E.: The inhomogeneous Dirichlet problem in Lipschitz domains. *J. Funct. Anal.* **130**, 161–219 (1995)
9. Kow, P.-Z., Larson, S., Salo, M., Shahgholian, H.: Quadrature domains for the Helmholtz equation with applications to non-scattering phenomena. *Potential Anal.* (2022). <https://doi.org/10.1017/s11118-022-10054-5>
10. Lee, J.M.: Introduction to Smooth Manifolds, 2nd edn. Graduate Texts in Mathematics, vol. 218. Springer, New York (2013)
11. Logunov, A., Malinnikova, E.: Review of Yau’s conjecture on zero sets of Laplace eigenfunctions. In: *Current Developments in Mathematics*, vol. 2018, pp. 179–212. Int. Press, Somerville, MA (2020)
12. Madsen, I., Tornehave, J.: From Calculus to Cohomology: de Rham Cohomology and Characteristic Classes. Cambridge University Press, Cambridge (1997)
13. Melrose, R.B.: Geometric Scattering Theory. Stanford Lectures. Cambridge University Press, Cambridge (1995)
14. Salo, M., Shahgholian, H.: Free boundary methods and non-scattering phenomena. *Res. Math. Sci.* **8**, 58 (2021)
15. Yafaev, D.R.: Mathematical Scattering Theory. Analytic Theory. Mathematical Surveys and Monographs, vol. 158. American Mathematical Society, Providence, RI (2010)

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