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FIXED ANGLE INVERSE SCATTERING FOR SOUND SPEEDS CLOSE TO CONSTANT

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ABSTRACT. We study the fixed angle inverse scattering problem of determining a sound speed from scattering measurements corresponding to a single incident wave. The main result shows that a sound speed close to constant can be stably determined by just one measurement. Our method is based on studying the linearized problem, which turns out to be related to the acoustic problem in photoacoustic imaging. We adapt the modified time-reversal method from [P. Stefanov and G. Uhlmann, *Thermoacoustic tomography with variable sound speed*, *Inverse Problems* **25** (2009), 075011] to solve the linearized problem in a stable way, and use this to give a local uniqueness result for the nonlinear inverse problem.

Keywords: Inverse scattering, wave equation, time-reversal, progressive wave expansion.

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1. INTRODUCTION AND MAIN RESULT

Inverse scattering problems appear in real-life phenomena and have applications in a wide range of fields such as radar, sonar, fault detection in fiber optics, geophysical exploration, medical imaging and nondestructive testing. In this work, we study the inverse acoustic scattering problem of recovering a sound speed from fixed angle scattering measurements. Let $n \geq 2$ and $\theta \in S^{n-1}$ be a fixed vector. Consider $\eta \in H^{s_0}(\mathbb{R}^n)$ with $s_0 > 0$ being large enough and so that $1 - \eta$ is compactly supported in Ω with $\Omega := \{x \in \mathbb{R}^n : |x| < 1\}$, i.e.

$$\eta(x) = 1, \quad |x| \geq 1 - \sigma, \tag{1.1}$$

for a fixed number $\sigma \in (0, 1)$. We also assume that for some $M > 1$

$$M^{-1} \leq \eta(x) \leq M, \quad x \in \mathbb{R}^n. \tag{1.2}$$

Now consider a plane wave solution \tilde{U} to the wave equation

$$(\eta(x) \partial_t^2 - \Delta) \tilde{U} = 0 \text{ in } \mathbb{R}^{n+1}, \quad \tilde{U}|_{\{t < -1\}} = \delta(t - x \cdot \theta). \tag{1.3}$$

This model describes the propagation of scattered sound waves in an inhomogeneous medium (whose properties are described by the coefficient η) produced by the interaction of an incident plane wave of the form $\delta(t - x \cdot \theta)$ with the medium. If $c(x)$ is the sound speed in the medium then $\eta(x) = c(x)^{-2}$, but for simplicity we will refer to η as the sound speed.

The aim of this paper is to prove that the sound speed η is uniquely determined by boundary measurements of a solution to (1.3) corresponding to a fixed direction $\theta \in S^{n-1}$. From the mathematical viewpoint, for $T > 1$, the measurements are encoded by the map

$$\tilde{\mathcal{A}} : \eta \mapsto \tilde{U}|_{\Sigma_T}, \tag{1.4}$$

where \tilde{U} is the solution to (1.3) and Σ_T stands for the lateral boundary of the space-time cylinder $\Omega \times (-T, T)$, that is

$$\Sigma_T := \partial\Omega \times (-T, T).$$

In Proposition 1.2 we prove that there is a unique solution \tilde{U} to (1.3) in a suitable Hilbert space so that the restriction on Σ_T , denoted by $\tilde{U}|_{\Sigma_T}$, is well defined.

In this framework, our first main result states that if $\tilde{\mathcal{A}}(\eta) = \tilde{\mathcal{A}}(1)$ where η satisfies (1.1)–(1.2) and is close to 1, then $\eta = 1$. Moreover, we also derive the corresponding quantification (stability estimate) with a modulus of continuity of Hölder type.

Theorem 1.1. *Let $n \geq 2$, $M > 1$ and $\theta \in S^{n-1}$. Fix some $s_0 \gg n/2 + 2$. There exist a small $\varrho > 0$, $\mu = \mu(n) \in (0, 1)$, $T = T(n) > 1$, $s_1 > s_0$, and $C > 0$ such that*

$$\|\eta - 1\|_{H^{s_0}(\mathbb{R}^n)} \leq C \|\tilde{\mathcal{A}}(\eta) - \tilde{\mathcal{A}}(1)\|_{H^{-5}((-T, T); H^{1/2}(S^{n-1}))|_{\{t > z\}}}^\mu$$

for all $\eta \in H^{s_0}(\mathbb{R}^n)$ satisfying (1.1)–(1.2) and $\|\eta - 1\|_{H^{s_0}(\mathbb{R}^n)} \leq \varrho$, $\|\eta\|_{H^{s_1}(\mathbb{R}^n)} \leq M$.

Theorem 1.1 is an instance of a fixed angle inverse scattering result, where one determines a sound speed from scattering measurements corresponding to a single incident plane wave. There are several results of this type for determining a time-independent potential q instead of a sound speed. The equation in this case is

$$(\partial_t^2 - \Delta + q(x))\mathcal{W} = 0 \text{ in } \mathbb{R}^{n+1}, \quad \mathcal{W}|_{\{t \ll 0\}} = \delta(t - x \cdot \theta). \quad (1.5)$$

One can alternatively work on the frequency side with the Schrödinger equation

$$\begin{cases} (-\Delta - k^2 + q(x))u = 0 \text{ in } \mathbb{R}^n, \\ u \text{ is outgoing.} \end{cases}$$

The fixed angle scattering problem consists of determining q from the knowledge of $\mathcal{W}|_{\Sigma_T}$, or equivalently from the scattering amplitude $a_q(k, \theta, \omega)$ for all $k > 0$ and $\omega \in S^{n-1}$. This equivalence is discussed in detail in [RS20a].

There are several known results related to recovering small or generic potentials and singularities from fixed angle measurements [BCL⁺20, BLM89, Mer18, Rui01, Ste92]. In the recent works [RS20a, RS20b] it was shown that a potential $q \in C_c^\infty(\mathbb{R}^n)$ is uniquely determined by measurements corresponding to two incident plane waves from opposite directions $\theta = \pm\theta_0$, or just a single incident plane wave if q satisfies some symmetry conditions. This result was extended in [MS21] to the case when $-\Delta$ is replaced by the Laplace-Beltrami operator $-\Delta_g$ in (1.5), with g being a known metric satisfying certain symmetry conditions. Using similar ideas, in [MPMS21] the authors proved analogous results in the case of time-independent first order coefficients. We also mention the recent work [KRS20], which studies fixed angle scattering for time-dependent coefficients also in the case of first-order perturbations. However, the problem of determining a general potential $q \in C_c^\infty(\mathbb{R}^n)$ from fixed angle measurements remains open and so does the corresponding inverse backscattering problem (see [RU14] for more information).

The purpose of the present article is to study the fixed angle problem for determining a sound speed η instead of a potential q . The method in [RS20b, RS20a], which is based on Carleman estimates and reflection arguments, requires symmetry and appears to break down for most nonconstant sound speeds. In this work we approach the problem for sound speeds by studying the linearized problem. The main observation is that the linearization of the fixed angle inverse problem for a sound speed has similar features as the acoustic problem in thermo/photoacoustic tomography [Ba1, SU13]. We then adapt the modified time-reversal method introduced in [SU09b] to our case and establish uniqueness, stability and reconstruction for the linearized problem at a constant sound speed. This can be used to prove local uniqueness and stability for the nonlinear inverse problem as in [SU09a], leading to Theorem 1.1.

Remark 1.1. It is possible that the methods in this work can be extended to deal with the linearized problem at a general sound speed, and to obtain a counterpart of Theorem 1.1 showing that $\|\eta_1 - \eta_2\| \leq C \|\tilde{\mathcal{A}}(\eta_1) - \tilde{\mathcal{A}}(\eta_2)\|^\mu$ in suitable norms when both η_1 and η_2

are close to some fixed nonconstant sound speed. These questions are more involved and will be left to a future work.

We now describe our method in more detail. Throughout this work we will assume that $\theta = e_n$, the n th vector of the standard basis in \mathbb{R}^n . If $x \in \mathbb{R}^n$ we write $x = (y, z)$ where $y \in \mathbb{R}^{n-1}$ and $z \in \mathbb{R}$. Thus $x \cdot \theta = z$. As mentioned above, the proof of Theorem 1.1 reduces to studying the injectivity and stability properties of the linearization of the map $\tilde{\mathcal{A}}$ at a constant sound speed and to using a general result in [SU09a]. Due to technical reasons, we consider a smoother initial value than the one in (1.3) and study instead the equation

$$(\eta(x) \partial_t^2 - \Delta) \mathcal{U} = 0 \text{ in } \mathbb{R}^{n+1}, \quad \mathcal{U}|_{\{t < -1\}} = H_1(t - z), \quad (1.6)$$

where $H_1(s) = s$ when $s \geq 0$ and zero otherwise. We now consider the fixed angle inverse scattering problem associated with (1.6), where the measurement operator is given by

$$\mathcal{A}(\eta) = \mathcal{U}|_{\Sigma_T}. \quad (1.7)$$

Since η is time-independent, one has $\tilde{\mathcal{U}} = \partial_t^2 \mathcal{U}$ and thus the fixed angle inverse scattering problems for \mathcal{A} and $\tilde{\mathcal{A}}$ are equivalent. The following result describes the precise function spaces involved (its proof is presented in Appendix A).

Proposition 1.2. *Let $n \geq 2$, $T > 1$ and $M > 1$. Fix some $s_0 > n/2 + 2$. Let $\eta \in H^{s_0}(\mathbb{R}^n)$ be a function satisfying*

$$\|\eta\|_{H^{s_0}(\mathbb{R}^n)} \leq M, \quad M^{-1} \leq \eta(x) \leq M \quad \text{a.e in } \mathbb{R}^n.$$

There exists a unique solution $\mathcal{U} \in H^{-1}((-T, T); H^1(\mathbb{R}^n))$ to (1.6) with

$$\|\mathcal{U}\|_{H^{-1}((-T, T); H^1(\mathbb{R}^n))} \leq C(n, T, M).$$

Moreover, $\tilde{\mathcal{U}} = \partial_t^2 \mathcal{U} \in H^{-3}((-T, T); H^1(\mathbb{R}^n))$ is the unique solution to (1.3), and hence it satisfies

$$\begin{aligned} \|\tilde{\mathcal{U}}|_{\Sigma_T}\|_{H^{-3}((-T, T); H^{1/2}(S^{n-1}))} &\lesssim \|\tilde{\mathcal{U}}\|_{H^{-3}((-T, T); H^1(\mathbb{R}^n))} \\ &\lesssim \|\mathcal{U}\|_{H^{-1}((-T, T); H^1(\mathbb{R}^n))} \leq C(n, T, M). \end{aligned}$$

Since our approach involves a linearization argument, we shall first study the Fréchet derivative of \mathcal{A} at the constant 1, denoted by A_1 . It is proved in Section 2 that A_1 is given by

$$A_1(f) = U|_{\Sigma_T},$$

where U solves

$$(\partial_t^2 - \Delta)U = -f(x)\delta(t - z) \text{ in } \mathbb{R}^{n+1}, \quad U|_{\{t < -1\}} = 0. \quad (1.8)$$

By employing the progressive wave expansion method, one can further show that any solution to (1.8) can be written as

$$U(t, x) = u(t, x)H(t - z),$$

where H stands for the Heaviside function and u is a C^2 function in the set $\{t \geq z\}$ solving the equation

$$\begin{cases} (\partial_t^2 - \Delta)u = 0 & \text{in } \{t > z\}, \\ (\partial_t + \partial_z)u = -\frac{1}{2}f & \text{on } \{t = z\}. \end{cases} \quad (1.9)$$

Thus the linearized inverse problem associated with A_1 amounts to determining the initial value $f(x)$ from the knowledge of $u|_{\Sigma_T \cap \{t > z\}}$.

It is worth mentioning that the linearized problem above is similar to the acoustic problem in thermo/photoacoustic tomography. There one needs to recover an initial condition $g(x)$ from the boundary measurement $v|_{\partial\Omega \times (0, T)}$, where v solves

$$\begin{cases} (\partial_t^2 - \Delta)v = 0 & \text{in } \mathbb{R}^n \times (0, T), \\ v = g & \text{on } \{t = 0\}, \\ \partial_t v = 0 & \text{on } \{t = 0\}. \end{cases}$$

This problem was studied in detail in [SU09b], also for general sound speeds, by using a modification of the time-reversal method. Error estimates, reconstruction formulae and numerical implementations in case of non-trapping sound speeds were obtained earlier in [HKN08] and [Hri09] by using the time-reversal method. For more details on thermo/photoacoustic tomography inverse problems we refer the readers to [KK08, KK11, SU13, AKK17] and the references therein. We will adapt the modified time-reversal method to show uniqueness, stability and reconstruction for the inverse problem associated with (1.9). The main difference with previous results is that the initial data is given on the characteristic set $\{t = z\}$ instead of the standard set $\{t = 0\}$. This creates various difficulties, and in order to overcome these we employ energy estimates in space-time domains that are adapted to this characteristic geometry.

The following result is the precise statement of uniqueness and stability in the linearized inverse problem with respect to suitable norms. We refer to Proposition 4.6 for a reconstruction formula involving a Neumann series. Denote $\Gamma := (\Omega \times (-T, T)) \cap \{t = z\}$.

Proposition 1.3. *Let $s_0 \gg n/2 + 2$, $M > 1$ and $T > 1$. There exists $C > 0$ so that*

$$\|f\|_{L^2(\Omega)} \leq C(\|A_1 f\|_{H^1(\Sigma_T \cap \{t > z\})} + \frac{T^{1/2}}{\sigma^{1/4}} \|A_1 f\|_{H^1(\Sigma_T \cap \Gamma)}),$$

for $f \in H_{\Omega}^{s_0}(\mathbb{R}^n)$ with $\text{supp } f \in \{x \in \mathbb{R}^n; |x| \leq 1 - \sigma\}$.

We outline the method for proving Proposition 1.3. Instead of A_1 , it is convenient to work with the closely related operator $A'_1 : F \mapsto u|_{\Sigma}$ where u solves

$$\begin{cases} (\partial_t^2 - \Delta)u = 0 & \text{in } \{t > z\}, \\ u = F & \text{on } \{t = z\}, \end{cases} \quad (1.10)$$

and $\Sigma := (\partial\Omega \times (-T, T)) \cap \{t > z\}$. Recall $\Gamma = (\Omega \times (-T, T)) \cap \{t = z\}$. Define the space

$$\mathcal{H} := \{F \in H^1(\{t = z\}) : \text{supp}(F) \subset \bar{\Gamma}\}$$

equipped with the H^1 -norm. In Section 3.3 we actually use a slightly different definition of \mathcal{H} , see (3.26) and (3.27), but here for illustration purpose we prefer to keep the definition simple. One would like to think of A'_1 as a bounded operator $\mathcal{H} \rightarrow H^1(\Sigma)$. In fact this holds in the standard case where the initial surface is $\{t = 0\}$ instead of $\{t = z\}$ since the trace of u is in H^1 [BS91, FR05]. We are not aware of such a result for our slanted case, so we will work with smooth functions instead and use norm estimates with uniform bounds in the energy spaces.

After some natural derivations, uniqueness and stability for A_1 reduces to uniqueness and stability for A'_1 . Now we use a time-reversal method as in [SU09b] and define an approximate inverse B for A'_1 as the map $B : h \mapsto v|_{\Gamma}$, where v solves the Dirichlet problem

$$\begin{cases} (\partial_t^2 - \Delta)v = 0 & \text{in } Q_T \cap \{t > z\}, \\ v = h & \text{on } \Sigma, \\ v = \phi_0, v_t = 0 & \text{in } Q_T \cap \{t = T\}, \\ \Delta\phi_0 = 0 & \text{in } \Omega, \quad \phi_0 = h(\cdot, T) \text{ on } \partial\Omega, \end{cases} \quad (1.11)$$

with $Q_T = \Omega \times [-T, T]$. We prove that B is a bounded operator $H^1(\Sigma) \rightarrow H^1(\Gamma)$. Note that in odd dimensions, the sharp Huygens' principle implies that any solution u of (1.10) with $F \in \mathcal{H}$ satisfies $u(\cdot, T)|_{\bar{\Omega}} = 0$ for T large enough. Hence letting $h = u|_{\Sigma}$ in (1.11) implies that $v = u$ and hence $BA'_1 F = F$, so that B is an exact inverse of A'_1 . In even dimensions this is no longer true. Instead, we will show that B is almost an inverse of A'_1 in the following parametrix sense

$$BA'_1 F = F - \tilde{K}F$$

where \tilde{K} is a bounded operator on \mathcal{H} with norm strictly less than 1 when T is large enough. In [SU09b] this was done by using a unique continuation property for the wave equation, but in our case due to a technical issue we use local energy decay instead. This argument allows us to invert A'_1 by a Neumann series and prove Proposition 1.3.

This paper is structured as follows. Section 2 is dedicated to proving Theorem 1.1 under the assumption that Proposition 1.3 related to linearized problem is known. Then the proof is reduced to finding suitable function spaces so that [SU09a, Theorem 2] can be applied. In Section 3 we study the basic properties of the linear map A_1 . In Section 4 we prove Proposition 1.3 as a consequence of a reconstruction formula stated in Proposition 4.6. In the Appendix we give some results on the well-posedness of the forward problem for wave equations in negative Sobolev spaces that are required in our arguments. The results are stated with finite regularity assumptions on the coefficients, and the dependence of the constants in norm estimates on different quantities is explicitly specified.

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2. THE NONLINEAR MAP. PROOF OF THEOREM 1.1

In this section we will prove Theorem 1.1 by using Propositions 1.2, 1.3 and 3.1 as well as Lemma A.5. The proofs of these results will be given in later sections. We shall use the following abstract local uniqueness and stability result from [SU09a, Theorem 2].

Proposition 2.1. *Let $\mathcal{X}_j, \mathcal{Y}_j$ with $j = 1, 2, 3$ be Banach spaces with $\mathcal{X}_3 \subset \mathcal{X}_1 \subset \mathcal{X}_2$ and $\mathcal{Y}_3 \subset \mathcal{Y}_2 \subset \mathcal{Y}_1$, such that the following interpolation estimates hold:*

$$\|f\|_{\mathcal{X}_1} \lesssim \|f\|_{\mathcal{X}_2}^{\mu_1} \|f\|_{\mathcal{X}_3}^{1-\mu_1}, \quad \|g\|_{\mathcal{Y}_2} \lesssim \|g\|_{\mathcal{Y}_1}^{\mu_2} \|g\|_{\mathcal{Y}_3}^{1-\mu_2}, \quad \mu_1, \mu_2 \in (0, 1], \quad \mu_1 \mu_2 > 1/2. \quad (2.1)$$

Let $\mathcal{A} : \mathcal{V}_1 \rightarrow \mathcal{Y}_1$ be a nonlinear map where $\mathcal{V}_1 \subset \mathcal{X}_1$ is an open subset of \mathcal{X}_1 . Consider $f_0 \in \mathcal{V}_1$ and assume that

$$\mathcal{A}(f) = \mathcal{A}(f_0) + A_{f_0}(f - f_0) + R_{f_0}(f), \quad \|R_{f_0}(f)\|_{\mathcal{Y}_1} \leq C(f_0) \|f - f_0\|_{\mathcal{X}_1}^2 \quad (2.2)$$

holds for all f in some neighbourhood of f_0 in \mathcal{V}_1 . Here A_{f_0} stands for the Fréchet derivative of \mathcal{A} at f_0 . In addition, suppose that

$$\|h\|_{\mathcal{X}_2} \leq C \|A_{f_0} h\|_{\mathcal{Y}_2}, \quad h \in \mathcal{X}_1. \quad (2.3)$$

Then for any $L > 0$ there exists $\epsilon > 0$, so that for any f with

$$\|f - f_0\|_{\mathcal{X}_1} \leq \epsilon, \quad \|f\|_{\mathcal{X}_3} \leq L, \quad (2.4)$$

one has the conditional stability estimate

$$\|f - f_0\|_{\mathcal{X}_1} \leq CL^{2-\mu_1-\mu_2} \|\mathcal{A}(f) - \mathcal{A}(f_0)\|_{\mathcal{Y}_1}^{\mu_1 \mu_2}.$$

In particular, if $\mathcal{A}(f) = \mathcal{A}(f_0)$ for some f satisfying (2.4), then $f = f_0$.

The task now is verifying that \mathcal{A} and A_1 from Section 1 satisfy all conditions of Proposition 2.1 in appropriate Banach spaces. In fact we will change the setup slightly and write $\eta = 1 + f$, and \mathcal{A} will be considered as the operator $f \mapsto \mathcal{U}|_{\Sigma_T \cap \{t > z\}}$ (note that $\mathcal{U} = 0$ when $t < z$). Similarly, instead of A_1 we consider the operator $A_0 : f \mapsto U|_{\Sigma_T \cap \{t > z\}}$. In this setting, we will apply Proposition 2.1 with $f_0 \equiv 0$.

We first introduce some useful notation. If $F \subset \mathbb{R}^n$ is closed, define for any $s_0 \in \mathbb{R}$ the set

$$H_F^{s_0}(\mathbb{R}^n) = \{f \in H^{s_0}(\mathbb{R}^n) : \text{supp}(f) \subset F\}.$$

We also write

$$\Omega_\sigma = \{x \in \mathbb{R}^n : |x| < 1 - \sigma\}$$

where $\sigma \in (0, 1)$ is fixed as in (1.1). Let $s_0 > n/2 + 2$ and $M > 1$. The map \mathcal{A} will be defined in the open subset

$$\mathcal{V}_M^{s_0}(\mathbb{R}^n) = \{f \in H_{\Omega_\sigma}^{s_0}(\mathbb{R}^n) : M^{-1} < 1 + f < M\}$$

of $H_{\Omega_\sigma}^{s_0}(\mathbb{R}^n)$ as the map

$$\mathcal{A} : \mathcal{V}_M^{s_0}(\mathbb{R}^n) \rightarrow H^{-1}((-T, T); H^{1/2}(S^{n-1}))|_{\{t > z\}}, \quad \mathcal{A}(f) = \mathcal{U}|_{\Sigma_T \cap \{t > z\}} \quad (2.5)$$

where $\eta = 1 + f$. This is a well defined map by Proposition 1.2. We will prove in Lemma 2.2 that the linearization of \mathcal{A} at 0 is given by

$$A_0 : H_{\Omega_\sigma}^{s_0}(\mathbb{R}^n) \rightarrow H^{-1}((-T, T); H^{1/2}(S^{n-1}))|_{\{t > z\}}, \quad A_0 f = U|_{\Sigma_T \cap \{t > z\}} \quad (2.6)$$

where U solves

$$\begin{cases} (\partial_t^2 - \Delta)U = -f(x) \delta(t - z) & \text{in } \mathbb{R}^{n+1}, \\ U|_{\{t < -1\}} = 0. \end{cases} \quad (2.7)$$

Let us verify one by one the conditions of Proposition 2.1.

Condition (2.2). We claim that the map \mathcal{A} , defined by (2.5), verifies the condition (2.2) with

$$\mathcal{X}_1 = H_{\Omega_\sigma}^{s_0}(\mathbb{R}^n), \quad \mathcal{Y}_1 = H^{-3}((-T, T); H^{1/2}(S^{n-1}))|_{\{t > z\}}.$$

Indeed, the proof is contained in the following result, where we write

$$\begin{aligned} \hat{\mathcal{A}} : \mathcal{V}_M^{s_0}(\mathbb{R}^n) &\rightarrow H^{-1}((-T, T); H^{1/2}(S^{n-1})), \quad \hat{\mathcal{A}}(f) = \mathcal{U}|_{\Sigma_T}, \\ \hat{A}_0 : H_{\Omega_\sigma}^{s_0}(\mathbb{R}^n) &\rightarrow H^{-1}((-T, T); H^{1/2}(S^{n-1})), \quad \hat{A}_0 f = U|_{\Sigma_T}. \end{aligned}$$

Thus \mathcal{A} and A_0 are the restrictions $\mathcal{A}(f) = \hat{\mathcal{A}}(f)|_{\{t > z\}}$ and $A_0 f = \hat{A}_0 f|_{\{t > z\}}$.

Lemma 2.2. *Let $s_0 \gg n/2 + 2$, $M > 1$ and $T > 1$. The map $\hat{\mathcal{A}}$ is well defined*

$$\hat{\mathcal{A}} : \mathcal{V}_M^{s_0}(\mathbb{R}^n) \subset H_{\Omega_\sigma}^{s_0}(\mathbb{R}^n) \rightarrow H^{-1}((-T, T); H^{1/2}(S^{n-1})).$$

Moreover, it is $C^{1,1}$ near 0 as a map $\mathcal{V}_M^{s_0}(\mathbb{R}^n) \rightarrow H^{-3}((-T, T); H^{1/2}(S^{n-1}))$, so that

$$\|\hat{\mathcal{A}}(f) - \hat{\mathcal{A}}(0) - \hat{A}_0(f)\|_{H^{-3}((-T, T); H^{1/2}(S^{n-1}))} \lesssim \|f\|_{H^{s_0}(\mathbb{R}^n)}^2, \quad (2.8)$$

for all $f \in \mathcal{V}_M^{s_0}(\mathbb{R}^n)$ near the origin.

Proof. The fact that $\hat{\mathcal{A}}$ maps $\mathcal{V}_M^{s_0}(\mathbb{R}^n)$ to $H^{-1}((-T, T); H^{1/2}(S^{n-1}))$ is an immediate consequence of Proposition 1.2. Let us move to prove the $C^{1,1}$ regularity near 0. Fix $\eta = 1 + f$ with $f \in \mathcal{V}_M^{s_0}(\mathbb{R}^n)$ near the origin. Let $\hat{A}_0 f = U|_{\Sigma_T}$ where U is as in (2.7), and let \mathcal{R} be defined by

$$\hat{\mathcal{A}}(f) = \hat{\mathcal{A}}(0) + \hat{A}_0(f) + \mathcal{R}|_{\Sigma_T}, \quad \mathcal{R}|_{\Sigma_T} := \hat{\mathcal{A}}(f) - \hat{\mathcal{A}}(0) - \hat{A}_0(f).$$

In order to prove (2.8), we will actually prove a stronger estimate

$$\|\mathcal{R}|_{\Sigma_T}\|_{H^{-3}((-T,T);H^{1/2}(S^{n-1}))} \lesssim \|f\|_{L^\infty(\mathbb{R}^n)}^2, \quad f \in H_{\Omega}^{s_0}(\mathbb{R}^n),$$

which trivially implies the required estimate by using Morrey's inequality since $s_0 > n/2$. To do that, we first set

$$\hat{A}(f) = \mathcal{U}|_{\Sigma_T}, \quad \hat{A}(0) = \mathcal{U}_0|_{\Sigma_T}, \quad \hat{A}_0(f) = U|_{\Sigma_T},$$

where the distributions \mathcal{U} , \mathcal{U}_0 and U satisfy

$$\begin{aligned} (\eta(x)\partial_t^2 - \Delta)\mathcal{U} &= 0, & \text{in } \mathbb{R}^{n+1}, \quad \mathcal{U}|_{\{t < -1\}} &= H_1(t-z), \\ (\partial_t^2 - \Delta)\mathcal{U}_0 &= 0, & \text{in } \mathbb{R}^{n+1}, \quad \mathcal{U}_0|_{\{t < -1\}} &= H_1(t-z), \\ (\partial_t^2 - \Delta)U &= -f\delta(t-z), & \text{in } \mathbb{R}^{n+1}, \quad U|_{\{t < -1\}} &= 0. \end{aligned}$$

By Proposition 1.2, we deduce that $\mathcal{U} \in H^{-1}((-T, T); H^1(\mathbb{R}^n))$. By uniqueness of distributional solutions, we have $\mathcal{U}_0(y, z, t) = H_1(t-z)$. Furthermore, Proposition 3.1 ensures that $U \in H^{-1}((-T, T); L^2(\mathbb{R}^n))$. On the other hand, a straightforward computation shows that $\mathcal{R} := \mathcal{U} - \mathcal{U}_0 - U$ satisfies in \mathbb{R}^{n+1}

$$(\partial_t^2 - \Delta)\mathcal{R} = -f\partial_t^2(\mathcal{U} - H_1), \quad \mathcal{R}|_{\{t < -1\}} = 0.$$

In addition, we also have in \mathbb{R}^{n+1}

$$(\partial_t^2 - \Delta)(\mathcal{U} - H_1) = -f\partial_t^2\mathcal{U}, \quad \mathcal{U} - H_1|_{\{t < -1\}} = 0. \quad (2.9)$$

Note that the sources on the right of above equations belong to $H^{-3}((-T, T); L^2(\mathbb{R}^n))$. Since $f = f(x)$ is independent of t , for all $\alpha \geq 0$ and $F \in H^{-\alpha}((-T, T); L^2(\mathbb{R}^n))$ one has

$$\|fF\|_{H^{-\alpha}((-T, T); L^2(\mathbb{R}^n))} \lesssim \|f\|_{L^\infty(\mathbb{R}^n)} \|F\|_{H^{-\alpha}((-T, T); L^2(\mathbb{R}^n))}.$$

These facts combined with Lemma A.4 give that

$$\|\mathcal{R}\|_{H^{-3}((-T, T); H^1(\mathbb{R}^n))} \lesssim \|f\|_{L^\infty(\mathbb{R}^n)} \|\partial_t^2(\mathcal{U} - H_1)\|_{H^{-3}((-T, T); L^2(\mathbb{R}^n))}. \quad (2.10)$$

We bound the norm on the left with the help of (2.9) and Lemma A.4 as follows

$$\begin{aligned} \|\partial_t^2(\mathcal{U} - H_1)\|_{H^{-3}((-T, T); L^2(\mathbb{R}^n))} &\leq \|\partial_t^2(\mathcal{U} - H_1)\|_{H^{-3}((-T, T); H^1(\mathbb{R}^n))} \\ &\leq \|\mathcal{U} - H_1\|_{H^{-1}((-T, T); H^1(\mathbb{R}^n))} \\ &\lesssim \|f\partial_t^2\mathcal{U}\|_{H^{-1}((-T, T); L^2(\mathbb{R}^n))} \\ &\lesssim \|f\|_{L^\infty(\mathbb{R}^n)} \|\partial_t^2\mathcal{U}\|_{H^{-1}((-T, T); L^2(\mathbb{R}^n))}. \end{aligned}$$

This estimate combined with (2.10) gives

$$\|\mathcal{R}\|_{H^{-3}((-T, T); H^1(\mathbb{R}^n))} \lesssim \|f\|_{L^\infty(\mathbb{R}^n)}^2 \|\partial_t^2\mathcal{U}\|_{H^{-1}((-T, T); L^2(\mathbb{R}^n))}. \quad (2.11)$$

Finally, by using the trace theorem, (2.11) gives the desired estimate for the remainder term $\mathcal{R}|_{\Sigma_T}$. This finishes the proof. \square

Condition (2.3). By Proposition 1.3, whose proof is presented in Section 4, we consider

$$\mathcal{X}_2 = L_{\Omega_\sigma}^2(\mathbb{R}^n), \quad \mathcal{Y}_2 = H^{3/2}(\Sigma_T \cap \{t > z\}).$$

Now for any $f \in \mathcal{X}_1 = H_{\Omega_\sigma}^{s_0}(\mathbb{R}^n)$, Proposition 1.3 and the trace theorem imply that

$$\|f\|_{\mathcal{X}_2} \leq C\|A_0f\|_{\mathcal{Y}_2}.$$

Thus condition (2.3) is satisfied.

Condition (2.1). We have to choose a pair of Banach spaces $\mathcal{X}_3 \subset \mathcal{X}_1$ and $\mathcal{Y}_3 \subset \mathcal{Y}_2$ so that for some $\mu_1, \mu_2 \in (0, 1)$ with $\mu_1\mu_2 > 1/2$ one has

$$\|f\|_{H^{s_0}(\mathbb{R}^n)} \lesssim \|f\|_{L^2(\mathbb{R}^n)}^{\mu_1} \|f\|_{\mathcal{X}_3}^{1-\mu_1}, \quad f \in \mathcal{X}_3, \quad (2.12)$$

and

$$\|g\|_{H^{3/2}(\Sigma_T \cap \{t > z\})} \lesssim \|g\|_{H^{-3}((-T, T); H^{1/2}(S^{n-1}))|_{\{t > z\}}}^{\mu_2} \|g\|_{\mathcal{Y}_3}^{1-\mu_2}, \quad g \in \mathcal{Y}_3. \quad (2.13)$$

Fix an arbitrary $\mu_1 \in (0, 1)$. Consider $s_1 > s_0$ satisfying $s_0 = 0(\mu_1) + s_1(1 - \mu_1)$. A standard interpolation argument yields (2.12) with

$$\mathcal{X}_3 = H_{\Omega_\sigma}^{s_1}(\mathbb{R}^n).$$

On the other hand, if we fix $\mu_2 \in (0, 1)$ and choose $s_2 > 3/2$ with $3/2 = (-3)(\mu_2) + s_2(1 - \mu_2)$, then interpolation gives that

$$\|g\|_{H^{3/2}(\Sigma_T \cap \{t > z\})} \lesssim \|g\|_{H^{-3}(\Sigma_T \cap \{t > z\})}^{\mu_2} \|g\|_{H^{s_2}(\Sigma_T \cap \{t > z\})}^{1-\mu_2}.$$

Thus choosing

$$\mathcal{Y}_3 = H^{s_2}(\Sigma_T \cap \{t > z\})$$

implies (2.13). Note that we can make $\mu_1\mu_2$ as close to 1 as we want by choosing s_1 and s_2 large enough. Thus, the condition $\mu_1\mu_2 > 1/2$ is satisfied.

We are now in a position to apply Proposition 2.1: there exist $\mu \in (0, 1)$, $C = C(L) > 0$ and $\varrho > 0$ small enough so that

$$\|f\|_{H^{s_0}(\mathbb{R}^n)} \leq C(L) \|\mathcal{A}(f) - \mathcal{A}(0)\|_{H^{-3}((-T, T); H^{1/2}(S^{n-1}))|_{\{t > z\}}}^\mu$$

whenever $\|f\|_{H^{s_0}(\mathbb{R}^n)} \leq \varrho$ and $\|f\|_{H^{s_1}(\mathbb{R}^n)} \leq L$ for $f \in H_{\Omega_\sigma}^{s_1}(\mathbb{R}^n)$. If we recall that $\mathcal{A}(f) = \hat{\mathcal{A}}(f)|_{\{t > z\}}$ and $\tilde{\mathcal{A}}(1 + f) = \partial_t^2 \hat{\mathcal{A}}(f)$, Lemma A.5 applied with $k = -5$ implies that

$$\|\eta - 1\|_{H^{s_0}(\mathbb{R}^n)} \leq C(L) \|\tilde{\mathcal{A}}(\eta) - \tilde{\mathcal{A}}(1)\|_{H^{-5}((-T, T); H^{1/2}(S^{n-1}))|_{\{t > z\}}}^\mu$$

when $\|\eta - 1\|_{H^{s_0}(\mathbb{R}^n)} \leq \varrho$ and $\|\eta\|_{H^{s_1}(\mathbb{R}^n)} \leq L$ for $f \in H_{\Omega_\sigma}^{s_1}(\mathbb{R}^n)$. This finishes the proof of Theorem 1.1.

3. THE LINEARIZED MAP

We now concentrate on studying the main properties of the linearization at 0 of the map \mathcal{A} given in (2.5). Recall that this linearization is denoted by A_0 and it is given by (2.6) and (2.7). The existence and uniqueness of solutions to (2.7) is provided by the following result. Its proof is provided in the Appendix.

Proposition 3.1. *Let $s_0 \gg n/2 + 2$ and $T > 1$. Consider $f \in H_{\Omega}^{s_0}(\mathbb{R}^n)$. There is a unique distributional solution $U(y, z, t)$ to (2.7), and it is supported in the region $\{t \geq z\}$. In particular, one has*

$$U(y, z, t) = u(y, z, t)H(t - z),$$

where u is a C^2 function in $\{t \geq z\}$ satisfying the IVP

$$\begin{cases} (\partial_t^2 - \Delta)u = 0 & \text{in } \{t > z\}, \\ (\partial_t + \partial_z)u = -\frac{1}{2}f & \text{on } \{t = z\}, \\ u|_{\{t < -1\}} = 0. \end{cases} \quad (3.1)$$

In addition, given any $K \geq 3$ we may arrange that u is C^K in the set $\{t \geq z\}$ by taking s_0 large enough. In particular, one always has $U \in H^{-1}((-T, T), L^2(\mathbb{R}^n))$.

Using standard ODE techniques, one has

$$u(y, z, t)|_{t=z} = -\frac{1}{2} \int_{-\infty}^0 f(y, s+z) ds := F(y, z, z) \quad (3.2)$$

It turns out that instead of (3.1) it is more convenient to consider the characteristic initial value problem

$$\begin{cases} (\partial_t^2 - \Delta)u = 0 & \text{in } \{t > z\}, \\ u = F & \text{on } \{t = z\}. \end{cases} \quad (3.3)$$

We will apply this with F given in (3.2), so that the initial value $F(y, z, z)$ satisfies

$$ZF(y, z, z) = -\frac{1}{2}f(y, z).$$

Here Z is defined as $ZF(y, z, z) := \partial_s(F(y, s, s))|_{s=z}$. Thanks to the chain rule, one can see that $Z := \partial_t + \partial_z$.

As explained in Section 2, the Fréchet derivative of \mathcal{A} at 0 is the map $A_0: f \mapsto U|_{\Sigma_T}$. However, it is technically easier to study the map

$$A'_0: F \mapsto u|_{\Sigma_T \cap \{t \geq z\}}, \quad (3.4)$$

where u solves (3.3). The original map A_0 can be recovered by $A_0 f = E_0(A'_0 F)$ where F is defined in (3.2) and E_0 denotes extension by zero from $\Sigma_T \cap \{t \geq z\}$ to Σ_T . Therefore in the linearized part we mainly focus on the inverse problem for A'_0 instead of A_0 . We shall adapt the modified time-reversal method proposed in [SU09b] to study A'_0 . Denote

$$\begin{cases} \Sigma := \{(y, z, t); z \leq t \leq T, (y, z) \in \partial\Omega\}, \\ \Sigma_- := \{(y, z, t); -1 \leq t \leq z, (y, z) \in \partial\Omega\}, \\ \Gamma := \{(y, z, z); (y, z) \in \Omega\}, \quad \Gamma_T := \{(y, z, T); (y, z) \in \Omega\}, \\ Q := \Omega \times [-1, T], \end{cases} \quad (3.5)$$

In the time reversal procedure, we generate an approximate inverse for A'_0 as the map $h \mapsto v|_{\Gamma}$, where v solves the problem

$$(\partial_t^2 - \Delta)v = 0 \text{ in } Q, \quad v = h \text{ on } \Sigma \cup \Sigma_-, \quad v = v_0, v_t = v_1 \text{ in } \Gamma_T,$$

with prescribed boundary data h and certain data (v_0, v_1) at the final time $t = T$. For the time reversal argument it would be natural to work with energy spaces. However, it is not obvious that $h \in H^1(\Sigma)$ would imply $v|_{\Gamma} \in H^1(\Gamma)$. We will prove this fact by using energy estimates. To that end, we first investigate an initial boundary value problem.

3.1. An initial boundary value problem. For a function h , we adopt the following convention:

$$h_t := \partial_t h, \quad h_z := \partial_z h, \quad h_y := \nabla_y h, \quad h_x := (\nabla_y h, \partial_z h), \quad h_\nu := \nu \cdot h_x.$$

We define a seminorm $\|\cdot\|_{\mathcal{H}}$ as follows,

$$\|h\|_{\mathcal{H}} = \left(\frac{1}{\sqrt{2}} \int_{\Gamma} (|\nabla_y h|^2 + |h_z + h_t|^2) dS \right)^{\frac{1}{2}}. \quad (3.6)$$

We introduce the following PDE,

$$\begin{cases} (\partial_t^2 - \Delta)v = G & \text{in } Q, \\ v = u & \text{on } \Sigma \cup \Sigma_-, \\ v = \phi_0, v_t = \phi_1 & \text{in } \Gamma_T. \end{cases} \quad (3.7)$$

We have the following a priori estimate.

Proposition 3.2. *Assume $v \in C^2(\bar{Q})$ solves (3.7) with $\phi_0 \in H^1(\Omega)$, $\phi_1 \in L^2(\Omega)$, $G \in L^2(Q)$ and $u|_{\Sigma \cup \Sigma_-} \in H^1(\Sigma \cup \Sigma_-)$. For $T \geq 1$, we have*

$$\|v\|_{\mathcal{H}} \leq C(\|\nabla\phi_0\|_{L^2(\Omega)} + \|\phi_1\|_{L^2(\Omega)} + \|G\|_{L^2(Q)} + \sqrt{ne}^{T/2}\|v\|_{H^1(\Sigma)}). \quad (3.8)$$

for some constant C independent of v , ϕ_0 , ϕ_1 , G , n and T .

Remark 3.1. The result in Proposition 3.2 is contained in [Ma21]. Though, for self-containedness of this paper, and also for preparing necessary ingredients for the following sections, we demonstrate a proof of Proposition 3.2 below.

To prove Proposition 3.2, we do some preparations first. Let us introduce the following notations:

$$\begin{cases} \tilde{Q}_\tau := \{(y, z, t); (y, z) \in \Omega, z \leq t \leq \tau\} \subset Q, \\ \tilde{\Sigma}_\tau := \{(y, z, t); (y, z) \in \partial\Omega, z \leq t \leq \tau\} \subset \Sigma, \\ \tilde{\Gamma}_\tau := \{(y, z, \tau); (y, z) \in \Omega\} \cap \tilde{Q}_\tau, \end{cases}$$

thus $\partial\tilde{Q}_\tau = \tilde{\Gamma}_\tau \cup \tilde{\Sigma}_\tau \cup \Gamma$ for $\tau \geq 1$. Note also that when $-1 \leq \tau < 1$ the set $\tilde{\Gamma}_\tau$ is a strict subset of $\Omega \times \{\tau\}$. To abbreviate, we denote $\square := \partial_t^2 - \Delta$.

Lemma 3.3. *Under the same assumptions as in Proposition 3.2, we have*

$$\|v\|_{\mathcal{H}}^2 \leq C(\|\nabla\phi_0\|_{L^2(\Omega)}^2 + \|\phi_1\|_{L^2(\Omega)}^2 + e^T\|G\|_{L^2(Q)}^2 + \|v_t\|_{L^2(\Sigma)}\|v_\nu\|_{L^2(\Sigma)})$$

for some constant C independent of u , ϕ_0 , ϕ_1 , G and T .

Proof. Integrating the identity

$$2 \operatorname{Re}\{\bar{v}_t \square v\} = \operatorname{Re} \operatorname{div}_{x,t} (-2\bar{v}_t \nabla_x v, |v_t|^2 + |\nabla_x v|^2) \quad (3.9)$$

over \tilde{Q}_τ when $\tau \geq 1$ and noticing that $\square v = G$ in \tilde{Q}_τ , we obtain that

$$\begin{aligned} \operatorname{Re} \int_{\tilde{Q}_\tau} 2\bar{v}_t G &= \int_{\tilde{\Gamma}_\tau} (|v_t|^2 + |\nabla_x v|^2) \, dS - 2 \operatorname{Re} \int_{\tilde{\Sigma}_\tau} \bar{v}_t v_\nu \, dS \\ &\quad - \frac{1}{\sqrt{2}} \int_{\Gamma} (|\nabla_y v|^2 + |v_z + v_t|^2) \, dS, \end{aligned} \quad (3.10)$$

Hence, (3.10) becomes

$$\operatorname{Re} \int_{\tilde{Q}_\tau} 2\bar{v}_t G = \int_{\tilde{\Gamma}_\tau} (v_t^2 + |\nabla_x v|^2) \, dS - \frac{1}{\sqrt{2}} \|v\|_{\mathcal{H}}^2 - 2 \operatorname{Re} \int_{\tilde{\Sigma}_\tau} \bar{v}_t v_\nu \, dS. \quad (3.11)$$

We set

$$e(\tau) := \int_{\tilde{\Gamma}_\tau} (|\nabla_x v|^2 + |v_t|^2) \, dS. \quad (3.12)$$

From (3.11) we deduce

$$e(\tau) \leq \int_0^\tau e(s) \, ds + \int_{\tilde{Q}_\tau} |G|^2 + \frac{1}{\sqrt{2}} \|v\|_{\mathcal{H}}^2 + 2 \int_{\tilde{\Sigma}_\tau} |v_t v_\nu| \, dS,$$

so the Gronwall's inequality gives

$$e(\tau) \leq e^\tau \left(\frac{1}{\sqrt{2}} \|v\|_{\mathcal{H}}^2 + \int_{\tilde{Q}_\tau} |G|^2 + 2 \int_{\tilde{\Sigma}_\tau} |v_t v_\nu| \, dS \right). \quad (3.13)$$

From (3.11) we also obtain

$$\begin{aligned} \frac{1}{\sqrt{2}} \|v\|_{\mathcal{H}}^2 &= \int_{\tilde{\Gamma}_\tau} (v_t^2 + |\nabla_x v|^2) \, dS - 2 \operatorname{Re} \int_{\tilde{\Sigma}_\tau} \bar{v}_t v_\nu \, dS - \operatorname{Re} \int_{\tilde{Q}_\tau} 2\bar{v}_t G \\ &\leq e(\tau) + 2[1 + \epsilon(e^\tau - 1)] \int_{\tilde{\Sigma}_\tau} |v_t v_\nu| \, dS + \left[\frac{1}{\epsilon} + \epsilon(e^\tau - 1) \right] \int_{\tilde{Q}_\tau} |G|^2 + \epsilon(e^\tau - 1) \frac{1}{\sqrt{2}} \|v\|_{\mathcal{H}}^2. \end{aligned}$$

By setting $\epsilon = [2(e^\tau - 1)]^{-1} > 0$ and absorbing $\|v\|_{\mathcal{H}}^2$ on the right hand side, we obtain

$$\frac{1}{2\sqrt{2}}\|v\|_{\mathcal{H}}^2 \leq e(\tau) + 3 \int_{\tilde{\Sigma}_\tau} |v_t v_\nu| dS + 4e^\tau \int_{\tilde{Q}_\tau} |G|^2. \quad (3.14)$$

Now setting $\tau = T$ and noting that $v(\cdot, T) = \phi_0$ and $v_t(\cdot, T) = \phi_1$, (3.14) gives

$$\|v\|_{\mathcal{H}}^2 \leq C(\|\nabla\phi_0\|_{L^2(\Omega)}^2 + \|\phi_1\|_{L^2(\Omega)}^2 + \|v_t\|_{L^2(\Sigma)}\|v_\nu\|_{L^2(\Sigma)} + e^T\|G\|_{L^2(Q)}^2),$$

for some constant C independent of T . We arrive at the conclusion. \square

The norm $\|v_\nu\|_{L^2(\Sigma)}$ can also be estimated. Following [RS20b, Lemma 3.3], we examine the quantity $(x \cdot \nabla v)\square v$ and use integration by parts.

Lemma 3.4. *Assume $v \in C^2(\bar{Q})$ solves (3.7) For $\tau \geq 1$, we have*

$$\|v_\nu\|_{L^2(\tilde{\Sigma}_\tau)}^2 \leq Cne^\tau(\|G\|_{L^2(\tilde{Q}_\tau)}^2 + \|v\|_{\mathcal{H}}^2 + ne^\tau\|v\|_{H^1(\tilde{\Sigma}_\tau)}^2),$$

for some constant C independent of v , G , τ and the dimension n .

Proof. Integrating the identity

$$2 \operatorname{Re}(x \cdot \nabla \bar{v})\square v = \operatorname{Re} \operatorname{div}_{x,t} (x(|\nabla v|^2 - |v_t|^2) - 2(x \cdot \nabla \bar{v})\nabla v, 2(x \cdot \nabla \bar{v})v_t) + n|v_t|^2 - (n-2)|\nabla v|^2$$

over \tilde{Q}_τ and noticing that $\square v = G$ in \tilde{Q}_τ , similar to (3.10) we now have

$$\begin{aligned} 2 \int_{\tilde{\Sigma}_\tau} |v_\nu|^2 &= \frac{1}{\sqrt{2}} \int_{\Gamma} [z|\nabla(v)|^2 - 2 \operatorname{Re}\{x \cdot \nabla(\bar{v})\partial_z(v)\}] + \operatorname{Re} \int_{\tilde{\Gamma}_\tau} 2(x \cdot \nabla v)v_t \\ &\quad - 2 \operatorname{Re} \int_{\tilde{Q}_\tau} (x \cdot \nabla \bar{v})G + \int_{\tilde{\Sigma}_\tau} (|v_\nu|^2 + \frac{1}{2} \sum_{i \neq j} |\Omega_{ij}v|^2 - |v_t|^2) \\ &\quad + \int_{\tilde{Q}_\tau} (n|v_t|^2 - (n-2)|\nabla v|^2), \end{aligned} \quad (3.15)$$

where we used (see [RS20b, (1.19)]) the fact

$$|\nabla\varphi(x)|^2 = |\varphi_\nu(x)|^2 + \frac{1}{2} \sum_{i \neq j} |\Omega_{ij}\varphi(x)|^2, \quad \text{where } |x| = 1, \quad \Omega_{ij} = x_i\partial_j - x_j\partial_i.$$

By moving the v_ν -term on the RHS of (3.15) to the left, we further obtain

$$\int_{\tilde{\Sigma}_\tau} |v_\nu|^2 \lesssim \|v\|_{\mathcal{H}}^2 + e(\tau) + \|v\|_{H^1(\tilde{\Sigma}_\tau)}^2 + n \int_0^\tau e(s) ds + n \int_{\tilde{Q}_\tau} |G|^2, \quad (3.16)$$

where $e(\tau)$ is defined in (3.12). Combining (3.16) and (3.13), we have

$$\int_{\tilde{\Sigma}_\tau} |v_\nu|^2 \lesssim (n+1)e^\tau(\|v\|_{\mathcal{H}}^2 + \int_{\tilde{Q}_\tau} |G|^2 + \frac{1}{\epsilon}\|v_t\|_{\tilde{\Sigma}_\tau}^2 + \epsilon\|v_\nu\|_{\tilde{\Sigma}_\tau}^2) + \|v\|_{\mathcal{H}}^2 + \|v\|_{H^1(\tilde{\Sigma}_\tau)}^2.$$

Setting $\epsilon = [2C(n+1)e^\tau]^{-1}$ for some constant C big enough, we have

$$\int_{\tilde{\Sigma}_\tau} |v_\nu|^2 \lesssim (n+1)e^\tau \int_{\tilde{Q}_\tau} |G|^2 + ne^\tau\|v\|_{\mathcal{H}}^2 + n^2e^{2\tau}\|v\|_{H^1(\tilde{\Sigma}_\tau)}^2.$$

We arrive at the conclusion. \square

Now we are ready to prove Proposition 3.2.

Proof of Proposition 3.2. Combining Lemma 3.3 and Lemma 3.4, we get

$$\begin{aligned} \|v\|_{\mathcal{H}}^2 &\leq C(\|\nabla\phi_0\|_{L^2(\Omega)}^2 + \|\phi_1\|_{L^2(\Omega)}^2) + \epsilon Cne^T\|G\|_{L^2(Q)}^2 + \left(\frac{C}{4\epsilon} + \epsilon Cn^2e^{2T}\right)\|v\|_{H^1(\Sigma)}^2 \\ &\quad + \epsilon Cne^T\|v\|_{\mathcal{H}}^2. \end{aligned}$$

By setting $\epsilon = (2Cne^T)^{-1}$ and absorbing the $\|v\|_{\mathcal{H}}^2$ term on the RHS by the LHS, we finally arrive at (3.8). The proof is complete. \square

3.2. The time reversal method. Recall the notation from (3.5). Given the data $\{u|_{\Sigma}\}$, we aim to construct an approximation of F in (3.3) by using a modified time-reversal method as in [SU09b]. We define a function v as the solution of the following system:

$$\begin{cases} (\partial_t^2 - \Delta)v = 0 & \text{in } Q, \\ v = \tilde{u} & \text{on } \Sigma \cup \Sigma_-, \\ v = \phi_0, v_t = 0 & \text{in } \Gamma_T, \\ \Delta\phi_0 = 0 & \text{in } \Omega, \quad \phi_0 = u(\cdot, T) \text{ on } \partial\Omega. \end{cases} \quad (3.17)$$

Here \tilde{u} is a suitable extension of u from Σ to Σ_- given in Lemma 3.6 below. The choice of this extension will have no influence on the analysis in the rest of the paper. However, as mentioned earlier, it is not obvious that $v|_{\Gamma}$ is in $H^1(\Gamma)$ if $u \in H^1(\Sigma)$. Hence, we shall define

$$v|_{\Gamma} := \lim_{\epsilon \rightarrow 0^+} v_{\epsilon}|_{\Gamma}$$

where v_{ϵ} solves

$$\begin{cases} (\partial_t^2 - \Delta)v_{\epsilon} = 0 & \text{in } Q, \\ v_{\epsilon} = u_{\epsilon} & \text{on } \Sigma \cup \Sigma_-, \\ v_{\epsilon} = \phi_0, \partial_t v_{\epsilon} = 0 & \text{in } \Gamma_T, \\ \Delta\phi_0 = 0 & \text{in } \Omega, \quad \phi_0 = u_{\epsilon}(\cdot, T) \text{ on } \partial\Omega. \end{cases} \quad (3.18)$$

Here u_{ϵ} is a modification of u on the lateral boundary given in (3.19) such that the compatibility requirements on $\partial\Gamma_T$ required by the existence of a smooth solution v_{ϵ} (when $u|_{\Sigma}$ is smooth) are satisfied. Then we know that $v_{\epsilon} \in C^2(\bar{Q})$ by [LLT86, Remark 2.10], and hence Proposition 3.2 can be applied to v_{ϵ} and the limit exists due to the estimate given in Proposition 3.2. For convenience we reproduce [LLT86, Remark 2.10] in the next lemma.

Lemma 3.5. *Let Φ be a solution of the system*

$$\begin{cases} (\partial_t^2 - \Delta)\Phi = F & \text{in } \Omega \times [0, T], \\ \Phi = g & \text{on } \Xi := \partial\Omega \times [0, T], \\ \Phi = \Phi_0, \Phi_t = \Phi_1 & \text{in } \Omega \times \{t = 0\}, \end{cases}$$

with (F, g, Φ_0, Φ_1) satisfying the regularity assumptions (m is a non-negative integer)

$$\begin{cases} F \in L^1(0, T; H^m(\Omega)), \quad \frac{d^m F}{dt^m} \in L^1(0, T; L^2(\Omega)), \\ \Phi_0 \in H^{m+1}(\Omega), \quad \Phi_1 \in H^m(\Omega), \\ g \in H^{m+1}(\Xi) := L^2(0, T; H^{m+1}(\Xi)) \cap H^{m+1}(0, T; L^2(\Xi)) \end{cases}$$

and satisfying all necessary compatibility conditions up to order m . Then

$$\Phi \in C([0, T]; H^{m+1}(\Omega)), \quad \frac{d^{(m+1)}\Phi}{dt^{(m+1)}} \in C([0, T]; L^2(\Omega)), \quad \text{and} \quad \frac{\partial\Phi}{\partial\nu} \in H^m(\Xi).$$

Note that in contrast with u which is defined in the infinite half plane $\{(y, z, t); t \geq z\}$, v is only defined in the finite cylinder $\Omega \times [-1, T]$.

Lemma 3.6. *If $u|_{\Sigma} \in C^K(\Sigma)$ where $K = \lceil n/2 \rceil + 2$, then there exists a unique solution $v \in C([-1, T]; H^1(\Omega))$ of the system (3.17), and $v|_{\Gamma} \in H^1(\Gamma)$.*

Proof. We extend u from Σ to Σ_- as follows. Inspired by the Taylor's expansion, we define the following extension of u ,

$$\tilde{u}(y, z, t) := \begin{cases} u(y, z, t), & \text{on } \Sigma, \\ \sum_{k=0}^K (\partial_t^k u)(y, z, z) \cdot (t - z)^k / k!, & \text{on } \Sigma_-. \end{cases}$$

Then $\tilde{u} \in C^K(\Sigma \cup \Sigma_-) \subset H^K(\Sigma \cup \Sigma_-)$. By the trace theorem $\tilde{u}|_{\partial\Omega \times \{T\}}$ is in $H^{K-\frac{1}{2}}(\partial\Omega \times \{T\})$, which implies that the harmonic function ϕ_0 in (3.17) is in $H^K(\Gamma_T)$.

Now we construct a series of approximate Dirichlet boundary data $\{u_\epsilon\}_{\epsilon>0}$ such that the compatibility conditions needed in Lemma 3.5 will be satisfied. To that end, we fix a cutoff function $\chi_0 \in C_c^\infty(\mathbb{R})$ satisfying $\chi_0(t) = 1$ when $|t| \leq 1$ and $\chi_0(t) = 0$ when $|t| \geq 2$, and we set

$$u_\epsilon(x, t) := \chi_0\left(\frac{T-t}{\epsilon}\right)\tilde{u}(x, T) + (1 - \chi_0\left(\frac{T-t}{\epsilon}\right))\tilde{u}(x, t). \quad (3.19)$$

We see $u_\epsilon = \tilde{u} = u$ on Γ_T . We denote the initial velocity of (3.18) as ϕ_1 , i.e. $u_t = \phi_1 = 0$ in Γ_T . The compatibility conditions for Lemma 3.5 up to the order K are the following (cf. e.g. [Eva10, §7.2 (62)]):

$$(\Delta^k \phi_0, \Delta^k \phi_1) = (\partial_t^{2k} u_\epsilon, \partial_t^{2k+1} u_\epsilon) \text{ on } \Gamma_T, \quad \forall k : 0 \leq k \leq \left\lfloor \frac{K}{2} \right\rfloor. \quad (3.20)$$

It can be checked that (3.20) is true, and this is simply because $\phi_0 = u_\epsilon(\cdot, T)$, $\Delta\phi_0 = 0$, $\phi_1 = 0$ and $\partial_t^j u_\epsilon(\cdot, T) = 0$ for $\forall j \geq 1$.

It is straightforward to check that

$$\|u - u_\epsilon\|_{L^2(\Sigma)} \rightarrow 0, \quad \|\nabla|_{\Sigma, x}(u - u_\epsilon)\|_{L^2(\Sigma)} \rightarrow 0, \quad \text{as } \epsilon \rightarrow 0^+, \quad (3.21)$$

where $\nabla|_{\Sigma, x}$ stands for the spacial component of the gradient on the manifold Σ . We aim to make the sequence $\{u_\epsilon\}_{\epsilon>0}$ converge to u in $H^1(\Sigma)$. According to (3.21), it is left to show $\|\partial_t(u - u_\epsilon)\|_{L^2(\Sigma)} \rightarrow 0$ as $\epsilon \rightarrow 0^+$. One can compute

$$\begin{aligned} \partial_t(u_\epsilon - \tilde{u})(x, t) &= \frac{1}{\epsilon} \chi_0'\left(\frac{T-t}{\epsilon}\right) [\tilde{u}(x, t) - \tilde{u}(x, T)] + (1 - \chi_0\left(\frac{T-t}{\epsilon}\right)) \tilde{u}_t(x, t) - \tilde{u}_t(x, t) \\ &= \frac{1}{\epsilon} \chi_0'\left(\frac{T-t}{\epsilon}\right) [(t-T)\tilde{u}_t(x, T) + \mathcal{O}(|t-T|^2)] - \chi_0\left(\frac{T-t}{\epsilon}\right) \tilde{u}_t(x, t) \\ &= -\frac{T-t}{\epsilon} \chi_0'\left(\frac{T-t}{\epsilon}\right) [\tilde{u}_t(x, T) + \mathcal{O}(|t-T|)] - \chi_0\left(\frac{T-t}{\epsilon}\right) \tilde{u}_t(x, t). \end{aligned}$$

Note that $\frac{T-t}{\epsilon} \chi_0'\left(\frac{T-t}{\epsilon}\right) \rightarrow 0$ point-wise in $[0, T]$, and $\chi_0\left(\frac{T-t}{\epsilon}\right) \rightarrow 0$ point-wise in $[0, T]$ as $\epsilon \rightarrow 0^+$, so $\partial_t(u_\epsilon - \tilde{u})(x, t) \rightarrow 0$ almost everywhere on $\bar{\Sigma}$, and hence

$$\|\partial_t(\tilde{u} - u_\epsilon)\|_{L^2(\Sigma)} \rightarrow 0, \quad \text{as } \epsilon \rightarrow 0^+. \quad (3.22)$$

Combining (3.21) with (3.22), we arrive at

$$\|\tilde{u} - u_\epsilon\|_{H^1(\Sigma)} \rightarrow 0, \quad \text{as } \epsilon \rightarrow 0^+. \quad (3.23)$$

The smoothness of \tilde{u} implies $u_\epsilon \in C^K(\Sigma \cup \Sigma_-) \subset H^K(\Sigma \cup \Sigma_-)$. For the system (3.18), the prerequisites of Lemma 3.5 are all satisfied now, especially the compatibility requirements (3.20), so we can conclude

$$v_\epsilon \in C([-1, T]; H^K(\Omega)), \quad \partial_t v_\epsilon \in C([-1, T]; H^{K-1}(\Omega)).$$

By the Sobolev embedding theorem we know $H^K(\Omega) \subset C^2(\bar{\Omega})$ when $K - n/2 \geq 2$, so we set $K = \lceil n/2 \rceil + 2$, and thus

$$v_\epsilon \in C([-1, T]; C^2(\bar{\Omega})), \quad \partial_t v_\epsilon \in C([-1, T]; C^1(\bar{\Omega})), \quad (3.24)$$

so $\Delta v_\epsilon(\cdot, t) \in C(\overline{\Omega})$ for each t . Therefore, the equation $(\partial_t^2 - \Delta)v_\epsilon = 0$ implies

$$\partial_t^2 v_\epsilon(x, t) \text{ is continuous w.r.t. } t. \quad (3.25)$$

By (3.24) and (3.25) we obtain $v_\epsilon \in C^2(\overline{Q})$, thus v_ϵ is well-defined on the slanted plane Γ and $v_\epsilon|_\Gamma \in H^1(\Gamma)$.

Finally, by Proposition 3.2 and the convergence in (3.23), we see $\{v_\epsilon|_\Gamma\}$ is a Cauchy sequence in $H^1(\Gamma)$, so we can define $v|_\Gamma$ as the limit of $v_\epsilon|_\Gamma$, i.e.,

$$v|_\Gamma := \lim_{\epsilon \rightarrow 0^+} v_\epsilon|_\Gamma,$$

and the estimate in Proposition 3.2 implies $v|_\Gamma \in H^1(\Gamma)$. The proof is done. \square

Remark 3.2. According to (3.3) and (3.18) we see that an $F \in H_{\overline{\Omega}}^{s_0}(\mathbb{R}^n)$ produces a u and a v_ϵ . By Proposition 3.1 and Lemma 3.6, we can conclude there exists a large enough integer s_0 such that $u \in C^{\lceil n/2 \rceil + 2}$ and $v_\epsilon \in C^2$, respectively.

3.3. The approximate inverse of A'_0 . To introduce the approximate inverse B of A'_0 , we introduce some function spaces first. Recall once more the notation given in (3.5). We wish to consider functions F in $\{t = z\}$ that satisfy $ZF = 0$ outside Γ , since this is true in (3.2) when $\text{supp}(f) \subset \overline{\Omega}$. Thus we consider functions h on $\{t = z\}$ that satisfy the conditions:

$$\text{supp}(h) \subset \{(y, z, z); |(y, z)| \leq 1 \text{ and } |y| \leq 1 - \frac{\sigma}{2}\}, \text{ and } Zh = 0 \text{ outside } \Gamma. \quad (3.26)$$

Later on, we will see that in (4.9), we have to make sure that the denominator $1 - |y|^2$ is positive away from zero, and due to this technical reason, we have to put the prerequisite “ $|y| \leq 1 - \frac{\sigma}{2}$ ” in (3.26), and for the same reason we also need “ $|x| \geq 1 - \sigma$ ” in (1.1). These are due to computations in (4.9)-(4.11) and we do not explain the details here.

We define the function spaces

$$\begin{cases} \mathcal{H}' := \{h \in H_{\text{loc}}^1(\{t = z\}); h \text{ satisfies (3.26)}\}, \\ \mathcal{H} := \{h \in \mathcal{H}'; \text{supp } h \subset \overline{\Gamma}\}, \end{cases} \quad (3.27)$$

Recall the seminorm $\|h\|_{\mathcal{H}}$ defined in (3.6),

$$\|h\|_{\mathcal{H}} = \left(\frac{1}{\sqrt{2}} \int_{\Gamma} (|\nabla_y h|^2 + |h_z + h_t|^2) dS \right)^{\frac{1}{2}}.$$

Note that $Z = \partial_z + \partial_t$ is tangential to Γ , and hence $\|h\|_{\mathcal{H}}$ corresponds to $\|\nabla_{\Gamma} h\|_{L^2(\Gamma)}$. Note also that by (3.26) $\|\cdot\|_{\mathcal{H}}$ is in fact a norm on \mathcal{H}' , so that $(\mathcal{H}', \|\cdot\|_{\mathcal{H}})$ is a Banach space and $(\mathcal{H}, \|\cdot\|_{\mathcal{H}})$ is a closed subspace. Moreover, by Poincaré inequality, the norm $\|h\|_{\mathcal{H}}$ is comparable to $\|h\|_{H^1(\Gamma)}$ for $h \in \mathcal{H}$. We shall use these facts several times in the following computations.

By Lemma 3.6 we can define an *approximate inverse* B of A'_0 as follows,

$$\begin{cases} B := \lim_{\epsilon \rightarrow 0^+} B_\epsilon, \text{ where } B_\epsilon: C^K(\Sigma) \subset H^1(\Sigma) \rightarrow \mathcal{H}', \quad u|_\Sigma \mapsto \tilde{v}_\epsilon, \\ \tilde{v}_\epsilon \text{ is the zero extension of } v_\epsilon|_\Gamma \text{ to } \{t = z\} \text{ where } v_\epsilon \text{ solves (3.18)}. \end{cases} \quad (3.28)$$

Remark 3.3. Due to the prerequisite “ $|y| \leq 1 - \frac{\sigma}{2}$ ” in the definition of \mathcal{H}' given in (3.26), the actual approximate inverse will be χB where χ is a function satisfying

$$\chi \in C^\infty(\{t = z\}), \quad |\chi| \leq 1, \quad \chi(y, z, z) = \begin{cases} 1, & \text{when } |y| \leq 1 - \sigma, \\ 0, & \text{when } |y| \geq 1 - \sigma/2. \end{cases} \quad (3.29)$$

From the proof of Lemma 3.6 we see $v|_{\Gamma} := \lim_{\epsilon \rightarrow 0^+} v_{\epsilon}|_{\Gamma} \in H^1(\Gamma)$, and \tilde{v} is the zero extension of v , so $\tilde{v} \in \mathcal{H}'$. Note that the restriction of u on $\{t = z\}$ belongs to \mathcal{H}' , so a solution of (3.17) will always have an extension in \mathcal{H}' . The map B shall be understood intuitively as a parametrix of A'_0 , and BA'_0F as an approximation of F for any $F \in \mathcal{H}$ but not for $F \in \mathcal{H}'$. This fact related to compact supports will be taken care of in Section 4). We will prove below in Proposition 3.7 that in fact $B : H^1(\Sigma) \rightarrow \mathcal{H}'$ is a bounded map.

3.4. Boundedness of the approximate inverse.

Proposition 3.7. *The map B extends as a bounded operator from $H^1(\Sigma)$ to \mathcal{H}' , and for $T \geq 1$, we have*

$$\|Bu\|_{\mathcal{H}} \leq Ce^{T/2}\|u\|_{H^1(\Sigma)}$$

for some constant C independent of u and T .

Proof. By density it is enough to prove the estimate when $u|_{\Sigma}$ is smooth. The equation (3.17) is a special case of (3.7) when we choose (ϕ_0, ϕ_1, G) according to

$$\begin{cases} \Delta\phi_0 = 0 \text{ in } \Omega, & \phi_0 = u(\cdot, T) \text{ on } \partial\Omega, \\ \phi_1 = 0 \text{ in } \Omega, & G = 0 \text{ in } Q. \end{cases}$$

Recall the definition of B in (3.28) and the proof of Lemma 3.6. From (3.8) we obtain

$$\|v_{\epsilon}\|_{\mathcal{H}}^2 \leq C(\|\nabla\phi_0\|_{L^2(\Omega)}^2 + ne^T\|v_{\epsilon}\|_{H^1(\Sigma)}^2),$$

and by taking the limit $\epsilon \rightarrow 0^+$ it gives

$$\|v\|_{\mathcal{H}}^2 = \|Bu\|_{\mathcal{H}}^2 = C(\|\nabla\phi_0\|_{L^2(\Omega)}^2 + ne^T\|u\|_{H^1(\Sigma)}^2). \quad (3.30)$$

It remains to bound the $\|\nabla\phi_0\|$ term by $\|u\|$. Since ϕ_0 is a harmonic function in Ω with Dirichlet data $u(\cdot, T)|_{\partial\Omega}$, so it follows from standard estimates for the Dirichlet problem and from the trace theorem on Σ that

$$\|\phi_0\|_{H^1(\Omega)} \leq C\|u(\cdot, T)\|_{H^{1/2}(\partial\Omega)} \leq C\|u\|_{H^1(\Sigma)}.$$

Then (3.30) becomes

$$\|Bu\|_{\mathcal{H}} \leq C(\|u\|_{H^1(\Sigma)} + \sqrt{ne^T}\|u\|_{H^1(\Sigma)}).$$

This is the required statement. \square

Readers may note that the constant $e^{T/2}$ in Lemma 3.4, Proposition 3.2 and Proposition 3.7 might not be optimal.

4. THE ERROR OPERATOR

In this section we recover the function F in (3.3) by employing an iterative algorithm, and the function space at the beginning of the algorithm is different from these in the rest of the algorithm. We define

$$\mathcal{H}_{\text{int}} := \{h \in H_{\text{loc}}^1(\{t = z\}); h \text{ satisfies (3.26) with } \frac{\sigma}{2} \text{ replaced by } \sigma\} \cap H_{\Gamma}^{s_0}(\{t = z\}).$$

Here s_0 shall be chosen large enough such that Proposition 3.1 can be applied. The subscript in \mathcal{H}_{int} stands for ‘‘initial’’, and \mathcal{H}_{int} is the function space at the beginning of the iteration. The construction of χ in (3.29) guarantees

$$\chi h = h \text{ in } \mathcal{H}_{\text{int}}. \quad (4.1)$$

Recall the operators A'_0 and B_{ϵ} defined in (3.4) and (3.28). Recall Remark 3.3. We wish to consider the operator $\chi BA'_0\chi$, where χ satisfies (3.29). We need χ because F is supported in $\{\chi = 1\}$ (recall (1.1)), so we want the approximate inverse to be also supported

in $\{\chi = 1\}$. By (4.1) and Proposition 3.1, we see that if $h \in \mathcal{H}_{\text{int}}$, we have $A'_0 \chi h \in H^1(\Sigma)$, and moreover $h = A'_0 \chi h$ on $\partial\Gamma$. Here $A'_0 \chi h$ stands for $A'_0(\chi h)$. And by (3.18), (3.19) and (3.28) we see $A'_0 \chi h = BA'_0 \chi h$ on $\partial\Gamma$. Hence, $h = BA'_0 \chi h$ on $\partial\Gamma$, and combining this with $h = \chi h$ on $\partial\Gamma$ induced by (4.1), one can conclude

$$h - \chi BA'_0 \chi h = 0 \text{ on } \partial\Gamma. \quad (4.2)$$

By Proposition 3.7 we have $\chi BA'_0 \chi h \in \mathcal{H}'$, so,

$$h - \chi BA'_0 \chi h \in \mathcal{H}'. \quad (4.3)$$

We introduce the second cutoff to preserve the vanishing trace of $h - \chi BA'_0 \chi h$. Indeed, by combining (4.2) and (4.3), we can conclude

$$\forall h \in \mathcal{H}_{\text{int}} \Rightarrow h - \chi BA'_0 \chi h \in \mathcal{H}.$$

Hence we can define an *error operator* K by the formula

$$K := \lim_{\epsilon \rightarrow 0^+} K_\epsilon, \quad \text{where } K_\epsilon: \mathcal{H}_{\text{int}} \rightarrow \mathcal{H}, \quad h \mapsto (I - \chi B_\epsilon A'_0 \chi)h. \quad (4.4)$$

There are two cutoff functions χ inside the operator $(I - \chi BA'_0 \chi)$. The first χ (on the right-hand-side of A'_0) is to make sure the argument after (3.26) stays valid. The second χ (on the left-hand-side of B) is to guarantee $(I - \chi BA'_0 \chi)$ maps \mathcal{H}_{int} into not only \mathcal{H}' but also, more importantly, into \mathcal{H} ; without the second χ , this will not be true.

In the linearized inverse problem for (3.4), KF indicates the difference between the original F and the approximation $\chi B(u|_\Sigma) = \chi BA'_0 \chi F$. For $F \in \mathcal{H}' \cap H^{s_0}(\{t = z\})$ with $\text{supp } F \subset \text{supp } \chi$, we have

$$(I - K)F = \chi B(A'_0 \chi F) = \chi B(A'_0 F) = \chi B(u|_\Sigma).$$

Note that $u|_\Sigma$, B and K are all known. Hence one would expect to recover F using the formula $(I - K)^{-1} \chi B(u|_\Sigma)$ provided that $I - K$ is invertible, which would hold in particular if the operator norm $\|K\|$ between suitable spaces is strictly less than 1.

Instead of (4.4), it would be more convenient if the domain and image of K would be the same. This is necessary, for instance, to perform a Neumann series argument. This leads us to considering the restriction of K to \mathcal{H} . We denote the restriction as \tilde{K} , i.e.,

$$\tilde{K}: \mathcal{H} \cap H^{s_0}(\{t = z\}) \rightarrow \mathcal{H}, \quad h \mapsto (I - \chi BA'_0 \chi)h.$$

We will prove in Proposition 4.4 that \tilde{K} extends as a bounded operator $\mathcal{H} \rightarrow \mathcal{H}$ with norm strictly less than 1 if T is large enough. Based on this, a reconstruction formula $F = (I + \sum_{j \geq 0} \tilde{K}^j K) \chi B(u|_\Sigma)$ is given in (4.22). After obtaining F , the f can be recovered by $f = ZF$. Based on this, a stability result is also given, see Proposition 4.6 for details. An illustration of the recovering procedure of F is given in Fig. 1.

$$\mathcal{H}_{\text{int}} \xrightarrow{BA'_0} \mathcal{H}', \quad \mathcal{H}_{\text{int}} \xrightarrow{K} \mathcal{H} \xrightarrow{\tilde{K}} \mathcal{H} \xrightarrow{\tilde{K}} \dots \xrightarrow{\tilde{K}} \mathcal{H} \xrightarrow{\tilde{K}} \dots$$

FIGURE 1. An illustration of the recovering procedure.

To elaborate how the error operator K works, we give an intuitive example. Assume $F \in \mathcal{H}'$, u and v satisfy the following systems,

$$\begin{cases} (\partial_t^2 - \Delta)u = 0 & \text{in } \{t > z\} \\ u = F & \text{on } \{t = z\}, \\ \text{supp } F \subset \text{supp } \chi \end{cases}, \quad \begin{cases} (\partial_t^2 - \Delta)v = 0 & \text{in } Q \\ v = u & \text{on } \Sigma \cup \Sigma_- \\ v = \phi_0, v_t = 0 & \text{in } \Gamma_T \end{cases}$$

with

$$\begin{cases} \Delta\phi_0 = 0 & \text{in } \Omega, \\ \phi_0 = u(\cdot, T) & \text{on } \partial\Omega. \end{cases}$$

Then, by ignoring the regularity issues we could say $KF = (u - \chi v)|_\Gamma$.

4.1. The energy estimates. For $\tau \in [0, T]$ and $h \in H^1(\mathbb{R}^n \times \{t = \tau\})$, we define the energy $E(\tau, h)$ as follows

$$E(\tau, h) := \int_{\mathbb{R}^n} (|\nabla_x h(x, \tau)|^2 + |h_t(x, \tau)|^2) dx.$$

We also define a functional $\|h\|_{\mathcal{H}_T}$ as follows,

$$\|h\|_{\mathcal{H}_T} := \left(\int_{\Omega} (|\nabla_x h(x, T)|^2 + |h_t(x, T)|^2) dx \right)^{1/2}. \quad (4.5)$$

The relations among $\|\cdot\|_{\mathcal{H}}$, $\|\cdot\|_{\mathcal{H}_T}$ and $E(T, \cdot)$ are described in the following lemmas.

Lemma 4.1. *Assume u is $C^2(\overline{Q})$ and u and v are associated to each other by means of (3.17), and u_ϵ and v_ϵ are defined as in the proof of Lemma 3.6. Set $w_\epsilon := u - v_\epsilon$. Then we have*

$$\lim_{\epsilon \rightarrow 0^+} \|w_\epsilon\|_{\mathcal{H}} = \lim_{\epsilon \rightarrow 0^+} \|w_\epsilon\|_{\mathcal{H}_T}.$$

Moreover, we have

$$\|u\|_{\mathcal{H}_T}^2 = \|u - \phi_0\|_{\mathcal{H}_T}^2 + \int_{\Omega} |\nabla\phi_0(x)|^2 dx.$$

Proof. According to the construction of v_ϵ , we know $\omega_\epsilon \in C^2(\overline{Q})$. Moreover, ω_ϵ satisfies

$$\begin{cases} (\partial_t^2 - \Delta)\omega_\epsilon = 0 & \text{in } Q, \\ v = u - u_\epsilon & \text{on } \Sigma \cup \Sigma_-. \end{cases}$$

Then according to (3.10) we have

$$0 = \|\omega_\epsilon\|_{\mathcal{H}_T}^2 - \|\omega_\epsilon\|_{\mathcal{H}}^2 - 2 \operatorname{Re} \int_{\overline{\Sigma}_T} \partial_t \overline{\omega_\epsilon} \partial_\nu \omega_\epsilon dS,$$

Note that $\omega_\epsilon \in C^2(\overline{Q})$. Combining this with Lemma 3.4, we can compute

$$\begin{aligned} \|\omega_\epsilon\|_{\mathcal{H}}^2 &\leq \|\omega_\epsilon\|_{\mathcal{H}_T}^2 + 2\|\omega_\epsilon\|_{H^1(\Sigma)} \|\partial_\nu \omega_\epsilon\|_{L^2(\Sigma)} \\ &\leq \|\omega_\epsilon\|_{\mathcal{H}_T}^2 + C\|\omega_\epsilon\|_{H^1(\Sigma)} (\|\omega_\epsilon\|_{\mathcal{H}} + C\|\omega_\epsilon\|_{H^1(\Sigma)}) \\ &= \|\omega_\epsilon\|_{\mathcal{H}_T}^2 + C\|\omega_\epsilon\|_{H^1(\Sigma)} \|\omega_\epsilon\|_{\mathcal{H}} + C^2\|\omega_\epsilon\|_{H^1(\Sigma)}^2 \\ &= \|\omega_\epsilon\|_{\mathcal{H}_T}^2 + 2C^2\|\omega_\epsilon\|_{H^1(\Sigma)}^2 + \frac{1}{2}\|\omega_\epsilon\|_{\mathcal{H}}^2 + C\|\omega_\epsilon\|_{H^1(\Sigma)}^2. \end{aligned} \quad (4.6)$$

We have $\|\omega_\epsilon\|_{H^1(\Sigma)} \rightarrow 0$ as $\epsilon \rightarrow 0^+$ due to (3.23). And, the function ω_ϵ ($:= u - v_\epsilon$) is independent of ϵ on Γ_T because v_ϵ is independent of ϵ on Γ_T , see (3.19). Hence, (4.6) implies $\|\omega_\epsilon\|_{\mathcal{H}}$ is uniformly bounded when $\epsilon \rightarrow 0^+$. Based on this fact, we re-compute (4.6) in the following way with the help of Lemma 3.4,

$$\|\|\omega_\epsilon\|_{\mathcal{H}}^2 - \|\omega_\epsilon\|_{\mathcal{H}_T}^2\| \leq 2\|\omega_\epsilon\|_{H^1(\Sigma)} \|\partial_\nu \omega_\epsilon\|_{L^2(\Sigma)} \leq C\|\omega_\epsilon\|_{H^1(\Sigma)} (\|\omega_\epsilon\|_{\mathcal{H}} + C\|\omega_\epsilon\|_{H^1(\Sigma)}). \quad (4.7)$$

By the boundedness of $\|\omega_\epsilon\|_{\mathcal{H}}$ w.r.t. ϵ as well as the fact that $\|\omega_\epsilon\|_{H^1(\Sigma)} \rightarrow 0$ as $\epsilon \rightarrow 0^+$, we can conclude from (4.7) that

$$\|\omega_\epsilon\|_{\mathcal{H}} - \|\omega_\epsilon\|_{\mathcal{H}_T} \rightarrow 0 \quad \text{as } \epsilon \rightarrow 0^+.$$

For the second claim, we have

$$\begin{aligned}
\|u - \phi_0\|_{\mathcal{H}_T}^2 &= \int_{\Omega} |\nabla_x(u(x, T) - \phi_0)|^2 dx + \int_{\Omega} |u_t|^2 dx \\
&= \|u\|_{\mathcal{H}_T}^2 - 2 \operatorname{Re} \int_{\Omega} \nabla_x(u) \cdot \nabla_x \overline{\phi_0} dx + \int_{\Omega} |\nabla_x \phi_0|^2 dx \\
&= \|u\|_{\mathcal{H}_T}^2 - 2 \operatorname{Re} \int_{\Omega} \nabla_x(u - \phi_0) \cdot \nabla_x \overline{\phi_0} dx - \int_{\Omega} |\nabla_x \phi_0|^2 dx \\
&= \|u\|_{\mathcal{H}_T}^2 - 2 \operatorname{Re} \int_{\partial\Gamma_T} (u - \phi_0) \partial_\nu \overline{\phi_0} dx + 2 \operatorname{Re} \int_{\Omega} (u - \phi_0) \Delta \phi_0 dx \\
&\quad - \int_{\Omega} |\nabla_x \phi_0|^2 dx \\
&= \|u\|_{\mathcal{H}_T}^2 - \int_{\Omega} |\nabla_x \phi_0|^2 dx.
\end{aligned}$$

The last equal sign is due to $u - \phi_0 = 0$ on $\partial\Gamma_T$ and $\Delta \phi_0 = 0$. The proof is complete. \square

Lemma 4.2. *Assume h satisfies $(\partial_t^2 - \Delta)h = 0$ in $\{t \geq z\}$ with the restriction $h|_{\Gamma}$ satisfying $h|_{\Gamma} \in \mathcal{H}_{int}$, then we have*

$$\|h\|_{\mathcal{H}_T}^2 \leq E(T, h) \leq \begin{cases} \|h\|_{\mathcal{H}}^2 + \frac{\sqrt{2}T}{\sqrt{\sigma}} \|h\|_{H^1(\partial\Gamma)}^2, & \text{if } h \in \mathcal{H}_{int}, \\ \|h\|_{\mathcal{H}}^2, & \text{if } h \in \mathcal{H} \cap H_{\Gamma}^{s_0}(\{t = z\}). \end{cases}$$

Proof. By the definition of the energy and the norm, it is obvious that $\|h\|_{\mathcal{H}_T}^2 \leq E(T, h)$.

Now we prove the second inequality. Recall Remark 3.2. Replacing the function v in identity (3.9) by h and integrating the identity in the region

$$\{(y, z, t) \in \mathbb{R}^{n+1}; y \in \mathbb{R}^{n-1}, z \leq t \leq T\},$$

we can obtain

$$E(T, h) = \frac{1}{\sqrt{2}} \int_{\{t=z\}} (|\nabla_y h|^2 + |\partial_z(h(y, z, z))|^2) dS. \quad (4.8)$$

The computation is similar to what has been done in Section 3.1, so we omit the details here.

When $h|_{\Gamma} \in \mathcal{H}' \cap H_{\Gamma}^{s_0}(\{t = z\})$, we know $h \in C^2(\{t \geq z\})$. Moreover, $h \in \mathcal{H}'$ gives $\partial_z(h(y, z, z)) = 0$ outside Γ and $\nabla_y h$ is constant along the curve $\gamma(z) := (y, z, z)$ for each y , so (4.8) becomes

$$\begin{aligned}
E(T, h) &= \frac{1}{\sqrt{2}} \int_{\{t=z; t \leq T\}} (|\nabla_y h|^2 + |\partial_z(h(y, z, z))|^2) dS \\
&= \|h\|_{\mathcal{H}}^2 + \frac{1}{\sqrt{2}} \int_{\{t=z; t \leq T\} \setminus \Gamma} |\nabla_y(h(y, z_y, z_y))|^2 dS \quad (\partial_z(h(y, z, z)) = 0) \\
&= \|h\|_{\mathcal{H}}^2 + \int_{\pi(\{t=z; t \leq T\} \setminus \Gamma)} |\nabla_y(h(y, z_y, z_y))|^2 dy dz \\
&\leq \|h\|_{\mathcal{H}}^2 + T \int_{|y| \leq 1} |\nabla_y(h(y, z_y, z_y))|^2 dy \\
&= \|h\|_{\mathcal{H}}^2 + T \int_{|y| \leq 1} \frac{1}{1 - |y|^2} |\vec{X}h(y, z_y, z_y)|^2 dy, \quad (4.9)
\end{aligned}$$

where $z_y := \sqrt{1 - |y|^2}$, $\pi: (x, t) \mapsto x$ is a projection map, and \vec{X} is a vector of vector fields defined as

$$\vec{X} := \sqrt{1 - |y|^2} \nabla_y - y \partial_z - y \partial_t.$$

The vector field \vec{X} is tangential to $\partial\Gamma_+ := \partial\Gamma \cap \{(y, z, t); z \geq 0\}$. To see this, we denote $S_1 := \{(y, z, t); |(y, z)|^2 = 1\}$ and $S_2 := \{(y, z, t); t = z\}$. Then $\partial\Gamma_+ = S_1 \cap S_2$, and $\vec{X}(|(y, z)|^2 - 1) = \vec{X}(t - z) = 0$ on $\partial\Gamma_+$. Therefore, \vec{X} is tangential to both S_1 and S_2 at $\partial\Gamma_+$, and hence tangential to $\partial\Gamma_+$.

Denote the volume form on $\partial\Gamma_+$ as dS , then it can be checked that

$$\frac{\sqrt{1 + |y|^2}}{\sqrt{1 - |y|^2}} dy = dS.$$

Hence, (4.9) implies

$$E(T, h) \leq \|h\|_{\mathcal{H}}^2 + T \int_{\partial\Gamma_+} \frac{1}{\sqrt{1 - |y|^4}} |\vec{X}h|^2 dS. \quad (4.10)$$

Recall that $\text{supp } h|_{\Gamma} \subset \{(y, z, t); |y| \leq 1 - \sigma/2\}$ for the fixed constant $\sigma \in (0, 1)$, so $\vec{X}h = 0$ when $|y| \geq 1 - \sigma/2$, see also [MPMS21, Proof of Lemma 2.9]. Therefore,

$$\begin{aligned} E(T, h) &\leq \|h\|_{\mathcal{H}}^2 + \frac{T}{\sqrt{1 - (1 - \sigma/2)^4}} \int_{\partial\Gamma_+} |\vec{X}h|^2 dS \\ &\leq \|h\|_{\mathcal{H}}^2 + \frac{\sqrt{2}T}{\sqrt{\sigma}} \|h\|_{H^1(\partial\Gamma)}^2. \end{aligned} \quad (4.11)$$

Hence, when $h \in C^2 \cap \mathcal{H}$, we know $h = 0$ on $\partial\Gamma$, so $E(T, h) \leq \|h\|_{\mathcal{H}}^2$. The proof is done. \square

4.2. Boundedness of the error operators. In the definition of K and \tilde{K} , the two operators are defined on certain subsets of \mathcal{H}' and \mathcal{H} , respectively. As we mentioned earlier, their domain of definitions can be extended to the whole \mathcal{H}' and \mathcal{H} . To that end, we first show some boundedness result of K and \tilde{K} in certain dense subsets of their domain. For simplicity we denote the operator norm of \tilde{K} as $\|\tilde{K}\|$, i.e., $\|\tilde{K}\| = \|\tilde{K}\|_{\mathcal{H} \rightarrow \mathcal{H}}$, and we also write

$$\|K\| := \sup \left\{ \frac{\|Kh\|_{\mathcal{H}}}{\|h\|_{\mathcal{H}} + \|h\|_{H^1(\partial\Gamma)}}; h \in \mathcal{H}_{\text{int}}, h \neq 0 \right\}.$$

Lemma 4.3. *Let $\sigma \in (0, 1)$ be as in (1.1). If $C > 0$ is the constant from Lemma 4.2 and s_0 is determined as in Remark 3.2, then we have*

$$\|Kh\|_{\mathcal{H}} \leq C \left(\|h\|_{\mathcal{H}} + \frac{T^{1/2}}{\sigma^{1/4}} \|h\|_{H^1(\partial\Gamma)} \right), \quad h \in \mathcal{H}_{\text{int}}, \quad (4.12)$$

$$\|\tilde{K}h\|_{\mathcal{H}} \leq C \|h\|_{\mathcal{H}}, \quad h \in \mathcal{H} \cap H_{\Gamma}^{s_0}(\{t = z\}). \quad (4.13)$$

Note that (4.12) implies $\|K\| \leq CT^{1/2}/\sigma^{1/4}$.

Proof. For any function $h \in \mathcal{H}' \cap H_{\Gamma}^{s_0}(\{t = z\})$, let u be the solution of system (3.3) driven by h so that $u = h$ on Γ , then according to Proposition 3.1 we can conclude $u \in C^{[n/2]+2}$ in $\{t \geq z\}$.

Assume v is associated to u by (3.17), and u_{ϵ} and v_{ϵ} are defined as in the proof of Lemma 3.6, and we denote $w_{\epsilon} = u - v_{\epsilon}$. By (4.1) we know $\chi h = h$. By (4.4) we see

$$K_{\epsilon}h = w_{\epsilon}|_{\Gamma} = h - \chi v_{\epsilon}|_{\Gamma} = \chi(h - v_{\epsilon}|_{\Gamma}) = \chi w_{\epsilon},$$

so we have

$$\|Kh\|_{\mathcal{H}}^2 = \lim_{\epsilon \rightarrow 0^+} \|K_\epsilon h\|_{\mathcal{H}}^2 = \lim_{\epsilon \rightarrow 0^+} \|\chi w_\epsilon\|_{\mathcal{H}}^2.$$

Note that

$$\begin{aligned} \|\chi w_\epsilon\|_{\mathcal{H}}^2 &\leq \|(\nabla_\Gamma \chi)w_\epsilon\|_{L^2(\Gamma)}^2 + \|\chi \nabla_\Gamma w_\epsilon\|_{L^2(\Gamma)}^2 \\ &\leq \|\nabla_\Gamma \chi\|_{L^\infty(\Gamma)}^2 \|w_\epsilon\|_{L^2(\Gamma)}^2 + \|\nabla_\Gamma w_\epsilon\|_{L^2(\Gamma)}^2 \\ &\leq C \|\nabla_\Gamma \chi\|_{L^\infty(\Gamma)}^2 \|\nabla_\Gamma w_\epsilon\|_{L^2(\Gamma)}^2 + \|\nabla_\Gamma w_\epsilon\|_{L^2(\Gamma)}^2 \\ &= (1 + C \|\nabla_\Gamma \chi\|_{L^\infty(\Gamma)}^2) \|w_\epsilon\|_{\mathcal{H}}^2, \end{aligned}$$

where we used Poincaré inequality with partial zero boundary trace, so we can continue

$$\|Kh\|_{\mathcal{H}}^2 \leq (1 + C \|\nabla_\Gamma \chi\|_{L^\infty(\Gamma)}^2) \lim_{\epsilon \rightarrow 0^+} \|w_\epsilon\|_{\mathcal{H}}^2 = (1 + C \|\nabla_\Gamma \chi\|_{L^\infty(\Gamma)}^2) \lim_{\epsilon \rightarrow 0^+} \|w_\epsilon\|_{\mathcal{H}_T}^2,$$

where the last equal sign is due to Lemma 4.1. Note that on Γ_T , $w_\epsilon = u - \phi_0$, which is independent of ϵ , so

$$\|Kh\|_{\mathcal{H}}^2 \leq \mathcal{C} \|u - \phi_0\|_{\mathcal{H}_T}^2, \quad (4.14)$$

where $\mathcal{C} = (1 + C \|\nabla_\Gamma \chi\|_{L^\infty(\Gamma)}^2)$. Now, combining (4.14) with Lemma 4.1 and Lemma 4.2 we have

$$\begin{aligned} \|Kh\|_{\mathcal{H}}^2 &\leq \mathcal{C} \|u\|_{\mathcal{H}_T}^2 && \text{(by Lemma 4.1)} \\ &\leq CE(T, u) \leq \mathcal{C} [\|u\|_{\mathcal{H}}^2 + T/\sqrt{\sigma} \|u\|_{H^1(\partial\Gamma)}^2] && \text{(by Lemma 4.2)} \\ &= \mathcal{C} [\|h\|_{\mathcal{H}}^2 + T/\sqrt{\sigma} \|h\|_{H^1(\partial\Gamma)}^2], \end{aligned} \quad (4.15)$$

where $\mathcal{C} = (1 + C \|\nabla_\Gamma \chi\|_{L^\infty(\Gamma)}^2)$. This is (4.12), and we obtain $\|K\| \leq \mathcal{C}$ from (4.12). When $h \in \mathcal{H} \cap H_\Gamma^{s_0}(\{t = z\})$, we know $h|_{\partial\Gamma} = 0$, so similar to (4.12) we have

$$\|\tilde{K}h\|_{\mathcal{H}}^2 \leq C \|u\|_{\mathcal{H}_T}^2 \leq CE(T, u) \leq C \|u\|_{\mathcal{H}}^2 = C \|h\|_{\mathcal{H}}^2, \quad (4.16)$$

where $C = (1 + \|\nabla_\Gamma \chi\|_{L^\infty})^2$. This gives (4.13). We arrive at the conclusion. \square

We are ready to extend \tilde{K} and to show its boundedness with norm strictly less than 1.

Proposition 4.4. *\tilde{K} can be extended to \mathcal{H} , and it is a bounded linear operator $\mathcal{H} \rightarrow \mathcal{H}$ with $\|\tilde{K}\| < 1$.*

Proof. The linearity of \tilde{K} can be easily seen from systems (3.3) and (3.7). Recall the s_0 in Remark 3.2. It can be checked that $\mathcal{H} \cap H_\Gamma^{s_0}(\{t = z\})$ is dense in \mathcal{H} . Assume $h \in \mathcal{H}$, we choose a sequence $\{h_j\}_{j \geq 1}$ in $\mathcal{H} \cap H_\Gamma^{s_0}(\{t = z\})$ such that $h_j \rightarrow h$ under the $\|\cdot\|_{\mathcal{H}}$ -norm. By (4.13) we know $\{\tilde{K}h_j\}_{j \geq 1}$ is a Cauchy sequence in \mathcal{H} , so $\lim_{j \rightarrow \infty} \tilde{K}h_j$ exists and the limit is unique, and we define

$$\tilde{K}h := \lim_{j \rightarrow \infty} \tilde{K}h_j. \quad (4.17)$$

The norm estimate for $\|\tilde{K}\|$ in Lemma 4.3 can be improved to $\|\tilde{K}\| < 1$ by using local energy decay. Indeed, fix a constant $T_0 \in (1, T)$, then when $t > T_0$, (3.3) gives

$$\begin{cases} (\partial_t^2 - \Delta)u = 0 & \text{in } \mathbb{R}^n \times [T_0, T], \\ u = u_0, \quad u_t = u_1 & \text{on } \{t = T_0\}, \end{cases} \quad (4.18)$$

where $u_0(x) = u(x, T_0)$ and $u_1(x) = u_t(x, T_0)$. By the local energy result in [Ike05] we see that when $t - T_0$ is large enough, we have

$$\|u\|_{\mathcal{H}_t}^2 \leq \frac{\lambda E(T_0, u)}{t - C}, \quad \forall t: \max\{T_0, C\} \leq t \leq T, \quad (4.19)$$

for some real numbers λ and C independent of u and t . Here the \mathcal{H}_t -norm is defined by changing T to t in (4.5). The authors remind that originally the result in [Ike05] is for wave equations in an exterior domain with a bounded Dirichlet obstacle contained inside. But the method in [Ike05] applies directly to the case where the obstacle is an empty set, i.e. applies to the free upper space case, e.g. (4.18). See also [Tam81, Vod04].

Now let $h \in \mathcal{H} \cap H_{\mathbb{F}}^{\geq 0}(\{t = z\})$ and let u be associated with h by (3.3). By (4.19) and Lemma 4.2, when $T - T_0$ is large enough, we see that

$$C\|u\|_{\mathcal{H}_T}^2 \leq \lambda' E(T_0, u) \leq \lambda' \|u\|_{\mathcal{H}}^2$$

of some constant $\lambda' < 1$, and the second equal sign is due to the fact that $h \in \mathcal{H}$. Hence, the chain of inequalities (4.16) can be further improved: for $h \in \mathcal{H} \cap H_{\mathbb{F}}^{\geq 0}(\{t = z\})$ we have

$$\|\tilde{K}h\|_{\mathcal{H}}^2 \leq C\|u\|_{\mathcal{H}_T}^2 \leq \lambda' \|u\|_{\mathcal{H}}^2 = \lambda' \|h\|_{\mathcal{H}}^2.$$

Therefore, we can conclude $\|\tilde{K}\| < 1$. The proof is complete. \square

Remark 4.1. The authors think the arguments in [SU09b, Proof of Theorem 1] in proving $\|\tilde{K}\| < 1$ does not directly apply to our case. This is mainly because \tilde{K} is defined by a limit process, see (4.17). In the limit, the condition $\|\tilde{K}\| < 1$ becomes $\|\tilde{K}\| \leq 1$. Consequently, it is not possible to use a Neumann series argument to invert an operator of the form $I - \tilde{K}$. See Section 4.3 for details.

4.3. Reconstruction and stability result of the linearized sound speed. By definition (4.4) we have $(I - K)F = \chi BA'_0 \chi F$. When $\text{supp } F \subset \{\chi = 1\}$, we have $(I - K)F = \chi BA'_0 \chi F$. Proposition 4.4 says the operator norm of \tilde{K} is less than 1, so we can use Neumann series to recover F . Before that, we need the following claim.

Claim 4.5. $(\chi BA'_0 \chi)K = K(\chi BA'_0 \chi)$.

Proof. We have

$$\begin{aligned} (\chi BA'_0 \chi)K &= (\chi BA'_0 \chi)(I - \chi BA'_0 \chi) = \chi BA'_0 \chi - (\chi BA'_0 \chi)(\chi BA'_0 \chi) \\ &= (I - \chi BA'_0 \chi)(\chi BA'_0 \chi) = K(\chi BA'_0 \chi). \end{aligned}$$

The proof is done. \square

Proposition 4.6. *Assume $F \in \mathcal{H}_{int}$, and f satisfies $ZF = f$. Let u be the solution of the system (3.3). Then*

$$F = (I + \sum_{j \geq 0} \tilde{K}^j K) \chi B(u|_{\Sigma}), \quad f = Z[(I + \sum_{j \geq 0} \tilde{K}^j K) \chi B(u|_{\Sigma})], \quad (4.20)$$

Moreover, we have the stability

$$\|f\|_{L^2(\Omega)} \leq \mathcal{C} (\|A_0 f\|_{H^1(\Sigma)} + \frac{T^{1/2}}{\sigma^{1/4}} \|A_0 f\|_{H^1(\partial\Gamma)}), \quad (4.21)$$

where $\mathcal{C} = C\sqrt{nT}e^{CT/2}[1 + \|K\|(1 - \|\tilde{K}\|)^{-1}]$ and the constant C is independent of f , K and T .

Remark 4.2. Note that $\|\tilde{K}^j K B u\|_{\mathcal{H}} \leq \|\tilde{K}\|^j \|K\| \|B u\|_{\mathcal{H}} \lesssim \|\tilde{K}\|^j \|K\| \|u\|_{H^1(\Sigma)}$ with $0 < \|\tilde{K}\| < 1$ by Proposition 4.4. The last inequality sign is due to Proposition 3.7. Hence the infinite sum in (4.20) converges.

Proof of Proposition 4.6. Note that (4.4) gives the identity $I - K = \chi BA'_0 \chi$, and \tilde{K} is a restriction of K , so $I - \tilde{K} = \chi BA'_0 \chi$ also holds when restricted to \mathcal{H} .

For $F \in \mathcal{H}_{\text{int}}$, we denote $F_1 := (I - \chi BA'_0 \chi)F = KF$, thus $F_1 \in \mathcal{H}$ and $(I - \tilde{K})F_1 = \chi BA'_0 \chi F_1$ because $I - \tilde{K} = \chi BA'_0 \chi$ on \mathcal{H} , so

$$F_1 = (I - \tilde{K})^{-1} \chi BA'_0 \chi F_1 = (I - \tilde{K})^{-1} \chi BA'_0 \chi KF,$$

and by Claim 4.5 we obtain $F_1 = (I - \tilde{K})^{-1} K \chi BA'_0 \chi F$. Because $F_1 = (I - \chi BA'_0 \chi)F$, we have $F - \chi BA'_0 \chi F = (I - \tilde{K})^{-1} K \chi BA'_0 \chi F$, which finally gives

$$\begin{aligned} F &= [I + (I - \tilde{K})^{-1} K] \chi BA'_0 \chi F = [I + (I - \tilde{K})^{-1} K] \chi BA'_0 F \\ &= (I + \sum_{j \geq 0} \tilde{K}^j K) \chi B(u|_{\Sigma}). \end{aligned} \quad (4.22)$$

We used the fact $\chi F = F$. Note that χ , K , \tilde{K} , B and $u|_{\Sigma}$ are known, thus F can be recovered using (4.22). The invertibility of $(I - \tilde{K})$ is guaranteed by Proposition 4.4. The image of K , i.e. \mathcal{H} , falls into the domain of $(I - \tilde{K})^{-1}$, so the RHS of (4.22) is well-defined. The identity (4.22) is an analogue of [SU09b, eq. (11)]. Recall $f = ZF$, so finally we obtained (4.20).

Combining Proposition 3.7, Lemma 4.3 and the fact $\|\tilde{K}\| < 1$, we have

$$\begin{aligned} \|F\|_{\mathcal{H}} &= \|(I + (I - \tilde{K})^{-1} K) \chi Bu\|_{\mathcal{H}} \leq \|\chi Bu\|_{\mathcal{H}} + (1 - \|\tilde{K}\|)^{-1} \|K \chi Bu\|_{\mathcal{H}} \\ &\leq C \|Bu\|_{\mathcal{H}} + C (1 - \|\tilde{K}\|)^{-1} \|K\| (\|\chi Bu\|_{\mathcal{H}} + \frac{T^{1/2}}{\sigma^{1/4}} \|\chi Bu\|_{H^1(\partial\Gamma)}) \\ &\leq C [\|Bu\|_{\mathcal{H}} + (1 - \|\tilde{K}\|)^{-1} \|K\| (\|Bu\|_{\mathcal{H}} + \frac{T^{1/2}}{\sigma^{1/4}} \|\chi u\|_{H^1(\partial\Gamma)})] \\ &\leq C [1 + \|K\| (1 - \|\tilde{K}\|)^{-1}] (\|Bu\|_{\mathcal{H}} + \frac{T^{1/2}}{\sigma^{1/4}} \|u\|_{H^1(\partial\Gamma)}) \\ &\leq C \sqrt{ne}^{T/2} [1 + \|K\| (1 - \|\tilde{K}\|)^{-1}] (\|u\|_{H^1(\Sigma)} + \frac{T^{1/2}}{\sigma^{1/4}} \|u\|_{H^1(\partial\Gamma)}). \end{aligned}$$

On the another side, the norm of F satisfies (see (3.6))

$$\begin{aligned} \|F\|_{\mathcal{H}}^2 &= \int_{\Omega} |\nabla_y F(y, z, z)|^2 + |ZF(y, z, z)|^2 dx \geq \int_{\Omega} |ZF(y, z, z)|^2 dx \\ &= \int_{\Omega} |f(y, z)|^2 dx. \end{aligned}$$

Hence,

$$\|f\|_{L^2(\Omega)} \leq C \sqrt{ne}^{T/2} [1 + \|K\| (1 - \|\tilde{K}\|)^{-1}] (\|u\|_{H^1(\Sigma)} + \frac{T^{1/2}}{\sigma^{1/4}} \|u\|_{H^1(\partial\Gamma)}).$$

Recall $u|_{\Sigma} = A_0 f$. The proof is done. \square

Remark 4.3. It is worth noting that, when the dimension is odd, and when the final time T is large enough, by the Huygens' principle we know $u(\cdot, T)$ will be zero in Ω (see e.g. [Eva10, §2. eq. (31)]), hence the final condition in (3.17) will be $v = 0$ and $v_t = 0$ in Γ_T , and this matches with the values of u and u_t in Γ_T because they are all zero functions at time T . In this case, the function v will be exactly the same as u , and hence \tilde{K} will be a zero map, and so the reconstruction formula (4.20) can be simplified as

$$F = (I + K) \chi B(u|_{\Sigma}), \quad f = Z[(I + K) \chi B(u|_{\Sigma})].$$

However, in even dimensions we do not have such a conclusion.

APPENDIX A. WELL-POSEDNESS OF THE LINEAR FORWARD PROBLEM

We wish to consider the well-posedness for the wave equation in Sobolev spaces that have a negative smoothness index with respect to time. For smooth coefficients such results may be found in [Hör76, Chapter 9]. Since we need to deal with coefficients having finite regularity, we will give the required results and proofs in this appendix.

Throughout this Appendix, the notation H^s with $s \in \mathbb{R}$ stands for $H^s(\mathbb{R}^n)$. For $s \in \mathbb{R}$, consider the norm $\|f\|_{H^s} = \|J^s f\|_{L^2}$ where $J^s = (1 - \Delta_x)^{s/2}$, and the corresponding inner product as $(f, g)_{H^s} = (J^s f, J^s g)_{L^2}$. We will consider wave equations of the form

$$\partial_t^2 u - \gamma^{-1} \partial_j (\gamma a^{jk} \partial_k u) + qu = 0$$

where γ and (a^{jk}) are positive. This form includes both the standard wave equation $(\partial_t^2 - c(x)^2 \Delta_x)u = 0$ as well as the Riemannian wave equation $(\partial_t^2 - \Delta_g)u = 0$. However, in the proofs we need to use the modified inner product $(\gamma u, v)_{L^2}$ in order to make the elliptic part symmetric.

We begin with an energy estimate for smooth functions.

Lemma A.1. *Consider the operator*

$$Au = -\gamma^{-1} \partial_j (\gamma a^{jk} \partial_k u) + qu$$

where $(a^{jk}(x))$ is a symmetric matrix and one has for some $M \geq 1$

$$M^{-1} |\xi|^2 \leq a^{jk}(x) \xi_j \xi_k \leq M |\xi|^2, \quad M^{-1} \leq \gamma(x) \leq M \quad \text{a.e. in } \mathbb{R}^n. \quad (\text{A.1})$$

Let $s \in \mathbb{R}$, and assume further the Sobolev multiplier properties

$$\|[J^s, \gamma^{-1} \partial_j (\gamma a^{jk} \partial_k)]\|_{H^{s+1} \rightarrow L^2} + \|q\|_{H^{s+1} \rightarrow H^s} \leq M. \quad (\text{A.2})$$

Let $T > 0$ and let $t_0 \in [0, T]$. Then for any $u \in C_c^\infty(\mathbb{R}^n \times \mathbb{R})$ one has

$$\sum_{j=0}^1 \|\partial_t^j u\|_{L^\infty((0, T), H^{s+1-j})} \leq C(\|(\partial_t^2 + A)u\|_{L^1((0, T), H^s)} + \|u(t_0)\|_{H^{s+1}} + \|\partial_t u(t_0)\|_{H^s}),$$

where $C = C(M, T) > 0$.

Proof. Write $F = (\partial_t^2 + A)u$ and $v = J^s u$, and define the energy

$$\begin{aligned} E(t) &= \|u(t)\|_{H^s}^2 + \|\nabla_x u(t)\|_{H^s}^2 + \|\partial_t u(t)\|_{H^s}^2 \\ &= \|v(t)\|_{L^2}^2 + \|\nabla_x v(t)\|_{L^2}^2 + \|\partial_t v(t)\|_{L^2}^2. \end{aligned}$$

Thus

$$M^{-2} E(t) \leq E_1(t) := (\gamma v(t), v(t))_{L^2} + (\gamma a^{jk} \partial_j v(t), \partial_k v(t))_{L^2} + (\gamma \partial_t v(t), \partial_t v(t))_{L^2}.$$

Since $(\partial_t^2 + A)u = F$, note that the function v solves the equation

$$\partial_t^2 v - \gamma^{-1} \partial_j (\gamma a^{jk} \partial_k v) = G,$$

where

$$G = J^s F + [J^s, \gamma^{-1} \partial_j (\gamma a^{jk} \partial_k)]u - J^s(qu).$$

Differentiating $E_1(t)$ and using the equation $\partial_t^2 v - \gamma^{-1} \partial_j (\gamma a^{jk} \partial_k v) = G$ gives

$$\begin{aligned} \frac{1}{2} E_1'(t) &= (\gamma \partial_t v(t), v(t))_{L^2} + (\gamma a^{jk} \partial_j v(t), \partial_k \partial_t v(t))_{L^2} + (\gamma \partial_t^2 v(t), \partial_t v(t))_{L^2} \\ &= (\gamma \partial_t v(t), v(t))_{L^2} + (\gamma G(t), \partial_t v(t))_{L^2}. \end{aligned}$$

Note that

$$\|\gamma^{1/2} G(t)\|_{L^2} \leq M^{1/2} (\|F(t)\|_{H^s} + M \|u(t)\|_{H^{s+1}}).$$

Thus we have

$$\begin{aligned} \frac{1}{2}E_1'(t) &\leq \|\gamma^{1/2}\partial_t v(t)\|_{L^2}(\|\gamma^{1/2}v(t)\|_{L^2} + M^{1/2}\|F(t)\|_{H^s} + M^{3/2}\sqrt{E(t)}) \\ &\leq \sqrt{E_1(t)}((1 + M^{5/2})\sqrt{E_1(t)} + M^{1/2}\|F(t)\|_{H^s}). \end{aligned}$$

Since $\sqrt{E_1(t)}$ may not be smooth, we study $\sqrt{\varepsilon + E_1(t)}$ for $\varepsilon > 0$. This satisfies

$$\frac{d}{dt}\sqrt{\varepsilon + E_1(t)} = \frac{E_1'(t)}{2\sqrt{\varepsilon + E_1(t)}} \leq C_1\sqrt{\varepsilon + E_1(t)} + M^{1/2}\|F(t)\|_{H^s}$$

where $C_1 = 1 + M^{5/2}$. By Gronwall's inequality we obtain for any $\tau \in [0, T]$ that

$$\sqrt{\varepsilon + E_1(\tau)} \leq e^{C_1 T}(\sqrt{\varepsilon + E_1(t_0)} + M^{1/2} \int_{t_0}^{\tau} \|F(t)\|_{H^s} e^{-C_1 t} dt).$$

Letting $\varepsilon \rightarrow 0$ proves the required statement. \square

We next discuss sufficient conditions for (A.2). We will use the following result.

Lemma A.2. *For any $a \in H^{\alpha_0}$ with $\alpha_0 > n/2$ one has the Sobolev multiplier property*

$$\|au\|_{H^\alpha} \leq C_{\alpha_0}\|a\|_{H^{\alpha_0}}\|u\|_{H^\alpha}, \quad |\alpha| \leq \alpha_0.$$

For any $a \in H^{s_0}$ with $s_0 > n/2 + 1$ one has the commutator estimate

$$\|J^s(au) - aJ^s u\|_{L^2} \leq C_{s_0}\|\nabla a\|_{H^{s_0-1}}\|u\|_{H^{s-1}}, \quad 1 \leq |s| \leq s_0.$$

Proof. The first part follows e.g. from [Tay11, Section 13.10]. The second part for $s = 1$ follows from Calderón's commutator formula [Ste93, Section VII.3.5], and for $n/2 + 1 < s \leq s_0$ it follows from the Kato-Ponce inequality [KP88]. Since we could not locate a precise reference for this result under the given conditions, we will give a proof following [Kat75].

The commutator estimate follows if we show that the operator $T = [J^s, a]J^{1-s}$ is bounded on L^2 . First let $1 \leq s \leq s_0$. Computing the Fourier transform of Tu gives

$$\begin{aligned} \widehat{Tu}(\xi) &= \langle \xi \rangle^s (\hat{a} * \langle \xi \rangle^{1-s} \hat{u}) - \hat{a} * \langle \xi \rangle \hat{u} \\ &= \int (\langle \xi \rangle^s - \langle \eta \rangle^s) \hat{a}(\xi - \eta) \langle \eta \rangle^{1-s} \hat{u}(\eta) d\eta. \end{aligned}$$

Since the function $f(t) = (1 + t^2)^{s/2}$ satisfies $f''(t) \geq 0$ for $s \geq 1$, we have

$$|\langle \xi \rangle^s - \langle \eta \rangle^s| \leq s(\langle \xi \rangle^{s-1} + \langle \eta \rangle^{s-1})|\xi - \eta|.$$

Let h be the function satisfying $\hat{h}(\xi) = |\xi| |\hat{a}(\xi)|$. It follows that $T = T_1 + T_2$ where

$$\begin{aligned} |\widehat{T_1 u}(\xi)| &\leq s \langle \xi \rangle^{s-1} \int \hat{h}(\xi - \eta) \langle \eta \rangle^{1-s} |\hat{u}(\eta)| d\eta, \\ |\widehat{T_2 u}(\xi)| &\leq s \int \hat{h}(\xi - \eta) |\hat{u}(\eta)| d\eta. \end{aligned}$$

Now one has

$$\begin{aligned} \|\widehat{T_1 u}\|_{L^2} &\leq s \|h \mathcal{F}^{-1}\{\langle \eta \rangle^{1-s} |\hat{u}(\eta)|\}\|_{H^{s-1}} \leq s \|h\|_{H^{s-1} \rightarrow H^{s-1}} \|\hat{u}\|_{L^2}, \\ \|\widehat{T_2 u}\|_{L^2} &\leq s \|h\|_{L^\infty} \|\mathcal{F}^{-1}\{|\hat{u}(\eta)|\}\|_{L^2} \leq s \|h\|_{L^\infty} \|\hat{u}\|_{L^2}. \end{aligned}$$

Moreover, since $s_0 > n/2 + 1$, we have the estimate $\|h\|_{L^\infty} \leq C \|h\|_{H^{s_0-1}} = C \|\nabla a\|_{H^{s-1}}$ and $\|h\|_{H^{s-1} \rightarrow H^{s-1}} \leq C \|h\|_{H^{s_0-1}} = C \|\nabla a\|_{H^{s_0-1}}$. The required commutator estimate for $1 \leq s \leq s_0$ follows by using the Plancherel theorem.

Now let $-s_0 \leq s \leq -1$. We use duality and compute

$$\|[J^s, a]u\|_{L^2} = \sup_{\|v\|_{L^2}=1} ([J^s, a]u, v)$$

$$\begin{aligned}
&= \sup_{\|v\|_{L^2}=1} (J^{s-1}u, J^{1-s}aJ^s v - J^1av) \\
&\leq \sup_{\|v\|_{L^2}=1} \|u\|_{H^{s-1}} \|J^1[J^{-s}, a]J^s v\|_{L^2}.
\end{aligned}$$

Note that

$$\|J^1[J^{-s}, a]J^s v\|_{L^2} \leq \|[J^{-s}, a]J^s v\|_{L^2} + \|[J^{-s}, \nabla a]J^s v\|_{L^2} + \|[J^{-s}, a]\nabla J^s v\|_{L^2}.$$

We use the commutator estimate for $[J^{-s}, a]$ to conclude that the first and last terms on the right are $\lesssim \|v\|_{L^2}$. For the middle term we use that

$$\|(J^{-s}(\nabla a) - (\nabla a)J^{-s})J^s v\|_{L^2} \leq \|(\nabla a)J^s v\|_{H^{-s}} + \|(\nabla a)v\|_{L^2} \lesssim \|v\|_{L^2}.$$

This concludes the proof. \square

Remark A.1. Using Lemma A.2, one can check that a sufficient condition for (A.2) is the following: if $s_0 > n/2$ and

$$\gamma \geq 1/E, \quad a^{jk}\xi_j\xi_k \geq |\xi|^2/E, \quad \|\gamma\|_{H^{s_0+1}} + \|a^{jk}\|_{H^{s_0+1}} + \|q\|_{H^{s_0}} \leq E, \quad (\text{A.3})$$

then (A.2) holds whenever $|s| \leq s_0$ for some $M = M(n, s, E)$.

We next give a simple solvability result which follows directly from the energy estimate in Lemma A.1 and duality. In fact under suitable regularity assumptions for the coefficients one has a unique solution in the class $C([0, T], H^{s+1}) \cap C^1([0, T], H^s)$, see e.g. [Smi98].

Lemma A.3. *Assume (A.1) and (A.3) for some $s_0 > n/2$, and let $-s_0 \leq s \leq s_0$. Given any $u_0 \in H^{s+1}$, $u_1 \in H^s$ and $F \in L^1((0, T), H^s)$, there is $u \in L^\infty((0, T), H^{s+1})$ which is a weak solution of*

$$(\partial_t^2 + A)u = F \text{ in } \mathbb{R}^n \times (0, T), \quad u(0) = u_0, \quad \partial_t u(0) = u_1,$$

in the sense that

$$(\gamma u, (\partial_t^2 + A)\varphi) = (\gamma F, \varphi) - (\gamma u_0, \partial_t \varphi(0)) + (\gamma u_1, \varphi(0))$$

for any $\varphi \in C_c^\infty(\mathbb{R}^n \times [0, T])$. This solution satisfies

$$\|u\|_{L^\infty((0, T), H^{s+1})} \leq C(\|F\|_{L^1((0, T), H^s)} + \|u_0\|_{H^{s+1}} + \|u_1\|_{H^s})$$

where $C = C(n, s, M, E, T)$.

Proof. Let $X = C_c^\infty(\mathbb{R}^n \times [0, T])$, i.e. any $\varphi \in X$ vanishes near $t = T$. Now if $\varphi_1, \varphi_2 \in X$ and $(\partial_t^2 + A)\varphi_1 = (\partial_t^2 + A)\varphi_2$, then Lemma A.1 with $t_0 = T$ implies that $\varphi_1 = \varphi_2$. We may thus define the linear functional

$$\ell : (\partial_t^2 + A)X \rightarrow \mathbb{R}, \quad \ell((\partial_t^2 + A)\varphi) = (\gamma F, \varphi) - (\gamma u_0, \partial_t \varphi(0)) + (\gamma u_1, \varphi(0)).$$

Applying Lemma A.1 with $t_0 = T$ again, one has

$$|\ell((\partial_t^2 + A)\varphi)| \lesssim (\|F\|_{L^1 H^s} + \|u_0\|_{H^{s+1}} + \|u_1\|_{H^s}) \|(\partial_t^2 + A)\varphi\|_{L^1 H^{-s-1}}.$$

By the Hahn-Banach theorem ℓ extends as a bounded linear functional on $L^1 H^{-s-1}$, and by duality it can be represented by a function $u \in L^\infty H^{s+1}$ satisfying

$$\ell((\partial_t^2 + A)\varphi) = (\gamma u, (\partial_t^2 + A)\varphi), \quad \varphi \in X,$$

and

$$\|u\|_{L^\infty H^{s+1}} \lesssim \|F\|_{L^1 H^s} + \|u_0\|_{H^{s+1}} + \|u_1\|_{H^s}.$$

This proves the result. \square

We next move to solutions that are in negative Sobolev spaces with respect to time. For any integer $k \geq 0$ and any $\alpha \in \mathbb{R}$, we let $H^{-k}((-T, T), H^\alpha(\mathbb{R}^n))$ be the dual of $H_0^k((-T, T), H^{-\alpha}(\mathbb{R}^n))$.

Lemma A.4. *Assume (A.1) and (A.3) for some $s_0 > n/2 + 2$. Let $k \geq 0$ be an integer, and let $-s_0 \leq s \leq s_0$. Given any $F \in H^{-k}((-T, T), H^s)$ with $F|_{\{t < 0\}} = 0$, there is a unique distributional solution $u \in H^{-k}((-T, T), H^{s+1})$ of*

$$(\partial_t^2 + A)u = F \text{ in } \mathbb{R}^n \times (-T, T), \quad u|_{\{t < 0\}} = 0. \quad (\text{A.4})$$

This solution satisfies

$$\|u\|_{H^{-k}((-T, T), H^{s+1})} \leq C \|F\|_{H^{-k}((-T, T), H^s)}$$

where $C = C(n, k, s, M, E, T)$.

Proof. Consider the space $X = C_c^\infty(\mathbb{R}^n \times [0, T])$ and the linear functional

$$\ell : (\partial_t^2 + A)X \rightarrow \mathbb{R}, \quad \ell((\partial_t^2 + A)\varphi) = (\gamma F, E\varphi)$$

where $E\varphi$ is any smooth extension of φ from $\mathbb{R}^n \times [0, T]$ to $\mathbb{R}^n \times (-T, T)$. The right hand side is well defined since $F|_{\{t < 0\}} = 0$. Choosing E to be a bounded extension operator on H^k , one has

$$|\ell((\partial_t^2 + A)\varphi)| \lesssim \|F\|_{H^{-k}((-T, T), H^s)} \|\varphi\|_{H^k((0, T), H^{-s})}.$$

Applying Lemma A.1 with $t_0 = T$ to $\varphi, \partial_t \varphi, \dots, \partial_t^k \varphi$ gives that

$$|\ell((\partial_t^2 + A)\varphi)| \lesssim \|F\|_{H^{-k}((-T, T), H^s)} \|(\partial_t^2 + A)\varphi\|_{H^k((0, T), H^{-s-1})}.$$

Using the Hahn-Banach theorem and duality, there is $u \in H^{-k}((-T, T), H^{s+1})$ satisfying

$$(\gamma u, (\partial_t^2 + A)\varphi) = \ell((\partial_t^2 + A)\varphi) = (\gamma F, E\varphi), \quad \varphi \in X,$$

and

$$\|u\|_{H^{-k}((-T, T), H^{s+1})} \lesssim \|F\|_{H^{-k}((-T, T), H^s)}.$$

Given any $\psi \in C_c^\infty(\mathbb{R}^n \times (-T, 0))$ we may find a solution of $(\partial_t^2 + A)\varphi = \psi$ with $\varphi(0) = \partial_t \varphi(0) = 0$. Then we have $(\gamma u, \psi) = (\gamma F, E\varphi)$ for any extension of φ , and choosing the zero extension gives $u = 0$ for $t < 0$. Thus u is a distributional solution of (A.4).

It remains to prove that any $u \in H^{-k}((-T, T), H^{s+1})$ satisfying

$$(\partial_t^2 + A)u = 0 \text{ in } \mathbb{R}^n \times (-T, T), \quad u|_{\{t < 0\}} = 0,$$

must be identically zero. Fix some $F \in C_c^\infty(\mathbb{R}^n \times (-T, T))$ and choose φ to be a solution of

$$(\partial_t^2 + A)\varphi = F \text{ in } \mathbb{R}^n \times (-T, T), \quad \varphi(T) = \partial_t \varphi(T) = 0.$$

Then $\varphi \in L^\infty H^{s_0+1}$. For any $\chi \in C^\infty([-T, T])$ with $\chi = 1$ near $[0, T]$ and $\chi = 0$ near $t = -T$ one has

$$(u, F) = (u, (\partial_t^2 + A)\varphi) = (u, (\partial_t^2 + A)(\chi(t)\varphi)).$$

Thanks to Assumption (A.3), we have $\gamma, a^{jk} \in C^{1,1}(\mathbb{R}^n)$ for $j, k = 1, 2, \dots, n$. By [Smi98, Theorem 4.7], we conclude that φ belongs to $L^2 H^{s_0+3}$. We claim that φ has more regularity in time. Indeed, since $\partial_t^2 \varphi = F - A\varphi$, we deduce $\varphi \in H^2 H^{s_0+1}$. We can iterate this argument so that by means of the Morrey's embedding one has $\varphi \in C^2((-T, T); \mathbb{R}^n)$ and hence $\chi(t)\varphi$ so. Since $\varphi(T) = \partial_t \varphi(T) = 0$ and $\chi(-T) = \chi'(-T) = 0$, one can thus integrate by parts and obtain that

$$(u, F) = ((\partial_t^2 + A)u, \chi(t)\varphi) = 0.$$

Since F was arbitrary we obtain $u = 0$ as required. \square

As a consequence, we prove the well-posedness of (1.3) and (1.6).

Proof of Proposition 1.2. We start by setting $\mathcal{U} := \mathcal{V} + H_1(t - z)$. Thus, \mathcal{U} is a solution to (1.6) if and only if \mathcal{V} is a solution to the IVP

$$(\eta(x) \partial_t^2 - \Delta)\mathcal{V} = -(1 - \eta)\delta(t - z) \text{ in } \mathbb{R}^{n+1}, \quad \mathcal{V}|_{\{t < -1\}} = 0. \quad (\text{A.5})$$

The source on the LHS belongs to $H_{loc}^{-1}(\mathbb{R}; L^2(\mathbb{R}^n))$. Hence, Lemma A.4 with the configuration

$$\gamma = \eta, \quad a^{jk} = \eta^{-1} \delta_{jk}, \quad s = 0, \quad k = -1 \quad \text{and} \quad s_0 = [n/2] + 3$$

ensures that there exists a unique distributional solution $\mathcal{V} \in H^{-1}((-T, T); H^1(\mathbb{R}^n))$ to (A.5) with

$$\|\mathcal{V}\|_{H^{-1}((-T, T); H^1(\mathbb{R}^n))} \lesssim \|(1 - \eta)\delta(t - z)\|_{H^{-1}((-T, T); L^2(\mathbb{R}^n))} \lesssim C(T, M).$$

In addition, one can verify that $\mathcal{U} := \mathcal{V} + H_1(t - z)$ is a solution to the IVP (1.6). Moreover, the above estimate remains valid with \mathcal{U} instead of \mathcal{V} . A straightforward computation shows that $\partial_t^2 \mathcal{U}$ satisfies (1.3). By uniqueness of distributional solutions we deduce $\tilde{\mathcal{U}} = \partial_t^2 \mathcal{U}$. Thanks to the trace theorem one has the following sequence of inequalities

$$\|\tilde{\mathcal{U}}|_{\Sigma_T}\|_{H^{-3}((-T, T); H^{1/2}(S^{n-1}))} \lesssim \|\mathcal{U}\|_{H^{-1}((-T, T); H^1(\mathbb{R}^n))} \lesssim C(T, M).$$

This finishes the proof of Proposition 1.2.

Proof of Proposition 3.1. We start by proving the uniqueness. It is reduced to proving that zero is the unique distributional solution to the homogeneous equation

$$(\partial_t^2 - \Delta)U = 0 \text{ in } \mathbb{R}^{n+1}, \quad U|_{\{t < -1\}} = 0,$$

which is true by [Hör83, Theorem 23.2.7]. The method we shall use to prove the existence is the so-called progressing wave expansion method, see e.g. [Shi85, Lemma 1] and [RU14, Theorem 1]. For any $j \geq 0$, define

$$s_+^j = \begin{cases} s^j, & s \geq 0, \\ 0, & s < 0. \end{cases}$$

Note that $s_+^0 = H(s)$ is the unidimensional Heaviside function at $s \in \mathbb{R}$. Let $f \in H_{\bar{\Omega}}^{s_0}(\mathbb{R}^n)$. Assume that $f \in C_c^m(\mathbb{R}^n)$, where $m \in \mathbb{N}$ (which will be fixed later) with $s_0 > n/2 + m$. This is always possible due to Morrey's inequality and the fact that $\text{supp}(f) \subset \bar{\Omega}$. Consider $N \in \mathbb{N}$ and suppose that the solutions to $(\partial_t^2 - \Delta)U = -f(x)\delta(t - z)$ have the following ansatz

$$U(x, t) = \sum_{j=0}^N a_j(x)(t - z)_+^j + R_N(x, t), \quad (\text{A.6})$$

where the coefficients a_j , $j = 0, 1, \dots, N$, and the remainder term R_N satisfy the initial value conditions

$$R_N|_{t < -1} = 0, \quad a_j|_{z < -1} = 0, \quad j = 0, 1, \dots, N.$$

A straightforward computation shows that the remainder term R_N must hold

$$\begin{aligned} (\partial_t^2 - \Delta)R_N(x, t) &= -(2\partial_z a_0 + f)\delta(t - z) + (\Delta a_N)(t - z)_+^N \\ &\quad - \sum_{j=0}^{N-1} (2(j+1)\partial_z a_{j+1} - \Delta a_j)(t - z)_+^j. \end{aligned} \quad (\text{A.7})$$

The task now is proving the existence of the coefficients $(a_j)_{j=0}^N$ and R_N , satisfying the recursive identity (A.7). One expects to get a smoother remainder term R_N when N is

large enough. This can be done by dropping most of the non-smooth terms on the right of (A.7). By standard ODE techniques, it is easy to see that if

$$a_0(y, z) = -\frac{1}{2} \int_{-\infty}^0 f(y, s + z) ds,$$

and for $k = 0, 1, \dots, N - 1$:

$$a_{k+1}(y, z) = \frac{1}{2(k+1)} \int_{-\infty}^0 (\Delta a_k)(y, s + z) ds,$$

then the remainder term R_N has to satisfy

$$(\partial_t^2 - \Delta)R_N(x, t) = (\Delta a_N)(t - z)_+^N, \quad R_N|_{\{t < -1\}} = 0. \quad (\text{A.8})$$

Note that $a_k \in C^{m-2k}(\mathbb{R}^n)$ and

$$\|a_k\|_{L^\infty(\mathbb{R}^n)} \lesssim \|f\|_{C^{m-2k}(\mathbb{R}^n)}.$$

Setting

$$\beta = \min\{m - 2N - 2, N - 1\},$$

we deduce that $r(y, z, t) := (\Delta a_N)(t - z)_+^N$ belongs to $C^\beta(\mathbb{R}^{n+1})$. In particular, r belongs to $H_{loc}^{\beta_1}(\mathbb{R}; H_{loc}^{\beta_2}(\mathbb{R}^n))$ with $\beta_j \geq 0$ and $\beta_1 + \beta_2 = \beta$. Here both β_1 and β_2 will be fixed later. Since $\partial_t^2 - \Delta$ is a strictly hyperbolic operator, [Hör76, Theorems 9.3.1 and 9.3.2] ensures that there exists a unique solution $R_N \in H_{loc}^{\beta_1+1}(\mathbb{R}; H_{loc}^{\beta_2}(\mathbb{R}^n))$ to (A.8) such that for any given $T > -1$ we have

$$\|R_N\|_{H^{\beta_1+1}((-1, T]; H^{\beta_2}(\mathbb{R}^n))} \lesssim \|(\Delta a_N)(t - z)_+^N\|_{H^{\beta_1}((-1, T]; H^{\beta_2}(\mathbb{R}^n))}.$$

We claim that $R_N \in C^2(\mathbb{R}^{n+1})$ by suitably choosing the parameters m , N , β_1 , and β_2 . Indeed, by [Hör83, Theorem B.2.8/Vol III], this follows if for instance

$$\frac{n+3}{2} < \beta_1 + \beta_2 = \beta \quad \text{and} \quad \frac{3}{2} < \beta_1,$$

and furthermore

$$\begin{aligned} \|R_N\|_{C^2((-1, T] \times \mathbb{R}^n)} &\lesssim \|R_N\|_{H^{\beta_1+1}((-1, T]; H^{\beta_2}(\mathbb{R}^n))} \\ &\lesssim \|(\Delta a_N)(t - z)_+^N\|_{H^{\beta_1}((-1, T]; H^{\beta_2}(\mathbb{R}^n))}. \end{aligned} \quad (\text{A.9})$$

Equating the parameters involved in the definition of β , that is, $m - 2N - 2 = N - 1$; allow us to choose $m = 3N + 1$, and hence $\beta = N - 1$. We distinguish two cases:

- When n is even we consider

$$N = \frac{n+6}{2}, \quad m = \frac{3}{2}n + 10, \quad \beta_1 = 2, \quad \beta_2 = \frac{n}{2}.$$

- When n is odd we consider

$$N = \frac{n+7}{2}, \quad m = \frac{3}{2}(n+1) + 10, \quad \beta_1 = 2, \quad \beta_2 = \frac{n+1}{2}.$$

The desired claim is proved by combining the above choices with (A.9). On the other hand, by (A.6) we deduce that U can be written as follows

$$U(x, t) = u(x, t)H(t - z),$$

where clearly

$$u(x, t) = \sum_{j=0}^N a_j(x)(t - z)_+^j + R_N(x, t)$$

is of class C^2 in the region $\{t \geq z\}$. This shows that u satisfies all the properties stated in Proposition 3.1. We point out that taking s_0 large enough we can make another choice of the parameters β_1, β_2, m and N so that u will be in C^K in the region $\{t \geq z\}$ for arbitrary $K \geq 3$. This finishes the proof of Proposition 3.1.

Relation between \mathcal{A} and $\tilde{\mathcal{A}}$. Recall that $\tilde{\mathcal{A}}(\eta) = \tilde{\mathcal{U}}|_{\Sigma_T}$ and $\mathcal{A}(\eta) = \mathcal{U}|_{\Sigma_T}$, where $\tilde{\mathcal{U}} = \partial_t^2 \mathcal{U}$ and both \mathcal{U} and $\tilde{\mathcal{U}}$ vanish for $t < -1$. The following result can be used to estimate \mathcal{A} terms of $\tilde{\mathcal{A}}$ and vice versa. This is quite standard, but one needs some care in the case of Sobolev spaces with negative smoothness index.

Lemma A.5. *Let \mathcal{M} be a compact smooth manifold without boundary. Let $T > 1$, let $k \in \mathbb{Z}$, and let $\alpha \in \mathbb{R}$. There is $C > 0$ such that*

$$C^{-1} \|u\|_{H^{k+2}((-T, T), H^\alpha(\mathcal{M}))} \leq \|\partial_t^2 u\|_{H^k((-T, T), H^\alpha(\mathcal{M}))} \leq C \|u\|_{H^{k+2}((-T, T), H^\alpha(\mathcal{M}))}$$

for any $u \in H^k((-T, T), H^\alpha(\mathcal{M}))$ satisfying $u|_{\{t < -1\}} = 0$.

The proof uses the following Poincaré type inequality.

Lemma A.6. *Let $a < a' < b$ and let $k \in \mathbb{Z}$. There is $C > 0$ so that*

$$\|u\|_{H^{k+1}((a, b))} \leq C \|\partial_t u\|_{H^k((a, b))}, \quad u \in C_c^\infty((a', b)).$$

Proof. We begin with a simple Poincaré inequality: for any $u \in C_c^\infty((a, b))$ one has

$$\int_a^b u^2 dt = \int_a^b \partial_t(t-b)u^2 dt = -2 \int_a^b (t-b)u \partial_t u dt \leq \frac{1}{2} \int_a^b u^2 dt + 2 \int_a^b (t-b)^2 (\partial_t u)^2 dt.$$

Absorbing one term to the left hand side gives the inequality

$$\|u\|_{L^2((a, b))} \leq C \|\partial_t u\|_{L^2((a, b))}, \quad u \in C_c^\infty((a, b)).$$

This shows that for any $k \geq 0$

$$\|u\|_{H^{k+1}((a, b))}^2 = \sum_{j=0}^{k+1} \|\partial_t^j u\|_{L^2((a, b))}^2 \leq C \|\partial_t u\|_{H^k((a, b))}^2.$$

Suppose now that $k \leq -1$. If $v \in C_c^\infty((a, b))$, define

$$w(t) = - \int_t^b v(s) ds.$$

Then $\partial_t w = v$. Using that $|w(t)| \leq C \|v\|_{L^2((a, b))}$ for $t \in [a, b]$, one has

$$\|w\|_{H^{-k}((a, b))} \leq C \|v\|_{H^{-k-1}((a, b))}.$$

For any $u \in C_c^\infty((a, b))$ and $v \in C_c^\infty((a, b))$, since $w(b) = 0$ we have

$$\int_a^b uv dt = \int_a^b u \partial_t w dt = - \int_a^b (\partial_t u) w dt.$$

Here w is in $H_0^{-k}([a, b])$ but not necessarily in $H_0^{-k}((a, b))$, so we cannot directly use duality to get a bound in terms of $\|\partial_t u\|_{H^k}$. However, if we use the fact that u vanishes for $a \leq t \leq a'$, we have

$$\int_a^b (\partial_t u) w dt = \int_a^b (t-a)^{-|k|} (\partial_t u) (t-a)^{|k|} w dt \leq \|(t-a)^{-|k|} \partial_t u\|_{H^k} \|(t-a)^{|k|} w\|_{H_0^{-k}}.$$

Now $\|(t-a)^{-|k|} \partial_t u\|_{H^k} \leq C \|\partial_t u\|_{H^k}$ and $\|(t-a)^{|k|} w\|_{H_0^{-k}} \leq C \|w\|_{H^{-k}} \leq C \|v\|_{H^{-k-1}}$. We have proved that

$$\int_a^b uv dt \leq C \|\partial_t u\|_{H^k} \|v\|_{H^{-k-1}}.$$

Since $H^{k+1}((a, b))$ is the dual of $H_0^{-k-1}((a, b))$, we obtain the desired inequality also for $k \leq -1$. \square

Proof of Lemma A.5. By density argument, it is enough to prove the statement is true for $u \in C^\infty([-T, T], C^\infty(\mathcal{M}))$ with $u|_{\{t < -1\}} = 0$. Clearly for $k \geq 0$

$$\|\partial_t^2 u\|_{H^k((-T, T), H^\alpha)}^2 \leq \|u\|_{H^{k+2}((-T, T), H^\alpha)}^2$$

and by duality for $k \leq -2$

$$\|\partial_t^2 u\|_{H^k((-T, T), H^\alpha)} = \sup_{\|v\|_{H_0^{-k}}=1} (\partial_t^2 u, v) = \sup_{\|v\|_{H_0^{-k}}=1} (u, \partial_t^2 v) \leq C \|u\|_{H^{k+2}((-T, T), H^\alpha)}.$$

For the remaining case $k = -1$, note that

$$\|\partial_t^2 u\|_{H^{-1}((-T, T), H^\alpha)} = \sup_{\|v\|_{H_0^1}=1} (\partial_t^2 u, v) = \sup_{\|v\|_{H_0^1}=1} (\partial_t u, \partial_t v) \leq C \|\partial_t u\|_{L^2((-T, T), H^\alpha)}$$

and use that $\|\partial_t u\|_{L^2((-T, T), H^\alpha)} \leq \|u\|_{H^1((-T, T), H^\alpha)}$.

To prove the converse inequality, let g be some Riemannian metric on \mathcal{M} and let $(\varphi_l)_{l=1}^\infty$ be an orthonormal basis of $L^2(\mathcal{M})$ consisting of eigenfunctions of $-\Delta_g$. Write $u_l(t) = (u(t), \varphi_l)_{L^2(\mathcal{M})}$ and $\langle l \rangle = (1 + l^2)^{1/2}$. If $m \geq 0$ one has

$$\begin{aligned} \|u\|_{H^m((-T, T), H^\alpha)}^2 &= \sum_{j=0}^m \int_{-T}^T \|\partial_t^j u(t)\|_{H^\alpha}^2 dt = \sum_{j=0}^m \sum_{l=1}^\infty \int_{-T}^T \langle l \rangle^{2\alpha} |\partial_t^j u_l(t)|^2 dt \\ &= \sum_{l=1}^\infty \langle l \rangle^{2\alpha} \|u_l\|_{H^m((-T, T))}^2. \end{aligned} \quad (\text{A.10})$$

For $m \leq -1$, using (A.10) to $\|v\|_{H^{-m}((-T, T), H^{-\alpha})}$ gives

$$\begin{aligned} \|u\|_{H^m((-T, T), H^\alpha)} &= \sup_{\|v\|_{H_0^{-m}((-T, T), H^{-\alpha})}=1} (u, v)_{L^2} \\ &= \sup_{v \in H_0^{-m}, \sum_l \langle l \rangle^{-2\alpha} \|v_l\|_{H^{-m}}^2 = 1} \int_{-T}^T \sum_{l=1}^\infty \langle l \rangle^\alpha u_l(t) \langle l \rangle^{-\alpha} v_l(t) dt \\ &\leq \sup_{v \in H_0^{-m}, \sum_l \langle l \rangle^{-2\alpha} \|v_l\|_{H^{-m}}^2 = 1} \sum_{l=1}^\infty \|\langle l \rangle^\alpha u_l\|_{H^m((-T, T))} \|\langle l \rangle^{-\alpha} v_l\|_{H_0^{-m}((-T, T))} \\ &\leq \left(\sum_{l=1}^\infty \langle l \rangle^{2\alpha} \|u_l\|_{H^m((-T, T))}^2 \right)^{1/2}. \end{aligned}$$

Choosing v_l so that

$$\|\langle l \rangle^{-\alpha} v_l\|_{H_0^{-m}((-T, T))} = \frac{\|\langle l \rangle^\alpha u_l\|_{H^m((-T, T))}}{\left(\sum_{l=1}^\infty \langle l \rangle^{2\alpha} \|u_l\|_{H^m((-T, T))}^2 \right)^{1/2}}$$

gives that

$$\|u\|_{H^m((-T, T), H^\alpha)}^2 = \sum_{l=1}^\infty \langle l \rangle^{2\alpha} \|u_l\|_{H^m((-T, T))}^2 \quad (\text{A.11})$$

for any $m \in \mathbb{Z}$.

Now let $k \in \mathbb{Z}$. Since $u_l(t) = 0$ for $t < -1$, we obtain from (A.11) and Lemma A.6 (used twice) that

$$\|u\|_{H^{k+2}((-T,T),H^\alpha)}^2 \leq C \sum_{l=1}^{\infty} \langle l \rangle^{2\alpha} \|\partial_t^2 u_l\|_{H^k((-T,T))}^2 \leq C \|\partial_t^2 u\|_{H^k((-T,T),H^\alpha)}^2.$$

This concludes the proof. \square

REFERENCES

- [AKK17] M. Agranovsky, P. Kuchment, and L. Kunyansky, *On reconstruction formulas and algorithms for the thermoacoustic tomography*, Photoacoustic imaging and spectroscopy, 2017, pp. 89–102.
- [BCL⁺20] J. A. Barceló, C. Castro, T. Luque, C. J. Meroño, A. Ruiz, and M. d. l. C. Vilela, *Uniqueness for the inverse fixed angle scattering problem*, J. Inverse Ill-Posed Probl. **28** (2020), no. 4, 465–470. MR4129366
- [BLM89] A. Bayliss, Y. Y. Li, and C. S. Morawetz, *Scattering by a potential using hyperbolic methods*, Math. Comp. **52** (1989), no. 186, 321–338. MR958869
- [BS91] G. Bao and W. W. Symes, *A trace theorem for solutions of linear partial differential equations*, Math. Methods Appl. Sci. **14** (1991), no. 8, 553–562. MR1129188
- [Eva10] L. C. Evans, *Partial differential equations*, Second, Graduate Studies in Mathematics, vol. 19, American Mathematical Society, Providence, RI, 2010. MR2597943
- [FR05] D. Finch and Rakesh, *Trace identities for solutions of the wave equation with initial data supported in a ball*, Math. Methods Appl. Sci. **28** (2005), no. 16, 1897–1917. MR2170771
- [HKN08] Y. Hristova, P. Kuchment, and L. Nguyen, *Reconstruction and time reversal in thermoacoustic tomography in acoustically homogeneous and inhomogeneous media*, Inverse problems **24** (2008), no. 5, 055006.
- [Hör76] L. Hörmander, *Linear partial differential operators: 4th printing*, Springer, 1976.
- [Hör83] L. Hörmander, *The analysis of linear partial differential operators, vol. i-iv*, Springer-Verlag, Berlin, 1983.
- [Hri09] Y. Hristova, *Time reversal in thermoacoustic tomography—an error estimate*, Inverse Problems **25** (2009), no. 5, 055008, 14. MR2501026
- [Ike05] R. Ikehata, *Local energy decay for linear wave equations with variable coefficients*, J. Math. Anal. Appl. **306** (2005), no. 1, 330–348. MR2132904
- [Kat75] T. Kato, *Quasi-linear equations of evolution, with applications to partial differential equations*, Spectral theory and differential equations (Proc. Sympos., Dundee, 1974; dedicated to Konrad Jörgens), 1975, pp. 25–70. Lecture Notes in Math., Vol. 448. MR0407477
- [KK08] P. Kuchment and L. Kunyansky, *Mathematics of thermoacoustic tomography*, European Journal of Applied Mathematics **19** (2008), no. 2, 191–224.
- [KK11] P. Kuchment and L. Kunyansky, *Mathematics of photoacoustic and thermoacoustic tomography*, Handbook of mathematical methods in imaging, 2011, pp. 817–865.
- [KP88] T. Kato and G. Ponce, *Commutator estimates and the Euler and Navier-Stokes equations*, Comm. Pure Appl. Math. **41** (1988), no. 7, 891–907. MR951744
- [KRS20] V. Krishnan, Rakesh, and S. Senapati, *Stability for a formally determined inverse problem for a hyperbolic pde with space and time dependent coefficients*, arXiv:2010.11726 (2020).
- [LLT86] I. Lasiecka, J.-L. Lions, and R. Triggiani, *Nonhomogeneous boundary value problems for second order hyperbolic operators*, J. Math. Pures Appl. (9) **65** (1986), no. 2, 149–192. MR867669
- [Ma21] S. Ma, *Regularity and energy of second order hyperbolic boundary value problems on non-timelike hypersurfaces*, arXiv:2107.09762 (2021).
- [Mer18] C. J. Meroño, *Recovery of singularities in inverse scattering*, Thesis (Ph.D.)—Universidad Autónoma de Madrid, Madrid, 2018.
- [MPMS21] C. J. Meroño, L. Potenciano-Machado, and M. Salo, *The fixed angle scattering problem with a first order perturbation*, Ann. Henri Poincaré (2021).
- [MS21] S. Ma and M. Salo, *Fixed angle inverse scattering in the presence of a riemannian metric*, J. Inverse Ill-Posed Probl. (2021).
- [RS20a] Rakesh and M. Salo, *Fixed angle inverse scattering for almost symmetric or controlled perturbations*, SIAM J. Math. Anal. **52** (2020), no. 6, 5467–5499. MR4170189
- [RS20b] Rakesh and M. Salo, *The fixed angle scattering problem and wave equation inverse problems with two measurements*, Inverse Problems **36** (2020), no. 3, 035005, 42. MR4068234

- [RU14] Rakesh and G. Uhlmann, *Uniqueness for the inverse backscattering problem for angularly controlled potentials*, Inverse Problems **30** (2014), no. 6, 065005, 24. MR3224125
- [Rui01] A. Ruiz, *Recovery of the singularities of a potential from fixed angle scattering data*, Comm. Partial Differential Equations **26** (2001), no. 9-10, 1721–1738. MR1865943
- [Shi85] T. Shiota, *An inverse problem for the wave equation with first order perturbation*, Amer. J. Math. **107** (1985), no. 1, 241–251. MR778095
- [Smi98] H. F. Smith, *A parametrix construction for wave equations with $C^{1,1}$ coefficients*, Ann. Inst. Fourier (Grenoble) **48** (1998), no. 3, 797–835. MR1644105
- [Ste92] P. Stefanov, *Generic uniqueness for two inverse problems in potential scattering*, Comm. Partial Differential Equations **17** (1992), no. 1-2, 55–68. MR1151256
- [Ste93] E. M. Stein, *Harmonic analysis: real-variable methods, orthogonality, and oscillatory integrals*, Princeton Mathematical Series, vol. 43, Princeton University Press, Princeton, NJ, 1993. With the assistance of Timothy S. Murphy, Monographs in Harmonic Analysis, III. MR1232192
- [SU09a] P. Stefanov and G. Uhlmann, *Linearizing non-linear inverse problems and an application to inverse backscattering*, J. Funct. Anal. **256** (2009), no. 9, 2842–2866. MR2502425
- [SU09b] P. Stefanov and G. Uhlmann, *Thermoacoustic tomography with variable sound speed*, Inverse Problems **25** (2009), no. 7, 075011, 16. MR2519863
- [SU13] P. Stefanov and G. Uhlmann, *Multiwave methods via ultrasound*, Inverse problems and applications: inside out. II, 2013, pp. 271–323. MR3098660
- [Tam81] H. Tamura, *On the decay of local energy for wave equations with time-dependent potentials*, J. Math. Soc. Japan **33** (1981), no. 4, 605–618. MR630627
- [Tay11] M. E. Taylor, *Partial differential equations III. Nonlinear equations*, Second, Applied Mathematical Sciences, vol. 117, Springer, New York, 2011. MR2744149
- [Vod04] G. Vodev, *Local energy decay of solutions to the wave equation for nontrapping metrics*, Ark. Mat. **42** (2004), no. 2, 379–397. MR2101395

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