

JYX



This is a self-archived version of an original article. This version may differ from the original in pagination and typographic details.

Author(s): Kow, Pu-Zhao; Uhlmann, Gunther; Wang, Jenn-Nan

Title: Optimality of Increasing Stability for an Inverse Boundary Value Problem

Year: 2021

Version: Accepted version (Final draft)

Copyright: © 2021 Society for Industrial and Applied Mathematics

Rights: In Copyright

Rights url: <http://rightsstatements.org/page/InC/1.0/?language=en>

Please cite the original version:

Kow, P.-Z., Uhlmann, G., & Wang, J.-N. (2021). Optimality of Increasing Stability for an Inverse Boundary Value Problem. *SIAM Journal on Mathematical Analysis*, 53(6), 7062-7080.

<https://doi.org/10.1137/21M1402169>

OPTIMALITY OF INCREASING STABILITY FOR AN INVERSE BOUNDARY VALUE PROBLEM

PU-ZHAO KOW, GUNTHER UHLMANN, AND JENN-NAN WANG

ABSTRACT. In this work we study the optimality of increasing stability of the inverse boundary value problem (IBVP) for the Schrödinger equation. The rigorous justification of increasing stability for the IBVP for the Schrödinger equation were established by Isakov [Isa11] and by Isakov, Nagayasu, Uhlmann, Wang of the paper [INUW14]. In [Isa11], [INUW14], the authors showed that the stability of this IBVP increases as the frequency increases in the sense that the stability estimate changes from a logarithmic type to a Hölder type. In this work, we prove that the instability changes from an exponential type to a Hölder type when the frequency increases. This result verifies that results in [Isa11], [INUW14] are optimal.

1. INTRODUCTION

In this paper, we study the instability phenomenon of the inverse boundary value problem (IBVP) for the Schrödinger equation with a frequency. This IBVP is notoriously ill-posed. Ignoring the effect of the frequency, a logarithmic type estimate were first derived in [Ale88] and this logarithmic estimate was shown be optimal in the form of exponential instability in [Man01]. Such exponential instability was also established for different inverse problems in [DR03, RS18, ZZ19]. However, by taking account of the frequency, it was observed numerically [CHP03] or [KW20] that the stability of some inverse problems will improve as the frequency increases. The increasing stability phenomena were rigorously proved in other situations [DI07, DI10, HI04, ILX20, INUW14, Isa07, Isa11, KU19, LLU19, NUW13], not only for inverse problems, but also for the unique continuation property.

In this work, we will study the counterpart of the increasing stability by investigating how the exponential instability is affected by the frequency. Here we consider the Schrödinger equation with a potential and a frequency

$$(1.1) \quad \left[\Delta + q(x) + \kappa^2 \right] u = 0 \quad \text{in } \Omega \subset \mathbb{R}^d$$

with $d \geq 2$. The Cauchy data corresponding to the Schrödinger equation (1.1) is defined by

$$\mathcal{C}_q := \left\{ \left(u|_{\partial\Omega}, \frac{\partial}{\partial\nu} u \Big|_{\partial\Omega} \right) \mid u \text{ is a solution to (1.1)} \right\}.$$

Let $\text{dist}(\mathcal{C}_{q_1}, \mathcal{C}_{q_2})$ be the Hausdorff distance between two Cauchy data with respect to q_1, q_2 . Under some appropriate assumptions, it was shown in [INUW14] that

$$(1.2) \quad \|\tilde{q}\|_{H^{-s}(\mathbb{R}^d)} \leq C \left(\kappa + \log \frac{1}{\text{dist}(\mathcal{C}_{q_1}, \mathcal{C}_{q_2})} \right)^{-2s-d} + C\kappa^4 \text{dist}(\mathcal{C}_{q_1}, \mathcal{C}_{q_2}),$$

2020 *Mathematics Subject Classification.* 35J15; 35R25; 35R30.

Key words and phrases. increasing stability phenomena; instability; inverse boundary value problem; Schrödinger equation.

where \tilde{q} is the zero extension of $q_1 - q_2$. The estimate (1.2) clearly indicates that the logarithmic part decreases as κ increases and the estimate changes from a logarithmic type to a Hölder type. Isakov in [Isa11], as well as Isaev and Novikov in [IN12], proved a similar estimate in terms of the Dirichlet-to-Neumann map.

We now briefly describe the problem considered in this work. For simplicity, we consider (1.1) in $\Omega = B_1$ and $\partial\Omega = \mathcal{S}^{d-1} = \partial B_1$. If $\kappa^2 \notin \text{spec}(-(\Delta + q))$, where $\text{spec}(-(\Delta + q))$ denotes the Dirichlet spectra of $-(\Delta + q)$ in Ω , then the Dirichlet-to-Neumann map of the Schrödinger equation (1.1), $\Lambda_q : u|_{\mathcal{S}^{d-1}} \mapsto \partial_r u|_{\mathcal{S}^{d-1}}$ is well-defined. It is well-known that

$$\Lambda_q \in \mathcal{L}(H^{\frac{1}{2}}(\mathcal{S}^{d-1}), H^{-\frac{1}{2}}(\mathcal{S}^{d-1})),$$

where $\mathcal{L}(\mathcal{X}, \mathcal{Y})$ denotes the space of bounded linear operators from \mathcal{X} to \mathcal{Y} . Clearly, given any $s \geq \frac{1}{2}$, we have

$$\Lambda_q \in \mathcal{L}(H^s(\mathcal{S}^{d-1}), H^{-s}(\mathcal{S}^{d-1})).$$

To simplify the notations, we simply denote

$$\|\bullet\|_{\mathcal{L}(H^s(\mathcal{S}^{d-1}), H^{-s}(\mathcal{S}^{d-1}))} = \|\bullet\|_{s \rightarrow -s}.$$

We now state the main results of this paper.

Theorem 1.1. *Let $\alpha > 0$, $R > 0$, and any*

$$(1.3) \quad \kappa^2 > 0$$

and

$$(1.4) \quad \text{perturbation } \theta \in (0, \min\{(2 + \log 2)^{-2\alpha}, R\}),$$

where R is related to the size of the potentials. Then there exist potentials $q_1, q_2 \in C^\alpha(B_1)$ such that

$$(1.5) \quad \|q_1 - q_2\|_{L^\infty(B_1)} \geq \theta$$

and

$$(1.6) \quad \|\Lambda_{q_1} - \Lambda_{q_2}\|_{\frac{d+4}{2} \rightarrow -\frac{d+4}{2}} \leq C_R(1 + \kappa^2) \left[\exp\left(-\frac{1 + \kappa^2}{3} \theta^{-\frac{1}{2\alpha}}\right) + 3\theta^{\frac{1}{2\alpha}} \right],$$

where C_R is a positive constant depends only on the parameter R .

Remark 1.2. *There are two reasons in using $H^{\frac{d+4}{2}}$ norm in the Dirichlet-to-Neumann map. One is to bound the $\|\cdot\|_{\frac{d+4}{2} \rightarrow -\frac{d+4}{2}}$ norm of the Dirichlet-to-Neumann map by the X_s of its corresponding discretization (see (3.21)). The other is to bound the cardinality of the δ -net Y (see Remark 3.7).*

Remark 1.3. *In the proof, we actually construct q_ℓ ($\ell = 1, 2$) with $q_\ell = \tilde{q}_\ell + i$ for some real-valued functions \tilde{q}_ℓ , where $i = \sqrt{-1}$. Therefore, $\kappa^2 \notin \cup_{\ell=1,2} \text{spec}(-(\Delta + q_\ell))$ and so Λ_{q_ℓ} are well-defined. We want to point out that the discrepancy parameter θ is independent of κ and the estimate on the right hand side of (1.6) is expressed explicitly in κ . Furthermore, C_R in (1.6) only depends on the size of the potentials and is independent of κ . The assumption of $\Im(q_\ell) = 1$ is to achieve this purpose with the help of explicit elliptic estimates. On the other hand, in the usual stability estimate of the inverse boundary value problem for the Schrödinger equation, potentials q_ℓ , $\ell = 1, 2$, are a priori given such that we can assume that κ^2 is not a Dirichlet eigenvalue of $-(\Delta + q_\ell)$. In which case, we can consider the Dirichlet-to-Neumann maps [Isa11]. Instead of Dirichlet-to-Neumann maps, we can also use the Cauchy*

data in the stability estimate [INUW14]. However, in the instability estimate, potentials q_ℓ are determined posteriorly. Therefore, either assuming that κ^2 is not a Dirichlet eigenvalue of $-(\Delta + q_\ell)$ or considering the Cauchy data are not feasible.

By the standard elliptic estimates, one can show that for any $q_1, q_2 \in L^\infty(B_1)$

$$(1.7) \quad \|\Lambda_{q_1} - \Lambda_{q_2}\|_{\frac{1}{2} \rightarrow -\frac{1}{2}} \leq C(q_1, q_2, \kappa) \|q_1 - q_2\|_{L^\infty(B_1)}$$

provided Λ_{q_1} and Λ_{q_2} are well-defined, where the constant $C(q_1, q_2, \kappa)$ depends on q_1, q_2 , and κ . Estimate (1.7) can be viewed as an well-posedness estimate of the mapping $q \rightarrow \Lambda_q$. However, it provides no information of the ill-posedness of $\Lambda_q \rightarrow q$.

Theorem 1.1 is valid for all frequencies κ . The estimate (1.6) apparently demonstrates that its right hand side is dominated by the Hölder component $\theta^{\frac{1}{2\alpha}}$ when κ is large. To further elucidate our result, it is interesting to compare Theorem 1.1 and a similar instability estimate proved by Isaev in [Isae13]. In our work, we first give an arbitrary wave number κ^2 and a perturbation θ (independent of κ), and then construct suitable potentials q_1, q_2 satisfying (1.5) and (1.6). In [Isae13], Isaev constructed some wave number κ^2 and some potentials $q \in \mathcal{C}^{d+1}(B_1)$ such that

$$(1.8) \quad \|q\|_{L^\infty(B_1)} > (1 + \kappa)\delta + (1 + \kappa)^{-(d+2)}(\ln(3 + \delta^{-1}))^{-\frac{d+2}{2}},$$

where $\delta = \|\Lambda_q - \Lambda_0\|_{L^\infty(\partial B_1) \rightarrow L^\infty(\partial B_1)}$. In other words, (1.8) validates the optimality of the increasing stability estimate obtained in [IN12] (similar to (1.2)) for potentials near zero. We want to point out that the potentials q_1, q_2 constructed in Theorem 1.1 are not necessarily small. Unlike the κ -independent discrepancy parameter θ in (1.5), the lower bound of $\|q\|_{L^\infty(B_1)} = \|q - 0\|_{L^\infty(B_1)}$ depends on κ in view of (1.8).

However, Theorem 1.1 is not optimal in the lower frequency. It is expected that the IVBP is exponentially unstable if κ is small. To verify this, we prove a similar estimate in the lower frequency.

Theorem 1.4. *Let $\alpha > 0$ and*

$$0 < \kappa^2 \leq \frac{1}{4}\kappa_1$$

and

$$\text{perturbation } \theta \in \left(0, \min \left\{ \left(\frac{\kappa_1}{2} + \log 2\right)^{-2\alpha}, \frac{1}{4}\kappa_1 \right\}\right),$$

where κ_1 is the first Dirichlet eigenvalue of $-\Delta$ on B_1 . Then there exist (real-valued) potentials $q_1, q_2 \in C^\alpha(B_1)$ such that (1.5) holds and

$$(1.9) \quad \|\Lambda_{q_1} - \Lambda_{q_2}\|_{\frac{d+4}{2} \rightarrow -\frac{d+4}{2}} \leq 16\sqrt{2} \exp\left(-\frac{1}{3}\theta^{-\frac{1}{2\alpha}}\right) + 24\sqrt{2}\kappa^2\theta^{\frac{1}{2\alpha}}.$$

Remark 1.5. *Remarkd as above, we construct real-valued potentials $q_1, q_2 \in L^\infty(B_1)$ with*

$$\|q_l\|_{L^\infty} + \kappa^2 < \kappa_1, \quad l = 1, 2.$$

Therefore, $\kappa^2 \notin \cup_{l=1,2} \text{spec}(-(\Delta + q_l))$ and the Dirichlet-to-Neumann map Λ_{q_l} is well-defined.

Estimate (1.6) shows that the instability changes from an exponential type to a Hölder type when κ^2 increases, conversely, (1.9) shows that the instability is an exponential type

when κ^2 is small as in [Man01]. In particular, for any perturbation θ , we have the following dichotomy:

$$\begin{cases} \text{exponential instability when } \kappa^2 < \theta_* & \left(\text{i.e. when } \kappa^2 \text{ is small} \right), \\ \text{Hölder instability when } \kappa^2 > \theta^* & \left(\text{i.e. when } \kappa^2 \text{ is large} \right), \end{cases}$$

where

$$\begin{aligned} \theta_* &= \min \left\{ 3\theta^{\frac{1}{2\alpha}} \log \left(\frac{1}{3}\theta^{-\frac{1}{2\alpha}} \right), \frac{2}{3}\theta^{\frac{1}{2\alpha}} \exp \left(-\frac{1}{3}\theta^{-\frac{1}{2\alpha}} \right) \right\}, \\ \theta^* &= \max \left\{ 3\theta^{\frac{1}{2\alpha}} \log \left(\frac{1}{3}\theta^{-\frac{1}{2\alpha}} \right), \frac{2}{3}\theta^{\frac{1}{2\alpha}} \exp \left(-\frac{1}{3}\theta^{-\frac{1}{2\alpha}} \right) \right\}. \end{aligned}$$

Such transition of instability was also proved for an inverse problem in the stationary radiative transport equation in [ZZ19]. Our study is also inspired by their result. We also want to comment that the stability estimate of determining a potential in the wave equation by the knowledge of the hyperbolic Dirichlet-to-Neumann map was shown to be a Hölder type, see [Sun90]. In other words, in the high frequency, the inverse boundary value for (1.1) is as stable as the inverse problem for the wave equation.

We can actually combine Theorem 1.1 and Theorem 1.4 into a single theorem.

Theorem 1.6. *Let $\alpha > 0$ and $R > 0$ be any given constants. There exists a positive constant $\theta_0 = \theta_0(\alpha, R)$ such that the following statement holds: For each $\kappa > 0$ and $0 < \theta < \theta_0$, there exist potentials $q_1, q_2 \in C^\alpha(B_1)$ such that (1.5) holds and*

$$(1.10) \quad \|\Lambda_{q_1} - \Lambda_{q_2}\|_{\frac{d+4}{2} \rightarrow -\frac{d+4}{2}} \leq C_R \max\{1, \kappa^2\} \exp \left(-c_0 \max\{1, \kappa^2\} \theta^{-\frac{1}{2\alpha}} \right) + C_R \kappa^2 \theta^{\frac{1}{2\alpha}},$$

for some constants C_R depending only on R and absolute constant c_0 .

This paper is organized as follows. We list some preliminary materials in Section 2. Section 3 is a collection of some useful estimates needed in the proofs of main theorems. We then prove Theorem 1.1 and Theorem 1.4 in Section 4. Like other works in the instability of the inverse problem, our proof is based on Kolmogorov's entropy theorem [KT61].

2. PRELIMINARIES

2.1. Existence of θ -discrete set for some neighborhood. Fixing $r_0 \in (0, 1)$. Given any $\alpha > 0$, $\theta > 0$, and $\beta > 0$, we consider the following set

$$\mathcal{N}_{\alpha\beta}^\theta(B_{r_0}) := \left\{ f \geq 0 \mid \text{supp}(f) \subset B_{r_0}, \|f\|_{L^\infty} \leq \theta, \|f\|_{C^\alpha} \leq \beta \right\},$$

where B_{r_0} denotes the ball of radius r_0 centered at the origin. The following proposition can be found in [ZZ19, Lemma 5.2] (or in [KT61] in a more abstract form), see also [Man01, Lemma 2] for a direct proof for the special case when $r_0 = \frac{1}{2}$.

Proposition 2.1. *There exists a constant $\mu > 0$ such that the following statement holds for all $\beta > 0$ and for all $\theta \in (0, \mu\beta)$:*

there exists a θ -discrete (i.e. θ -distinguishable) subset Z of $(\mathcal{N}_{\alpha\beta}^\theta(B_{r_0}), \|\bullet\|_{L^\infty})$,

that is, $\|f_1 - f_2\|_{L^\infty} \geq \theta$ for all $f_1, f_2 \in Z \subset \mathcal{N}_{\alpha\beta}^\theta(K)$. Moreover, the cardinality of Z , denoted by $|Z|$, is bounded below by

$$|Z| \geq \exp \left[2^{-(d+1)} \left(\frac{\mu\beta}{\theta} \right)^{\frac{d}{\alpha}} \right].$$

2.2. Matrix representation via spherical harmonics. As in [Man01, ZZ19], we will use the set of d -dimensional spherical harmonics:

$$\mathbb{H}^d := \left\{ Y_{mj} \mid m \geq 0, 1 \leq j \leq p_m \right\},$$

where

$$(2.1) \quad p_m := \binom{m+d-1}{d-1} - \binom{m+d-3}{d-1} \leq 2(1+m)^{d-2},$$

which is a complete orthogonal set in $L^2(\mathcal{S}^{d-1})$. Indeed, $\|\bullet\|_{H^s(\mathcal{S}^{d-1})}$ is also equivalent to the following norm:

$$(2.2) \quad \left\| \sum_{m,j} a_{mj} Y_{mj} \right\|_{H^s(\mathcal{S}^{d-1})}^2 := \sum_{m,j} (1+m)^{2s} |a_{mj}|^2,$$

see e.g. [Man01].

Given any bounded linear operator $\mathcal{A} : H^s(\mathcal{S}^{d-1}) \rightarrow H^{-s}(\mathcal{S}^{d-1})$, we define $a_{mjnk} := \langle \mathcal{A} Y_{mj}, Y_{nk} \rangle$ and consider the Banach space:

$$X_s := \left\{ (a_{mjnk}) \mid \|(a_{mjnk})\|_{X_s} := \sup_{mjnk} (1 + \max\{m, n\})^{\frac{d}{2}-s} |a_{mjnk}| \right\}.$$

(see [Man01, ZZ19]). The following proposition can be found in [ZZ19, Lemma 5.3], which is crucial in our work.

Proposition 2.2. *If $s \geq \frac{d}{2}$, then*

$$\|\mathcal{A}\|_{s \rightarrow -s} \leq 4\sqrt{2} \|(a_{mjnk})\|_{X_s}.$$

Proposition 2.2 suggests that one can interpret $a_{mjnk} := \langle \mathcal{A} Y_{mj}, Y_{nk} \rangle$ as the matrix representation of the bounded linear operator $\mathcal{A} : H^s(\mathcal{S}^{d-1}) \rightarrow H^{-s}(\mathcal{S}^{d-1})$.

3. SOME USEFUL ESTIMATES

In this section, we would like to derive some useful estimates that are needed in the proofs of main theorems.

3.1. Elliptic estimate.

Lemma 3.1. *Fixing $\lambda > 0$. Let $\kappa^2 > 0$, $q \in L^\infty(B_1)$, $f \in L^2(B_1)$, and $v \in H_0^1(B_1)$ be a solution to*

$$(3.1) \quad \left[\Delta + q(x) + \kappa^2 \right] v = f \quad \text{in } B_1.$$

If

$$(3.2) \quad \mathfrak{F}(q)(x) \geq \lambda \quad \text{or} \quad \mathfrak{F}(q)(x) \leq -\lambda \quad \text{a.e.},$$

then

$$(3.3) \quad \|v\|_{L^2(B_1)} \leq \frac{1}{\lambda} \|f\|_{L^2(B_1)}.$$

Proof. Multiplying (3.1) by \bar{v} and integrating over B_1 , we have

$$(3.4) \quad - \int_{B_1} |\nabla v|^2 dx + \int_{B_1} q(x)|v|^2 dx + \kappa^2 \int_{B_1} |v|^2 dx = \int_{B_1} f\bar{v} dx.$$

Taking the imaginary part of (3.4), we have

$$\lambda \|v\|_{L^2(B_1)}^2 \leq \left| \int_{B_1} \Im(q(x))|v|^2 dx \right| = \left| \Im \int_{B_1} f\bar{v} dx \right| \leq \frac{1}{2\lambda} \|f\|_{L^2(B_1)}^2 + \frac{\lambda}{2} \|v\|_{L^2(B_1)}^2,$$

which implies (3.3). \square

The following elliptic estimate is useful in our proof.

Proposition 3.2. *Let $\lambda > 0$, $\kappa^2 > 0$, $q \in L^\infty(B_1)$, $\phi \in H^{\frac{3}{2}}(B_1)$, and $u \in H^1(B_1)$ be a solution to*

$$(3.5) \quad \begin{cases} \left[\Delta + q(x) + \kappa^2 \right] u = 0 & \text{in } B_1, \\ u = \phi & \text{on } \mathcal{S}^{d-1}. \end{cases}$$

Assume that (3.2) holds. Then there exists a positive constant C_d such that

$$(3.6) \quad \|u\|_{L^2(B_1)} \leq C_d \left(\frac{1}{\lambda} (1 + \|q\|_{L^\infty(B_1)} + \kappa^2) + 1 \right) \|\phi\|_{H^{\frac{3}{2}}(\mathcal{S}^{d-1})}.$$

Here and after, C_d (and C'_d, C''_d) are general constants depending only on dimension d .

Proof. Let $\tilde{\phi} \in H^2(B_1)$ be an extension of ϕ satisfying the inequality

$$(3.7) \quad \|\tilde{\phi}\|_{H^2(B_1)} \leq C_d \|\phi\|_{H^{\frac{3}{2}}(\mathcal{S}^{d-1})}$$

for some positive constant C_d . By letting $v = u - \tilde{\phi} \in H_0^1(B_1)$, we have

$$(3.8) \quad \left[\Delta + q(x) + \kappa^2 \right] v = f \quad \text{in } B_1,$$

where $f = - \left[\Delta + q(x) + \kappa^2 \right] \tilde{\phi} \in L^2(B_1)$. From (3.3), we have

$$\|v\|_{L^2(B_1)} \leq \frac{1}{\lambda} \|f\|_{L^2(B_1)} \leq \frac{(1 + \|q\|_{L^\infty(B_1)} + \kappa^2)}{\lambda} \|\tilde{\phi}\|_{H^2(B_1)}.$$

Since $\|u\|_{L^2(B_1)} \leq \|v\|_{L^2(B_1)} + \|\tilde{\phi}\|_{L^2(B_1)}$, together with (3.7), we have

$$\begin{aligned} \|u\|_{L^2(B_1)} &\leq \|v\|_{L^2(B_1)} + \|\tilde{\phi}\|_{L^2(B_1)} \\ &\leq \left(\frac{1}{\lambda} (1 + \|q\|_{L^\infty(B_1)} + \kappa^2) + 1 \right) \|\tilde{\phi}\|_{H^2(B_1)} \\ &\leq C_d \left(\frac{1}{\lambda} (1 + \|q\|_{L^\infty(B_1)} + \kappa^2) + 1 \right) \|\phi\|_{H^{\frac{3}{2}}(\mathcal{S}^{d-1})}, \end{aligned}$$

which implies (3.6). \square

3.2. Matrix representation. Let $q_{\text{ref}} \in L^\infty(B_1)$ satisfy (3.2). For $R > 0$, and $0 < r_0 < 1$, we define

$$B_{+,R}^{\infty,r_0} := \left\{ q \text{ is real-valued} \mid \text{supp}(q) \subset \overline{B_{r_0}}, 0 \leq q \leq R \text{ a.e.} \right\}.$$

For each $\kappa^2 > 0$, we consider the mapping

$$\Gamma(q; q_{\text{ref}}) := \Lambda_{q+q_{\text{ref}}} - \Lambda_{q_{\text{ref}}} \quad \text{for all } q \in B_{+,R}^{\infty,r_0}.$$

Let $u_{mj} \in H^1(B_1)$ be the unique solution to

$$\begin{cases} (\Delta + q_{\text{ref}}(x) + q(x) + \kappa^2)u_{mj} = 0 & \text{in } B_1, \\ u_{mj} = Y_{mj} & \text{on } \mathcal{S}^{d-1}. \end{cases}$$

Similarly, let $\dot{u}_{mj} \in H^1(B_1)$ be the unique solution to

$$\begin{cases} (\Delta + q_{\text{ref}}(x) + \kappa^2)\dot{u}_{mj} = 0 & \text{in } B_1, \\ \dot{u}_{mj} = Y_{mj} & \text{on } \mathcal{S}^{d-1}. \end{cases}$$

We also define $v_{mj} \in H^1(B_1)$ the unique solution to

$$\begin{cases} (\Delta + \overline{q_{\text{ref}}(x)} + q(x) + \kappa^2)v_{mj} = 0 & \text{in } B_1, \\ v_{mj} = Y_{mj} & \text{on } \mathcal{S}^{d-1}, \end{cases}$$

and let $\dot{v}_{mj} \in H^1(B_1)$ be the unique solution to

$$\begin{cases} (\Delta + \overline{q_{\text{ref}}(x)} + \kappa^2)\dot{v}_{mj} = 0 & \text{in } B_1, \\ \dot{v}_{mj} = Y_{mj} & \text{on } \mathcal{S}^{d-1}. \end{cases}$$

By (2.2) and (3.6), we can derive

$$(3.9a) \quad \|u_{mj}\|_{L^2(B_1)} \leq C'_d \left(\frac{1}{\lambda} (1 + R + \lambda + \kappa^2) + 1 \right) (1+m)^{\frac{3}{2}},$$

$$(3.9b) \quad \|\dot{u}_{mj}\|_{L^2(B_1)} \leq C'_d \left(\frac{1}{\lambda} (1 + R + \lambda + \kappa^2) + 1 \right) (1+m)^{\frac{3}{2}},$$

$$(3.9c) \quad \|v_{mj}\|_{L^2(B_1)} \leq C'_d \left(\frac{1}{\lambda} (1 + R + \lambda + \kappa^2) + 1 \right) (1+m)^{\frac{3}{2}},$$

$$(3.9d) \quad \|\dot{v}_{mj}\|_{L^2(B_1)} \leq C'_d \left(\frac{1}{\lambda} (1 + R + \lambda + \kappa^2) + 1 \right) (1+m)^{\frac{3}{2}}.$$

Write $\tilde{Y}_{mj}(x) := |x|^m Y_{mj}(x/|x|)$, which is harmonic in B_1 . We can prove

Lemma 3.3. *Let $q_{\text{ref}} \in L^\infty(B_1)$ satisfy (3.2) and $q \in B_{+,R}^{\infty,r_0}$. Then*

$$(3.10) \quad \begin{aligned} \langle \Gamma(q; q_{\text{ref}})Y_{mj}, Y_{nk} \rangle &= I_{mjnk}^{(1)}(q) + I_{mjnk}^{(2)}(q_{\text{ref}}) \\ &= L_{mjnk}^{(1)}(q) + L_{mjnk}^{(2)}(q_{\text{ref}}), \end{aligned}$$

where

$$\begin{aligned}
I_{mjnk}^{(1)}(q) &= - \int_{B_{r_0}} q(x) u_{mj} \overline{\tilde{Y}_{nk}} dx, \\
I_{mjnk}^{(2)}(q_{\text{ref}}) &= - \int_{B_1} (q_{\text{ref}}(x) + \kappa^2) (u_{mj} - \dot{u}_{mj}) \overline{\tilde{Y}_{nk}} dx, \\
L_{mjnk}^{(1)}(q) &= - \int_{B_{r_0}} q(x) \overline{v_{nk}} \tilde{Y}_{mj} dx, \\
L_{mjnk}^{(2)}(q_{\text{ref}}) &= - \int_{B_1} (\overline{q_{\text{ref}}(x)} + \kappa^2) \overline{(v_{nk} - \dot{v}_{nk})} \tilde{Y}_{mj} dx.
\end{aligned}$$

Proof. We begin to prove

$$(3.11) \quad \langle \Gamma(q; q_{\text{ref}}) Y_{mj}, Y_{nk} \rangle = I_{mjnk}^{(1)}(q; q_{\text{ref}}) + I_{mjnk}^{(2)}(q; q_{\text{ref}}).$$

Using integration by parts, we have

$$\begin{aligned}
(3.12) \quad \langle \Gamma(q; q_{\text{ref}}) Y_{mj}, Y_{nk} \rangle &= \int_{\mathcal{S}^{d-1}} \partial_r (u_{mj} - \dot{u}_{mj}) \overline{\tilde{Y}_{nk}} ds \\
&= \int_{B_1} \Delta (u_{mj} - \dot{u}_{mj}) \overline{\tilde{Y}_{nk}} dx.
\end{aligned}$$

In the computation of (3.12), we used the fact $u_{mj} = \dot{u}_{mj}$ on \mathcal{S}^{d-1} and \tilde{Y}_{nk} is harmonic in B_1 . Since

$$(\Delta + q_{\text{ref}}(x) + q(x) + \kappa^2) u_{mj} = 0 \quad \text{and} \quad (\Delta + q_{\text{ref}}(x) + \kappa^2) \dot{u}_{mj} = 0,$$

and $\text{supp}(q) \subset B_{r_0}$, (3.12) implies

$$\begin{aligned}
&\langle \Gamma(q; q_{\text{ref}}) Y_{mj}, Y_{nk} \rangle \\
&= - \int_{B_{r_0}} q(x) u_{mj} \overline{\tilde{Y}_{nk}} dx - \int_{B_1} (q_{\text{ref}}(x) + \kappa^2) (u_{mj} - \dot{u}_{mj}) \overline{\tilde{Y}_{nk}} dx,
\end{aligned}$$

which is exactly (3.11).

Next, we want to establish

$$(3.13) \quad \langle \Gamma(q; q_{\text{ref}}) Y_{mj}, Y_{nk} \rangle = L_{mjnk}^{(1)}(q; q_{\text{ref}}) + L_{mjnk}^{(2)}(q; q_{\text{ref}}).$$

Since the adjoint operator of $\Gamma(q; q_{\text{ref}})$ is $\Gamma(q; \overline{q_{\text{ref}}})$, similar to the derivation of (3.12), we have

$$\begin{aligned}
\langle \Gamma(q; q_{\text{ref}}) Y_{mj}, Y_{nk} \rangle &= \langle Y_{mj}, \Gamma(q; \overline{q_{\text{ref}}}) Y_{nk} \rangle \\
&= \int_{B_1} \overline{\Delta (v_{nk} - \dot{v}_{nk})} \tilde{Y}_{mj} dx.
\end{aligned}$$

Since

$$(\Delta + \overline{q_{\text{ref}}(x)} + q(x) + \kappa^2) v_{mj} = 0 \quad \text{and} \quad (\Delta + \overline{q_{\text{ref}}(x)} + \kappa^2) \dot{v}_{mj} = 0,$$

and $\text{supp}(q) \subset B_{r_0}$, we obtain

$$\begin{aligned}
\langle \Gamma(q; q_{\text{ref}}) Y_{mj}, Y_{nk} \rangle &= \langle Y_{mj}, \Gamma(q; \overline{q_{\text{ref}}}) Y_{nk} \rangle \\
&= - \int_{B_{r_0}} q(x) \overline{v_{nk}} \tilde{Y}_{mj} dx - \int_{B_1} (\overline{q_{\text{ref}}(x)} + \kappa^2) \overline{(v_{nk} - \dot{v}_{nk})} \tilde{Y}_{mj} dx
\end{aligned}$$

and hence (3.13). \square

Next, we define

$$(3.14) \quad M_{mjnk}^{(1)}(q) := \begin{cases} I_{mjnk}^{(1)}(q) & \text{if } n \geq m, \\ L_{mjnk}^{(1)}(q) & \text{if } n < m, \end{cases} \quad M_{mjnk}^{(2)}(q_{\text{ref}}) := \begin{cases} I_{mjnk}^{(2)}(q_{\text{ref}}) & \text{if } n \geq m, \\ L_{mjnk}^{(2)}(q_{\text{ref}}) & \text{if } n < m. \end{cases}$$

Therefore, from (3.10), we have

$$(3.15) \quad |\langle \Gamma(q; q_{\text{ref}})Y_{mj}, Y_{nk} \rangle| \leq |M_{mjnk}^{(1)}(q)| + |M_{mjnk}^{(2)}(q_{\text{ref}})|.$$

Lemma 3.4. *Let $q_{\text{ref}} \in L^\infty(B_1)$ be such that (3.2) holds. Let $\ell = \max\{m, n\}$, $\kappa^2 > 0$, and $q \in B_{+,R}^{\infty,r_0}$. Then*

$$(3.16) \quad |M_{mjnk}^{(1)}(q)| \leq C_d'' \Phi(R, \lambda, \kappa)(1 + \ell)r_0^\ell,$$

$$(3.17) \quad |M_{mjnk}^{(2)}(q_{\text{ref}})| \leq C_d'' \Phi(R, \lambda, \kappa)(\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2)(1 + \ell)$$

with $C_d'' > 1$, where

$$(3.18) \quad \Phi(R, \lambda, \kappa) = (R + 1) \left(\frac{1}{\lambda} (1 + R + \lambda + \kappa^2) + 1 \right).$$

Remark 3.5. *Denote $\tau := s - \frac{d+2}{2}$. It follows from (3.16) and (3.17) that*

$$(3.19) \quad (1 + \ell)^{\frac{d}{2}-s} |M_{mjnk}^{(1)}(q)| \leq C_d'' \Phi(R, \lambda, \kappa)(1 + \ell)^{-\tau} r_0^\ell,$$

$$(3.20) \quad (1 + \ell)^{\frac{d}{2}-s} |M_{mjnk}^{(2)}(q_{\text{ref}})| \leq C_d'' \Phi(R, \lambda, \kappa)(\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2)(1 + \ell)^{-\tau}.$$

Therefore, if $s \geq \frac{d+2}{2}$ (i.e. $\tau \geq 0$), we know that

$$\|(\langle \Gamma(q; q_{\text{ref}})Y_{mj}, Y_{nk} \rangle)\|_{X_s} = \sup_{mjnk} (1 + \ell)^{\frac{d}{2}-s} |\langle \Gamma(q; q_{\text{ref}})Y_{mj}, Y_{nk} \rangle| < \infty.$$

If we define $\Gamma_{mjnk}^{q_{\text{ref}}}(q) := \langle \Gamma(q; q_{\text{ref}})Y_{mj}, Y_{nk} \rangle$, we have

$$(3.21) \quad (\Gamma_{mjnk}^{q_{\text{ref}}}(B_{+,R}^{\infty,r_0})) \subset X_s \quad \text{for all } s \geq \frac{d+2}{2}.$$

Proof of Lemma 3.4. Like in [ZZ19, (16)], using (3.9a) yields

$$\begin{aligned} |I_{mjnk}^{(1)}(q)| &= \left| \int_{B_{r_0}} q(x) |x|^n \overline{Y_{nk}(x/|x|)} u_{mj}(x) dx \right| \\ &\leq R \left(\int_0^{r_0} r^{2n+d-1} dr \right)^{\frac{1}{2}} \|u_{mj}\|_{L^2(B_1)} \\ &\leq C_d' R \left(\frac{1}{\lambda} (1 + R + \lambda + \kappa^2) + 1 \right) \frac{(1+m)^{\frac{3}{2}}}{\sqrt{2n+d}} r_0^{n+\frac{d}{2}} \\ &\leq C_d' R \left(\frac{1}{\lambda} (1 + R + \lambda + \kappa^2) + 1 \right) \frac{(1+m)^{\frac{3}{2}}}{\sqrt{2n+d}} r_0^n. \end{aligned}$$

Similarly, from (3.9c), we have

$$|L_{mjnk}^{(1)}(q)| \leq C_d' R \left(\frac{1}{\lambda} (1 + R + \lambda + \kappa^2) + 1 \right) \frac{(1+n)^{\frac{3}{2}}}{\sqrt{2m+d}} r_0^m.$$

Therefore,

$$\begin{aligned} |M_{mjnk}^{(1)}(q)| &\leq C'_d R \left(\frac{1}{\lambda} (1 + R + \lambda + \kappa^2) + 1 \right) \frac{(1+n)^{\frac{3}{2}}}{\sqrt{2n+d}} r_0^n \quad \text{if } n \geq m, \\ |M_{mjnk}^{(1)}(q)| &\leq C'_d R \left(\frac{1}{\lambda} (1 + R + \lambda + \kappa^2) + 1 \right) \frac{(1+m)^{\frac{3}{2}}}{\sqrt{2m+d}} r_0^m \quad \text{if } n < m, \end{aligned}$$

which gives (3.16).

Now, we turn to (3.17). Following [ZZ19, (16)] again, combining (3.9a) and (3.9b) implies

$$\begin{aligned} |I_{mjnk}^{(2)}(q_{\text{ref}})| &= \left| \int_{B_1} (q_{\text{ref}}(x) + \kappa^2) (u_{mj} - \dot{u}_{mj}) |x|^n \overline{Y_{nk}(x/|x|)} dx \right| \\ &\leq \left(\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2 \right) \left(\int_0^1 r^{2n+d-1} dr \right)^{\frac{1}{2}} \left(\|u_{mj}\|_{L^2(B_1)} + \|\dot{u}_{mj}\|_{L^2(B_1)} \right) \\ &\leq 2C'_d \left(\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2 \right) \left(\frac{1}{\lambda} (1 + R + \lambda + \kappa^2) + 1 \right) \frac{(1+m)^{\frac{3}{2}}}{\sqrt{2n+d}}. \end{aligned}$$

Likewise, by (3.9c) and (3.9d), we have

$$|L_{mjnk}^{(2)}(q_{\text{ref}})| \leq 2C'_d \left(\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2 \right) \left(\frac{1}{\lambda} (1 + R + \lambda + \kappa^2) + 1 \right) \frac{(1+n)^{\frac{3}{2}}}{\sqrt{2m+d}}.$$

Then we can derive (3.17) as above. \square

3.3. Construction of a net. In view of (3.21) in Remark 3.5, we want to construct a δ -net Y of $((\Gamma_{mjnk}^{q_{\text{ref}}}(B_{+,R}^{\infty,r_0})), \|\bullet\|_{X_s})$, which is not too-large. Precisely, we want to obtain the following proposition.

Proposition 3.6. *Assume that $q_{\text{ref}} \in L^\infty(B_1)$ satisfies (3.2). Let $s > \frac{d+2}{2}$ and $\kappa^2 > 0$. Define $\tau := s - \frac{d+2}{2}$. Given any $0 < \delta < \Phi$, where $\Phi = \Phi(R, \lambda, \kappa)$ is the function given in (3.18), there exists a δ -net Y of $((\Gamma_{mjnk}^{q_{\text{ref}}}(B_{+,R}^{\infty,r_0})), \|\bullet\|_{X_s})$ such that*

$$(3.22) \quad \log |Y| \leq \eta \left[1 + \log \left(1 + \frac{\Phi}{\delta} \right) + (\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2) \frac{\Phi}{\delta} + \left((\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2) \frac{\Phi}{\delta} \right)^{\frac{1}{\tau}} \right]^{2d}$$

for some constant $\eta = \eta(d, s, r_0)$.

Remark 3.7. *In view of Remark 3.5, here we exclude the case $s = \frac{d+2}{2}$ (i.e. $s > \frac{d+2}{2}$), to ensure that the quantity $1/\tau$ in (3.22) is well-defined.*

Proof. Step 1: Initialization. Suggested by (3.19) and (3.20), let ℓ_1 and ℓ_2 be the solution to

$$(1 + \ell_1)^{-\tau} r_0^{\ell_1} = \frac{\delta}{8\sqrt{2}C''_d \Phi} \quad \text{and} \quad (1 + \ell_2)^{-\tau} (\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2) = \frac{\delta}{8\sqrt{2}C''_d \Phi},$$

respectively. This is valid because $0 < \delta < \Phi$ and $C''_d > 1$. Let ℓ_* be the smallest nonnegative integer such that

$$(3.23) \quad (1 + \ell)^{-\tau} \left(r_0^\ell + \|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2 \right) \leq \frac{\delta}{4\sqrt{2}C''_d \Phi} \quad \text{for all } \ell \geq \ell_*.$$

Observe that

$$\begin{aligned}
& (1 + (\ell_1 + \ell_2))^{-\tau} \left(r_0^{\ell_1 + \ell_2} + \|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2 \right) \\
&= r_0^{\ell_2} (1 + (\ell_1 + \ell_2))^{-\tau} r_0^{\ell_1} + (1 + (\ell_1 + \ell_2))^{-\tau} \left(\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2 \right) \\
&\leq (1 + \ell_1)^{-\tau} r_0^{\ell_1} + (1 + \ell_2)^{-\tau} (\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2) \leq \frac{\delta}{4\sqrt{2}C_d''\Phi},
\end{aligned}$$

that is, $\ell_1 + \ell_2$ satisfies (3.23). Since ℓ_* is the smallest nonnegative integer such that (3.23) holds, then

$$(3.24) \quad \ell_* \leq \ell_1 + \ell_2.$$

Note that $\tau \log(1 + \ell_1) \geq 0$. Thus, we have

$$\ell_1 \log r_0 \geq -\tau \log(1 + \ell_1) + \ell_1 \log r_0 = \log \left((1 + \ell_1)^{-\tau} r_0^{\ell_1} \right) = \log \left(\frac{\delta}{8\sqrt{2}C_d''\Phi} \right).$$

Since $0 < r_0 < 1$, i.e., $\log r_0 < 0$, it yields

$$(3.25) \quad \ell_1 \leq \frac{1}{\log r_0} \log \left(\frac{\delta}{8\sqrt{2}C_d''\Phi} \right) = \frac{1}{-\log r_0} \log \left(\frac{8\sqrt{2}C_d''\Phi}{\delta} \right).$$

We also have

$$(3.26) \quad \ell_2 = \left(\frac{\delta}{8\sqrt{2}C_d''\Phi(\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2)} \right)^{-\frac{1}{\tau}} - 1 = (8\sqrt{2}C_d'')^{\frac{1}{\tau}} \left((\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2) \frac{\Phi}{\delta} \right)^{\frac{1}{\tau}} - 1.$$

Combining (3.24), (3.25), and (3.26), leads to

$$\begin{aligned}
\ell_* + 1 &\leq \frac{1}{-\log r_0} \log \left(\frac{8\sqrt{2}C_d''\Phi}{\delta} \right) + (8\sqrt{2}C_d'')^{\frac{1}{\tau}} \left((\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2) \frac{\Phi}{\delta} \right)^{\frac{1}{\tau}} \\
&= \frac{1}{-\log r_0} \left(\log(8\sqrt{2}C_d'') + \log \left(\frac{\Phi}{\delta} \right) \right) + (8\sqrt{2}C_d'')^{\frac{1}{\tau}} \left((\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2) \frac{\Phi}{\delta} \right)^{\frac{1}{\tau}} \\
(3.27) \quad &\leq C_1 \left(1 + \log \left(1 + \frac{\Phi}{\delta} \right) + \left((\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2) \frac{\Phi}{\delta} \right)^{\frac{1}{\tau}} \right).
\end{aligned}$$

for some positive constant $C_1 = C_1(d, r_0, s)$.

Step 2: Construction of sets. For each 4-tuple (m, j, n, k) with $\ell := \max\{m, n\} \leq \ell_*$, (3.16) and (3.17) imply

$$\begin{aligned}
|M_{mjnk}^{(1)}(q) &\leq C_d'' C_2 \Phi(R, \lambda, \kappa), \\
|M_{mjnk}^{(2)}(q_{\text{ref}}) &\leq C_d'' \Phi(R, \lambda, \kappa) (\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2) (1 + \ell_*),
\end{aligned}$$

where $C_2 = C_2(d, r_0) = \sup_{\ell} (1 + \ell) r_0^\ell < \infty$. Let $\delta' = \frac{\delta}{8\sqrt{2}}$ and define

$$\begin{aligned}
Y_1' &:= \left\{ a = a_1 + ia_2 \in \frac{\delta'}{\sqrt{2}}\mathbb{Z} + i\frac{\delta'}{\sqrt{2}}\mathbb{Z} \mid |a_1|, |a_2| \leq C_d'' C_2 \Phi(R, \lambda, \kappa) \right\}, \\
Y_2' &:= \left\{ a = a_1 + ia_2 \in \frac{\delta'}{\sqrt{2}}\mathbb{Z} + i\frac{\delta'}{\sqrt{2}}\mathbb{Z} \mid |a_1|, |a_2| \leq C_d'' \Phi(R, \lambda, \kappa) (\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2) (1 + \ell_*) \right\},
\end{aligned}$$

where \mathbb{Z} denotes the set of integers. Then,

$$(3.28a) \quad |Y'_1| = \left(1 + 2 \left[\frac{C_d'' C_2 \Phi}{\delta' / \sqrt{2}} \right] \right)^2 \leq \left(1 + \frac{32 C_d'' C_2 \Phi}{\delta}\right)^2,$$

$$(3.28b) \quad |Y'_2| \leq \left(1 + \frac{32 C_d'' \Phi (\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2)(1 + \ell_*)}{\delta}\right)^2.$$

Define

$$Y_1 := \left\{ (b_{mjnk}) \left| \begin{array}{l} b_{mjnk} \in Y'_1, \text{ if } \max\{m, n\} \leq \ell_*; \\ b_{mjnk} = 0, \text{ otherwise} \end{array} \right. \right\},$$

$$Y_2 := \left\{ (c_{mjnk}) \left| \begin{array}{l} c_{mjnk} \in Y'_2, \text{ if } \max\{m, n\} \leq \ell_*; \\ c_{mjnk} = 0, \text{ otherwise} \end{array} \right. \right\},$$

and let $Y := Y_1 + Y_2$.

Step 3: Verifying Y is a δ -net of $((\Gamma_{mjnk}^{q_{\text{ref}}}(B_{+,R}^{\infty,r_0})), \|\bullet\|_{X_s})$. Our goal is to construct

$$\begin{cases} (b_{mjnk}) \in Y_1 & \text{which is an approximation of } (M_{mjnk}^{(1)}(q)), \\ (c_{mjnk}) \in Y_2 & \text{which is an approximation of } (M_{mjnk}^{(2)}(q_{\text{ref}})). \end{cases}$$

- If $\max\{m, n\} \leq \ell_*$, we take

$$\begin{cases} b'_{mjnk} \in Y'_1 & \text{the closest element to } M_{mjnk}^{(1)}(q), \\ c'_{mjnk} \in Y'_2 & \text{the closest element to } M_{mjnk}^{(2)}(q_{\text{ref}}). \end{cases}$$

Then

$$|b'_{mjnk} - M_{mjnk}^{(1)}(q)| \leq \delta' \quad \text{and} \quad |c'_{mjnk} - M_{mjnk}^{(2)}(q_{\text{ref}})| \leq \delta'.$$

By $s > \frac{d}{2}$, we then have

$$\begin{aligned} & 4\sqrt{2}(1 + \max\{m, n\})^{\frac{d}{2}-s} (|b'_{mjnk} - M_{mjnk}^{(1)}(q)| + |c'_{mjnk} - M_{mjnk}^{(2)}(q_{\text{ref}})|) \\ & \leq 8\sqrt{2}\delta' = \delta. \end{aligned}$$

- If $\ell = \max\{m, n\} > \ell_*$, we simply choose $b'_{mjnk} = c'_{mjnk} = 0$. Hence, we know that

$$\begin{aligned} & 4\sqrt{2}(1 + \max\{m, n\})^{\frac{d}{2}-s} (|b'_{mjnk} - M_{mjnk}^{(1)}(q)| + |c'_{mjnk} - M_{mjnk}^{(2)}(q_{\text{ref}})|) \\ & = 4\sqrt{2}(1 + \ell)^{\frac{d}{2}-s} (|M_{mjnk}^{(1)}(q)| + |M_{mjnk}^{(2)}(q_{\text{ref}})|) \\ & \leq 4\sqrt{2}C_d''\Phi(1 + \ell)^{-\tau} \left(r_0^\ell + \|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2 \right) \quad (\text{by (3.19), (3.20)}) \\ & \leq \delta \quad (\text{using (3.23)}). \end{aligned}$$

Combining these two cases, we conclude that Y is a δ -net of $((\Gamma_{mjnk}^{q_{\text{ref}}}(B_{+,R}^{\infty,r_0})), \|\bullet\|_{X_s})$

Step 4: Calculating the cardinality of Y . Let n_ℓ be the number of 4-tuples (m, j, n, k) be such that $\max\{m, n\} = \ell$, and let $n_* := \sum_{\ell=0}^{\ell_*} n_\ell$. It is not hard to check that

$$|Y| = |Y_1||Y_2| = (|Y'_1||Y'_2|)^{n_*}.$$

Also, by (2.1), we have

$$n_* \leq 8(1 + \ell_*)^{2d-2}.$$

Therefore, from (3.28a) and (3.28b), it follows

$$\begin{aligned} \log |Y| \leq & 8(1 + \ell_*)^{2d-2} \left[2 \log \left(1 + \frac{32C_d'' C_2 \Phi}{\delta} \right) \right. \\ & \left. + 2 \log \left(1 + \frac{32C_d'' \Phi (\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2)(1 + \ell_*)}{\delta} \right) \right]. \end{aligned}$$

Since $\log(1 + t) \leq t$ for all $t \geq 0$, we obtain

$$\begin{aligned} \log |Y| & \leq 16(1 + \ell_*)^{2d-2} \left[\log \left(1 + \frac{32C_d'' C_2 \Phi}{\delta} \right) + \frac{32C_d'' \Phi (\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2)(1 + \ell_*)}{\delta} \right] \\ & \leq 16(1 + \ell_*)^{2d-1} \left[\log \left(1 + \frac{32C_d'' C_2 \Phi}{\delta} \right) + \frac{32C_d'' \Phi (\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2)}{\delta} \right] \\ (3.29) \quad & \leq \eta' (1 + \ell_*)^{2d-1} \left[1 + \log \left(1 + \frac{\Phi}{\delta} \right) + (\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2) \frac{\Phi}{\delta} \right] \end{aligned}$$

for some constant $\eta' = \eta'(d, s, r_0)$. Combining (3.29) with (3.27) yields

$$\begin{aligned} \log |Y| \leq & \eta' \left[C_1 \left(1 + \log \left(1 + \frac{\Phi}{\delta} \right) + \left((\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2) \frac{\Phi}{\delta} \right)^{\frac{1}{\tau}} \right) \right]^{2d-1} \times \\ & \times \left[1 + \log \left(1 + \frac{\Phi}{\delta} \right) + (\|q_{\text{ref}}\|_{L^\infty(B_1)} + \kappa^2) \frac{\Phi}{\delta} \right], \end{aligned}$$

which implies (3.22). □

4. PROOFS OF THE MAIN RESULTS

In this section, we would like to present detailed proofs of Theorem 1.1 and Theorem 1.4. Before proving our results, we first give some observations. Taking $q_{\text{ref}} = i$ (hence $\lambda = 1$) and $s = \frac{d+4}{2}$ (i.e. $\tau = 1$), see also Remark 1.3, we see that (3.22) becomes

$$(4.1) \quad \log |Y| \leq \eta \left[1 + \log \left(1 + \frac{\Phi}{\delta} \right) + 2(1 + \kappa^2) \frac{\Phi}{\delta} \right]^{2d},$$

where

$$(4.2) \quad \Phi(R, \kappa) = (R + 1)(3 + R + \kappa^2).$$

Note that

$$\inf_{0 < \delta < \Phi} \frac{1}{(1 + \kappa^2)} \left[\log \left(1 + \frac{\Phi}{\delta} \right) + 2(1 + \kappa^2) \frac{\Phi}{\delta} \right] = \frac{\log 2}{1 + \kappa^2} + 2 < 2 + \log 2$$

for all $\kappa^2 > 0$. Therefore, given any θ satisfies

$$(4.3) \quad 2 + \log 2 < \theta^{-\frac{1}{2\alpha}} < \infty \quad \left(\text{equivalently, } \theta \in (0, (2 + \log 2)^{-2\alpha}) \right),$$

there exists a unique $\delta \in (0, \Phi)$ such that

$$(4.4) \quad \theta^{-\frac{1}{2\alpha}} = \frac{1}{1 + \kappa^2} \left[\log \left(1 + \frac{\Phi}{\delta} \right) + 2(1 + \kappa^2) \frac{\Phi}{\delta} \right],$$

and (4.1) is reduced to

$$(4.5) \quad |Y| \leq \exp \left[\eta \left(1 + (1 + \kappa^2) \theta^{-\frac{1}{2\alpha}} \right)^{2d} \right]$$

In view of these observations, we are ready to prove Theorem 1.1, the case of higher frequency.

Proof of Theorem 1.1. By Proposition 2.1, there exists a constant $\mu > 0$ such that the following statement holds for all $\beta > 0$ and for all $\theta \in (0, \mu\beta)$:

there exists a θ -discrete subset Z of $(\mathcal{N}_{\alpha\beta}^\theta(B_{r_0}), \|\bullet\|_{L^\infty})$

with

$$|Z| \geq \exp \left[2^{-(d+1)} \left(\frac{\mu\beta}{\theta} \right)^{\frac{d}{\alpha}} \right].$$

It suffices to restrict $0 < \theta < \min\{(2 + \log 2)^{-2\alpha}, R, \mu\beta\}$. Therefore, we can choose $0 < \delta < \Phi$ such that (4.4) is satisfied.

Let $0 < r_0 < 1$. Since $0 < \theta < R$, we see that $Z \subset \mathcal{N}_{\alpha\beta}^\theta(B_{r_0}) \subset B_{+,R}^{\infty,r_0}$. Recall that we have chosen $q_{\text{ref}} = i$. By Proposition 3.6, there exists a δ -net Y of $((\Gamma_{mjnk}^i(B_{+,R}^{\infty,r_0}), \|\bullet\|_{X_s})$ with $s = \frac{d+4}{2}$ such that (4.5) holds. Of course, Y is also a δ -net of $((\Gamma_{mjnk}^i(Z), \|\bullet\|_{X_s})$.

We now choose $\beta > R$ sufficiently large (depends on d, μ, θ, κ) satisfying

$$|Z| \geq \exp \left[2^{-(d+1)} \left(\frac{\mu\beta}{\theta} \right)^{\frac{d}{\alpha}} \right] > \exp \left[\eta \left(1 + (1 + \kappa^2) \theta^{-\frac{1}{2\alpha}} \right)^{2d} \right] \geq |Y|.$$

Therefore, we can choose two different $\tilde{q}_1, \tilde{q}_2 \in Z \subset \mathcal{N}_{\alpha\beta}^\theta(B_{r_0}) \subset L^\infty(B_1)$ such that there exists $y \in Y$ such that

$$\begin{aligned} \|(\Gamma_{mjnk}^i(\tilde{q}_1) - y_{mjnk})\|_{X_s} &\leq \delta, \\ \|(\Gamma_{mjnk}^i(\tilde{q}_2) - y_{mjnk})\|_{X_s} &\leq \delta. \end{aligned}$$

Combining with Proposition 2.2, we conclude that

$$\|\Lambda_{q_1} - \Lambda_{q_2}\|_{\frac{d+4}{2} \rightarrow -\frac{d+4}{2}} \leq 8\sqrt{2}\delta,$$

where $q_1 = \tilde{q}_1 + i$ and $q_2 = \tilde{q}_2 + i$. Now we estimate δ by θ .

- **Case 1.** If $(1 + \kappa^2) \frac{\Phi}{\delta} \leq \log(1 + \frac{\Phi}{\delta})$, then (4.4) implies

$$\theta^{-\frac{1}{2\alpha}} \leq \frac{3}{1 + \kappa^2} \log \left(1 + \frac{\Phi}{\delta} \right),$$

that is,

$$\delta \leq (R + 1)(3 + R + \kappa^2) \exp \left(-\frac{1 + \kappa^2}{3} \theta^{-\frac{1}{2\alpha}} \right).$$

- **Case 2.** If $(1 + \kappa^2) \frac{\Phi}{\delta} \geq \log(1 + \frac{\Phi}{\delta})$, the (4.4) implies

$$\theta^{-\frac{1}{2\alpha}} \leq \frac{1}{1 + \kappa^2} \left[3(1 + \kappa^2) \frac{\Phi}{\delta} \right] = \frac{3\Phi}{\delta} = 3(R + 1)(3 + R + \kappa^2) \delta^{-1},$$

that is,

$$\delta \leq 3(R + 1)(3 + R + \kappa^2) \theta^{\frac{1}{2\alpha}}.$$

Combining these two cases yields

$$\delta \leq (R+1)(3+R+\kappa^2) \left[\exp\left(-\frac{1+\kappa^2}{3}\theta^{-\frac{1}{2\alpha}}\right) + 3\theta^{\frac{1}{2\alpha}} \right],$$

which gives (1.6). Since Z is θ -discrete subset with respect to the norm $\|\bullet\|_{L^\infty(B_1)}$, (1.5) follows immediately, which is our desired result. \square

Now, we want to prove Theorem 1.4 in which the frequency is small. First of all, we improve Proposition 3.2 without the restriction (3.2).

Proposition 4.1. *Let $\kappa^2 > 0$, $q \in L^\infty(B_1)$, $\phi \in H^{\frac{3}{2}}(B_1)$, and $u \in H^1(B_1)$ be the solution to (3.5). If*

$$(4.6) \quad \|q\|_{L^\infty(B_1)} \leq \frac{1}{4}\kappa_1 \quad \text{and} \quad \kappa^2 \leq \frac{1}{4}\kappa_1,$$

where κ_1 is the first Dirichlet eigenvalue of $-\Delta$ on B_1 (which depends only on dimension d), then

$$(4.7) \quad \|u\|_{L^2(B_1)} \leq C_d \|\phi\|_{H^{\frac{3}{2}}(\mathcal{S}^{d-1})}$$

for some constant C_d depending only on dimension d .

Proof. Let $\tilde{\phi} \in H^2(B_1)$ as in (3.7) and consider $v = u - \tilde{\phi} \in H_0^1(B_1)$. In this case, we write (3.8) as

$$(4.8) \quad \Delta v = f - q(x)v + \kappa^2 v,$$

where $f = -(\Delta + q(x) + \kappa^2)\tilde{\phi}$. Multiplying (4.8) by $-\bar{v}$ and using the integration by parts, we have

$$\begin{aligned} \int_{B_1} |\nabla v|^2 dx &= \int_{B_1} f\bar{v} dx + \int_{B_1} q(x)|v|^2 dx + \kappa^2 \int_{B_1} |v|^2 dx \\ &\leq \frac{1}{4\epsilon} \|f\|_{L^2(B_1)}^2 + \left(\epsilon + \|q\|_{L^\infty(B_1)} + \kappa^2 \right) \|v\|_{L^2(B_1)}^2 \end{aligned}$$

for $\epsilon > 0$. Poincaré's inequality implies that we have $\kappa_1 \|v\|_{L^2(B_1)}^2 \leq \|\nabla v\|_{L^2(B_1)}^2$ for all $v \in H_0^1(B_1)$. From (4.6), it follows that

$$\begin{aligned} \kappa_1 \|v\|_{L^2(B_1)}^2 &\leq \frac{1}{4\epsilon} \|f\|_{L^2(B_1)}^2 + \left(\epsilon + \|q\|_{L^\infty(B_1)} + \kappa^2 \right) \|v\|_{L^2(B_1)}^2 \\ &\leq \frac{1}{4\epsilon} \|f\|_{L^2(B_1)}^2 + \left(\epsilon + \frac{1}{2}\kappa_1 \right) \|v\|_{L^2(B_1)}^2. \end{aligned}$$

Choosing $\epsilon = \frac{1}{4}\kappa_1$, we have

$$\|v\|_{L^2(B_1)}^2 \leq \frac{4}{\kappa_1^2} \|f\|_{L^2(B_1)}^2.$$

Using (4.6) again gives

$$\begin{aligned} \|v\|_{L^2(B_1)} &\leq 2\kappa_1^{-1} \|f\|_{L^2(B_1)} \leq 2\kappa_1^{-1} (1 + \|q\|_{L^\infty(B_1)} + \kappa^2) \|\tilde{\phi}\|_{H^2(B_1)} \\ &\leq 2\kappa_1^{-1} \left(1 + \frac{1}{2}\kappa_1 \right) \|\tilde{\phi}\|_{H^2(B_1)}. \end{aligned}$$

Combining $\|u\|_{L^2(B_1)} \leq \|v\|_{L^2(B_1)} + \|\tilde{\phi}\|_{L^2(B_1)}$ and (3.7) implies (4.7). \square

Let $u_{mj} \in H^1(B_1)$, $\dot{u}_{mj} \in H^1(B_1)$, $v_{mj} \in H^1(B_1)$, $\dot{v}_{mj} \in H^1(B_1)$ as in Subsection 3.2. In this case, by using (4.7) rather than (3.6), we can improve (3.9a), (3.9b), (3.9c), and (3.9d) as follows:

$$(4.9a) \quad \|u_{mj}\|_{L^2(B_1)} \leq C'_d(1+m)^{\frac{3}{2}},$$

$$(4.9b) \quad \|\dot{u}_{mj}\|_{L^2(B_1)} \leq C'_d(1+m)^{\frac{3}{2}},$$

$$(4.9c) \quad \|v_{mj}\|_{L^2(B_1)} \leq C'_d(1+m)^{\frac{3}{2}},$$

$$(4.9d) \quad \|\dot{v}_{mj}\|_{L^2(B_1)} \leq C'_d(1+m)^{\frac{3}{2}},$$

for some constant C'_d .

Now, choosing $q_{\text{ref}} \equiv 0$ in Lemma 3.3, we can prove the following lemma.

Lemma 4.2. *For each $q \in B_{+,R}^{\infty,r_0}$, we have*

$$\begin{aligned} \langle \Gamma(q; 0)Y_{mj}, Y_{nk} \rangle &= I_{mjnk}^{(1)}(q) + I_{mjnk}^{(2)}(0), \\ &= L_{mjnk}^{(1)}(q) + L_{mjnk}^{(2)}(0), \end{aligned}$$

where

$$\begin{aligned} I_{mjnk}^{(1)}(q) &= - \int_{B_{r_0}} q(x) u_{mj} \overline{\tilde{Y}_{nk}} dx, \\ I_{mjnk}^{(2)}(0) &= -\kappa^2 \int_{B_{r_0}} (u_{mj} - \dot{u}_{mj}) \overline{\tilde{Y}_{nk}} dx, \\ L_{mjnk}^{(1)}(q) &= - \int_{B_{r_0}} q(x) \overline{v_{nk}} \tilde{Y}_{mj} dx, \\ L_{mjnk}^{(2)}(0) &= -\kappa^2 \int_{B_{r_0}} \overline{(v_{nk} - \dot{v}_{nk})} \tilde{Y}_{mj} dx. \end{aligned}$$

Remark 4.3. *Since the potential q is real-valued, $u_{mj} = v_{mj}$ and $\dot{u}_{mj} = \dot{v}_{mj}$.*

Let $M_{mjnk}^{(l)}$ ($l = 1, 2$) be defined in (3.14), then we again obtain (3.15), that is,

$$|\langle \Gamma(q; 0)Y_{mj}, Y_{nk} \rangle| \leq |M_{mjnk}^{(1)}(q)| + |M_{mjnk}^{(2)}(0)|.$$

Following exactly the same arguments in the proof of Lemma 3.4, we have

$$\begin{cases} |M_{mjnk}^{(1)}(q)| \leq C''_d(1+\ell)r_0^\ell & \text{for all } q \in B_{+, \frac{1}{4}\kappa_1}^{\infty, r_0} \\ |M_{mjnk}^{(1)}(q)| \leq C''_d\kappa^2(1+\ell) \end{cases}$$

Consequently, we can replace the function Φ in (3.16) and (3.17) by the constant 1. The observation significantly improves our estimates later.

Here, we remark that due to (4.6), we choose $R = \frac{1}{4}\kappa_1$. Repeating the same arguments as in the proof of Proposition 3.6, we can prove

Proposition 4.4. *Let $s > \frac{d+2}{2}$ and $0 < \kappa^2 \leq \frac{1}{4}\kappa_1$. Define $\tau = s - \frac{d+2}{2}$. Given any $0 < \delta < 1$, there exists a δ -net Y of $((\Gamma_{mjnk}^0(B_{+, \frac{1}{4}\kappa_1}^{\infty, r_0})), \|\bullet\|_{X_s})$ with*

$$(4.10) \quad \log |Y| \leq \eta \left[1 + \log \left(1 + \frac{1}{\delta} \right) + \frac{\kappa^2}{\delta} + \left(\frac{\kappa^2}{\delta} \right)^{\frac{1}{\tau}} \right]^{2d}$$

for some constant $\eta = \eta(d, s, r_0)$.

As above, we choose $s = \frac{d+4}{2}$, i.e. $\tau = 1$, and (4.10) becomes

$$(4.11) \quad \log |Y| \leq \eta \left[1 + \log \left(1 + \frac{1}{\delta} \right) + \frac{2\kappa^2}{\delta} \right]^{2d}.$$

Observe that

$$\inf_{0 < \delta < 1} \left[\log \left(1 + \frac{1}{\delta} \right) + \frac{2\kappa^2}{\delta} \right] \leq \log 2 + \frac{\kappa_1}{2}$$

Therefore, given any θ satisfies

$$0 < \theta < \left(\log 2 + \frac{\kappa_1}{2} \right)^{-2\alpha},$$

there exists a unique $\tilde{\delta} \in (0, 1)$ such that

$$(4.12) \quad \theta^{-\frac{1}{2\alpha}} = \log \left(1 + \frac{1}{\tilde{\delta}} \right) + \frac{2\kappa^2}{\tilde{\delta}}.$$

We see that (4.12) is similar to (4.4), except that the multiplier $(1 + \kappa^2)^{-1}$ in front of the right-hand-side of (4.4) is removed. In view of (4.12), (4.11) is reduced to

$$(4.13) \quad \log |Y| \leq \eta (1 + \theta^{-\frac{1}{2\alpha}})^{2d}.$$

Based on these observations, we are ready to prove Theorem 1.4.

Proof of Theorem 1.4. By Proposition 2.1, there exists a constant $\mu > 0$ such that for all $\beta > 0$ and for all $\theta \in (0, \mu\beta)$:

there exists a θ -discrete subset Z of $(\mathcal{N}_{\alpha, \beta}^\theta(B_{r_0}), \|\bullet\|_{L^\infty})$

with

$$|Z| \geq \exp \left[2^{-(d+1)} \left(\frac{\mu\beta}{\theta} \right)^{\frac{d}{\alpha}} \right].$$

It suffices to restrict $0 < \theta < \min\{\frac{1}{2}\kappa_1 + \log 2\}^{-2\alpha}, \frac{1}{4}\kappa_1, \mu\beta\}$. Therefore, we can find $0 < \tilde{\delta} < 1$ such that (4.12) is satisfied. Since $0 < \theta < \frac{1}{4}\kappa_1$, it is clear that $Z \subset \mathcal{N}_{\alpha, \beta}^\theta(B_{r_0}) \subset B_{+, \frac{1}{4}\kappa_1}^{\infty, r_0}$. Recall that we have chosen $q_{\text{ref}} = 0$. From Proposition 4.4, there exists a $\tilde{\delta}$ -net Y of $((\Gamma_{m j n k}^0(B_{+, \frac{1}{4}\kappa_1}^{\infty, r_0})), \|\bullet\|_{X_s})$ with $s = \frac{d+4}{2}$ such that (4.13) holds. Clearly, Y is also a $\tilde{\delta}$ -net of $((\Gamma_{m j n k}^0(Z)), \|\bullet\|_{X_s})$.

We now choose $\beta > \frac{1}{4}\kappa_1$ sufficiently large, depending on d, μ, θ, κ , such that

$$|Z| \geq \exp \left[2^{-(d+1)} \left(\frac{\mu\beta}{\theta} \right)^{\frac{d}{\alpha}} \right] > \exp \left[\eta (1 + \theta^{-\frac{1}{2\alpha}})^{2d} \right] \geq |Y|.$$

Therefore, we can choose two different $q_1, q_2 \in Z \subset \mathcal{N}_{\alpha, \beta}^\theta(B_{r_0}) \subset L^\infty(B_1)$ such that there exists $y \in Y$ such that

$$\begin{aligned} \|\Gamma_{m j n k}(q_1; 0) - y_{m j n k}\|_{X_s} &\leq \tilde{\delta}, \\ \|\Gamma_{m j n k}(q_2; 0) - y_{m j n k}\|_{X_s} &\leq \tilde{\delta}. \end{aligned}$$

It follows from Proposition 2.2 that

$$\|\Lambda_{q_1} - \Lambda_{q_2}\|_{\frac{d+4}{2} \rightarrow -\frac{d+4}{2}} \leq 8\sqrt{2}\tilde{\delta}.$$

Again, we estimate $\tilde{\delta}$ by θ .

- **Case 1.** If $\frac{2\kappa^2}{\tilde{\delta}} \leq \log(1 + \frac{1}{\tilde{\delta}})$, then (4.12) yields

$$\theta^{-\frac{1}{2\alpha}} \leq 3 \log \left(1 + \frac{1}{\tilde{\delta}} \right) \leq 3 \log \left(\frac{2}{\tilde{\delta}} \right),$$

namely,

$$\tilde{\delta} \leq 2 \exp \left(-\frac{1}{3} \theta^{-\frac{1}{2\alpha}} \right).$$

- **Case 2.** If $\frac{2\kappa^2}{\tilde{\delta}} \geq \log(1 + \frac{1}{\tilde{\delta}})$, then (4.12) implies

$$\theta^{-\frac{1}{2\alpha}} \leq \frac{3\kappa^2}{\tilde{\delta}}, \quad \text{i.e.,} \quad \tilde{\delta} \leq 3\kappa^2 \theta^{\frac{1}{2\alpha}}.$$

Putting these two cases together, we obtain

$$\tilde{\delta} \leq 2 \exp \left(-\frac{1}{3} \theta^{-\frac{1}{2\alpha}} \right) + 3\kappa^2 \theta^{\frac{1}{2\alpha}},$$

which gives (1.9). By the fact that Z is θ -discrete subset with respect to the norm $\|\bullet\|_{L^\infty(B_1)}$, we then obtain (1.5). The proof is now completed. \square

ACKNOWLEDGEMENTS

Kow and Wang were partly supported by MOST 108-2115-M-002-002-MY3 and 109-2115-M-002-001-MY3. Uhlmann was partly supported by NSF, a Walker Family Endowed Professorship at UW and a Si-Yuan Professorship at IAS, HKUST.

REFERENCES

- [Ale88] G. Alessandrini, *Stable determination of conductivity by boundary measurements*, Appl. Anal. **27** (1988), 153–172, [MR0922775](#), [doi:10.1080/00036818808839730](#).
- [CHP03] D. Colton, H. Haddar, and M. Piana, *The linear sampling method in inverse electromagnetic scattering theory*, Special section on imaging. Inverse Problems **19** (2003), no. 6, S105–S137, [MR2036524](#), [doi:10.1088/0266-5611/19/6/057](#).
- [DI07] Deepak Aralumallige Subbarayappa and V. Isakov, *On increased stability in the continuation of the Helmholtz equation*, Inverse Problems **23** (2007), no. 4, 1689–1697, [MR2348729](#), [doi:10.1088/0266-5611/23/4/019](#).
- [DI10] Deepak Aralumallige Subbarayappa and V. Isakov, *Increasing stability of the continuation for the Maxwell system*, Inverse Problems **26** (2010), no. 7, 074005, 14 pp, [MR2644032](#), [doi:10.1088/0266-5611/26/7/074005](#).
- [DR03] M. Di Cristo and L. Rondi, *Examples of exponential instability for inverse inclusion and scattering problems*, Inverse Problems **19** (2003), no. 3, 685–701, [MR1984884](#), [doi:10.1088/0266-5611/19/3/313](#).
- [FSU10] J. Feldman, M. Salo and G. Uhlmann, *The Calderón problem - an introduction to inverse problems, a partially completed preview version of the textbook* (2010).
- [HI04] T. Hrycak and V. Isakov, *Increased stability in the continuation of solutions to the Helmholtz equation*, Inverse Problems **20** (2004), no. 3, 697–712, [MR2067496](#), [doi:10.1088/0266-5611/20/3/004](#).
- [ILX20] V. Isakov, S. Lu, and B. Xu, *Linearized inverse Schrödinger potential problem at a large wavenumber*, SIAM J. Appl. Math. **80** (2020), no. 1, 338–358, [MR4062806](#), [doi:10.1137/18M1226932](#), [arXiv:1812.05011](#).
- [IN12] M. I. Isaev and R. G. Novikov, *Energy and regularity dependent stability estimates for the Gel'fand inverse problem in multidimensions*. J. Inverse Ill-Posed Probl. **20** (2012), no. 3, 313–325, [MR2984491](#), [doi:10.1515/jip-2012-0024](#).

- [INUW14] V. Isakov, S. Nagayasu, G. Uhlmann, and J.-N. Wang, *Increasing stability of the inverse boundary value problem for the Schrödinger equation*, Inverse problems and applications, 131–141, Contemp. Math., **615**, Amer. Math. Soc., Providence, RI (2014), [MR3221602](#), [doi:10.1090/conm/615](#), [arXiv:1302.0940](#).
- [Isa07] V. Isakov, *Increased stability in the continuation for the Helmholtz equation with variable coefficient*, Control methods in PDE-dynamical systems, 255–267, Contemp. Math. **426**, Amer. Math. Soc., Providence, RI (2007), [MR2311529](#), [doi:10.1090/conm/426](#).
- [Isa11] V. Isakov, *Increasing stability for the Schrödinger potential from the Dirichlet-to-Neumann map*, Discrete Contin. Dyn. Syst. Ser. S, **4** (2011), no. 3, 631–640, [MR2746425](#), [doi:10.3934/dcdss.2011.4.631](#).
- [Isae13] M. I. Isaev, *Instability in the Gel'fand inverse problem at high energies*. Appl. Anal. **92** (2013), no. 11, 2262–2274, [MR3169162](#), [doi:10.1080/00036811.2012.731501](#), [arXiv:1206.2328](#).
- [KT61] A. Kolmogorov and V. Tikhomirov, ε -entropy and ε -capacity in functional spaces. Mat. Nauk **14** (1959), pp. 3–86; English transl.: Amer. Math. Soc. Transl. **17** (1961), 277–364, [MR0124720](#), [MR0112032](#), [doi:10.1007/978-94-017-2973-4_7](#).
- [KW20] P.-Z. Kow and J.-N. Wang, *Reconstruction of an impenetrable obstacle in anisotropic inhomogeneous background*, IMA J. Appl. Math. **86** (2021), no. 2, 320–348, [MR4246858](#), [doi:10.1093/imamat/hxab002](#).
- [KU19] K. Krupchyk and G. Uhlmann, *Stability estimates for partial data inverse problems for Schrödinger operators in the high frequency limit*, J. Math. Pures Appl. (9) **126** (2019), 273–291, [MR3950020](#), [doi:10.1016/j.matpur.2019.02.017](#), [arXiv:1712.06200](#).
- [LLU19] R.-Y. Lai, Q. Li, and G. Uhlmann, *Inverse problems for the stationary transport equation in the diffusion scaling*, SIAM J. Appl. Math. **79** (2019), no. 6, 2340–2358, [MR4039537](#), [doi:10.1137/18M1207582](#), [arXiv:1808.02071](#).
- [Man01] N. Mandache, *Exponential instability in an inverse problem for the Schrödinger equation*, Inverse Problems **17** (2001), no. 5, 1435–1444, [MR1862200](#), [doi:10.1088/0266-5611/17/5/313](#).
- [NSU88] A. Nachman, J. Sylvester, and G. Uhlmann, *An n -dimensional Borg-Levinson theorem*, Comm. Math. Phys. **115** (1988), no. 4, 595–605, [MR0933457](#), [doi:10.1007/BF01224129](#).
- [NUW13] S. Nagayasu, G. Uhlmann, and J.-N. Wang, *Increasing stability in an inverse problem for the acoustic equation*, Inverse Problems **29** (2013), no. 2, 025012, 11 pp, [MR3020433](#), [doi:10.1088/0266-5611/29/2/025012](#), [arXiv:1110.5145](#).
- [RS18] A. Rüland and M. Salo, *Exponential instability in the fractional Calderón problem*, Inverse Problems **34** (2018), no. 4, 045003, 21 pp, [MR3774704](#), [doi:10.1088/1361-6420/aaac5a](#), [arXiv:1711.04799](#).
- [Sun90] Z. Sun, *On continuous dependence for an inverse initial boundary value problem for the wave equation*, J. Math. Anal. Appl. **149** (1990), 188–204, [MR1059582](#), [doi:10.1016/0022-247X\(90\)90207-V](#).
- [ZZ19] H. Zhao and Y. Zhong, *Instability of an inverse problem for the stationary radiative transport near the diffusion limit*, SIAM J. Math. Anal. **51** (2019), no. 5, 3750–3768, [MR4007614](#), [doi:10.1137/18M1222582](#), [arXiv:1809.01790](#).

DEPARTMENT OF MATHEMATICS, NATIONAL TAIWAN UNIVERSITY, TAIPEI 106, TAIWAN.

Email address: d07221005@ntu.edu.tw

DEPARTMENT OF MATHEMATICS, UNIVERSITY OF WASHINGTON, BOX 354350, SEATTLE, WA 98195-4350, USA, AND INSTITUTE FOR ADVANCED STUDY OF THE HONG KONG UNIVERSITY OF SCIENCE AND TECHNOLOGY, HONG KONG.

Email address: gunther@math.washington.edu

INSTITUTE OF APPLIED MATHEMATICAL SCIENCES, NATIONAL TAIWAN UNIVERSITY, TAIPEI 106, TAIWAN.

Email address: jnwang@math.ntu.edu.tw