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Article

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Charge Transfer Plasmons in Dimeric Electron Clusters

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Abstract

The tunability of the optical response of dimers of metal clusters and nanoparticles makes them ideal for many applications from sensing and imaging to inducing chemical reactions. We have studied charge transfer plasmons in separate and linked dimers of closed-shell electron clusters of 8 and 138 electrons using time-dependent density functional theory. The simple model clusters enable the systematic study of the charge transfer phenomenon from the electronic perspective. To identify the charge transfer plasmons, we have developed an index, the Charge Transfer Ratio, for quantifying the charge transfer nature of the excitations. In addition, we analyze the induced transition density and the electron transitions contributing to the dipole moment at the charge transfer plasmon energies. Our results show that the optical response of the dimers is very sensitive to changes in the inter-cluster separation and in the width of the linking channel, with charge transfer plasmon peaks appearing at low energies for dimers with linking or sufficient electron cloud overlap.

Introduction

The optical response of metal clusters and nanoparticles is highly dependent on the shape, size, and dielectric environment of the cluster.¹ The response of two adjacent particles depends also on the separation.² This tunability makes clusters and cluster assemblies or arrays ideal for many applications, such as surface-enhanced Raman spectroscopy for detecting single molecules,³ the molecular ruler for measuring distances in the nanoscale,⁴ cancer imaging and treatment,⁵ and driving chemical reactions.⁶

Plasmons are collective excitations of the conduction electrons, where the electron density of the cluster oscillates in resonance with the external electromagnetic field. In plasmonic metal clusters, the main excitation is the localized surface plasmon resonance (LSPR), where the electron density oscillations happen mostly on the surface of the cluster.⁷ There have been several approaches to understand and classify different optical excitations in metal

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3 clusters and nanoparticles, especially to recognize the LSPR peaks from the optical spectra.
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5 The distribution of the induced density^{8,9} and collectivity of the excitation^{10,11} have both
6
7 been used as criteria for plasmonic resonances. Different plasmonic indices based on the
8
9 strength of the induced potential have also been developed.^{12,13} Recently, Gieseck et al.¹⁴
10
11 proposed a method to identify plasmonic resonances based on three criteria - the collectivity
12
13 of the excitations, the superatomic character of the excitations, and the additivity of the
14
15 contributions from the single-particle excitations.
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18 In dimers, the plasmons of the two clusters can couple capacitively to form a bonding
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20 dipolar plasmon (BDP).¹⁵ The energy of the excitation is lowered, causing the BDP peak to
21
22 be redshifted in comparison to the LSPR peak of the single clusters.^{2,15} For short separations
23
24 and chemically linked dimers there can also be charge transfer plasmons (CTPs).^{16,17} In
25
26 these CTPs, the electron density oscillates between the two clusters, making one cluster
27
28 momentarily positively and the other negatively charged. The CTPs happen at energies lower
29
30 than BDPs, and thus add a new energy range for the applications of plasmons. Depending on
31
32 the chemical coupling between the clusters, the CTPs can be caused by electrons tunneling
33
34 or flowing between the clusters.¹⁷

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36 Charge transfer plasmons have been experimentally observed in many dimeric nanos-
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38 tructures, such as nanogaps between two gold coated atomic force microscope tips,¹⁸ gold
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40 nanorods connected with a gold nanojunction,¹⁹ and dimers of various nanoparticles, such
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42 as gold nanoshells,¹⁶ silver spheres,²⁰ and gold nanoblocks.²¹ For atomistically precise thiol-
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44 stabilized (SR) gold nanoclusters, Lahtinen et al. observed an absorption peak that they
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46 assigned to a tunneling CTP in covalently linked $\text{Au}_{\sim 250}(\text{SR})_n$ dimers.²²

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48 Large nanoparticles and dimers with a separation of several nanometers can be treated
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50 with classical methods solving the Maxwell equations, and good agreement with experiments
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52 have been found for nanoparticle dimers in such cases.²³ However, going to sizes of a few
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54 nanometers²⁴ or to sub-nanometer separations,^{17,25} the quantum effects cannot be ignored.
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56 Thus, to model dimer systems with small, atom-precise clusters with more accuracy, time-

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3 dependent density functional theory (TDDFT) has been employed in multiple studies.^{26–31}
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5 The drawback of TDDFT is that it is computationally heavy. Therefore, in the region
6 where quantum mechanical effects are important, but the size of the system makes atomistic
7 TDDFT calculations unfeasible, some simplifications have to be made. Xiang et al. have
8 employed the time-dependent orbital-free density functional theory (TD-OFDFT) to study
9 dimers and trimers of sodium clusters consisting of 55 or 1415 atoms in order to reduce the
10 computational cost.³² The quantum-corrected model (QCM)^{33,34} includes quantum effects
11 in the optical response of nanoparticles with a small gap by modeling electron tunneling by a
12 fictitious conductive material with the local permittivity reflecting the tunneling probability.
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21 One approach is to combine TDDFT with the simple jellium model to represent the
22 cluster dimers.^{17,35–38} There, the nuclei and core electrons are represented as a uniform,
23 positive background charge, and only the valence electrons are modeled explicitly. For
24 elements where the d-electron screening plays a role, such as gold and silver, a dielectric
25 background can be added to model this effect. The simple jellium model with a constant
26 background density is most suitable for alkali metals, such as sodium, for which the s-type
27 valence electrons are strongly delocalized. Different elements of this group can be represented
28 by tuning the background density to correspond to the experimental density.
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37 Since the jellium model allows for electron tunneling and the spilling of the conductance
38 electrons, it enables charge transfer excitations for small gap sizes, and can thus be used to
39 understand the principles of charge transfer plasmons in dimeric structures. To accurately
40 reproduce experimental results quantitatively, the background charge distribution at the edge
41 of the cluster should be adjusted. In close distances, the atomic structure has been shown
42 to affect the electric field enhancement and the charge transfer current.^{24,39,40} Zhang et al.³⁹
43 have shown that for spherical jellium clusters of about 300 atoms, the jellium model gives the
44 same trend for the BDP as atomistic models using bcc and icosahedral packing. The CTP
45 peak is also reproduced by the jellium model, but only at about a separation of 10 Å, whereas
46 for the two atomic arrangements a charge transfer mode was seen already at separations of
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3 30–40 Å. The results indicate that the main trends in the spectra are reproduced by the
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5 jellium model, and at sufficiently small separations, the jellium model is able to describe the
6
7 CTP.
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9 Here, we have employed the jellium model with TDDFT for dimers of spherical clusters
10
11 to study charge transfer plasmons from the electronic perspective. The simplicity of the
12
13 background charge potential eases the investigation of the electronic phenomena, since the
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15 ground state electron density has a very simple structure, and the changes in electron density
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17 during excitations are therefore easier to distinguish and analyze as compared to atomic
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19 structures. Since we are using the same model for all the systems, we can follow the trends
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21 in the shape of the spectra and the peak energies even if the jellium model causes a shift
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23 compared to atomic clusters.
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25 We studied homo- and heterodimers of two 8-electron or 138-electron clusters and varied
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27 the inter-cluster separation between the clusters. We also constructed some linked dimers
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29 by connecting the clusters with a cylindrical tunnel of the same jellium background density
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31 to gain understanding on the effect of conductive linking to the excitations. The conductive
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33 jellium linker establishes a coupling of metallic nature between the clusters, allowing electrons
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35 to flow between the clusters during a CTP excitation. These linked clusters serve as simple
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37 models for clusters that are connected by a conductive linker, such as one or more conductive
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39 ligands or a chain of metal atoms. It is possible to modify the background density of the
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41 linker to study the effect of the linker conductance. However, here we use the same density
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43 as for the clusters to have a linker with the same metallic character. For atomistic clusters,
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45 this kind of linking has been studied by Rossi et al.,²⁹ who studied CTPs in a system with
46
47 a narrow atomic contact between to metal clusters.
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49 Following the methods of our previous study of the birth of LSPR in jellium clusters,⁴¹
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51 we use the induced density and analysis of the contributions of the single-particle excitations
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53 to study the nature of the excitations in the dimer systems. In addition, we have developed
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55 an index to quantify the charge transfer nature of the excitations, enabling us to recognize
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3 the excitations with charge transfer character from the spectra and follow the development
4 of the CTPs as the separation gets smaller.
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8 9 Theory and methods 10

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12 The GPAW^{42,43} program, which performs the calculations using a uniform real-space grid,
13 was used for the density functional theory (DFT) calculations. The local density approxima-
14 tion (LDA) after Perdew and Wang⁴⁴ was employed for the exchange-correlation functional.
15 Typically, LDA is considered to yield inappropriate excitation energies due to lack of di-
16 vergence of exchange-correlation kernel at charge transfer excitations. In the supplementary
17 information (section 1), we discuss why charge transfer plasmons differ from this picture, and
18 why already LDA yields a balanced description of the intra and inter particle Coulomb forces.
19 The optical absorption spectra were calculated with linear response time dependent DFT
20 (lr-TDDFT) as implemented in GPAW.⁴⁵ Some of the peaks in the spectra were then further
21 analyzed using time-dependent density functional perturbation theory (TD-DFPT),⁴⁶ which
22 allows us to solve the contributions of the Kohn-Sham (K-S) electron-hole (e-h) transitions
23 contributing to the excitation and to calculate the induced transition density. The K-S con-
24 tributions were visualized using the dipole transition contribution map (DTCM) scheme,^{47,48}
25 where the relative strengths of the contributions to the transition dipole moment from all
26 the K-S transitions can be seen at the same time. To study the excitations with charge
27 transfer between the two clusters, the DTCM analysis is done for light polarized parallel to
28 the dimer axis for all the systems.
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46 Since the clusters of the dimers are spherical, the symmetries of the superatom-like elec-
47 tron orbitals were identified by doing the Y_{lm} analysis,⁴⁹ where the part of the wavefunction
48 inside a certain projection sphere is projected into spherical harmonics. This analysis was
49 done separately for the two clusters in one dimer, summing the contributions for the pro-
50 jected density of states (PDOS) figures. Thus these symmetries are the symmetries for one
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3 cluster, not the whole dimer system, where the orbitals are bonding and antibonding states
4 of the single-cluster orbitals. Introducing the conductive pathway breaks the local spherical
5 symmetry already at ground state level and thus allows the charge transfer excitations (also
6 discussed in supplementary section 1).
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11 In the simple jellium model used here, the edges of the positive uniform charge distribu-
12 tion were sharp. The jellium background of a linked dimer can be seen in figure 1 a). The
13 radius of the clusters r , the separation between the jellium edges S , and the radius of the
14 cylindrical linker a are shown. The electron density experiences spill-out from the positive
15 core, and this spill-out is relatively larger for smaller clusters. The electron densities of the
16 studied systems along the dimer axis can be seen in figure S1.
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23 For the separate spheres, the density used for the positive background was approximately
24 the density of sodium, calculated from the Wigner-Seitz radius $r_{ws} = 2.08 \text{ \AA}$. For the linked
25 spheres, the volume of the spheres was kept the same, so the volume of the dimer was changed
26 by the volume of the linker. Thus the jellium density is a little smaller for the linked systems,
27 maximum 5.6 % for the 8-electron sphere system with the wider linker.
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33 Because of the smooth jellium background, the grid spacing could be as big as 0.4 \AA
34 for the dimers of 8-electron clusters and 0.5 \AA for the dimers of 138-electron clusters. The
35 convergence of the energies relative to the grid spacing was tested. The energy of the
36 systems was converged with the accuracy of 0.5 meV per electron. In the optical spectrum
37 calculations, the energy cut-off was 5 eV or larger, meaning that at least all the e-h transitions
38 with a smaller energy difference than this were included. All the spectra are plotted for light
39 polarized along the dimer axis.
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47 To analyze the CT nature of the excitations, we introduce the *Charge Transfer Ratio*
48 (CTR) which is defined in equation 1, and which can be calculated for any energy in the
49 spectrum:
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$$\text{CTR} = \left| \frac{D \int_A \rho_{\text{ind}}(\mathbf{r}) d\mathbf{r}}{\int_{A+B} x \rho_{\text{ind}}(\mathbf{r}) d\mathbf{r}} \right| \quad (1)$$

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56 Here ρ_{ind} is the induced transition density at the studied energy, x the coordinate on the
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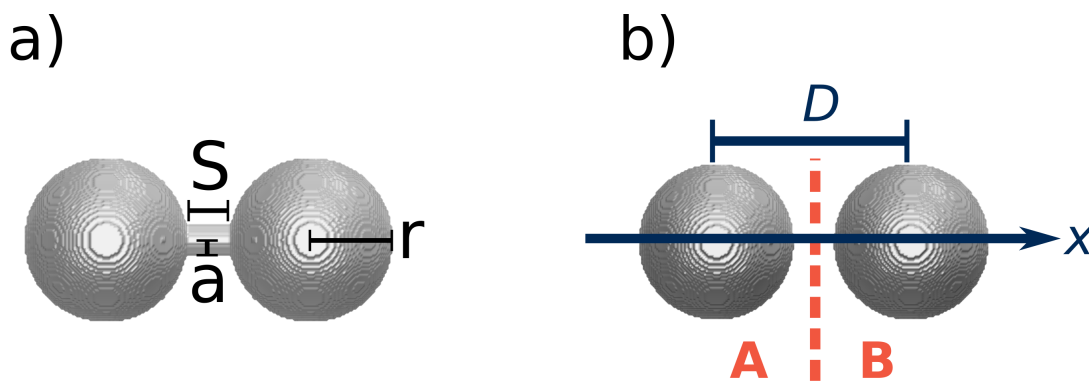


Figure 1: a) A coupled dimer with a radius of the spheres r , separation S , and radius of the linking channel a . Shown is the positive jellium background. b) Explanation of the symbols used in equation 1. A and B are the left and right sides of the calculation box, D is the distance between the centers of the clusters, and x is the direction of the dimer axis.

dipole axis and D the distance between the centers of the two clusters. The calculation box is divided into two equal sides, A and B , where the division is in the middle of the gap between the clusters. See figure 1 b) for a visual guide to the variables.

The denominator of equation 1 is the dipole moment of the system along the dimer axis, and the numerator is the dipole moment of a system where all the induced charge on side A is concentrated in the center of cluster A , and similarly for side B . The numerator is thus a measure of the charge separation between sides A and B , and as a result also between the two clusters. The numerator acts as normalization. The CTR compares the 'CT moment', the expression in the numerator, to the dipole moment: the higher the number is, the bigger the CT nature of the excitation is. If both clusters are neutral, i.e. the electron density oscillates separately inside each cluster, the integral of the induced density on side A , and thus the CTR value, is zero. Depending on the distribution of the charges, the ratio can also be negative or greater than one, but the absolute value is the measure of the charge transfer characteristic. With these numbers we can then plot a charge transfer spectrum showing how the charge transfer characteristics of the excitations evolve as a function of energy. Further discussion on the details of the CTR is given in the SI text.

Results and discussion

The studied dimer systems are named **1–12**. The parameters for each system are shown in table 1. The separations are given in terms of the radii of the systems, $r \approx 4.2 \text{ \AA}$ for the 8-electron clusters and $R \approx 10.7 \text{ \AA}$ for the 138-electron cluster.

Table 1: The parameters of the studied dimer systems. S is the separation between the jellium edges and a is the radius of the cylindrical linking channel in the case of linked dimers. $r \approx 4.2 \text{ \AA}$ and $R \approx 10.7 \text{ \AA}$ refer to the radii of the 8-electron and 138-electron clusters, respectively. See figure 1 a) for a visual guide to the parameters.

system	number of el.	S	a (\AA)
1	8+8	$2r$	-
2	8+8	$0.5r$	-
3	8+8	0	-
4	8+8	$2r$	1.0
5	8+8	$0.5r$	1.0
6	8+8	$0.5r$	2.1
7	138+138	$0.5R$	-
8	138+138	$0.25R$	-
9	138+138	$0.1R$	-
10	138+138	$0.5R$	1.0
11	138+138+2	$0.5R$	2.1
12	8+138	$0.5R$	-

Dimers of 8-electron clusters. The absorption spectra and the CTR spectra for the 8-electron cluster dimers can be seen in figure 2. The spectra are for light polarized in the direction of the dimer axis. The CTR values are calculated at the energies of the electronic excitations seen in the oscillation strength spectra without broadening. The upper row is for separate clusters and the lower for linked.

For dimer **1** with a separation of $S/r = 2$, there is virtually no electron-cloud overlap, as can be seen from figure S1 a). The absorption spectrum is almost identical to that of a single cluster.⁴¹ Here, the CTR is also very close to zero along the whole spectrum, indicating that there is no charge transfer between the clusters at any energies.

When the clusters are brought closer together, however, the increasing overlap causes

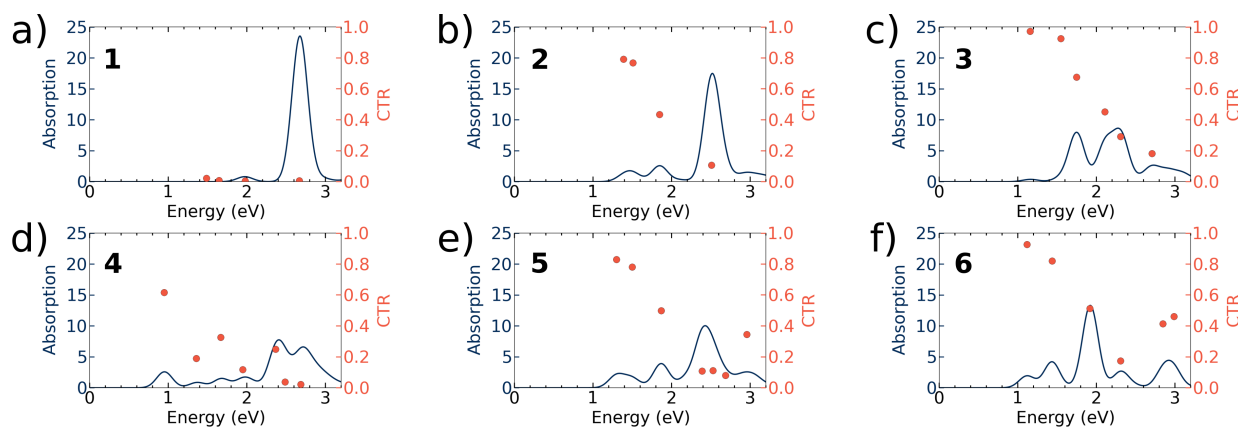


Figure 2: The absorption spectrum and the CTR for the dimers **1–6** with two 8-electron clusters.

coupling: although the BDP peak is still clearly the strongest peak for dimer **2** with the separation-radius ratio $S/r = 0.5$, some of the strength has already been transferred to lower energies. In the absorption spectrum this is seen as a red shift of the main peak and the appearance of an extra peak at energies less than 2 eV. The lowest-energy excitations have some charge-transfer nature, as proved by the CTR analysis discussed later. For dimer **3** with touching cluster, $S/r = 0$, the shape of the spectrum is already very different from separate cluster, with two larger peaks, and a significant CTR value even near 2 eV.

Going from larger to smaller separation, we can see the transfer of oscillation strength from the BDP peak to lower energies and excitations of CT nature. The spectrum is very sensitive to even small changes of a couple of Ångströms in the separation. This can be explained by changes in the energies and shapes of the electron orbitals, and by the drastic changes in the electron density between the clusters, as seen in figures S1 a) –b).

We can also couple the clusters by connecting them with a conducting linker, as for dimers **4–6**. From figures S1 d)–e) we can see that some of the electron density is now located in the linker. For the systems with the narrower linker (1.0 Å), **4** and **5**, the electron density between the clusters is only moderately bigger than for the separate clusters. For dimer **6** with the stronger linker, however, the electron density in the linker is comparable to that of inside the clusters, allowing the electrons to flow freely between the two clusters.

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3 Coupling the clusters with the separation of $S/r = 2$ even with the narrow linker, we can
4 see a drastic change in the spectrum of dimer **4** compared to the separate clusters with the
5 same inter-cluster distance, dimer **1**. For dimer **4** there is not anymore one main peak, and
6 there are many new peaks at energies lower from the BDP peak. All the peaks below 2 eV
7 have non-zero CTR values.
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13 Bringing the clusters closer together, to $S/r = 0.5$, but keeping the linker radius constant,
14 the spectrum becomes again more compressed, and less peaks can be seen in the spectrum
15 of dimer **5**. However, the CTR is higher for the smallest peaks for the smaller separation. If
16 we, conversely, keep the distance constant $S/r = 0.5$ but increase the linker width to 2.1 Å,
17 the change in the positions of the clearest peaks is not very drastic, but the relative strengths
18 of the peaks change, as is evident from the spectrum of dimer **6**.
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23 For the systems **3–6** with a strong coupling, the BDP peak is fragmented. Also some of
24 these fragments have some charge transfer nature, but the CTR value goes to minimum in
25 this energy range, indicating that the electron density oscillations happen mostly separately
26 inside each cluster.
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33 We will now study in more detail one absorption peak with a large CTR value for three of
34 the systems. In figure 3 are shown the DTCMs and the induced densities for one prominent
35 CTP peak in systems **2**, **3**, and **4** each. In all of the systems, the induced densities of the
36 strongest CTP resonances have a similar shape when looked at an appropriate isosurface
37 value. Here, the isosurface value is the same for each system, which makes the weaker
38 plasmon for system **2** appear to have a different shape.
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45 The main contributions to the dipole moment come from P→D transitions. The coupling
46 between the two clusters has caused some splitting in the orbitals for each system, the most
47 in system **3** with the touching clusters. This hybridization of orbitals due to coupling is
48 an absolute requirement for charge transfer excitations. The first 20 orbitals (1S, 1P, 1D,
49 and 2S symmetries) for dimers **2–4** are visualized in figures S1 – S4. For systems **2** and **3**,
50 the lowest-energy 1P orbital is the bonding state of the individual P orbitals, aligned in the
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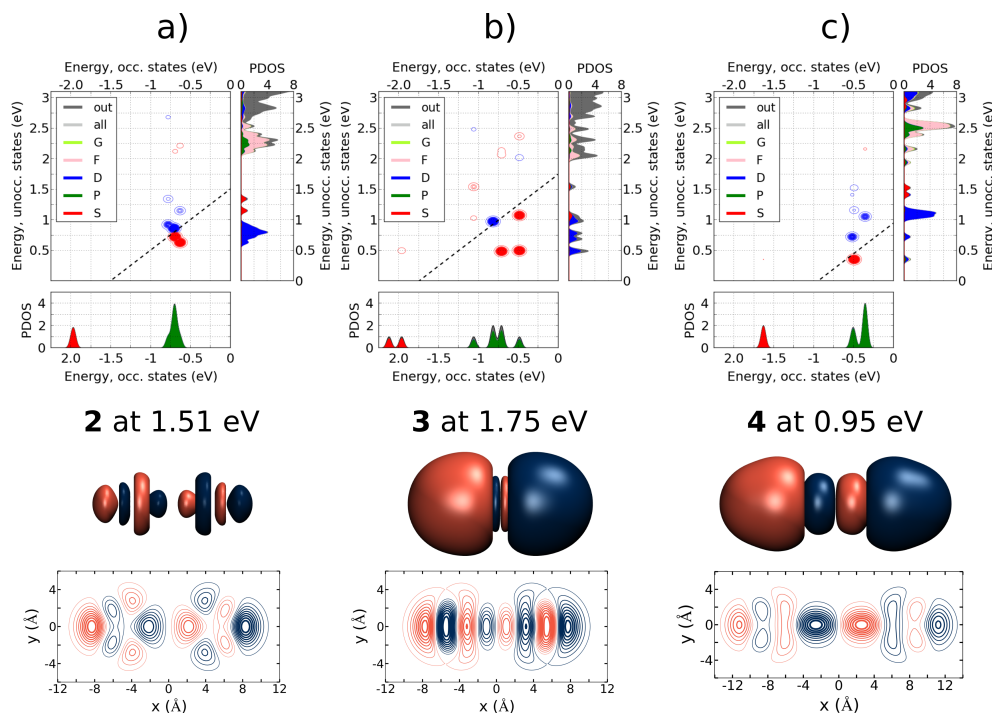


Figure 3: The DTCMs and induced densities for the $2 \cdot 8$ electron dimers **2**, **3**, and **4** for one CTP peak at energies indicated in the figure. The contour plot shows the relative strength of the contributions of each e-h transition to the induced dipole moment, read indicating positive and blue negative contribution. The colors of the projected density of states in the DTCM figures indicate the symmetries of the orbitals, as projected for spherical harmonics. The symmetry analysis is done for both clusters separately. The dashed lines are for $E_{ex} = \epsilon_{unocc} - \epsilon_{occ}$, where E_{ex} is the energy of the excitation, ϵ_{unocc} is the energy of the unoccupied K-S state, and ϵ_{occ} is the energy of the occupied K-S state. The induced densities are visualized as isosurfaces with a certain positive and negative value of the same absolute value in the middle row. The bottom row shows a contour plot of the induced density at $z = 0$, along the dimer plane. The isosurface values are the same for each system.

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3 dimer axis direction (σ orbital). This is followed by two bonding (π) and after those two
4 antibonding (π^*) states of orbitals aligned in different directions, and finally HOMO is the
5 antibonding state of the individual P orbitals aligned along the dimer axis (σ^*).
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9 For the two separate dimers **2** and **3**, the HOMO-LUMO transition is one of the strongest
10 positive contributions. LUMO in both systems is a bonding state of D orbitals, but for **2** of
11 the D_z^2 orbitals, and for **3** of the other type of D orbitals. For **3**, another positive contribution
12 is from P symmetry π^* to LUMO. For **2**, the other positive contribution is P symmetry π^*
13 to a bonding state of D orbitals. For **3**, there is a contribution from π^* to LUMO and from
14 HOMO to the antibonding state of 2S orbitals. Both systems have at least one screening
15 contribution from π to antibonding states of 1D.
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23 For the linked dimer, system **4**, the order of the 1P type orbitals is different: from lower
24 to higher energy, the states are σ , σ^* , $2 \times \pi$ and $2 \times \pi^*$, so the orbitals with more electron
25 density in the linker are lower in energy. Of these, the σ and σ^* orbitals are very close to
26 each other in energy, and the π and π^* orbitals are almost degenerate. Here, the positive
27 contribution comes from transition from the P symmetry σ^* orbital to the bonding state of
28 the D_z^2 orbitals aligned along the dimer axis. The main negative contributions are from P
29 type σ to the antibonding state of D_z^2 orbitals aligned along the dimer axis and from the
30 different σ or σ^* orbitals to different D orbital states.
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39 **Dimers of 138-electron clusters.** The absorption spectra and CTR spectra for the
40 138-electron cluster dimers can be seen in figure 4. In the upper row are again separate
41 clusters, from larger separation to smaller, and in the bottom row two linked systems with
42 the same separation but different linker width. For these larger clusters, the density of states
43 is bigger than for the 8-electron clusters, and thus the spectrum is more continuous, leading
44 to clearer trends in the response of the spectrum to changes in geometry. Going to smaller
45 separations, more and more absorption strength is transferred to smaller energies, which
46 correspond to excitations with a large CTR. The CTR value has a minimum near the BDP
47 peak. Linking the clusters has a similar effect, except the BDP peak is better preserved.
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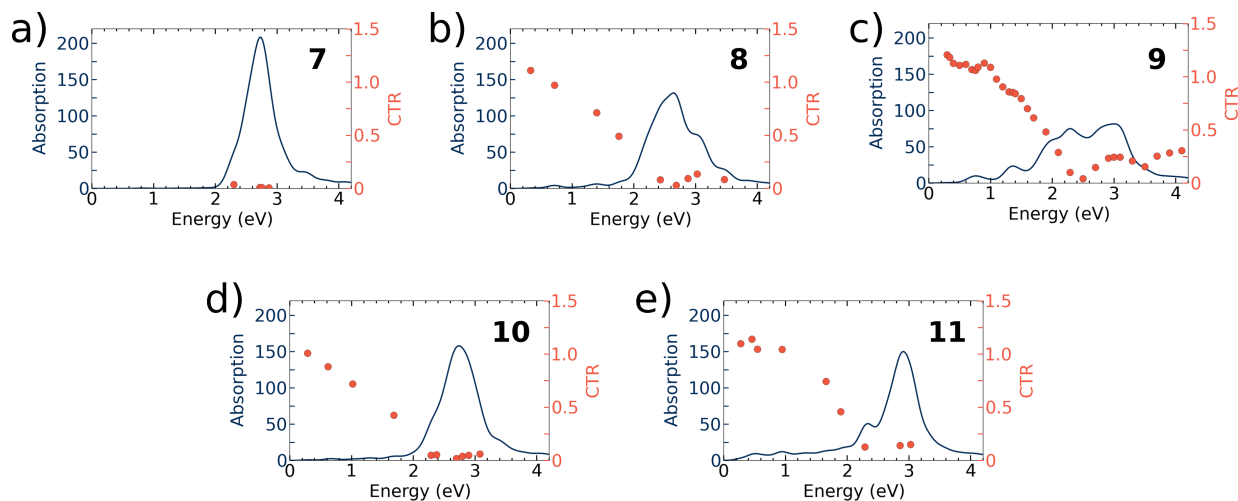


Figure 4: The absorption spectrum and the CTR for the dimers **7–11** with two 138-electron clusters.

One should notice that since the linker widths are the same here as they were for the smaller clusters, the linkers for these larger systems are relatively narrower compared to the cluster size.

The energies of the BDP and CTP peaks agree approximately with those obtained by Rossi et al.²⁹ using TDDFT for dimeric sodium systems with a total of 261 atoms. They varied the linking between the clusters in the dimer by starting from an uniform nanorod and stretching it to first form a narrow metallic bridge between two thicker ends and then finally breaking the system to two parts. The dimers with a narrow linker were found to have two CTP peaks between 0 and 1 eV and one BDP peak between 2 and 2.5 eV. The completely separate clusters supported a second BDP peak between 1 and 1.5 eV. Here, the highest CTR values are for absorption peaks below 1.5 eV, and the BDP peak, although strongly fragmented for systems **8** and **9**, is found between 2 and 3 eV. Our results are also qualitatively similar to experimental studies of dimeric silver²⁰ and gold nanoparticles^{19,21} and covalently linked gold clusters,²² with a CTP peak appearing at energies smaller than the BDP peak for small separations or touching or linked clusters.

If $2 \cdot 138$ electrons are placed in the potential of the positive jellium background with the wider linker (2.1 Å, the Wigner-Seitz radius used for the jellium background), the dimer

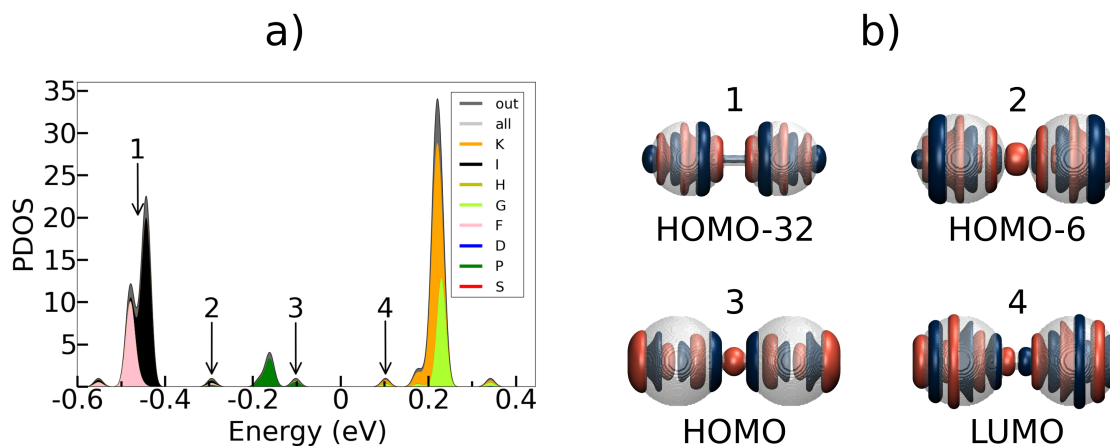


Figure 5: a) The projected density of states near the HOMO-LUMO gap and b) selected orbitals for dimer **11**. The arrows and numbers in a) indicate the orbitals that are shown in b). The jellium surface of the dimer is shown in transparent.

is not anymore a closed-shell system, with the HOMO-LUMO gap in the middle of the 3P states. This is caused by an extra electron orbital, HOMO-32, forming between the 2G and 1I orbitals. The orbital has some G and I symmetry inside the two clusters separately, but it is clearly asymmetric inside one cluster, and symmetric only with respect to the whole system. The positive and negative isosurfaces of this orbital can be seen in figure 5 b), where it is orbital number 1. To keep all the studied systems closed-shell, we then added two electrons and calculated the properties of this system. The dimer with the wider linker studied here has thus $2 \cdot 138 + 2$ electrons, two electrons more than the rest of the bigger dimers. The order and symmetries of the electronic states stayed the same after the addition of these extra electrons.

Closer inspection of the PDOS and the orbitals show that the orbitals aligned in the direction of the dimer axis are markedly separated in energy from the other orbitals, causing splitting in the DOS - the energy of the bonding combination of these orbitals is lowered, and the energy of the antibonding combination raised. The shapes of these orbitals are also distorted from the combination of the orbitals of two non-interacting clusters, making them quite asymmetric inside one cluster, and only symmetric regarding the whole dimer. Three of these orbitals are visualized in figure 5, where they are labeled as 2–4. In total, there is a

considerable amount of electron density in the linker, as can be seen from figure S1 k). As a result, the electrons can freely flow between the two clusters.

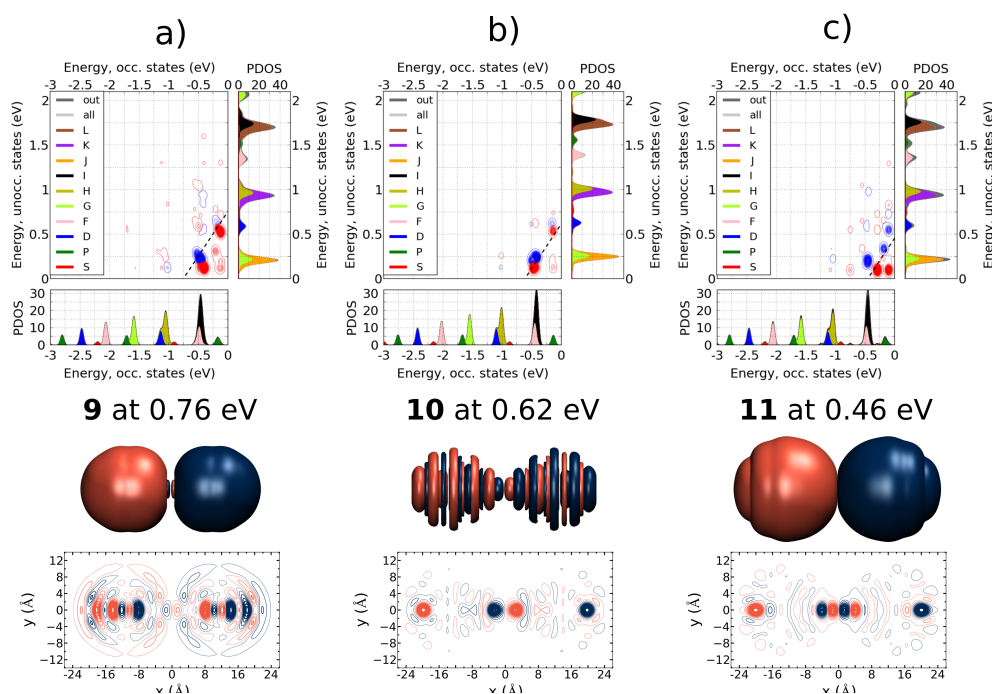


Figure 6: The DTCMs and induced densities for the $2 \cdot 138$ electron dimer systems **9**, **10**, and **11** for one CTP peak at energies indicated in the figure. See the caption of figure 3 for details.

The DTCMs and the induced densities for one peak with a high CTR value are shown for systems **9**, **10** and **11** in figure 6. In the DTCM, the difference between CTP and BDP excitations is that in the former there are stronger contributions from the screening transitions. The transitions are still mainly those that are allowed for a single spherical cluster, except of the transitions of the HOMO orbitals to the LUMO in system 11. These orbitals are the ones aligned along the dimer axis and having significant electron density in the linker.

As the CTR values in figure 4 show, for all of the energies analyzed in figure 6, in total one of the clusters is negatively and the other positively charged. However, there are also induced density variations inside the clusters. For the separate clusters, in system **9**, the induced density inside each cluster still somewhat resembles the induced density for the

LSPR peak for a single cluster,⁴¹ reflecting the shell structure of the electron density, which can be seen in figures S1 g)–i). Here, the highest density changes are near the center of the clusters. For the linked clusters, the picture changes. For system **10**, with the narrower linker, the biggest values for the induced density along the dimer axis are found at both ends of the clusters, forming two temporary dipoles inside each cluster. For the dimer with the wider linker, system **11**, additionally an extra dipole is formed in the linker.

From the DTCMs in figure 5 we can see that for systems **9** and **10**, the contributions come from transitions from 1I to 1K, 2F to 2G, and 3P to 3D. The main positive contributions for system **11** are from transitions from the HOMO-6 and HOMO orbitals to LUMO, as visualized in figure 5. The strongest negative contribution comes from transitions from the 1I orbitals to 1J orbitals, although there are some clearly weaker contributions from 3P to 3D and 3P to 1J/2G transitions.

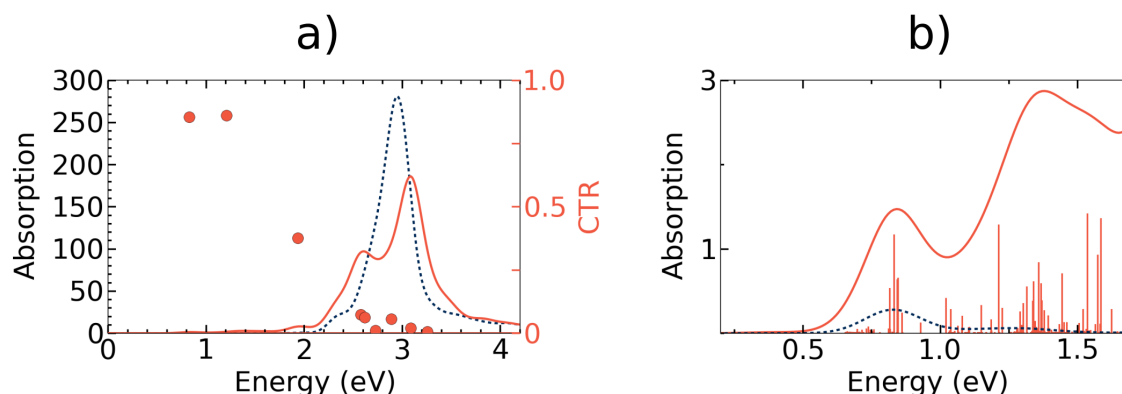


Figure 7: a) The absorption spectrum of the dimer **12** with 8- and 138-electron clusters (solid orange line) and the summed absorption spectra of the separate clusters with 8 and 138 electrons (dashed blue line). The CTR values are for the dimer. b) A zoom of the spectra between 0.2 eV and 1.8 eV, along with the oscillation strengths for the dimer.

Heterodimer. The heterodimer (dimer **12**) has two clusters of different sizes: the 138-electron cluster with a radius of 10.7 Å and the 8-electron cluster with a radius of 4.2 Å. The separation between the jellium edges was 2.7 Å, which is 1/4 of the radius of the larger cluster. The electron density overlap between the clusters, as shown in figure S1 l), is similar to that of the dimer of the 138-electron clusters with the same separation, seen in S1 h). In

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3 figure 7 a), the spectrum of the heterodimer is plotted in orange and the sum of the spectra
4 of the two separate clusters in blue. The CTR values are calculated for the dimer. The sides
5 A and B of the calculation box for equation 1 are again chosen so that the division goes
6 halfway between the jellium edges of the cluster. In the ground state, this division gives a
7 total of eight electrons on the side of the smaller cluster, and 138 on the side of the larger
8 cluster when integrating over the electron density. Figure 7 b) shows a zoom of the region
9 between 0.2 and 1.8 eV, with the oscillation strengths for the heterodimer visible. In figure
10 8 are shown the DTCMs and induced densities for the heterodimer at three energies.

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13 The individual 138-electron cluster has its LSPR peak at 2.95 eV, and the 8-electron
14 cluster at 2.72 eV.⁴¹ For the heterodimer, the two largest peaks for light polarized in the
15 dimer axis direction are found at 2.61 eV and 3.08 eV. The CTR value is close to zero for
16 both of these peaks, indicating little or no charge transfer between the cluster. Also the
17 other analysed excitations at energies 2.6 eV – 3.1 eV have small CTR values.

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20 For the first of the largest peaks, that is analyzed in figure 8 b), the induced density
21 inside the smaller cluster is similar to that of the individual 8-electron cluster at the LSPR
22 energy. From the contour plot we can see that most of the induced density is concentrated
23 around the smaller cluster. The main transition of the isolated smaller cluster, $1P \rightarrow 1D$, is
24 contributing to the dipole moment also in this system, but many transitions inside the larger
25 cluster are participating in the excitation too.

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28 The shape of the induced density distribution for the largest peak is similar to the induced
29 density of the LSPR peak for the individual 138-electron cluster. Also the main contributions
30 to the dipole moment come from the same e-h transitions as in the single cluster. However,
31 there are also changes in density inside the smaller cluster, and the $1P \rightarrow 1D$ transitions have
32 some screening contribution. Here the density oscillations inside the bigger cluster are almost
33 symmetric, as can be seen from the contour plot of the induced density, but very asymmetric
34 inside the smaller cluster.

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37 The latter peak around 3 eV is clearly smaller in the heterodimer than the LSPR peak

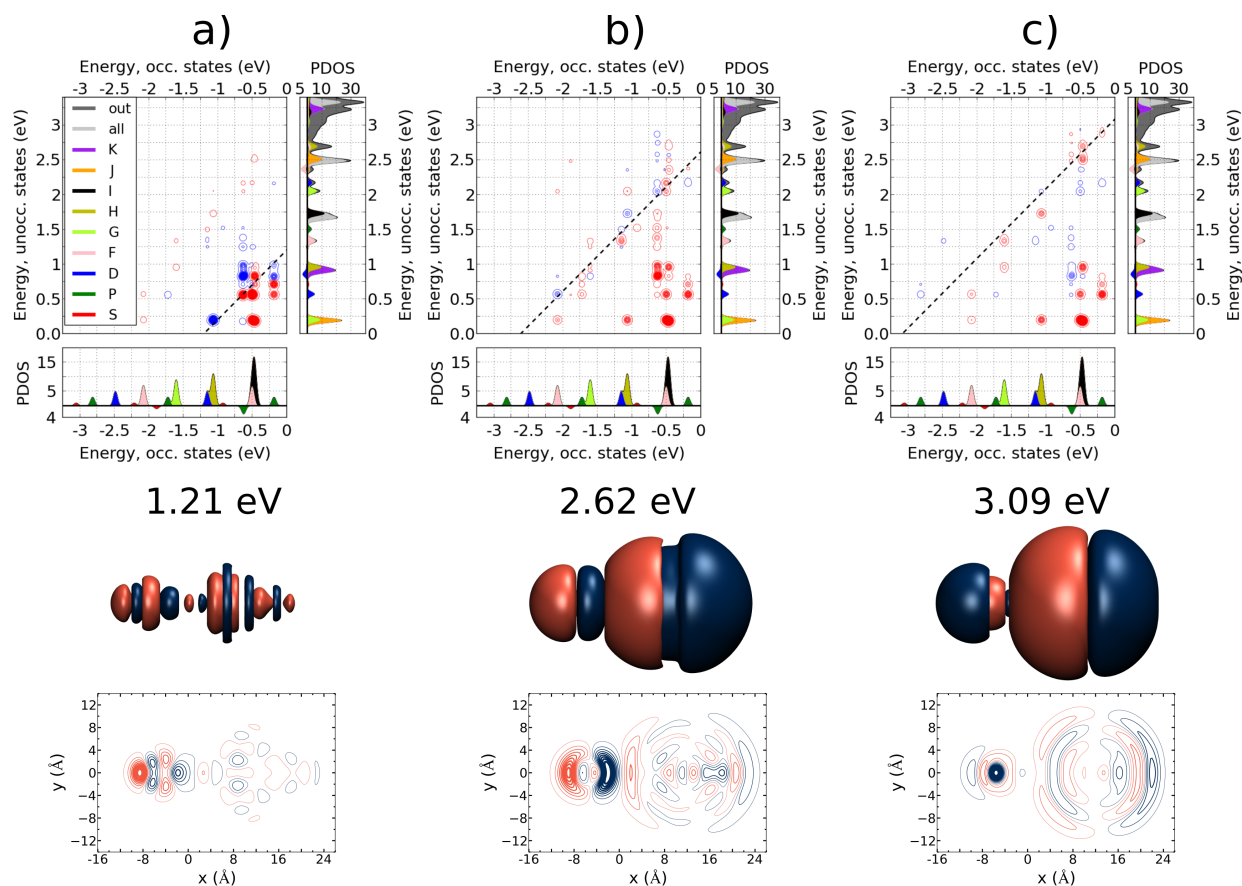


Figure 8: The DTCMs and induced densities for three different peaks for the heterodimer at energies indicated in the figure. Here, the PDOS analysis is plotted separately for the smaller cluster (down) and for the larger cluster (up). See further details in the caption of figure 3.

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3 for the individual 138-electron cluster, and the peak at 2.62 eV is larger than the LSPR peak
4 for the 8-electron cluster. A clear coupling is thus present between the clusters, and the
5 larger cluster 'borrows' some oscillation strength to the smaller one. Since the shape, the
6 energy of the largest peak, and the maximum absorption strength of the spectrum is changing
7 drastically from the single 138-electron cluster to the heterodimer, the larger cluster can also
8 be thought of as a sensor, as the presence of the smaller cluster can be detected from the
9 optical response of the system. The amplification of the main peak for the smaller cluster
10 also helps in this detection.

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12 The CTR values are larger for the excitations below 2 eV in energy. One of these ex-
13 citations, at 1.21 eV, is analyzed in figure 8 a). At this energy, there are no excitations in
14 the absorption spectra of the individual clusters, as can be seen from figure 7 b). For this
15 excitation, the biggest changes in the electron density are around the smaller cluster. The
16 induced density distribution is also quite asymmetric inside both of the clusters. The differ-
17 ence in the DTCM compared to the higher energy peaks is the presence of many screening
18 transitions (blue color), one of which is the 1P→1D transition of the smaller cluster.

35 Conclusions

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37 We have studied the coupling of plasmons and the emergence of the charge transfer plasmon
38 in dimers of metal clusters from the electronic perspective using the simple jellium model.
39 Here all the systems had cylindrical symmetry, but the real-space grid based GPAW code
40 enables the study of clusters and cluster assemblies of any kind of geometries. The high
41 scalability of the code made possible the lr-TDDFT calculations of the systems with the
42 larger dimers of 276 or 278 electrons, where the number of electron-hole pairs included in
43 the calculation was approximately 90 000.

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45 The Charge Transfer Ratio index, calculated post-processing from the induced densities,
46 makes it possible to indentify the excitations with the most charge transfer nature. The CTR
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3 spectra show that the excitations visible at small energies in the absorption spectra of dimers
4 with electron cloud overlap or linking all have charge transfer nature. These transitions are
5 not present in the spectrum of dimers with a large separation,
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9 For the smaller clusters, the effect of the linker on the optical spectra is quite drastic.
10 Here, the radius of the narrower linker is about $1/4$, and the wider linker $1/2$ of the radius
11 of the spheres, which changes the electron density profile and the energies of the orbitals
12 significantly from that of separate clusters with the same separation. For the larger clusters,
13 the ratio of the radius of the linker and the radius of the sphere is about 0.20 for the wider
14 linker. For these dimers, the shape of the spectra changes less with linking, but especially for
15 the wider linker, a significant amount of the oscillation strength is transferred to low-energy
16 charge transfer nature excitations. The wider linker, the radius of which corresponds to
17 the Wigner-Seitz radius used for the jellium density, also enables the emergence of a new
18 occupied K-S orbital.
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21 One heterodimer system with an 8-electron and 138-electron cluster was also studied. The
22 larger cluster is observed to borrow oscillation strength to the smaller, so that the LSPR
23 peak of the larger cluster is diminished while the LSPR peak for the smaller cluster gains
24 strength. Some excitations with charge transfer nature are also present at lower energies.
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28 Experimental systems such as cluster dimers linked with ligands, colloidal cluster assem-
29 blies and arrays of clusters on a substrate all require knowledge of the effect of the coupling
30 of the electronic excitations to the optical properties. The jellium linker used here can act
31 as a simplified model for one or several conducting linker molecules, or metal atom chain
32 between clusters, showing the trend for the absorption spectrum when conductive linker is
33 introduced. When modeling the complex ligand-protected clusters, the results of these sim-
34 ple jellium clusters can help to sort the different kind of excitations and to understand the
35 underlying principles of the electronic behavior.
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Notes

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Supporting Information Available

- The electron densities along the dimer axis for all the dimers, the 20 first electron orbitals for dimers **2–4**, a discussion about the choice of the LDA functional, and derivation of the CTR equation. This information is available free of charge at ...

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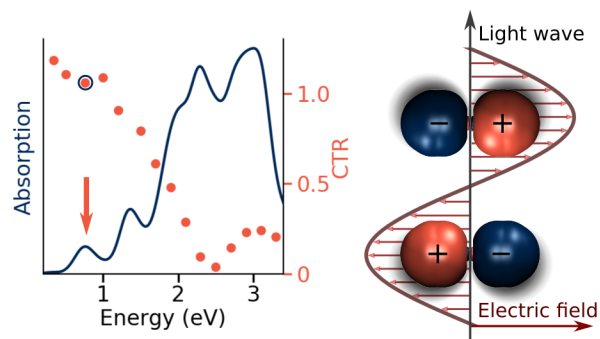
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