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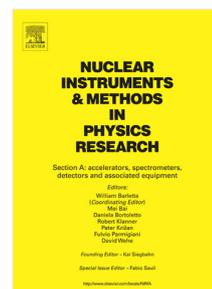
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Characterization and performance of the DTAS detector

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Abstract

DTAS is a segmented total absorption γ -ray spectrometer developed for the DESPEC experiment at FAIR. It is composed of up to eighteen NaI(Tl) crystals. In this work we study the performance of this detector with laboratory sources and also under real experimental conditions. We present a procedure to reconstruct offline the sum of the energy deposited in all the crystals of the spectrometer, which is complicated by the effect of NaI(Tl) light-yield non-proportionality. The use of a system to correct for time variations of the gain in individual detector modules, based on a light pulse generator, is demonstrated. We describe also an event-based method to evaluate the summing-pileup electronic distortion in segmented spectrometers. All of this allows a careful characterization of the detector with Monte Carlo simulations that is needed to calculate the response function for the analysis of total absorption γ -ray spectroscopy data. Special attention was paid to the interaction of neutrons with the spectrometer, since they are a source of contamination in studies of β -delayed neutron emitting nuclei.

Keywords: β decay, total absorption γ -ray spectrometer, exotic nuclei, NaI(Tl) detector, non-proportional scintillation light yield, Monte Carlo simulations

1. Introduction

Decay studies of exotic nuclear species at the focal plane of the FAIR-NUSTAR Super Fragment Separator in the DESPEC experiment [1] will provide information on the nuclear structure and the astrophysics impact of exotic nuclei. Far from stability, the Q_β values are very large, and the corre-

sponding increase in level density implies, on the one hand, the fragmentation of the β feeding into many levels populated in the decay and, on the other hand, the fragmentation of the γ intensity between many possible cascades. Total Absorption γ -Ray Spectroscopy (TAGS) has been shown to be an accurate tool to determine β -decay intensity distributions for such nuclei far from the valley of β stability. This technique avoids the so-called *Pandemonium* effect [2], related to the relatively poor efficiency of HPGe detectors. Instead of detecting individual γ rays as in high-resolution exper-

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20 iments with HPGe detectors, TAGS aims to detect the full β -delayed electromagnetic cascade. This is achieved with spectrometers made of large scintillator crystals covering a solid angle of $\sim 4\pi$. To extract the β -intensity distributions a de-convolution procedure is applied to the measured energy spectrum using the detector response, as will be explained later.

For this reason, a new spectrometer has been designed and constructed for the DESPEC experiment [3]. The Decay Total Absorption γ -ray Spectrometer (DTAS) is a segmented detector that consists of a maximum of eighteen NaI(Tl) crystals with dimensions 150 mm \times 150 mm \times 250 mm [3]. The choice of the material, as well as the geometry, comes from a careful study where several goals and constraints were taken into account: the efficiency, the amount of dead material, the energy and time resolution, the neutron sensitivity, the sensitivity limit for high-lying β -intensity, the coupling to ancillary detectors, and the cost [3]. The advantage of the segmentation in this case is threefold: the possibility to extract information from the module-multiplicity spectra, as will be explained later, the possibility of using the individual modules as single γ detectors, and the mechanical flexibility of the set-up. In fact, we consider two main configurations for DTAS: a sixteen-module configuration designed for experiments at fragmentation facilities, as proposed in [3], and an eighteen-module configuration for experiments at ISOL-type facilities. Both configurations without shielding can be seen in Fig. 1. In the eighteen-module configuration side holes can be made by moving away the modules of the horizontal central plane, thus allowing access from both sides of the detector, as shown in Fig. 1 bottom. In this way DTAS can be combined with ancillary detectors and it is possible to position a beam pipe in the centre of the spectrometer. This configuration has recently been commissioned at IGISOL [4], with holes of 10 cm used to place a HPGe detector from one side and the beam pipe with a β detector from the other side. The two central modules were separated by 16 cm instead of 10 cm in order to lower their counting rate, so that it was comparable to the external modules. The configuration foreseen for FAIR [3], with sixteen modules, will be coupled to the Advanced Implantation Detector Array (AIDA) [5]. In order to place AIDA in the center of DTAS, the two central modules in the eighteen-module configuration are removed and the two modules above the central hole are supported

by a specially designed aluminium frame with external dimensions identical to a module, as shown in Fig. 1 upper panel.

The shielding surrounding DTAS is composed of stainless steel sheets, lead bricks and aluminium, and it served to reduce the background counting rate by one order-of-magnitude in the measurements of this work. The allocation of individual modules to positions in the arrangement was done according to their resolutions, ranging from 7% to 9% at 661.7 keV, so that the positions associated with the lowest counting rates (the eight corners of the assembly shown in Fig. 1) were occupied by the modules with the poorest resolution.



Figure 1: DTAS detector in the sixteen-module configuration (top) and in the eighteen-module configuration (bottom) without radiation shielding.

The outline of the article is the following: in section 2 we will describe the procedure to reconstruct the full energy deposited in the detector from the signals of the individual modules. In section 3 a method to evaluate the summing-pileup contamination will be explained, and its validation with calibration sources will be discussed. Finally, the Monte Carlo (MC) response function of the detector will be described in section 4, and the repro-

duction of several calibration sources and the neutron contamination coming from β -delayed neutron emitters will be discussed.

2. Total energy reconstruction: hardware sum and software sum

In this section we will describe the electronic chain employed to process the signals from the individual modules of DTAS, and the procedure to reconstruct the total energy deposited in the detector. In particular, two methods to calculate the total energy sum will be discussed: the hardware sum and the software sum.

2.1. Signal processing

In order to analyse data from DTAS we have to reconstruct accurately, for each event, the energy deposited in the full spectrometer and its module-multiplicity, M_m (number of modules that fire above the threshold, often known as fold). The full energy released in the spectrometer is obtained by summing the energy deposited in the individual modules, either electronically or via software. The electronic chain to process the signals from the modules was designed with this idea in mind, and it is represented in Fig. 2.

We use Mesytec MSI-8p preamplifiers [6] for both anode and dynode signals from the photomultiplier tubes (PMTs). MSI-8p are custom adapted units with fixed gain, preamplifier constants optimized for these PMTs and without timing filter amplifier. After the preamplifier, dynode signals are split into two branches; one branch is sent to a CAEN N625 Quad Linear FAN-in FAN-out [7], and the other to Mesytec MSCF-16 shapers. The N625 module acts as an analog signal adder and one of the outgoing signals is processed in an ORTEC 671 amplifier [8] to produce the sum energy signal (hardware sum) sent to the analog to digital converter (ADC), a CAEN V785 module, of the data acquisition system (DACQ). Another output from the N625 module is used to construct a common stop signal sent to a time to digital converter (TDC), CAEN V775, using an ORTEC 474 Timing Filter Amplifier and an ORTEC 584 Constant Fraction Discriminator. The MSCF-16 shapers provide individual energy and timing output signals that are sent to the individual channels of the ADC and TDC modules respectively. The anode signals after the preamplifier are sent to sampling digitizers of a second digi-

tal DACQ, running in self-triggered mode, which is not discussed in this publication.

In order to carry out the hardware sum properly we need to match the gains of the different PMTs by adjusting the high voltage (HV) applied to them, so that the signals of individual modules are aligned. Note that aligned here means having the same amplitude for the same energy deposited.

The software sum is reconstructed offline from the individual signals processed with the MSCF-16 shapers. In the following subsections we will show a method of correcting possible changes in the gain of the modules, as well as the way to perform properly the alignment and determine the software sum of these signals.

2.2. Gain correction system

A system to correct changes in the gain of individual modules has been developed, as planned in [3]. These changes may be due to temperature variations [9], drift of the PMT current and fluctuations in the HV supply. In this system the gain of each module is monitored checking the position of the peak produced by a pulsed light source. An additional external reference detector, with a weak ^{137}Cs radioactive source, is used to monitor the stability of the light pulse generator. The following elements are employed in this system:

- An external reference well-type NaI(Tl) detector of 3" diameter \times 3" length manufactured by Saint Gobain [10]. The well has 15 mm diameter and 40 mm depth. The crystal is mounted on a 3" diameter ETI 9305 PMT as shown in Fig. 3.
- A 490 nm light pulse generator model 6010 from BNC [11]. The generator is triggered with an external 100 Hz clock signal.
- A 2 m long bundle of borosilicate glass fibres split into 20 bundles of 2 mm diameter, manufactured by FiberTech Optica [12]. The fibres are terminated with SMA type connectors.
- A weak ^{137}Cs source of ~ 300 Bq.

The fibre bundle splitter is used to distribute the light pulse from the generator to the reference detector and to each of the eighteen modules. The ^{137}Cs source is placed inside the well of the reference detector. The reference detector is surrounded by lead shielding and is placed close to DTAS. Since

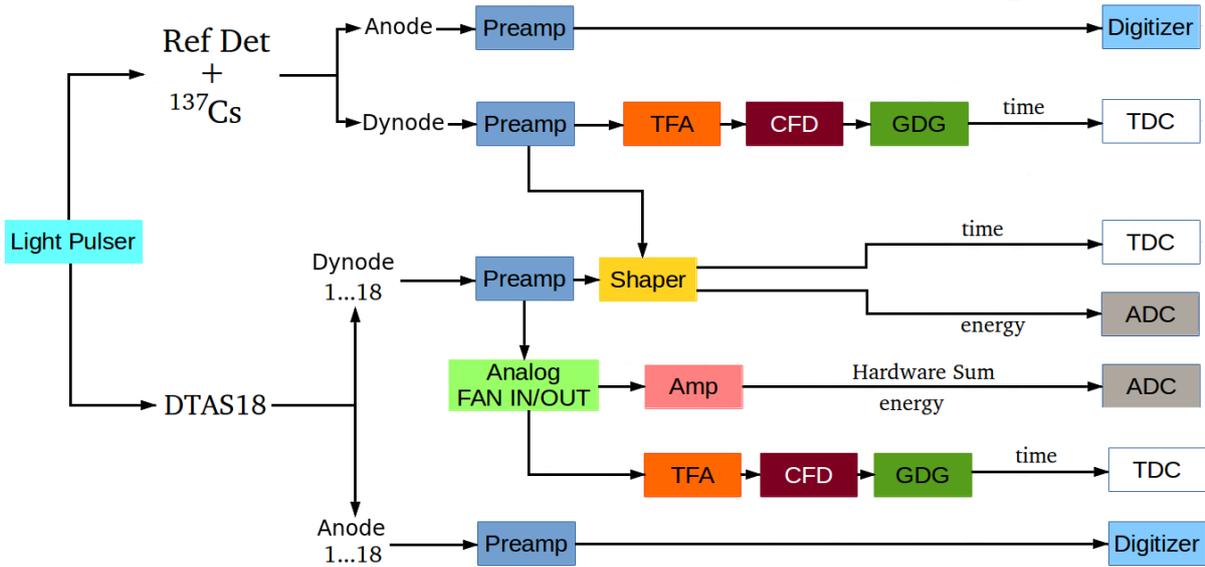


Figure 2: Schematic diagram of the electronic chain. The labels correspond to: Preamplifier (Preamp), Spectroscopic Amplifier (Amp), Timing Filter Amplifier (TFA), Constant Fraction Discriminator (CFD), Gate/Delay Generator (GDG), Time to Digital Converter (TDC), Analog to Digital Converter (ADC).

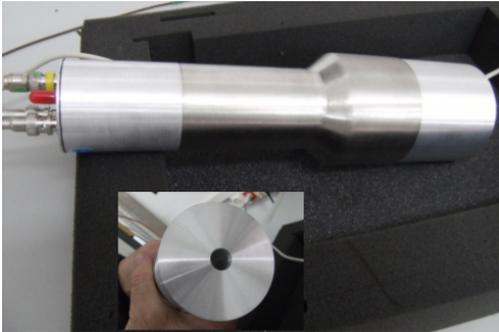


Figure 3: NaI(Tl) reference well type detector. Inset: a view of the front face with the hole where the weak ^{137}Cs source is placed.

190 both the reference detector and DTAS have shielding, this weak source does not affect the DTAS measurements. The position of the 661.7 keV peak in the well detector provides a reference for possible changes in the gain of this detector. Comparing the position of the light pulser peak with this peak we can determine if there are variations of the intensity of the light source. With this information we can separate in each module variations in the gain from variations in the light source intensity. The gain correction is calculated for short time inter-

vals, and the procedure will be detailed in the next subsection. An example of the spectra of the reference detector and one individual module of DTAS showing the light pulser peaks can be seen in Fig. 4.

In order not to disturb the measured individual spectra, the peak due to the light pulser has to be located beyond the energy region of interest, see Fig. 6 as an example. When choosing an optical fibre bundle for each module, we took into account that each of the 20 bundles does not transport the same amount of light, and the individual modules do not convert the same amount of incident light into the same signal amplitude in the PMT. For both reasons, in order to minimize the difference in position of light pulser peaks between modules we assigned the bundles that transport more light with the worst modules in terms of light conversion.

Apart from applying the gain correction offline, the gain correction system could also be used for maintaining the alignment of the signals of the modules during the measurement by applying periodic HV corrections to the PMTs. This requires information about the dependence of the gain with the HV for each module. Although we have tested this online correction method, it was not used in the actual measurements.

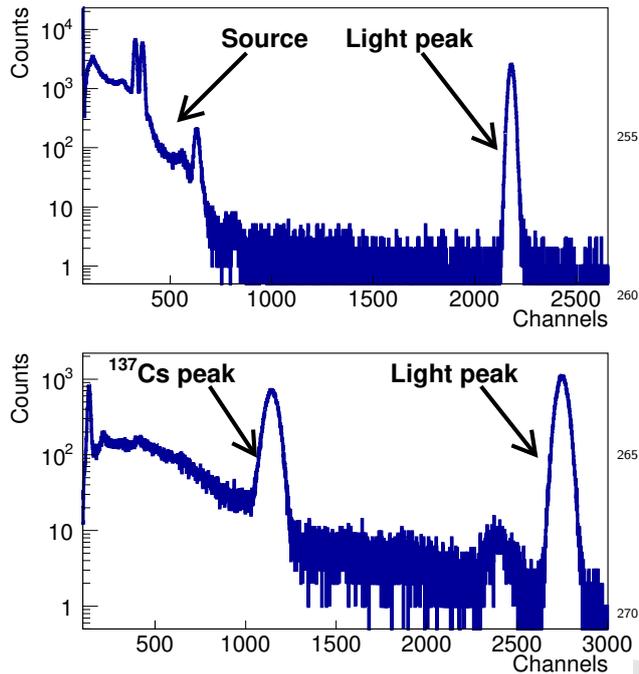


Figure 4: Individual DTAS detector spectrum with the light pulser peak in a ^{60}Co measurement (top). Reference detector spectrum with the 661.7 keV peak from the weak ^{137}Cs source, and the light pulser peak (bottom).

2.3. Software sum

Just as in the case of the hardware sum, before performing the software sum the amplitude of the signals stored for each event must be properly aligned. Although signal amplitudes were gain-matched before the FAN-in FAN-out for the hardware sum, and even though the gains of the shapers are set to a common value, the stored amplitude information needs to be corrected due to slight variations in gain and offset of the individual electronic channels.

The first idea for making this alignment was to convert signal amplitude (proportional to light collected) into energy for each of the individual channels. This conversion between light collected and deposited energy is what we will call energy calibration. A solution like this has been successfully adopted for a 12-fold segmented BaF_2 spectrometer in previous works [13, 14, 15, 16]. Nevertheless, we soon realized that it can not be applied in the case of a segmented detector made of $\text{NaI}(\text{Tl})$ because of the non-proportionality of the light yield in this material [17, 18]. The reason is related to what was

pointed out in [19], explaining the shift of the position of full energy peaks due to γ -ray cascades with respect to single γ -ray peaks of the same energy. For every primary electron created by γ -ray interactions there is a shift of about 10 keV in the apparent energy. Since γ -rays of several hundreds of keV to a few MeV typically require of the order of three interactions (two Compton, one photoelectric) to deposit the full energy this explains why for a γ -cascade of two γ -rays (γ -multiplicity, $M_\gamma=2$) the shift is approximately 30 keV, while for $M_\gamma=3$ it is 60 keV and so on. In the case of a segmented detector the situation is more complicated, and the shift depends not only on the γ -multiplicity, M_γ , but also on the number of modules where the energy is deposited, M_m , which determines the distribution of the number of primary electrons in each module. Taking into account the different ways that $3 \times M_\gamma$ electrons can be distributed in M_m modules one can determine that the apparent energy shifts follow approximately the numbers in Table 1. The first row in the table corresponds to the behaviour of a single $\text{NaI}(\text{Tl})$ crystal spectrometer like LUCRECIA at ISOLDE [20] or the LBNL spectrometer used at GSI [21].

M_γ	1	2	3	4
M_m				
1	0	+30	+60	+90
2	-30	0	+30	+60
3	-60	-30	0	+30
4	-90	-60	-30	0

Table 1: Shift in keV of the sum peak position due to the non-proportionality of the light yield in a segmented $\text{NaI}(\text{Tl})$ spectrometer when the individual modules are calibrated in energy before the software sum.

For a single crystal spectrometer the non-proportionality is not a problem as far as this effect is included in the MC simulations in the way detailed in [19]. Likewise, it does not present any problem for the hardware sum of a segmented $\text{NaI}(\text{Tl})$ spectrometer, as long as the PMTs are gain-matched. However, the consequence of applying an energy calibration to individual modules before summing, is that the resolution of the sum peaks is worsened due to the displacement of the different M_m module-multiplicities contributing to the sum. The non-proportionality of the light yield in $\text{NaI}(\text{Tl})$ is known to have an important contribution to the resolution of single crystal

290 detectors [22, 23], and by extension of the hardware sum of multi-crystal detectors. Here we are referring to an additional effect that must be taken into account when reconstructing the software sum for multi-crystal detectors. In Fig. 5 these shifts are shown for a measurement of ^{22}Na ($M_\gamma=3$) and for the corresponding MC simulation of this source that includes the non-proportionality of the light yield as in [19]. In both cases an energy calibration has been applied to all the individual modules before summing. The vertical black line corresponds to 2296.5 keV, the sum of the energies of the three γ -rays involved: 511 keV, 511 keV and 1274.5 keV. The sum peaks of the different M_m module-multiplicities are not aligned, showing a displacement in agreement with Table 1. Only $M_m=3$ is aligned with the nominal sum, since it corresponds to a 0 keV shift in Table 1, with three γ -rays detected in three crystals. Note that the experimental spectra are not background subtracted, whereas the MC is only widened by the light function from [19], without taking into account additional contributions to the resolution.

In summary, in order to maintain the resolution, we have to align the stored amplitudes of the individual modules, thus reproducing with our software sum the same behaviour of the hardware sum, and, equivalently, of a single crystal detector. In addition there are other effects that may worsen the resolution, like changes in the gain of the PMTs and the electronic chain. The correction we have to apply to counteract these effects is applied sequentially on a reduced number of events. The number of events should be sufficiently large to determine the peak positions accurately and sufficiently small to limit the effect of gain variations during the acquisition time. We have verified that 1 million events (that corresponds to approximately 4 minutes for a typical counting rate of 4-5 kHz in DTAS), fulfill this condition.

330 The stored amplitude is represented by the bin number in the histogram accumulated by each ADC channel (detector module) and the first step is to determine the offset and the gain. To determine the ADC offset for each channel we use the position of the peak due to the electronic noise. The gain is obtained from the position of the two peaks from a calibration run with a ^{22}Na source (511 keV and 1274.5 keV). With the offset and gain so obtained the alignment of the first million events is performed choosing one arbitrary module as a reference.

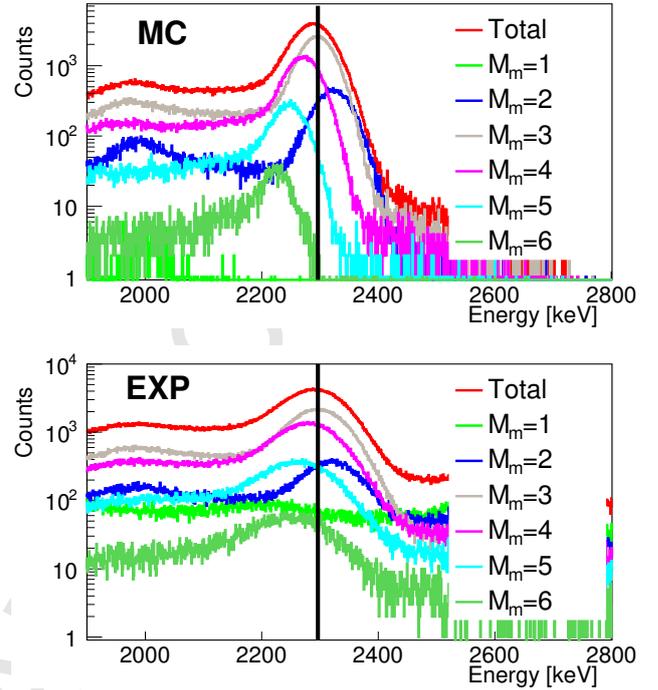


Figure 5: ^{22}Na software sum of the light produced in the individual detectors calibrated in energy. Both the MC (top) and the experimental measurement (bottom) show the shifts from Table 1. The vertical black line corresponds to the sum of the energies of the three γ -rays involved. (For interpretation of the references to color in this figure caption, the reader is referred to the web version of this paper.)

After the alignment, the reference values of the parameters involved in the gain correction procedure are determined. The ADC offset is represented by a_j , with $j = 0 \dots 18$, with $j = 0$ being the well detector and $j = 1 \dots 18$ the DTAS modules. The reference position of the light pulser peak for each module, L_j , is obtained by peak fitting. Analogously, the ^{137}Cs peak and the light pulser peak reference positions for the well detector, P_0 and L_0 respectively, are determined.

The next group of one million events is then processed. We define L'_j as the new light pulser peak position of module j , b_j as the gain change factor of module j , and C as the change factor in the light source intensity. The procedure described below is followed in order to calculate the gain corrections and sum the amplitudes of all modules stored in each event:

- The new position of the ^{137}Cs peak for the well detector, P'_0 , is determined, as well as the po-

sition of the light pulser peak L'_0 .

- The change in the gain of the PMT of the well detector is calculated:

$$b_0 = \frac{P_0 - a_0}{P'_0 - a_0} \quad (1)$$

- The change in the light produced by the light pulse generator, C , is calculated:

$$C = \frac{L_0 - a_0}{b_0(L'_0 - a_0)} \quad (2)$$

- L'_j is determined for each of the DTAS modules, and with this value the gain change factor is calculated taking into account the change of intensity of the light source from Equation 2:

$$b_j = \frac{L_j - a_j}{C(L'_j - a_j)} \quad (3)$$

Once the parameters b_j are determined we re-process the same group of events applying the gain correction factor in order to align the amplitudes of all modules to the first group of events used as a reference. As a result of applying this procedure, the software sum can be performed properly, as can be seen in Figure 6 for a ^{60}Co source. At this point an energy calibration can be applied (a conversion between light collected and energy) by using single peaks ($M_\gamma=1$), as in the case of the hardware sum.

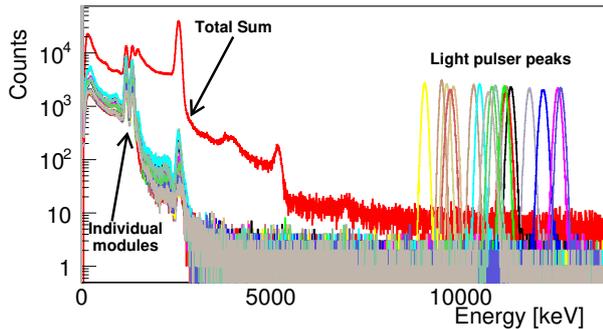


Figure 6: Software sum spectrum of a ^{60}Co source measurement and individual module spectra showing the alignment. The peaks in individual modules above 8 MeV are due to the light pulser.

The software sum reconstructed in this way exhibits the same behaviour as the hardware sum in

terms of the non-proportionality of the light yield, as seen in Fig. 7 for two calibration sources. In both cases the segmented detector behaves as a single crystal detector in terms of the position of the sum peak. The main differences are related to a slightly better resolution in the software sum with respect to the hardware sum due to the gain corrections, and a different shape in the pileup region that will be commented on next section. The final resolution achieved for the software sum of the eighteen modules is 8.6% at 661.7 keV, worse than the 7.7% value obtained for the single prototype module in [3].

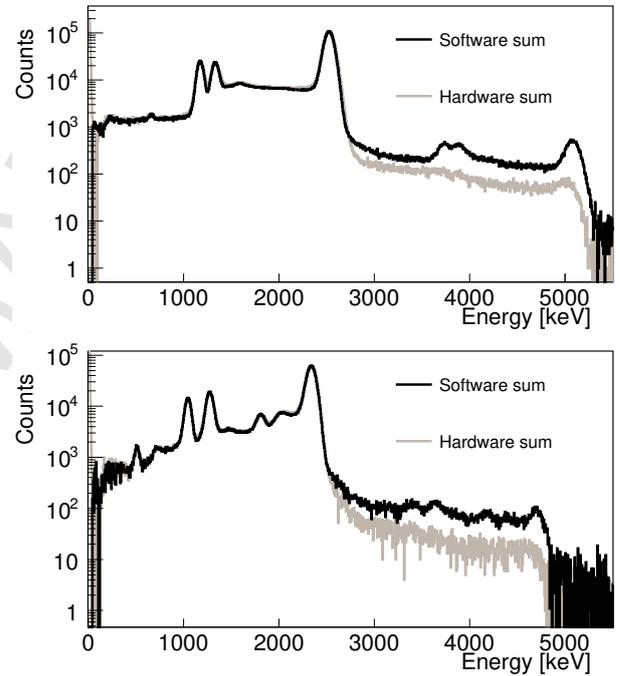


Figure 7: Comparison of the hardware sum (grey) and the software sum (black) in DTAS for a ^{60}Co source (top) and a ^{22}Na source (bottom).

In order to ensure that the treatment of the non-proportionality is correct, we can check the spectra of the M_m module-multiplicities after this process. In Fig. 8, we show the good alignment of the different M_m module-multiplicity spectra achieved with this method for a ^{22}Na source ($M_\gamma=3$) and a ^{60}Co source ($M_\gamma=2$), in contrast with the results shown in Fig. 5. The vertical black lines correspond to the sum peak positions calculated using the shift associated with single crystals according to [19] (first row of Table 1): 2296.5 keV+60 keV for the ^{22}Na

410 source, and 2505.7 keV+30 keV for the case of ^{60}Co .

We should point out that the procedure followed here solves the misalignment problems between simulation and experiment encountered in the calibration of the Modular Total Absorption Spectrometer (MTAS) [24], as will be shown in Section 4.

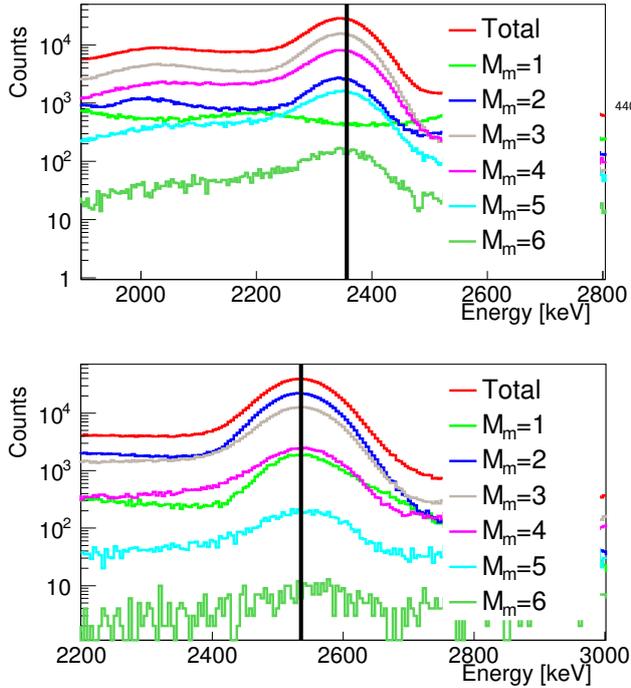


Figure 8: Multiplicity spectra of a ^{22}Na source (top) and a ^{60}Co source (bottom). Vertical black lines show the corresponding energy of the sum peak for a single crystal detector. (For interpretation of the references to color in this figure caption, the reader is referred to the web version of this paper.)

3. Summing-pileup calculation

An important source of distortion in the measured spectra is the random superposition of electronic signals within the time length of the ADC gate, due to the statistical nature of the decays. This superposition affects the pulse shape of a single detector leading to so-called pulse-pileup [25]. This applies to individual crystals as well as to the hardware sum of a multi-crystal detector. In the software sum of a segmented detector the distortion due to the superposition of events within the ADC gate takes an additional form; namely,

the sum of the signals detected in different modules and corresponding to different decays that are, however, stored in the same event. Thus to calculate the distortion of the final spectrum both processes must be taken into account, the pulse-pileup (that we will simply call pileup) and the random summing (that will simply be called summing, but has to be distinguished from the traditional use of this term in spectroscopy). We have developed a method to treat the distortion of spectra due to summing-pileup that was already used in previous works [13, 14, 15, 16], and will be detailed here. The quality of the reproduction of this type of spectrum distortion for a set of calibration sources, listed in Table 2, has been studied.

Source	Rate [kHz]
^{22}Na	4 and 5
^{60}Co	7
^{24}Na	14
^{137}Cs	19
^{152}Eu - ^{133}Ba	44
Background	3

Table 2: Set of sources used in the study of the summing-pileup, and their counting rates in DTAS, together with the environmental background rate. All these measurements were performed with shielding.

3.1. Procedure

The evaluation of the summing-pileup contamination is based on the event structure of the experimental data, and on the true electronic pulse shape of the individual modules after the MSCF-16 shapers. For the first order summing-pileup calculation, two arbitrary random events are read from the list-mode event file and the time difference between them is sampled randomly within the ADC gate length. If an individual detector has fired in both events, two pulses with their corresponding amplitude are summed, and the maximum within our effective ADC gate $\tau = 5.6 \mu\text{s}$ (the ADC gate minus the peaking time of the individual signals) is taken, according to [25]. If, on the contrary, the individual detector has only fired in one event, it contributes to the summing evaluation. The total summing-pileup is the sum of all contributions, as depicted in Fig. 9. This procedure assumes implicitly that the distortion of measured events is small. This approximation is valid if the rate is below 10 kHz. For higher rates a similar procedure,

but based on simulated data, is used as explained in the next subsection.

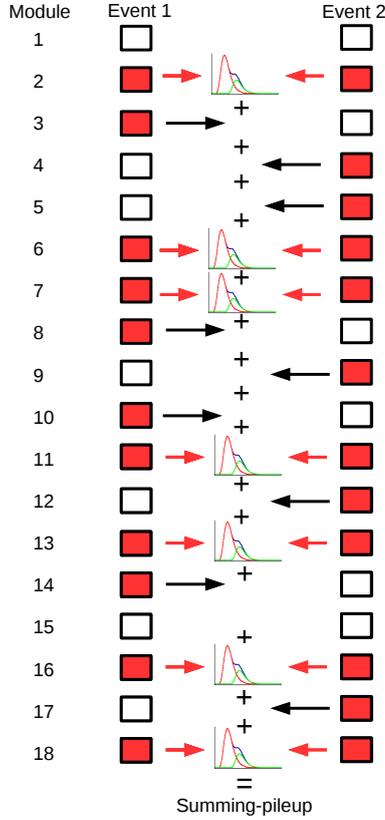


Figure 9: First order summing-pileup scheme. The red squares represent a signal over the threshold stored in the ADC for a module in a given event. When the same module fires in the two events used for the summing-pileup reconstruction, it is processed as a pulse pileup. In any other case, the signals are added giving rise to the summing contribution.

It is worth mentioning that, in our measurements, the majority of the summing-pileup events are coming from the summing contribution, as shown in Fig. 10 for a ^{60}Co source, where the total summing-pileup contains around 87% events with only summing, $\sim 1\%$ with only pulse pileup, and $\sim 12\%$ where both contribute.

The normalization factors needed to compare the calculated summing-pileup contribution with experimental spectra are obtained from the theoretical expression of Eq. 4, which is based on the expression used for pileup order n in [25], but adapted to a segmented detector. Here α_i are the individual

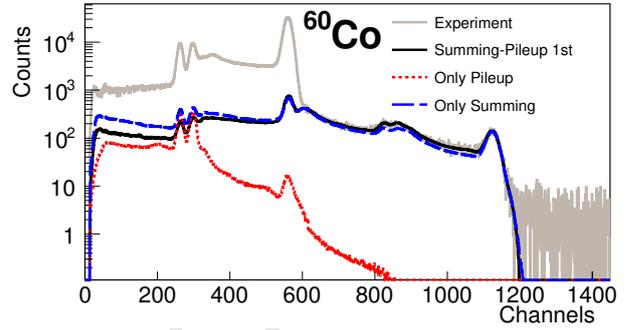


Figure 10: First order summing-pileup for the ^{60}Co source. The contributions of the pulse pileup (dotted red) and the random summing (dashed blue) are separated to show the region where they are affecting the spectrum. (For interpretation of the references to color in this figure caption, the reader is referred to the web version of this paper.)

counting rates of the 18 crystals and τ is the length of the effective ADC gate.

$$N_{theo}^n = \sum_{i=1}^{18} e^{-\alpha_i \tau} (1 - e^{-\alpha_i \tau})^n \quad (4)$$

When the counting rate is high (approximately above 10 kHz), second order summing-pileup contributions must be evaluated. This is the case for the ^{24}Na source and the ^{137}Cs source, whereas for the ^{152}Eu - ^{133}Ba source even the third order contribution was needed in order to reproduce the measured spectrum. The procedure in those cases represents just an extension of the method already described. In the second order contribution, for example, three events are taken each time, instead of two. The quality of the reproduction of this contamination in the set of calibration sources of Table 2 can be seen in Figs. 10 and 11.

3.2. Calculation with MC simulated data

We encountered difficulties to reproduce the shape of the summing-pileup contribution for the sources with high counting rates, ^{137}Cs and ^{152}Eu - ^{133}Ba . This is because in these sources, a large fraction of the events that are used for the calculation are already distorted by summing-pileup. However we are unable to distinguish if a measured event is distorted or not. The way out of this dilemma is to use realistic simulated data. For this purpose, we simulated the decay of sources with Geant4 [26] as will be explained in Section 4 and we stored the

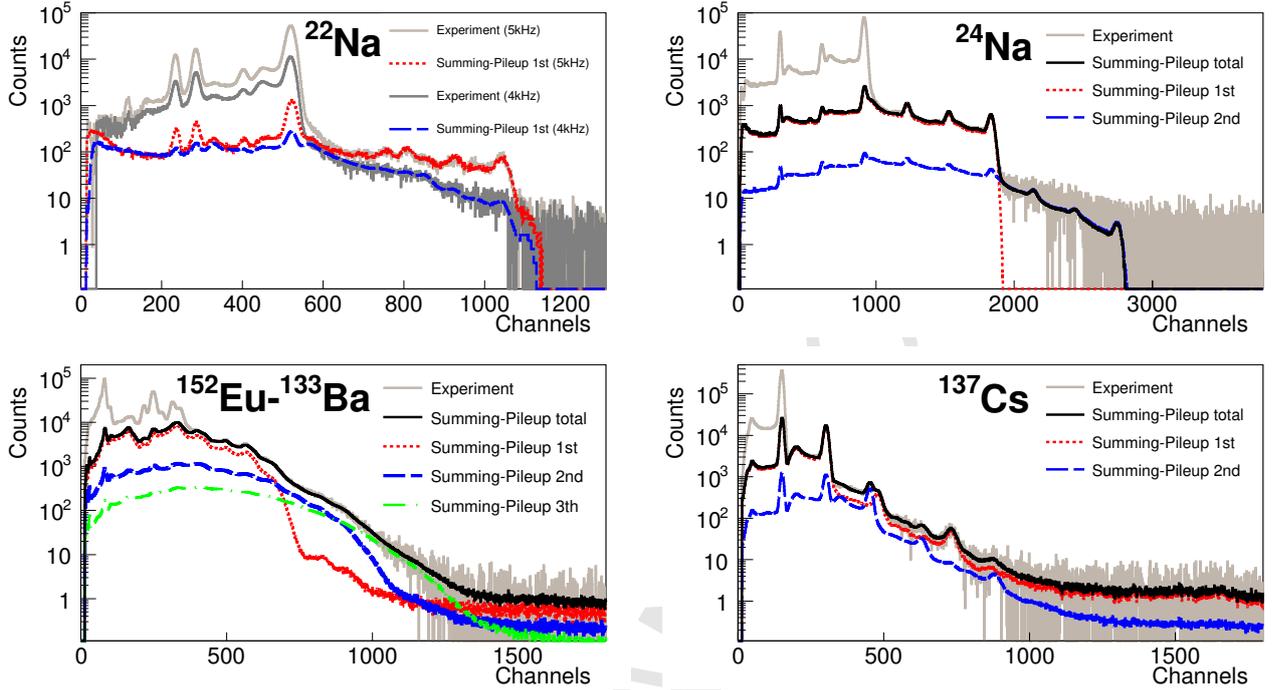


Figure 11: Summing-pileup contributions for a set of calibration sources. Two sources of ²²Na with different counting rates are compared, normalized by the time of the measurement. For the ¹³⁷Cs and the ¹⁵²Eu-¹³³Ba sources the calculated summing-pileup contribution is based on simulated data (see text for details).

relevant information for modules fired in each decay event in a format similar to experimental data. In the simulation the deposited energy is converted into light and the experimental resolution is introduced. The proper light to experimental amplitude calibration is then applied.

For the calculation of the summing-pileup contribution, the same procedure explained in the previous section is used with small modifications. In particular, we have to supplement the simulated data file with a real background data file. We assume that the summing-pileup distortion in background events is small. Consequently only the source-source and source-background summing-pileup contributions are calculated. For this reason the first event is always chosen from the pure source (MC simulation) and the second is taken either from the source, or from the background experimental file. The proportion between source and background for the second event is roughly fixed by the counting rates, and it is a parameter that can be adjusted by looking at the resulting spectra.

The use of MC data files to reconstruct the summing-pileup contribution has proven to be suc-

cessful, and the summing-pileup of ¹³⁷Cs and ¹⁵²Eu-¹³³Ba shown in Fig. 11 has been reconstructed by using MC simulated data, instead of using the experimental source file, thus validating this method for high counting rates.

4. Validation of MC simulations

The aim of the TAGS technique is to determine a β -intensity distribution from an experimental measured spectrum by solving the inverse problem represented by:

$$d_i = \sum_j R_{ij} f_j + c_i \quad (5)$$

where d_i is the number of counts in channel i of the experimental spectrum, f_j is the number of events that feed level j in the daughter nucleus, and R_{ij} is the response function of the detector that represents the probability that feeding to the level j gives a count in channel i of the spectrum. The sum of all contaminants in channel i is represented by c_i .

In order to perform this de-convolution and obtain the feeding distribution, a method was developed by the group of Valencia [27] which has been successfully applied to a large number of cases. An essential ingredient of this process is the determination of the response function, that is unique to each detector and to each decay scheme, and has to be calculated by means of MC codes. For this reason, a mandatory step in the characterization of the detector is to validate the MC simulation. This is achieved by comparison of simulations with measured calibration sources to verify that the best possible agreement is reached. The package Geant4 [26] has been used for this purpose, and the geometry of DTAS has been included in great detail, as shown in Fig. 12. In addition the relevant physics processes involved in particle detection have been incorporated. In particular, the non-proportional light yield in NaI(Tl) has been taken into account according to the parametrization and the procedure detailed in [19]. In the next subsection we present the results of such a comparison.

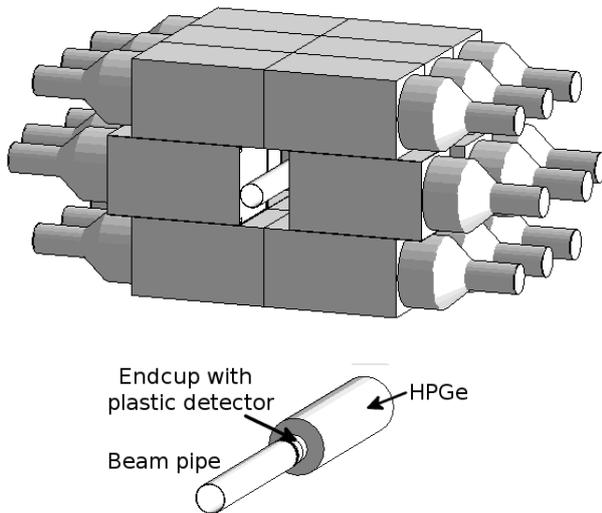


Figure 12: Geometry of the set-up implemented in the MC. The DTAS detector in the eighteen-module configuration (top) is represented. Inside we considered the beam pipe, a plastic β detector and a HPGe detector (bottom).

The efficiency of DTAS for γ -rays and β particles can be obtained from the MC simulations once the geometry and the physics have been validated. In Fig. 13 we show the calculated efficiencies for the complete set-up used in the commissioning with

radioactive beams performed at IGISOL with the eighteen-module configuration [4]. The beam pipe, a 3 mm thick plastic scintillator β detector with its PMT, and a HPGe detector, all inserted in DTAS, are included in the geometry. The efficiencies shown in Fig. 13 are calculated after applying a realistic energy threshold of 80 keV to the individual modules before reconstructing the sum energy. The total efficiency is above 80% over the whole range, while the peak efficiency at 1 MeV is 65%. When we consider the individual modules in the array, the peak γ efficiency at 1 MeV is 50%.

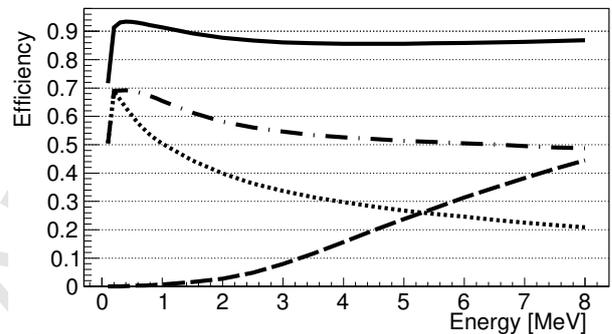


Figure 13: Simulated efficiencies as a function of the energy of the DTAS detector in the eighteen-module configuration. Total γ efficiency (solid line) and peak γ efficiency (dashed-dotted line) are shown. The peak γ efficiency for the array of individual modules is shown as a dotted line. The efficiency of DTAS for β particles as a function of the end-point energy is shown as a dashed line. Ancillary detectors are included in the simulation (see text for details).

The total and peak efficiencies are limited by the solid angle covered and the amount of both sensitive and dead material. The solid angle covered can be increased and the dead material decreased if we remove the HPGe detector and we close the gap between the modules that it occupies. In this case the efficiencies will increase to $\varepsilon_{\gamma}^T = 93\%$ and $\varepsilon_{\gamma}^P = 68\%$ respectively at 1 MeV.

In comparison with the Lucrecia spectrometer at ISOLDE [20] and the LBNL spectrometer at GSI [21], both single crystal spectrometers, the efficiency of the present eighteen module configuration is similar. When compared with the recently built MTAS spectrometer [28] our peak efficiency is 9% smaller at 1 MeV and 25% smaller at 3 MeV. This difference is a consequence of the much larger NaI(Tl) volume in MTAS, which is a factor of 2.5 larger than DTAS. It should be noted that the smaller efficiency of DTAS does not affect its perfor-

mance as a total absorption spectrometer. A nice example is provided by the decay of ^{137}I which has been measured by DTAS [4] (see also subsection 4.2) and MTAS [29]. Figure 4 in [29] compares the spectrum measured with the full MTAS with the spectrum measured with the sub-detector consisting of the 7 most central modules. This sub-detector is equivalent to DTAS in volume (about 100 litres) and efficiency. As can be observed the differences are minimal except for the contamination induced by the interaction of delayed neutrons emitted in the decay, which is much larger in the full MTAS. The reason why the much larger volume brings a seemingly small effect is to be found in the complex de-excitation pattern, with relatively large cascade γ -multiplicities, and the effect of β penetration which tends to wash out the increase in single γ -ray peak efficiencies. Such a consideration was taken into account during the design of DTAS when choosing the detector size [3].

4.1. Reproduction of the calibration sources

In this subsection we compare the results of the MC simulations with measurements for the calibration sources in Table 2. In the comparison, the different sources of contamination in the measured spectra are taken into account. The environmental background is subtracted from the measured spectra. However we choose to show explicitly the summing-pileup contribution, calculated as described in the previous section, adding it to the MC simulated spectra for the comparison. We use the DECAYGEN event generator [30] to generate the primary particles in the MC simulations. As can be observed in Fig. 14, we obtain an excellent reproduction of the experimental spectra.

One of the key features of DTAS is its segmentation. This allows one to obtain much richer information, provided by the energy spectra of the individual modules and more importantly by the sum energy spectra gated with different conditions on the number of modules that fired (M_m). These additional spectra are sensitive to the details of the de-excitation cascades (energies and module-multiplicities M_γ).

In the case of laboratory sources with known decay schemes the M_m module-multiplicity information provides a more stringent test of the accuracy of the MC simulation, both of the geometry and the physical processes included. As can be seen in Fig. 15 for the ^{22}Na source an excellent agreement is obtained proving that we have the MC sim-

ulations under good control. It should be noted that all calculated spectra shown in Fig. 15 are obtained simultaneously using the same energy calibration, a common normalization factor and the same summing-pileup calculation.

4.2. Reproduction of neutron interactions

The emission of β -delayed neutrons in the decay of exotic neutron-rich nuclei is a source of background for total absorption spectrometers like DTAS. Neutrons interact with detector materials producing γ -rays, either in an inelastic reaction or after neutron capture. When detected these γ -rays are indistinguishable from β delayed γ -rays. The ability to reproduce this type of background correctly with MC simulations is fundamental for the analysis of TAGS spectra from β -delayed neutron emitters. The issues related to the simulation of neutron interactions in inorganic scintillators, and in particular the use of the Geant4 simulation tool, have been discussed before [31, 3]. One important item is the quality of the information in nuclear data bases concerning reaction cross sections for all the materials encountered by the neutrons. In the present simulations we used the library ENDF-VII.0, which gave good results before. The data from this library were converted into the G4NDL data format [32]. Another important item is the description of the γ -ray de-excitations of excited states resulting from neutron interactions. We replace the standard capture cascade generator of Geant4, which is rather schematic, with a generator that uses the statistical model to describe realistically the multiplicity and energy distribution of the cascades. In the case of inelastic scattering we use the standard Geant4 PhotonEvaporation data base which relies on evaluated spectroscopic data [33].

We have studied the decay of two β -delayed neutrons emitters measured at IGISOL: ^{137}I and ^{95}Rb . Preliminary results for ^{137}I were presented in [4]. Relevant decay information for these nuclei is well established in the data bases. The values from ENSDF [33] of Q_β , neutron separation energy in the daughter S_n and neutron emission probability P_n are given in Table 3. We recently performed an accurate P_n measurement for both isotopes using the BELEN neutron counter [34] which gave 9.08(14)% for ^{95}Rb (compatible with the value in the table) and 7.76(14)% for ^{137}I (8% larger than the value in the table).

We compare with the simulation the sum energy spectra gated with β particles detected in a thin

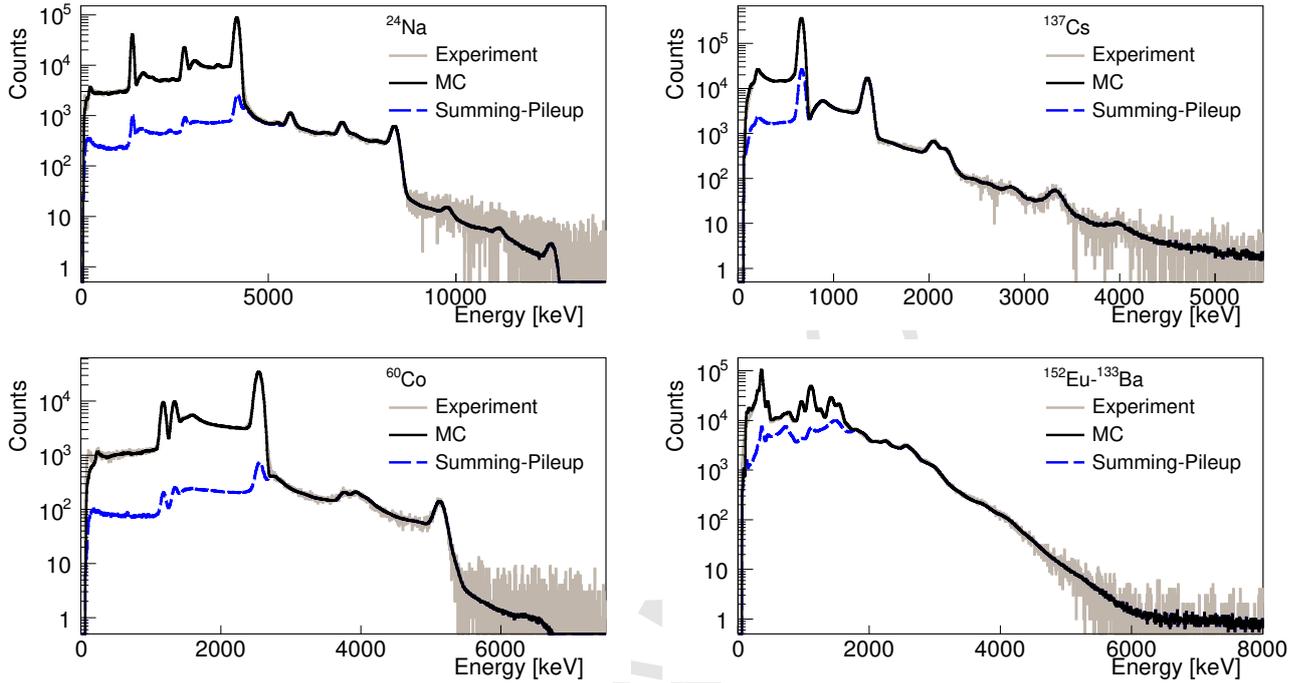


Figure 14: Experimental spectra of calibration sources after subtracting the environmental background (solid grey) compared with the MC simulations (solid black) taking into account the summing-pileup contamination (dashed blue).

Isotope	Q_β [MeV]	S_n [MeV]	P_n [%]
^{137}I	6.027(9)	4.02556(10)	7.14(23)
^{95}Rb	9.228(21)	4.348(7)	8.7(3)

Table 3: Properties of β -delayed neutron emitters used to test MC simulations.

plastic scintillator. These spectra are free from environmental background and are affected by the end-point energy dependence of the β efficiency which suppresses the β decays to states close to Q_β . In order to take this effect properly into account an event generator was implemented [15] that reproduces the known sequence of β -neutron- γ emission in the decay. It reproduces also the measured neutron spectra obtained from the ENDF/B VII.1 database, based on the work in [35]. The generator requires the reconstruction of the β intensity distribution from the measured neutron spectra using the information on neutron branchings to the excited levels in the final nucleus, I_n . The generator uses the associated γ branchings, I_γ , as well. Both I_n and I_γ data were retrieved from the ENSDF database [33].

In the simulations a time window for accumulation of the energy deposited after multiple neutron interactions was applied. This window takes into account the existence of a delay between neutron-induced γ -rays and the prompt β signal. A window of 500 ns was employed in accordance with the experimental coincidence time window between DTAS and the β plastic detector. This window ensures the collection of all the energy deposited. Figure 16 shows the comparison of measured and simulated spectra for the two β -delayed neutron emitters studied. As can be observed the reproduction of the gross structure above 6.8 MeV, mainly due to neutron capture in the iodine (^{127}I) in the crystal, is very good. In the case of ^{137}I the shape of the structure depends on the β -delayed neutron energy spectrum. It should be noted however that in the case of ^{95}Rb this structure includes partial summing of capture γ -rays with γ -rays emitted from excited states populated after neutron emission. The strongest of these γ -rays is also visible as a peak in the simulated spectra at 837 keV superimposed on the γ -ray background from neutron inelastic collisions. We found that the shape of the gross struc-

ture is quite sensitive to the de-excitation pattern after neutron emission. For ^{95}Rb , the use of the evaluated decay scheme available in ENSDF produced a wrong shape for the spectrum. However when we use the de-excitation scheme in ^{94}Sr measured by Kratz et al. [36] a good reproduction was obtained as can be seen in Fig. 16. Thus except when the neutron emission proceeds entirely to the ground state, the effect of the neutron energy distribution on the shape of the capture peak is obscured by the final nucleus γ spectra. Since the latter is often unknown or poorly known, it seems difficult to obtain reliable information about the shape of the β -delayed neutron spectrum from TAGS spectra in the general case.

Ideally the normalization of the β -delayed neutron contribution to the total spectrum should be determined by the P_n value, as we did in previous works with a BaF_2 spectrometer [13, 14, 15, 16]. However, in BaF_2 there is no clear peak to check the accuracy of this normalization. Recently it was suggested that precise P_n values could be obtained from TAGS measurements with a $\text{NaI}(\text{Tl})$ spectrometer [29]. In the present case, from the normalization of the simulated and measured counts in the capture bump at 6.8 MeV we have obtained P_n values of 6.8% and 10.9% for ^{137}I and ^{95}Rb respectively, after taking into account all the contaminants (summing-pileup and activity of the descendants). Note that in comparison with the numbers given in Table 3, we found a 5% smaller value for ^{137}I , while for ^{95}Rb a 25% larger value is obtained. Compared to our recently measured values [34] the differences are -12% and +20% respectively. We studied the dependence of the extracted P_n on the length of the time window applied to the experiment (coincidence gate) and to the MC simulation. We found that in the range 300-500 ns the results were stable within 3%. It should be noted that the value we obtain is 14% lower than the value of 7.9(4)% obtained by a similar procedure with MTAS [29]. We suspect that the differences with respect to the reference numbers given in Table 3 are due to uncertainties in the MC simulations of the interaction of neutrons. In view of this discrepancy and the fact that for ^{95}Rb we also obtain a large difference but of opposite sign, we conclude that further investigations are needed before deciding on the reliability of P_n extraction from TAGS spectra [29].

In any case, the key point for us is the reproducibility of the shape of the spectra of the β -

delayed neutron contamination that affects the extraction of $I_\beta(E_x)$ from the analysis of TAGS spectra. A proper determination of this background component is particularly relevant when extracting an accurate value for the β intensity above the neutron separation energy that proceeds by γ emission, P_γ [15]. The investigation of γ /neutron competition from neutron unbound states is a topic of current active research. The importance of the correction of the background due to β -delayed neutrons for the determination of P_γ using TAGS spectrometers made of $\text{NaI}(\text{Tl})$ can not be overlooked. This material has a large capture cross-section resulting in large β -delayed neutron detection efficiencies, of the order of 40%. We note, in particular, that this correction has been ignored in a recent measurement of P_γ for ^{70}Co decay with the SuN spectrometer [37] and might change their result significantly.

The sensitivity of the MC simulation of the β -n decay contamination to the knowledge of the decay (neutron spectrum and γ -cascades after neutron emission) represents a challenge for very neutron-rich nuclei in the general case where this information is poorly known or not known at all. Given that γ -rays produced by neutron interactions are delayed with respect to β -particle emission one can use timing information to discriminate between these signals [3]. We have tested this idea for ^{137}I and ^{95}Rb with reasonable results, as will be shown in a forthcoming publication. However, this type of time discrimination cannot be applied for the γ -ray de-excitation in the final nucleus after neutron emission since they are prompt with respect to the β -particles. The best option here seems to use the spectrometer itself to obtain information about this type of contamination as was suggested in [28]. The modularity of DTAS helps here, since there will be a certain degree of spatial separation between γ -rays coming from the final nucleus and those coming from neutron interactions. This can be exploited to tag β -delayed neutron events by setting a coincidence gate on the neutron capture "peak" observed, for example, in one half of the spectrometer and looking at the spectra in the other half of the spectrometer. Work to demonstrate the feasibility of this approach is in progress.

5. Conclusions

The characterization of the DTAS detector has been carried out. A gain stabilization system based

on a light pulse generator has been tested successfully. The non-proportionality of the light yield effects in a NaI(Tl) multi-crystal spectrometer were taken into account to reconstruct properly the sum of the total energy deposited in the spectrometer. The summing-pileup distortion of the spectrum was successfully computed using a revision of a method previously developed, and for high-rate measurements an improvement in this method has been introduced with the help of MC simulated data. A careful Geant4 MC simulation of the DTAS detector response to β -decays has been performed. The quality of the response function, needed for any TAGS analysis, has been validated after obtaining excellent agreement when comparisons were made with measurements of calibration sources. This includes in particular a good agreement of M_m module-multiplicity gated spectra. A nice agreement between the measured and simulated shape of the β -delayed neutron background was also obtained for two well known neutron emitters.

6. Acknowledgements

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References

- [1] B. Rubio, Int. J. Modern Phys. E 15 (2006) 1979.
- [2] J. Hardy, et al., Phys. Lett. B 71 (2) (1977) 307.
- [3] J. L. Tain, et al., Nucl. Instrum. and Methods A 803 (2015) 36.
- [4] V. Guadilla, et al., Nucl. Instrum. and Methods B 376 (2016) 334.
- [5] T. Davinson, et al., Technical Report for the Design, Construction and Commissioning of the Advanced Implantation Detector Array (AIDA)[link].
URL http://www.fair-center.eu/fileadmin/fair/publications_exp/aida_TDR.pdf
- [6] [link].
URL <http://www.mesytec.com/>
- [7] [link].
URL <http://www.caen.it/>
- [8] [link].
URL <http://www.ortec-online.com/>
- [9] M. Moszynski, et al., Nucl. Instrum. and Methods A 568 (2006) 739.
- [10] [link].
URL <http://www.crystals.saint-gobain.com/>
- [11] [link].
URL <http://www.berkeleyneucleonics.com/>
- [12] [link].
URL <https://fibertech-optica.com/>
- [13] J. L. Tain, et al., Phys. Rev. Lett. 115 (2015) 062502.
- [14] A.-A. Zakari-Issoufou, et al., Phys. Rev. Lett. 115 (2015) 102503.
- [15] E. Valencia, et al., Phys. Rev. C 95 (2017) 024320.
- [16] S. Rice, et al., Phys. Rev. C 96 (2017) 014320.
- [17] D. Engelkemeir, Rev. Sci. Instr. 27 (1956) 589.
- [18] S. Payne, et al., IEEE Transactions on Nuclear Science 56 (2009) 4.
- [19] D. Cano-Ott, et al., Nucl. Instrum. and Methods A 430 (1999) 333.
- [20] B. Rubio, et al., J. Phys. G: Nucl. Part. Phys. 44 (2017) 084004.
- [21] M. Karny, et al., Nucl. Instrum. and Methods B 126 (1997) 411.
- [22] P. Dorenbos, et al., IEEE Transactions on Nuclear Science 42 (1995) 6.
- [23] J. Valentine, et al., IEEE Transactions on Nuclear Science 45 (1998) 512.
- [24] B. Rasco, et al., Nucl. Instrum. and Methods A 788 (2015) 137.
- [25] D. Cano-Ott, et al., Nucl. Instrum. and Methods A 430 (1999) 488.
- [26] S. Agostinelli, et al., Nucl. Instrum. and Methods A 506 (2003) 250.
- [27] J. L. Tain, D. Cano-Ott, Nucl. Instrum. and Methods A 571 (2007) 728.
- [28] M. Karny, et al., Nucl. Instrum. and Methods A 836 (2016) 83.
- [29] B. C. Rasco, et al., Phys. Rev. C 95 (2017) 054328.
- [30] J. L. Tain, D. Cano-Ott, Nucl. Instrum. and Methods A 571 (2007) 719.
- [31] J. L. Tain, et al., Nucl. Instrum. and Methods A 774 (2015) 17.
- [32] E. Mendoza, et al., IEEE Transactions on Nuclear Science 61 (2014) 4.
- [33] Evaluated Nuclear Structure Data File, <http://www.nndc.bnl.gov/ensdf/>.
- [34] J. Agramunt, et al., Nucl. Instrum. and Methods A 807 (2016) 69.
- [35] M. Brady, Evaluation and Application of Delayed Neutron Precursor Data, Los Alamos National Laboratory, 1989.
- [36] K.-L. Kratz, et al., Z. Phys. A 306 (1982) 239.
- [37] A. Spyrou, et al., Phys. Rev. Lett. 117 (2016) 142701.

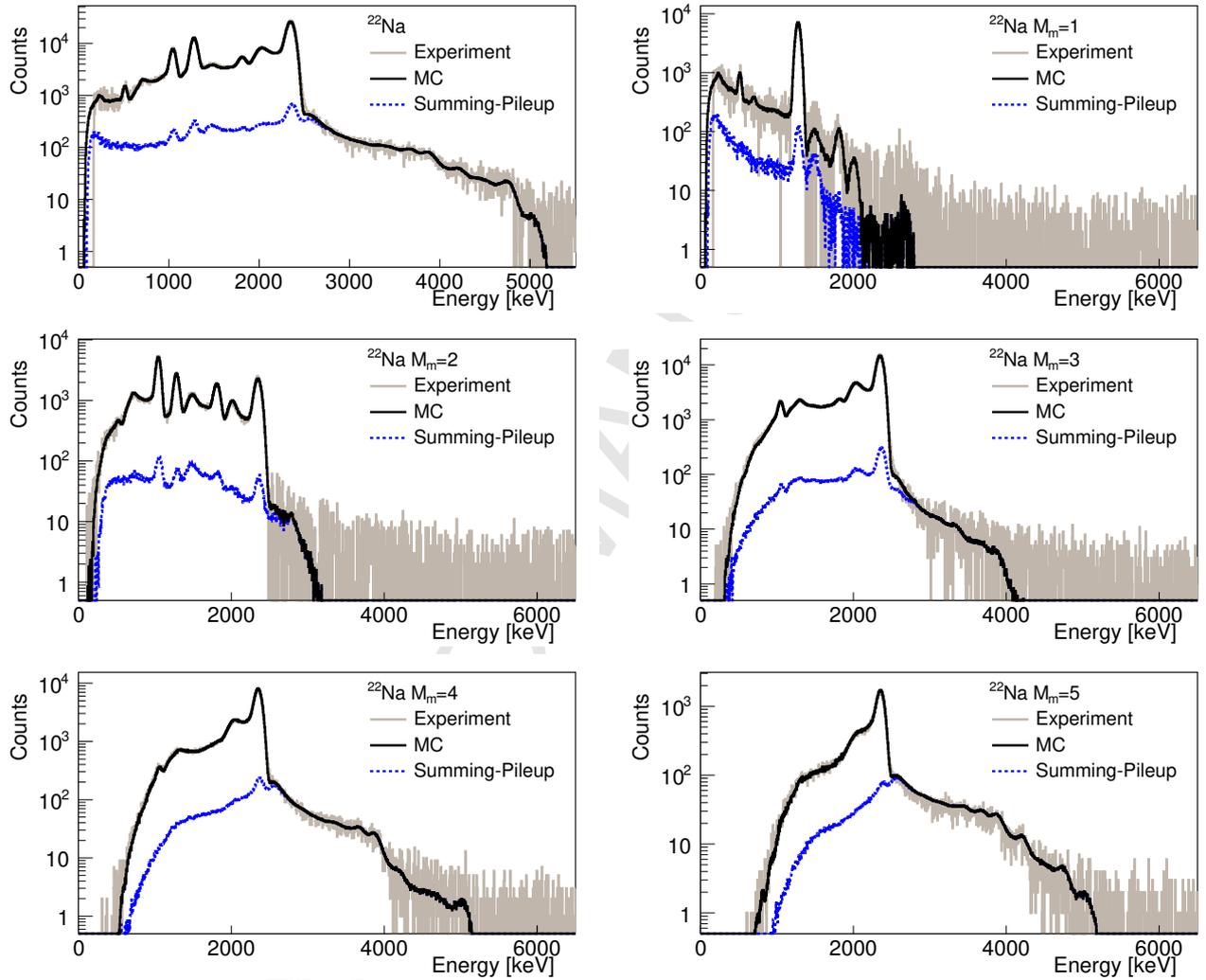


Figure 15: ^{22}Na experimental spectra after subtracting the environmental background (solid grey) compared with the MC simulations (solid black) taking into account the summing-pileup contamination (dashed blue). The sum energy spectrum without conditions and with a condition on module-multiplicity M_m from 1 to 5 is shown.

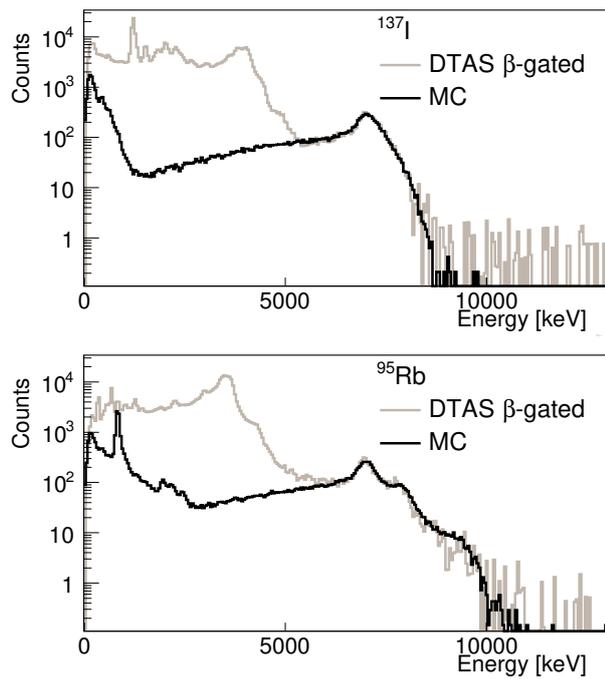


Figure 16: Simulation of the β -delayed neutron decay branch for neutron emitters measured in the commissioning of DTAS at IGISOL: ^{137}I (top) and ^{95}Rb (bottom). Experimental β -gated spectra (in grey) are compared to simulations (in black).