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MAPPINGS OF FINITE DISTORTION: BOUNDARY EXTENSIONS IN UNIFORM DOMAINS

TUOMO ÄKKINEN AND CHANG-YU GUO

ABSTRACT. In this paper, we consider mappings on uniform domains with exponentially integrable distortion whose Jacobian determinants are integrable. We show that such mappings can be extended to the boundary and moreover these extensions are exponentially integrable with quantitative bounds. This extends previous results of Chang and Marshall [7] on analytic functions, Poggi-Corradini and Rajala [20] and Äkkinen and Rajala [2] on mappings of bounded and finite distortion.

1. INTRODUCTION

A mapping $f: \Omega \rightarrow \mathbb{R}^n$, on a domain $\Omega \subset \mathbb{R}^n$ has finite distortion if the following conditions are fulfilled:

- (a) $f \in W_{\text{loc}}^{1,1}(\Omega, \mathbb{R}^n)$,
- (b) $J_f = \det(Df) \in L_{\text{loc}}^1(\Omega)$,
- (c) there exists a measurable $K_f: \Omega \rightarrow [1, \infty)$, so that for almost every $x \in \Omega$ we have

$$|Df(x)|^n \leq K_f(x)J_f(x),$$

where $|\cdot|$ is the operator norm. If $K_f \leq K < \infty$ almost everywhere, we say that f is K -quasiregular. If $n = 2$ and $K = 1$, we recover complex analytic functions. See [21], [22] and [23] for the theory of quasiregular mappings, and [11], [13] for the theory of mappings of finite distortion.

Let Ω be a bounded subdomain of \mathbb{R}^n with a distinguished point $x_0 \in \Omega$. We will consider the collection of mappings of finite distortion $f: \Omega \rightarrow \mathbb{R}^n$ satisfying $f(x_0) = 0$ and

$$(1.1) \quad \int_{\Omega} J_f(x) dx \leq |\Omega|,$$

where $|\Omega|$ denotes the Lebesgue n -measure of Ω . Moreover, we assume that there are constants $\lambda, \mathcal{K} > 0$ such that

$$(1.2) \quad \int_{\Omega} \exp(\lambda K_f) dx \leq \mathcal{K}.$$

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Let $\mathcal{F}_{\lambda, \mathcal{K}}(\Omega)$ denote the class of mappings satisfying (1.1) and (1.2). The class of K -quasiregular mappings with (1.1) will be denoted by $\mathcal{F}_K(\Omega)$.

If \mathcal{F} is the class of analytic functions on the unit disc \mathbb{D} with $f(0) = 0$ and

$$\int_{\mathbb{D}} |f'(z)|^2 dA(z) \leq \pi.$$

In [7, Corollary 1] Chang and Marshall proved the following sharp extension of an earlier result by Beurling [3], that is

$$(1.3) \quad \sup_{f \in \mathcal{F}} \int_0^{2\pi} \exp(|\bar{f}(e^{i\theta})|^2) d\theta < \infty.$$

This result is sharp, since for Beurling functions $B_a: \mathbb{D} \rightarrow \mathbb{C}$, $0 < a < 1$,

$$B_a(z) = \log \left(\frac{1}{1-az} \right) \log^{-1/2} \left(\frac{1}{1-a^2} \right)$$

one can show that

$$\lim_{a \rightarrow 0} \int_0^{2\pi} \exp(\gamma |B_a(e^{i\theta})|^2) d\theta = \infty$$

for every $\gamma > 1$. After this many natural generalizations of (1.3) have been proved. First, Marshall [17] gave another proof for (1.3) using different techniques than in [7]. Poggi-Corradini [18] was able to shorten the proof of Marshall [17] using so called ‘‘egg-yolk principle’’. Poggi-Corradini and Rajala proved the ‘‘egg-yolk principle’’ for quasiregular mappings in space and with this they showed in [20, Theorem 1.1] that for $n \geq 2$ we have

$$(1.4) \quad \sup_{f \in \mathcal{F}_K(\mathbb{B}^n)} \int_{\mathbb{S}^{n-1}} \exp(\alpha |\bar{f}(\xi)|^{\frac{n}{n-1}}) d\mathcal{H}^{n-1}(\xi) < \infty,$$

where the constant

$$\alpha = (n-1) \left(\frac{n}{2K} \right)^{\frac{1}{n-1}}$$

is sharp in dimension $n = 2$. Here \mathbb{B}^n and \mathbb{S}^{n-1} denote the unit ball and its boundary, respectively. Note that mappings $f \in \mathcal{F}_{\lambda, \mathcal{K}}(\mathbb{B}^n)$ can be extended to $\partial\mathbb{B}^n = \mathbb{S}^{n-1}$ using the fact that f has radial limits a.e. on \mathbb{S}^{n-1} . Recently it was shown by Äkkinen and Rajala [2] that (1.4) extends to the class $\mathcal{F}_{\lambda, \mathcal{K}}(\mathbb{B}^n)$ in the form

$$(1.5) \quad \sup_{f \in \mathcal{F}_{\lambda, \mathcal{K}}(\mathbb{B}^n)} \int_{\mathbb{S}^{n-1}} \exp(\beta |\bar{f}(\xi)|) d\mathcal{H}^{n-1}(\xi) < \infty,$$

where $\beta = \beta(n, \mathcal{K}, \lambda) > 0$. This result is sharp in terms of the exponent 1 but the sharp constant β is not known even in dimension $n = 2$, see discussion in [2, Section 4].

The purpose of this paper is to prove (1.4) and (1.5) for more general domains than the unit ball. We will consider domains Ω that are uniform, the justification for this restriction is given in Section 3. The boundary of a uniform domain is known to have Hausdorff dimension strictly less than

n but for technical reasons we will assume more regularity than this. If ν is a positive Borel measure we say that $\partial\Omega$ is upper s -Ahlfors regular with respect to ν if there exist constants $s \in (0, n)$ and $C > 0$ such that for any $x \in \Omega$

$$(1.6) \quad \nu(\partial\Omega \cap B(x, r)) \leq Cr^s \quad \text{for } 0 < r \leq \text{diam}(\Omega).$$

Our main theorems are the following:

Theorem 1.1. *Let $\Omega \subset \mathbb{R}^n$, $n \geq 2$, be a c -uniform domain with center x_0 and ν any positive Borel measure satisfying (1.6). Then there exists a constant $\alpha = \alpha(\Omega, n, K, c, s) > 0$ such that*

$$\sup_{f \in \mathcal{F}_K(\Omega)} \int_{\partial\Omega} \exp\left(\alpha |\bar{f}(w)|^{n/(n-1)}\right) d\nu(w) < \infty,$$

where \bar{f} is the extension of f to $\partial\Omega$ given in Definition 3.13.

Theorem 1.2. *Let $\Omega \subset \mathbb{R}^n$, $n \geq 2$, be a c -uniform domain with center x_0 and ν any positive Borel measure satisfying (1.6). Then there exists a constant $\beta = \beta(\Omega, n, \mathcal{K}, \lambda, c, s) > 0$ such that*

$$\sup_{f \in \mathcal{F}_{\lambda, \mathcal{K}}(\Omega)} \int_{\partial\Omega} \exp(\beta |\bar{f}(w)|) d\nu(w) < \infty,$$

where \bar{f} is the extension of f to $\partial\Omega$ given in Definition 3.13.

These results are sharp with respect to the exponent inside the exponential map since \mathbb{B}^n is a particular example of a uniform domain. The sharp constants α and β remain open, we would need new techniques for this problem. The general idea behind the proof is similar to the one used in the previous papers [2, 20], essentially relying on a version of a Theorem of Beurling [3] originally stated for analytic functions. Allowing more general domains produces a two-fold problem: first we have to find a way to extend our mapping to the boundary and also we need a new version of the Beurling Theorem.

It could be possible to deduce Theorem 1.2 from earlier works by Cianchi [5] and [6] as follows: one first uses the Orlicz-Sobolev extension in uniform domains for each of the component functions f_i , $i = 1, \dots, n$, to obtain a global Orlicz-Sobolev function $f'_i : \mathbb{R}^n \rightarrow \mathbb{R}$; then we can choose a quasicontinuous representative of f'_i , which means outside a set of small capacity, f'_i is continuous. This in return gives us a natural way to extend our original function f_i to the boundary of Ω as $\bar{f}_i : \partial\Omega \rightarrow \mathbb{R}$. Now extending results for Orlicz-Sobolev functions in [6] to similar ones given for Lorentz-Sobolev functions in [5, Theorems 2.4 and 2.5] and using above arguments one can show that Theorem 1.2 is valid for each component function \bar{f}_i . So Theorem 1.2 follows once we set $\bar{f} = (\bar{f}_1, \dots, \bar{f}_n)$.

In this way, one could prove Theorem 1.2 with a constant α independent of λ and \mathcal{K} and thus cannot obtain the sharp constant since the sharp constant depends on the distortion K ; see [20]. We believe that the techniques

used here can be further extended to more irregular domain than uniform domains since most of our results can be generalized to domains satisfying *quasihyperbolic boundary condition* (cf. [16]).

This paper is organized as follows. Section 2 fixes the notation. In section 3 we will show that f can be extended to the boundary of a uniform domain along John-curves and that this extension is well-defined with a small exceptional set on $\partial\Omega$ in terms of Hausdorff gauges. Section 4 contains the main ingredients of the proofs of Theorems 1.1 and 1.2, which is the capacity estimate given in Theorem 4.1. Finally in Section 5 we will prove our main Theorems.

2. NOTATION AND PRELIMINARIES

A curve γ in Ω is a continuous mapping $\gamma: [0, 1] \rightarrow \Omega$. The trace $\gamma([0, 1])$ of a curve γ is denoted by $|\gamma|$. A curve γ is said to connect points $x, y \in \Omega$ if $\gamma(0) = y$ and $\gamma(1) = x$ and points $x \in \Omega$, $y \in \partial\Omega$ if $\gamma(1) = x$, $\gamma((0, 1]) \subset \Omega$ and

$$\lim_{t \rightarrow 0^+} \gamma(t) = y.$$

We also use the notation $l(\gamma)$ to denote the (euclidean) length of a curve γ .

The quasihyperbolic metric k_Ω in a domain $\Omega \subsetneq \mathbb{R}^n$ is defined as

$$k_\Omega(x, y) = \inf_{\gamma} \int_{\gamma} \frac{ds}{d(z, \partial\Omega)},$$

where the infimum is taken over all rectifiable curves γ in Ω which join x to y . This metric was introduced by Gehring and Palka in [9]. A curve γ connecting x to y for which

$$k_\Omega(x, y) = \int_{\gamma} \frac{ds}{d(z, \partial\Omega)}$$

is called a quasihyperbolic geodesic. Quasihyperbolic geodesics joining any two points of a proper subdomain of \mathbb{R}^n always exists; see [8, Lemma 1]. Given two points $x, y \in \Omega$, we denote by $[x, y]$ any quasihyperbolic geodesic that joins x and y .

Let $\Omega \subset \mathbb{R}^n$ be a domain. We will denote by $\mathscr{W} := \mathscr{W}(\Omega)$ a Whitney type decomposition of Ω , where \mathscr{W} is a collection of closed cubes $Q \subset \Omega$ with pairwise disjoint interiors and having edges parallel to coordinate axis, such that

$$\Omega = \bigcup_{Q \in \mathscr{W}} Q.$$

Moreover, there exists $\lambda_0 > 1$ such that the cubes dilated by factor λ_0 have uniformly bounded overlap, the diameters of $Q \in \mathscr{W}$ take values in $\{2^{-j} : j \in \mathbb{Z}\}$ and satisfy condition

$$\text{diam } Q \leq d(Q, \partial\Omega) \leq 4 \text{ diam } Q.$$

We will also use the notation

$$\mathscr{W}_j = \{Q \in \mathscr{W} : \text{diam } Q = 2^{-j}\}.$$

3. BOUNDARY EXTENSIONS

In the previous papers [17], [20], [19] and [2] the boundary extension of $f \in \mathcal{F}_{\lambda, \mathcal{K}}(\mathbb{B}^n)$ was defined via the radial limits of f that exist at almost every point in \mathbb{S}^{n-1} by absolute continuity on lines, (1.1) and (1.2). In this section we prove similar results for more general class of domains, namely for uniform domains.

3.1. Existence of limits along curves. When considering general bounded domains Ω there might be big parts of the boundary $\partial\Omega$ that are not accessible by a rectifiable curve inside Ω , see Figure 1.



FIGURE 1. Topologist comb

To avoid such situations, we restrict our considerations to John domains.

Definition 3.1. A bounded domain $\Omega \subset \mathbb{R}^n$ is a c -John domain with distinguished point $x_0 \in \Omega$ (center) if every point $x \in \Omega$ can be joined to x_0 by a c -John curve. Recall also that a rectifiable curve $\gamma: [0, l] \rightarrow \Omega$ parametrized by arc length is called a c -John curve joining x to x_0 if it satisfies $\gamma(0) = x$, $\gamma(l) = x_0$ and

$$(3.1) \quad d(\gamma(t), \partial\Omega) \geq \frac{1}{c}t$$

for every $t \in [0, l]$.

It follows from the lower semicontinuity of $l(\gamma)$ and Ascoli-Arzelà Theorem that each boundary point can be connected to x_0 by a curve satisfying (3.1). Thus every point on the boundary is accessible, even in a non-tangential manner. Notice that typically there will be uncountably many c -John curves between $x_0 \in \Omega$ and $y \in \partial\Omega$. For the rest of this subsection the standing assumptions are: Let Ω be a c_0 -John domain with center x_0 , \mathscr{W} a Whitney decomposition of Ω and $c \geq c_0$.

If $\xi \in \partial\Omega$ then define $I_c(\xi, x_0)$ as the collection of all c -John curves connecting ξ to x_0 and

$$P_c(\xi) = \{Q \in \mathscr{W} : Q \cap \gamma \neq \emptyset \text{ for some } \gamma \in I_c(\xi, x_0)\}.$$

For $Q \in \mathscr{W}$ and $E \subset \partial\Omega$, we define the c -shadow of Q on E by

$$S_E^c(Q) = \{\xi \in E : Q \in P_c(\xi)\}.$$

When $E = \partial\Omega$, we write $S^c(Q)$ instead of $S_{\partial\Omega}^c(Q)$. We will omit the sub/superscripts c in the above notation if $c = c_0$.

The following two lemmas are well-known and the proofs can be found, for instance, in [10, Section 4].

Lemma 3.2. *Let $Q \in \mathscr{W}$ and $\xi \in \partial\Omega$. Then $S^c(Q)$ is closed and there exists a constant $C(c) > 0$ such that*

$$\text{diam } S^c(Q) \leq C(c) \text{diam } Q.$$

Also there exists a constant $C(n, c) > 0$ such that for any $j \in \mathbb{Z}$

$$\#\{Q \in \mathscr{W}_j : \xi \in S^c(Q)\} \leq C(n, c).$$

Lemma 3.3. *Assume that μ is a Borel measure on $\partial\Omega$ and $E \subset \partial\Omega$ measurable. Then there exists a constant $C(n, c) > 0$ such that, for each $j \in \mathbb{Z}$ we have*

$$\sum_{Q \in \mathscr{W}_j} \mu(S_E^c(Q)) \leq C\mu(E).$$

Our aim is to prove that mappings in the classes $\mathcal{F}_K(\Omega)$ and $\mathcal{F}_{\lambda, K}(\Omega)$ have limits along c_0 -John curves emanating from the center x_0 . This will be true for big pieces of the boundary in terms of Hausdorff dimension. For this we introduce a notion of discrete length of a curve: Assume that $\xi \in \partial\Omega$ and $\gamma \in \mathbf{I}(\xi, x_0)$. Given a continuous mapping $f: \Omega \rightarrow \mathbb{R}^n$ we define the discrete length of $f(\gamma)$ by

$$\ell_d[f(\gamma)] = \sum_{\substack{Q \in \mathscr{W} \\ Q \cap \gamma \neq \emptyset}} \text{diam } f(Q).$$

It follows from the definition and Lemma 3.2 that if $\ell_d[f(\gamma)] < \infty$ then the limit

$$\lim_{t \rightarrow 0^+} f(\gamma(t)) \in \mathbb{R}^n$$

exists.

We also need oscillation estimates for mappings of finite distortion $f: \Omega \rightarrow \mathbb{R}^n$. The proof of the next lemma is contained in [15], see also [1].

Lemma 3.4. *Let $f: \Omega \rightarrow \mathbb{R}^n$ be a mapping of finite distortion and assume $\sigma > 1$. Then we have the following:*

i) If f is quasiregular, then

$$\text{diam } f(B) \leq C \left(\int_{\sigma B} J_f(x) dx \right)^{1/n}$$

for every $B = B(x, r)$ for which $\sigma B \subset\subset \Omega$. Here constant $C > 0$ depends only on n and K .

ii) If f has exponentially integrable distortion, then

$$\text{diam } f(B) \leq C \left(\int_{\sigma B} J_f(x) dx \right)^{1/n} \log(1/\text{diam } B)^{1/n}$$

for every $B = B(x, r)$ for which $\sigma B \subset\subset \Omega$. Here constant $C > 0$ depends only on n , λ and $\int_{\Omega} \exp(\lambda K) dx$.

Now we are ready to state a theorem which tells that a mapping of finite distortion can be extended to $\partial\Omega$ along John curves with a small exceptional set.

Theorem 3.5. *Assume that $f: \Omega \rightarrow \mathbb{R}^n$ is a mapping of finite distortion with $J_f \in L^1(\Omega)$. Let E_f be the set of points $w \in \partial\Omega$ for which there exists a curve $\gamma \in I(w, x_0)$ so that f does not have a limit along γ . Then we have the following:*

i) If f is quasiregular and h is a doubling gauge function satisfying

$$(3.2) \quad \int_0^1 h(t)^{1/(n-1)} \frac{dt}{t} < \infty$$

then $\mathcal{H}^h(E_f) = 0$.

ii) If f has exponentially integrable distortion and h is a doubling gauge function satisfying

$$(3.3) \quad \int_0^1 (h(t) \log(1/t))^{1/(n-1)} \frac{dt}{t} < \infty$$

then $\mathcal{H}^h(E_f) = 0$.

Proof. We will follow the proof of Theorem 1.1 in [15] and actually prove that $\mathcal{H}^h(A_\infty) = 0$, where A_∞ is the set of points in $\xi \in \partial\Omega$ for which there is a curve $\gamma \in I(\xi, x_0)$ so that $f(\gamma)$ has infinite discrete length. Note that $E_f \subset A_\infty$. On the contrary, assume that $\mathcal{H}^h(A_\infty) > 0$. Then $\mathcal{H}^h(A_k) > 0$, where A_k is the set of points in $\xi \in \partial\Omega$ for which there exist $\gamma \in I(\xi, x_0)$ so that $\ell_d[f(\gamma)] \geq k$. Then by Frostman's lemma there exists a Borel measure μ supported in A_k so that

$$\mu(B^n(x, r)) \leq h(r)$$

for every $B^n(x, r) \subset \mathbb{R}^n$ and

$$\mu(E_k) \simeq \mathcal{H}_\infty^h(A_k) \geq \mathcal{H}_\infty^h(A_\infty) > 0.$$

By the definition of A_k

$$\begin{aligned} \mu(A_k)k &\leq \int_{A_k} \ell_d[f(\gamma_w)] d\mu(w) \\ &\leq \int_{A_k} \sum_{\substack{Q \in \mathcal{W} \\ Q \cap \gamma_w \neq \emptyset}} \text{diam } f(Q) d\mu(w). \end{aligned}$$

From now on, we assume that f has exponentially integrable distortion. Let $\lambda_0 > 1$ be the constant in the definition of Whitney decomposition \mathscr{W} for which the cubes $q \in \mathscr{W}$ dilated by a factor λ_0 have uniformly bounded overlap. If $Q \in \mathscr{W}$ then set $B_Q = B^n(x_Q, r_Q)$, where x_Q is the center of Q and $r_Q = \text{diam } Q/2$. Using Lemma 3.4 for balls B_Q we have the following chain of inequalities:

$$\begin{aligned}
\mu(A_k)k &\leq \int_{A_k} \sum_{\substack{Q \in \mathscr{W} \\ Q \cap \gamma_w \neq \emptyset}} \text{diam } f(B_Q) d\mu(w) \\
&\leq \sum_{Q \in \mathscr{W}} \int_{A_k} \chi_{S(Q)}(w) d\mu(w) \text{diam } f(B_Q) \\
&\leq \sum_{Q \in \mathscr{W}} \mu(S_{A_k}(Q)) \text{diam } f(B_Q) \\
&\leq C \sum_{Q \in \mathscr{W}} \mu(S_{A_k}(Q)) \log^{\frac{1}{n}} \left(\frac{1}{\text{diam } B_Q} \right) \\
&\quad \times \left(\int_{\lambda_0 B_Q} J_f(x) dx \right)^{1/n} \\
&\leq C \left(\sum_{Q \in \mathscr{W}} \mu(S_{A_k}(Q))^{\frac{n}{n-1}} \log^{\frac{1}{n-1}} \left(\frac{1}{\text{diam } B_Q} \right) \right)^{(n-1)/n} \\
&\quad \times \left(\sum_{Q \in \mathscr{W}} \int_{\lambda_0 B_Q} J_f(x) dx \right)^{1/n} \\
&\leq C \left(\sum_{j=1}^{\infty} \sum_{Q \in \mathscr{W}_j} \mu(S_{A_k}(Q))^{n/(n-1)} j^{1/(n-1)} \right)^{(n-1)/n}.
\end{aligned}$$

Above we also used the fact that $\text{diam } B_Q \simeq 2^{-j}$ for $Q \in \mathscr{W}_j$ and also the property that $\lambda_0 Q$'s have uniformly bounded overlap. The left hand side is approximated using Lemmas 3.2 and 3.3 together with the doubling property of gauge function h

$$\begin{aligned}
\sum_{Q \in \mathscr{W}_j} \mu(S_{A_k}(Q))^{n/(n-1)} &\leq \max_{Q \in \mathscr{W}_k} \mu(S_{A_k}(Q))^{1/(n-1)} \sum_{Q \in \mathscr{W}_j} \mu(S_{A_k}(Q)) \\
&\leq C(n, c) \max_{Q \in \mathscr{W}_k} \mu(S(Q))^{1/(n-1)} \mu(A_k) \\
&\leq C(n, c, D) h(2^{-j})^{1/(n-1)} \mu(A_k).
\end{aligned}$$

Putting these estimates together gives

$$(3.4) \quad \mu(A_k)k \leq C\mu(A_k)^{(n-1)/n} \left(\sum_{j=1}^{\infty} h(2^{-j})^{1/(n-1)} j^{1/(n-1)} \right)^{(n-1)/n}.$$

Since h satisfies (3.3) we know that the sum on the right hand side is finite (and independent of k) and thus $\mu^{1/n}(A_k)k \leq C$ for every $k \in \mathbb{N}$ and thus $\mu(A_k) \rightarrow 0$ as $k \rightarrow \infty$, which is a contradiction. Thus we have $\mathcal{H}^h(A_\infty) = 0$ and the claim follows. For quasiregular f we do not get the term $j^{1/(n-1)}$ in (3.4) and thus a gauge satisfying (3.2) gives the desired conclusion. \square

Remark 3.6. As a corollary of Theorem 3.5 we get that the set E_f has Hausdorff dimension zero as well as the set A_∞ . Moreover, for a quasiregular f we have $\mathcal{H}^{h_\varepsilon}(E_f) = 0$ for every $\varepsilon > 0$, where $h_\varepsilon(t) = 1/\log^{n-1+\varepsilon}(1/t)$. Also, if f has exponentially integrable distortion, then $\mathcal{H}^{h_\varepsilon}(E_f) = 0$ for every $\varepsilon > 1$. In the quasiregular case this means that E_f is of conformal capacity zero in terms of gauge functions, see [25].

3.2. Uniqueness of limits along John curves. Assume that Ω is a c_0 -John domain. By the previous section we have a natural candidate for the extension of $f \in \mathcal{F}_{\lambda, \mathcal{K}}(\Omega)$ to $\partial\Omega$, that is $\bar{f}: \partial\Omega \rightarrow \mathbb{R}^n$ so that

$$(3.5) \quad \bar{f}(w) = \lim_{t \rightarrow 0^+} f(\gamma_w(t)),$$

where γ_w is a c_0 -John-curve connecting x_0 to w . The problem with this extension is that it might not be well-defined even for a smooth mapping $f: \Omega \rightarrow \mathbb{R}^n$.

Example 3.7. Let $\Omega = \mathbb{D} \setminus L$, where L is the (euclidean) line segment connecting $-\frac{1}{2}e_1$ to $\frac{1}{2}e_1$. Then $\Omega \subset \mathbb{R}^2$ is a c -John domain with $x_0 = -\frac{3}{4}e_1$. Set $A = B(0, 1/4) \cap \{x_2 > 0\}$ and $B = B(0, 1/4) \cap \{x_2 < 0\}$. There is a smooth mapping $f: \Omega \rightarrow \mathbb{R}^2$ such that $f(A) = A + i$ and $f(B) = B - i$ and thus the extension \bar{f} is not well-defined for any point in $(t, 0) \subset \partial\Omega$, where $t \in (-1/4, 1/4)$.

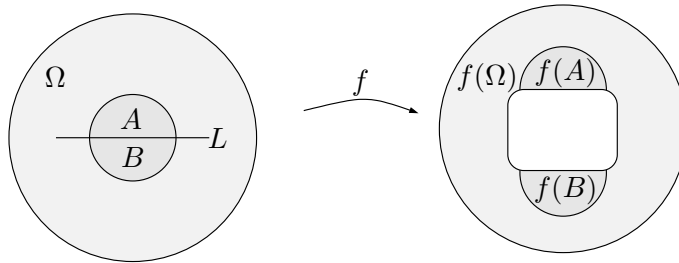


FIGURE 2. John is not sufficient

In order to exclude such domains we have to impose more conditions on Ω . Examples like this do not work if we strengthen our assumption from John condition to Ω being uniform.

Definition 3.8. A bounded domain $\Omega \subset \mathbb{R}^n$ is a c -uniform domain if there exists a constant $c > 0$ such that each pair of points $x_1, x_2 \in \Omega$ can be joined by a c -uniform curve γ in Ω , namely a curve γ with the following two properties:

$$(3.6) \quad l(\gamma) \leq c|x_1 - x_2|,$$

$$(3.7) \quad \min_{i=1,2} l(\gamma(x_i, x)) \leq cd(x, \partial\Omega) \quad \text{for all } x \in \gamma.$$

These domains first appeared in the work of Gehring and Osgood [8]. Later these domains were used in many different contexts such as Sobolev extension domains and uniformizing Gromov hyperbolic spaces, for these see [14, 4]. It should be noticed that quasihyperbolic geodesics between two points in a c -uniform domain are c_1 -uniform curves with $c_1 = c_1(n, c)$ [8]. We say that Ω is a John (or uniform) domain if it is c -John (or c -uniform) for some $c > 0$.

We may assign for each c -uniform domain a center point x_0 . Indeed, if Ω is a c -uniform domain, then we may fix a point $x_0 \in \Omega$ such that Ω is a c_1 -John domain with center x_0 . Moreover, the John constant c_1 depends only on c [24, section 2.17]. Since we are studying boundary extension of mappings, it is convenient to introduce the following concept.

Definition 3.9 (Uniform domain with center). We say that Ω is a c_0 -uniform domain with center x_0 if Ω is a uniform domain and as a John domain, it is c_0 -John with center x_0 .

As was observed in Section 3.1, for a given point $x \in \partial\Omega$ in a John domain Ω , there are typically several John curves connecting x to x_0 . Thus it might happen that $f \in \mathcal{F}_{\lambda, \mathcal{K}}(\Omega)$ has different limits at $w \in \partial\Omega$ along different John curves in $I(w, x_0)$; see for instance Figure 2. In the rest of this section, we will show that if Ω is c_0 -uniform domain with center x_0 , then mappings in $\mathcal{F}_{\lambda, \mathcal{K}}(\Omega)$ can be extended to $\partial\Omega$ in a unique way along curves in $I(w, x_0)$.

Theorem 3.10. *Let $\Omega \subset \mathbb{R}^n$ be a c_0 -uniform domain with center x_0 , $f \in \mathcal{F}_{\lambda, \mathcal{K}}(\Omega)$ and h a doubling gauge function satisfying (3.3). Then f has a unique limit along curves $\gamma \in I(w, x_0)$ for \mathcal{H}^h -a.e. $\xi \in \partial\Omega$, i.e. if $\gamma, \eta \in I(\xi, x_0)$ so that*

$$\lim_{t \rightarrow 0^+} f(\gamma(t)) = a \quad \text{and} \quad \lim_{t \rightarrow 0^+} f(\eta(t)) = b.$$

Then $a = b$.

We need the following two lemmas to prove Theorem 3.10.

Lemma 3.11. *Let Ω be a c_0 -John domain with center x_0 and $c \geq c_0$. If f has exponentially integrable distortion and F_f is the set of points $\xi \in \partial\Omega$ for which there exists a curve $\gamma \in \mathbf{I}_c(\xi, x_0)$ so that*

$$\sum_{i=1}^{\infty} a_i = \infty,$$

then $\mathcal{H}^h(F_f) = 0$, where h is as in (3.3). Above, a_i is defined in the following way: for a fixed $\gamma \in \mathbf{I}_c(\xi, x_0)$ connecting x_0 to $\xi \in \partial\Omega$, we denote by $Q_i \in \mathcal{W}$ those cubes that intersect γ and labeled in the order from x_0 to ξ , and set

$$a_i = \left(\int_{\lambda_0 B_{Q_i}} J_f(x) dx \right)^{1/n} \log(1/\text{diam } B_{Q_i})^{1/n}.$$

Proof. Notice that if $\xi \in F_f$, then

$$\sum_{i=1}^{\infty} \left(\int_{\lambda_0 B_{Q_i}} J_f(x) dx \right)^{1/n} \log(1/\text{diam } B_{Q_i})^{1/n} = \infty.$$

Let A_k be the collection of points $\xi \in \partial\Omega$ for which there exists $\gamma_\xi \in \mathbf{I}_c(\xi, x_0)$ so that labeling as in the lemma we have

$$\sum_{i=1}^{\infty} \left(\int_{\lambda_0 B_{Q_i}} J_f(x) dx \right)^{1/n} \log(1/\text{diam } B_{Q_i})^{1/n} \geq k.$$

Then $F_f \subset A_\infty$ and thus it is enough to show that $\mathcal{H}^h(A_\infty) = 0$. Towards a contradiction assume the opposite $\mathcal{H}^h(A_\infty) > 0$ and denote by μ the Frostman measure whose support lies in A_k , which satisfies $\mu(A_k) \simeq \mathcal{H}_\infty^h(A_k) \geq \mathcal{H}_\infty^h(A_\infty)$ and

$$\mu(B(x, r)) \leq Ch(r)$$

for every ball $B(x, r)$. Then we may follow the proof of Theorem 3.5 to obtain

$$\mu(A_k)k \leq C$$

with a universal constant $C > 0$. Thus we have a contradiction and the lemma is proved. \square

Lemma 3.12. *Let $\Omega \subset \mathbb{R}^n$ be a c_0 -uniform domain with center x_0 and \mathcal{W} a Whitney decomposition of Ω . Then there exists a constant C such that for any $s > 0$ and pair $x, \tilde{x} \in \Omega$ with $d(\tilde{x}, \partial\Omega) \geq s$, $d(x, \partial\Omega) \geq s$ and $|x - \tilde{x}| \leq c_1 s$ there exists a chain of Whitney cubes $\{Q_k\}_{k=1}^N$ connecting points x and \tilde{x} such that the number of cubes is uniformly bounded with respect to s , i.e. $N \leq C$, where C depends only on n , c_0 and c_1 but not on s . Moreover if $\xi \in S(Q)$ for some $Q \in \cup_{i=1}^N Q_i$ then $\xi \in S^{c'}(\tilde{Q})$ for any \tilde{Q} from this collection with $c' = c'(C, n, c_0) \geq c_0$.*

Proof. Let $s > 0$ and fix $x, \tilde{x} \in \Omega$ with $d(\tilde{x}, \partial\Omega) \geq s$, $d(x, \partial\Omega) \geq s$ and $|x - \tilde{x}| \leq c_1 s$. Connect x to \tilde{x} by a quasihyperbolic geodesic $[x, \tilde{x}]_k$ and enumerate the cubes $Q \in \mathcal{W}$, that intersect $[x, \tilde{x}]_k$, in a such way that $x \in Q_1$, $\tilde{x} \in Q_N$ and $Q_i \cap Q_{i+1} \neq \emptyset$ for every i . By [16, page 421] we know that the number of cubes in this chain is comparable to $k_\Omega(x, \tilde{x})$ and moreover, since Ω is uniform, by [4, Lemma 2.13] we have a constant $c'(n, c_0) > 0$ such that

$$(3.8) \quad k_\Omega(x, \tilde{x}) \leq c' \log \left(1 + \frac{|x - \tilde{x}|}{\min\{d(x, \partial\Omega), d(\tilde{x}, \partial\Omega)\}} \right).$$

The first part of the lemma follows.

Next we prove the latter part. Since $\xi \in S(Q)$ for some $Q \in \cup_{i=1}^N Q_i$ there exists a curve $\gamma \in I(\xi, x_0)$ such that $\gamma \cap Q \neq \emptyset$. Let $y \in Q \cap \gamma$. Let \tilde{Q} be any other Whitney cubes from the collection and fix $z \in \tilde{Q}$. Let γ_2 denote the piecewise linear curve from y to z s.t. $|\gamma_2| \subset \cup_{i=1}^N Q_i$. We define curve $\tilde{\gamma}$ as follows. We first travel from ξ along γ to y and denote this by γ_1 . Next, follow γ_2 from y to z and denote this by γ_2^+ , then go back from z to y along γ_2 and denote this by γ_2^- . Finally, follow γ from y to x_0 and denote this by γ_3 . In this way we obtained curve $\tilde{\gamma} = \gamma_1 \star \gamma_2^+ \star \gamma_2^- \star \gamma_3$ connecting ξ to x_0 . To illustrate that $\tilde{\gamma}$ is a c' -John curve with $c' = c'(C, n, c_0) \geq c_0$, we assume that $\tilde{\gamma} : [0, l'] \rightarrow \Omega$ is parametrized by arc length and we have to verify (3.1) with c' .

For $t \in [0, l(\gamma_1)]$ we have

$$t \leq c_0 d(\gamma(t), \partial\Omega),$$

since $\gamma \in I(\xi, x_0)$. Next assume that $t \in [l(\gamma_1), l(\gamma_1) + 2l(\gamma_2)]$. First note that all the Whitney cubes in the collection have uniformly comparable size since the number N is bounded from above uniformly by C and adjacent Whitney cubes have comparable size and thus we may estimate

$$l(\gamma_2) \leq c(n, C) d(\tilde{\gamma}(t), \partial\Omega) \quad \text{and} \quad d(y, \partial\Omega) \leq c(n, C) d(\tilde{\gamma}(t), \partial\Omega).$$

Using the above estimates, we have

$$t \leq l(\gamma_1) + 2l(\gamma_2) \leq c_0 d(y, \partial\Omega) + 2l(\gamma_2) \leq c(C, n, c_0) d(\tilde{\gamma}(t), \partial\Omega).$$

Last, assume that $t \in [l(\gamma_1) + 2l(\gamma_2), l(\tilde{\gamma})]$. Now we have $|y - \tilde{\gamma}(t)| \leq c_0 d(\tilde{\gamma}(t), \partial\Omega)$, which imply

$$l(\gamma_2) \leq c(n, C) d(y, \partial\Omega) \leq c(n, C, c_0) d(\tilde{\gamma}(t), \partial\Omega)$$

and thus

$$t \leq c_0 d(\tilde{\gamma}(t), \partial\Omega) + 2l(\gamma_2) \leq c(n, C, c_0) d(\tilde{\gamma}(t), \partial\Omega).$$

□

Proof of Theorem 3.10. Let F_f denote the set defined in Lemma 3.11. We know that $\mathcal{H}^h(F_f) = 0$. Let $\xi \in \partial\Omega \setminus F_f$ and $\gamma \in I(\xi, x_0)$ be such that

$$\lim_{t \rightarrow 0^+} f(\gamma(t)) = a \in \mathbb{R}^n.$$

We will show that f has the same limit along any curve $\eta \in I(\xi, x_0)$. Fix such a curve η . Given $r > 0$, let $t_r, \tilde{t}_r \in [0, 1]$ be such that $|\gamma(t_r) - w| = r$ and $|\eta(\tilde{t}_r) - w| = r$. Since γ, η are both c_0 -John curves this implies that

$$d(\gamma(t_r), \partial\Omega) \geq \frac{r}{c_0} \quad \text{and} \quad d(\eta(\tilde{t}_r), \partial\Omega) \geq \frac{r}{c_0}.$$

We now apply Lemma 3.12 with $s = r/c_0$ and for points $\gamma(t_r)$ and $\eta(\tilde{t}_r)$ to find a finite sequence of Whitney cubes $\{Q_i\}_{i=1}^N$ with $N \leq C$. Choose $x_i \in Q_i$ so that $x_1 = \gamma(t_r)$ and $x_N = \eta(\tilde{t}_r)$. Then by Lemma 3.4 we have

$$\begin{aligned} |f(\gamma(t_r)) - f(\eta(\tilde{t}_r))| &\leq \sum_{i=1}^N |f(x_{i+1}) - f(x_i)| \\ &\leq C_1 \sum_{i=1}^N \left(\int_{\lambda_0 B_{Q_i}} J_f(x) dx \right)^{1/n} \log(1/\text{diam } Q_i)^{1/n} \\ (3.9) \quad &\leq c(C, C_1) \max_{1 \leq i \leq N} \left(\int_{\lambda_0 B_{Q_i}} J_f(x) dx \right)^{1/n} \log(1/\text{diam } Q_i)^{1/n}. \end{aligned}$$

By Lemma 3.12 there is a constant $c' \geq c_0$ s.t. $\xi \in S^{c'}(Q_i)$ for every $i = 1, \dots, N$. Note that N does not depend on r . It follows from Lemma 3.11 that we may choose $r > 0$ so small that

$$\max_{1 \leq i \leq N} \left(\int_{\sigma Q_i} J_f(x) dx \right)^{1/n} \log(1/\text{diam } Q_i)^{1/n} < \varepsilon.$$

This together with (3.9) shows that

$$|f(\gamma(t_r)) - f(\eta(\tilde{t}_r))| \rightarrow 0$$

as $r \rightarrow 0$. From this we may conclude that

$$\lim_{t \rightarrow 0^+} f(\gamma(t)) = \lim_{t \rightarrow 0^+} f(\eta(t))$$

and the proof is complete. \square

Based on Theorem 3.10, we introduce the following definition.

Definition 3.13 (Boundary extension). Let Ω be a c_0 -uniform domain with center x_0 and $f \in \mathcal{F}_{\lambda, \mathcal{K}}(\Omega)$. We denote by $\bar{f}(\xi)$ the unique limit of f at $\xi \in \partial\Omega$ along c_0 -John curves.

Notice that by Theorem 3.5 and Theorem 3.10, \bar{f} is well-defined \mathcal{H}^h -a.e. in $\partial\Omega$. Note that the unit ball $\mathbb{B} \subset \mathbb{R}^n$ is a 1-uniform domain with center at the origin. According to Definition 3.13, $\bar{f}(w)$ is the unique limit of f along the unique 1-John curve, which is just the line segment \overline{ow} , thus in this case our Definition 3.13 coincides with the radial extension.

4. WEIGHTED CAPACITY ESTIMATES IN JOHN DOMAINS

In this section, we prove the following weighted capacity estimates in John domains. The techniques used here are rather standard and have been used effectively in [16, 10]. On the other hand, they give a new proof of the key level set estimate, namely Theorem 1.2 in [2], where Ω is the unit ball \mathbb{B} .

Recall that for disjoint compact sets E and F in $\bar{\Omega}$, we denote by $\text{Cap}_\omega(E, F, \Omega)$ the ω -weighted capacity of the pair (E, F) :

$$\text{Cap}_\omega(E, F, \Omega) = \inf_u \int_\Omega |\nabla u(x)|^n dx,$$

where the infimum is taken over all continuous functions $u \in W^{1,n}(\Omega)$ which satisfy $u(x) \leq 0$ for $x \in E$ and $u(x) \geq 1$ for $x \in F$.

Theorem 4.1. *Let $\Omega \subset \mathbb{R}^n$ be a c_0 -John domain with diameter one. Let $E \subset \partial\Omega$ be a compact set satisfying $\mathcal{H}^h(E) > 0$ with $h(t) = \log^{-n-1}(1/t)$. Then*

$$(4.1) \quad \text{Cap}(E, Q_0, \Omega) \geq C \left(\log \frac{1}{\mathcal{H}_\infty^1(E)} \right)^{1-n},$$

where Q_0 is the central Whitney cube containing the John center x_0 . Moreover, if $\omega : \Omega \rightarrow [0, \infty]$ be a weight which is exponentially integrable, then

$$(4.2) \quad \text{Cap}_{1/\omega}(E, Q_0, \Omega) \geq C \left(\log \frac{1}{\mathcal{H}_\infty^1(E)} \right)^{-n}.$$

The constant C in (4.1) and (4.2) depends only on the data associated to Ω and ω .

Proof. We only prove the weighted estimate (4.2), since the proof of the unweighted case is similar (and simpler).

Let $u \in W^{1,n}(\Omega) \cap C(\bar{\Omega})$ be a test function for the $\frac{1}{\omega}$ -weighted n -capacity of the pair (Q_0, E) . For each $x \in E$, we may fix a c_0 -John curve γ joining x to x_0 in Ω and define $P(x)$ to be the collection of Whitney cubes that intersect γ . We define a subcollection $P'(x) \subset P(Q(x))$ as follows: $P'(x) = \{Q_1, Q_2, \dots\}$ consists of a chain of Whitney cubes, which continue along $P(x)$ until it reaches the last cube Q_1 for which $\text{diam } Q_1 \geq \frac{1}{5} \mathcal{H}_\infty^1(E)$. Since adjacent Whitney cubes Q and Q' have $\text{diam } Q \leq 5 \text{diam } Q'$, we must have $\text{diam } Q \leq \mathcal{H}_\infty^1(E)$ for all $Q \in P'(x)$.

We first assume that for every $x \in E$ we have $u_{Q_1} \leq \frac{1}{2}$ for $Q_1 \in P'(x)$. We claim that

$$(4.3) \quad \int_\Omega |\nabla u(x)|^n \frac{1}{\omega(x)} dx \geq C \left(\log \frac{1}{\mathcal{H}_\infty^1(E)} \right)^{-n}.$$

If $x \in E$ and $P'(x) = \{Q_1, Q_2, \dots\}$ then $u(x) = 1$ and thus by Lebesgue differentiation Theorem we have $\lim_{i \rightarrow \infty} u_{Q_i} = 1$. The standard chaining argument involving Poincaré inequality (see for instance [16, 12]) gives us

the estimate

$$(4.4) \quad 1 \lesssim \sum_{Q \in P'(x)} \text{diam } Q \int_Q |\nabla u(y)| dy.$$

Note that for each $Q \subset \Omega$, by Hölder's inequality,

$$\begin{aligned} \int_Q |\nabla u| dy &= \int_Q |\nabla u| \omega^{-\frac{1}{n}} \omega^{\frac{1}{n}} dy \\ &\leq \left(\int_Q |\nabla u|^n \frac{1}{\omega} dy \right)^{\frac{1}{n}} \left(\int_Q \omega^{\frac{1}{n-1}} dy \right)^{\frac{n-1}{n}}. \end{aligned}$$

Notice that the function $t \mapsto \exp(\lambda t^{n-1})$ is convex and we may use Jensen's inequality to deduce that

$$\begin{aligned} \left(\int_Q \omega^{\frac{1}{n-1}} dy \right)^{\frac{n-1}{n}} &\leq \left(\lambda^{-1} \log \left(\int_Q \exp(\lambda \omega) dy \right) \right)^{\frac{1}{n}} \\ &\leq C(\lambda, L) \left(\log \frac{1}{\text{diam } Q} \right)^{\frac{1}{n}}, \end{aligned}$$

where $L = \int_{\Omega} \exp(\lambda \omega) dx$. Plugging the above estimate in (4.4), we obtain

$$(4.5) \quad 1 \leq C(\lambda, n) \sum_{Q \in P'(x)} \left(\int_Q |\nabla u|^n \frac{1}{\omega} dy \right)^{\frac{1}{n}} \log^{\frac{1}{n}} \frac{1}{\text{diam } Q}.$$

By our assumption $\mathcal{H}^h(E) > 0$ and thus we can choose a Frostman measure μ on the compact set E with growth function $h(r) = \log^{-n-1}(1/r)$ i.e. μ is a Borel measure on \mathbb{R}^n which satisfies the growth condition

$$\mu(E \cap B(x, r)) \leq h(r)$$

for every ball $B(x, r)$. Integrating (4.5) over the set E with respect to the Frostman measure μ , we see that

$$\mu(E) \leq C \int_E \sum_{Q \in P'(x)} \log^{\frac{1}{n}} \frac{1}{\text{diam } Q} \left(\int_Q |\nabla u|^n \frac{1}{\omega} dy \right)^{1/n} d\mu(x).$$

We now interchange the order of summation and integration to deduce that

$$\mu(E) \leq C \sum_{\substack{Q \in \mathcal{W} \\ \text{diam } Q \leq \mathcal{H}_{\infty}^1(E)}} \mu(S_E(Q)) \log^{\frac{1}{n}} \frac{1}{\text{diam } Q} \left(\int_Q |\nabla u|^n \frac{1}{\omega} dy \right)^{1/n}.$$

Applying Hölder's inequality to the right hand side leads to the following upper bound for $\mu(E)$

$$\begin{aligned} & \left(\sum_{\substack{Q \in \mathscr{W} \\ \text{diam } Q \leq \mathcal{H}_\infty^1(E)}} \mu(S_E(Q))^{\frac{n}{n-1}} \log^{\frac{1}{n-1}} \frac{1}{\text{diam } Q} \right)^{\frac{n-1}{n}} \left(\sum_{Q \in \mathscr{W}} \int_Q |\nabla u|^n \frac{1}{\omega} dy \right)^{1/n} \\ & \leq \left(\sum_{\substack{Q \in \mathscr{W} \\ \text{diam } Q \leq \mathcal{H}_\infty^1(E)}} \mu(S_E(Q))^{\frac{n}{n-1}} \log^{\frac{1}{n-1}} \frac{1}{\text{diam } Q} \right)^{\frac{n-1}{n}} \left(\int_\Omega |\nabla u|^n \frac{1}{\omega} dy \right)^{1/n}. \end{aligned}$$

The sum in the previous estimate can be dealt with in the following way: Choose $j_0 \in \mathbb{N}$ with $j_0 \geq C \log(1/\mathcal{H}_\infty^1(E))$ so that $\text{diam } Q \leq \mathcal{H}_\infty^1(E)$ implies that $Q \in \mathscr{W}_j$ for some $j \geq j_0$. We claim that

$$(4.6) \quad \sum_{j=j_0}^{\infty} \sum_{Q \in \mathscr{W}_j} \mu(S_E(Q))^{1+\frac{1}{n-1}} \log^{\frac{1}{n-1}} \frac{1}{\text{diam}(Q)} \leq C \mu(E) j_0^{-\frac{n}{n-1}}.$$

The growth condition on μ and Lemmas 3.2 and 3.3 imply that

$$\begin{aligned} & \sum_{j=j_0}^{\infty} \sum_{Q \in \mathscr{W}_j} \mu(S_E(Q))^{1+\frac{1}{n-1}} \log^{\frac{1}{n-1}} \frac{1}{\text{diam}(Q)} \\ & \leq \sum_{j=j_0}^{\infty} \sum_{Q \in \mathscr{W}_j} \mu(S_E(Q)) \mu(S_E(Q))^{\frac{1}{n-1}} \log^{\frac{1}{n-1}} \frac{1}{\text{diam}(Q)} \\ & \leq C \sum_{j=j_0}^{\infty} \sum_{Q \in \mathscr{W}_j} \mu(S_E(Q)) h(\text{diam } Q)^{\frac{1}{n-1}} \log^{\frac{1}{n-1}} \frac{1}{\text{diam}(Q)} \\ & \leq C \sum_{j=j_0}^{\infty} \sum_{Q \in \mathscr{W}_j} \mu(S_E(Q)) \log^{-\frac{n}{n-1}} \frac{1}{\text{diam } Q} \\ & \leq C \sum_{j=j_0}^{\infty} j^{-\frac{n}{n-1}} \sum_{Q \in \mathscr{W}_j} \mu(S_E(Q)) \leq C \mu(E) \sum_{j=j_0}^{\infty} j^{-\frac{n}{n-1}}. \end{aligned}$$

Inequality (4.6) follows from this since the sum on the right hand side converges. Putting together all the estimates above yields

$$\mu(E)^n \leq C \mu(E)^{n-1} \left(\log \frac{1}{\mathcal{H}_\infty^1(E)} \right)^{-n} \int_\Omega |\nabla u|^n \frac{1}{\omega} dy.$$

It follows that

$$\int_\Omega |\nabla u|^n \frac{1}{\omega} \geq C \mu(E) \log^n \frac{1}{\mathcal{H}_\infty^1(E)}.$$

To get (4.3) from here we just note that by concavity of h

$$\mu(E) \geq C\mathcal{H}_\infty^h(E) \geq C \log^{-1-n} \left(\frac{1}{\mathcal{H}_\infty^1(E)} \right)$$

and thus we have

$$\int_\Omega |\nabla u|^n \frac{1}{\omega} \geq C \left(\log \frac{1}{\mathcal{H}_\infty^1(E)} \right)^{-n}.$$

We are left with the case that there is $x \in E$ so that $u_{Q_1} \geq \frac{1}{2}$. In this case, we have $u_{Q_0} = 0$ and $u_{Q_1} \geq \frac{1}{2}$. By following the quasihyperbolic geodesic from x_1 to x_0 , we can find a finite chain of Whitney cubes $\{\tilde{Q}_j\}_{j=0}^l$ which joins Q_0 to Q_f , namely, cubes $\{\tilde{Q}_j\}_{j=1}^l$ such that $\tilde{Q}_0 = Q_0$ is the central cube, $\tilde{Q}_l = Q_1$. Moreover,

$$\text{diam } Q_1 \lesssim \min\{\text{diam } \tilde{Q}_j\} \quad \text{and} \quad l \lesssim \log \frac{1}{\text{diam } Q_1},$$

where the last inequality follows from the fact that the length of a quasihyperbolic geodesic is comparable with the number of Whitney cubes it intersects. Then repeating the chaining argument for finite chain of cubes $\{Q_0, \dots, Q_1\}$ gives the estimate

$$\begin{aligned} 1 &\lesssim \sum_{i=1}^l \text{diam } \tilde{Q}_i \int_{\tilde{Q}_i} |\nabla u| dy \lesssim \sum_{i=1}^l \left(\int_{\tilde{Q}_i} |\nabla u|^n \frac{1}{\omega} dy \right)^{\frac{1}{n}} \left(\log \frac{1}{\text{diam } \tilde{Q}_i} \right)^{\frac{1}{n}} \\ &\lesssim \left(\log \frac{1}{\text{diam } Q_1} \right)^{\frac{1}{n}} \sum_{i=1}^l \left(\int_{\tilde{Q}_i} |\nabla u|^n \frac{1}{\omega} dy \right)^{\frac{1}{n}} \\ &\lesssim \left(\log \frac{1}{\text{diam } Q_1} \right)^{\frac{1}{n}} \left(\sum_{i=1}^l \int_{\tilde{Q}_i} |\nabla u|^n \frac{1}{\omega} dy \right)^{\frac{1}{n}} l^{(n-1)/n}, \end{aligned}$$

from which we obtain

$$\int_\Omega |\nabla u(x)|^n \omega^{-1}(x) dx \geq C \left(\log \frac{1}{\text{diam } Q_1} \right)^{-n} \geq C \left(\log \frac{1}{\mathcal{H}_\infty^1(E)} \right)^{-n}$$

since $\text{diam } Q_1 \leq \mathcal{H}_\infty^1(E)$. □

In Theorem 4.1 we have assumed that Ω is a John domain. However, the proof works for more general domains. More precisely, it holds if Ω satisfies the quasihyperbolic boundary condition [16]. The only difference is that in Theorem 3.5, one has to consider limits along, instead of John curves, quasihyperbolic geodesics to the boundary.

Remark 4.2. Note that, in Theorem 4.1, we have assumed that $\text{diam } \Omega = 1$. This assumption was only used to ensure that $\log \frac{1}{\mathcal{H}_\infty^1(E)}$ is positive. For

domains with different diameter, one just needs to replace the right-hand side of (4.1) with

$$C \left(\log \frac{2 \operatorname{diam} \Omega}{\mathcal{H}_\infty^1(E)} \right)^{1-n},$$

where the constant C in front now depends additionally on Ω as well. Similar changes applies to (4.2).

Another observation one has to notice is that if we replace the central Whitney cube Q_0 by a ball $B_0 = B(x_0, r_0)$ in the capacity estimates in Theorem 4.1, then the estimates (4.1) and (4.2) become

$$(4.7) \quad \operatorname{Cap}(E, B_0, \Omega) \geq C \left(\log \frac{2 \operatorname{diam} \Omega}{\mathcal{H}_\infty^1(E)} \right)^{1-n},$$

and

$$(4.8) \quad \operatorname{Cap}(E, B_0, \Omega) \geq C \left(\log \frac{2 \operatorname{diam} \Omega}{\mathcal{H}_\infty^1(E)} \right)^{1-n},$$

respectively, where the constant C in both (4.7) and (4.8) now depends additionally on Ω and r_0 . Indeed, notice that if $u_{B_0} = 0$, then $u_{Q_0} \leq m_0 = m(r_0) < 1$. So we may repeat the proof of Theorem 4.1 by considering two cases:

- For every $x \in E$, $u_{Q_1} \leq \frac{1+m_0}{2} < 1$ for $Q_1 \in P'(x)$;
- There exists a point $x \in E$ so that $u_{Q_1} \geq \frac{1+m_0}{2} > m_0$.

The proof of the first case proceeds identically as before since $\frac{1+m_0}{2} < 1$. In the second case, since $u_{Q_1} \geq \frac{1+m_0}{2} > m_0$ and $u_{Q_0} \leq m_0$, we may proceed again as in the proof of Theorem 4.1 using the same finite chaining argument.

5. PROOFS OF THEOREMS 1.1 AND 1.2

We are ready to prove the main results. The proof is similar than the one given in [2, Section 2].

Proof of Theorems 1.1 and 1.2. Let Ω be a c_0 uniform domain with center x_0 . We will first prove the case $f \in \mathcal{F}_{\lambda, \kappa}(\Omega)$, the proof for $f \in \mathcal{F}_K(\Omega)$ is analogous. Assume that ν is a positive Borel measure satisfying (1.6) for $s \in (0, n)$. Since $\partial\Omega$ is compact it follows from (1.6) that

$$\nu(\partial\Omega) < \infty.$$

Furthermore, (1.6) together with the finiteness of $\nu(\partial\Omega)$ imply that for any Borel set $E \subset \partial\Omega$ we have

$$(5.1) \quad \nu(E) \leq C_\Omega \mathcal{H}_\infty^s(E).$$

We want to show that there are constants $\alpha > 0$ and $C < \infty$ independent of f such that

$$\int_{\partial\Omega} \exp(\alpha |\bar{f}(w)|) d\nu(w) \leq C,$$

where \bar{f} is the boundary extension given in Definition 3.13. By Fubini's theorem

$$\int_{\partial\Omega} \exp(\alpha|\bar{f}(w)|) d\nu(w) = \nu(\partial\Omega) + \int_1^\infty \nu(F_u) e^{\alpha u} du,$$

where $F_u = \{w \in \partial\Omega : |\bar{f}(w)| \geq u\}$. Thus by (5.1) it is enough to bound

$$\int_1^\infty \mathcal{H}_\infty^s(F_u) e^{\alpha u} du.$$

Define

$$E_t = \{x \in \Omega : |f(x)| = t\}$$

and

$$\mathcal{A}_{n-1}f(E_t) = \int_{S^{n-1}(0,t)} \text{card}f^{-1}(y) d\mathcal{H}^{n-1}(y).$$

Since the functions in $\mathcal{F}_{\lambda,\mathcal{K}}(\Omega)$ are equicontinuous, we can find $r_0 > 0$ such that $|f(x)| \leq 1$ for any $x \in B(x_0, r_0) \subset \Omega$ and $f \in \mathcal{F}_{\lambda,\mathcal{K}}$. Here r_0 depends only on n, λ, \mathcal{K} . Let Γ_{F_u} denote the collection of curves $\gamma: [0, 1] \rightarrow \mathbb{R}^n$ that connect points in $B(x_0, r_0)$ to points in F_u . Then by [22, Proposition 10.2]

$$\text{Cap}_{1/K}(B(x_0, r_0), F_u, \Omega) \leq \text{Mod}_{1/K}(\Gamma_{F_u}).$$

Similarly as in [2, 20] using the test function

$$\rho(x) = \left(\int_1^s \frac{dt}{(\mathcal{A}_{n-1}f(E_t))^{\frac{1}{n-1}}} \right)^{-1} \frac{|Df(x)|}{(\mathcal{A}_{n-1}f(E_t))^{\frac{1}{n-1}}}$$

when $|f(x)| = t \in (1, s)$ and $\rho(x) = 0$ otherwise for the weighted modulus we have the following upper bound

$$\text{Mod}_{1/K}(\Gamma_{F_u}) \leq \left(\int_1^s \frac{dt}{(\mathcal{A}_{n-1}f(E_t))^{\frac{1}{n-1}}} \right)^{n-1}.$$

From Theorem 4.1 we have a lower bound

$$\text{Cap}_{1/K}(B(x_0, r_0), F_s, \Omega) \geq C_2 \left(\log \frac{C_1}{\mathcal{H}_\infty^s(F_s)} \right)^{-n},$$

see also Remark 4.2. These together with the estimate

$$(s-1)^{n/(n-1)} \leq \omega_n^{1/(n-1)} \int_1^s \frac{dt}{(\mathcal{A}_{n-1}f(E_t))^{\frac{1}{n-1}}}$$

following from Hölder's inequality, see [2, 20], imply that

$$\mathcal{H}_\infty^s(F_s) \leq C_1 \exp\left(C_2 \omega_n^{1/n} (s-1)\right),$$

from which we conclude that for $\alpha < C_2 \omega_n^{1/n}$ we have

$$\int_1^\infty \mathcal{H}_\infty^s(E_s) e^{\alpha s} ds \leq C,$$

where the upper bound does not depend on f . □

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