Anomalous Evolution of the Near-Side Jet Peak Shape in Pb-Pb Collisions at $\sqrt{s_{NN}} = 2.76$ TeV


All material supplied via JYX is protected by copyright and other intellectual property rights, and duplication or sale of all or part of any of the repository collections is not permitted, except that material may be duplicated by you for your research use or educational purposes in electronic or print form. You must obtain permission for any other use. Electronic or print copies may not be offered, whether for sale or otherwise to anyone who is not an authorised user.
In elementary interactions with large momentum transfer ($Q^2 \gg \Lambda_{QCD}^2$), partons with high transverse momentum ($p_T$) are produced. They evolve from high to low virtuality producing parton showers and eventually hadronizing into a spray of collimated hadrons called jets. In interactions between heavy ions, such high-$p_T$ partons are produced at the early stages of the collisions. They propagate through the dense and hot medium created in these collisions and are expected to lose energy due to medium-induced gluon radiation and elastic scatterings, a process commonly referred to as jet quenching. Correspondingly, an inclusive jet suppression has been observed at the LHC [1–3] together with a large dijet asymmetry [4,5], while studies of the momentum and angular distributions of jet fragments show only a small modification of the jet core [6–8], and an excess of soft particles radiated to large angles from the jet axis [9]. Semi-inclusive hadron-jet correlations show a suppression of recoil jet yield, with no in-medium modification of transverse jet structure observed [10].

Dihadron angular correlations represent a powerful complementary tool to study jet modifications on a statistical basis in an energy region where jets cannot be identified event by event over the fluctuating background. Such studies involve measuring the distributions of the relative azimuthal angle $\Delta \phi$ and pseudorapidity $\Delta \eta$ between particle pairs consisting of a trigger particle in a certain transverse momentum $p_{T,\text{trig}}$ interval and an associated particle in a $p_{T,\text{assoc}}$ interval. In these correlations, jet production manifests itself as a peak centered around $(\Delta \phi = 0, \Delta \eta = 0)$ (near-side peak) and a structure elongated in $\Delta \eta$ at $\Delta \phi = \pi$ (the away side or recoil region). At low $p_T$, resonance decays as well as femtoscopic correlations also contribute to the near-side peak. The advantage of using dihadron correlations is that an event-averaged subtraction of the background from particles uncorrelated to the jet can be performed. This advantage is shared with the analysis of hadron-jet correlations recently reported in Refs. [9,10].

At RHIC, the near-side particle yield and peak shape of dihadron correlations have been studied for different systems and collision energies [11–13]. Small modifications of the yields with respect to a $pp$ reference from PYTHIA are observed and there is remarkably little dependence on the collision system at the center-of-mass energies of $\sqrt{s_{NN}} = 62.4$ and 200 GeV. An exception is the measurement in central Au-Au collisions at $\sqrt{s_{NN}} = 200$ GeV where the jetlike correlation is substantially broader and the momentum spectrum softer than in peripheral collisions and than those in collisions of other systems in this kinematic regime. In Ref. [12], the broadening observed in central Au-Au collisions at $\sqrt{s_{NN}} = 200$ GeV is seen as an indication of a modified jet fragmentation function. At the LHC, the measurement of the yield of particles associated with a high-$p_T$ trigger particle (8–15 GeV/$c$) in central Pb-Pb collisions
to the $pp$ reference at $p_{T,\text{assoc}} > 3$ GeV/$c$ shows a suppression on the away side and a moderate enhancement on the near side, indicating that medium-induced modifications can also be expected on the near side [14]. Much stronger modification is observed for lower trigger and associated particle $p_T$ ($3 < p_{T,\text{trig}} < 3.5$ GeV/$c$ and $1 < p_{T,\text{assoc}} < 1.5$ GeV/$c$) [15,16]. In the most central Pb-Pb collisions, the near-side yield is enhanced by a factor of 1.7. The present Letter expands these studies at the LHC to the characterization of the angular distribution of the associated particles with respect to the trigger particle. The angular distribution is sensitive to the broadening of the jet due to its energy loss and the distribution of radiated energy. Moreover, possible interactions of the parton shower with the collective longitudinal expansion [17–19] or with turbulent color fields [20] in the medium would result in near-side peak shapes that are broader in the $\Delta \eta$ than in the $\Delta \phi$ direction. Results from the study of the near-side peak shape of charged particles as a function of centrality and for different combinations of trigger and associated particle $p_T$ are discussed.

The data presented in this Letter were taken by the ALICE detector, of which a detailed description can be found in Ref. [21]. The main subsystems used in the present analysis are the Inner Tracking System (ITS), and the Time Projection Chamber (TPC). These have a common acceptance of $|\eta| < 0.9$. The ITS consists of six layers of silicon detectors for vertex finding, tracking, and triggering. The TPC is the main tracking detector measuring up to 159 space points per track. The V0 detector, two arrays of 32 scintillator tiles each, covering $2.8 < \eta < 5.1$ (V0-A) and $-3.7 < \eta < -1.7$ (V0-C), was used for triggering and centrality determination. All these detector systems have full azimuthal coverage.

Data from the 2010 and 2011 Pb-Pb runs of the LHC at $\sqrt{s_{NN}} = 2.76$ TeV are combined in the present analysis and compared with the 2011 $pp$ run at the same energy. In total, about 39 million Pb-Pb and 30 million $pp$ events are used. Details about the trigger and event selection in Pb-Pb ($pp$) collisions can be found in Ref. [22] (Ref. [23]), while the centrality determination is described in Ref. [24].

The collision-vertex position is determined with tracks reconstructed in the ITS and TPC [25], and its value in the beam direction ($z_{\text{vtx}}$) is required to be within 7 cm of the detector center. The Pb-Pb analysis is performed in the centrality classes 0%–10% (most central), 10%–20%, 20%–30%, 30%–50%, and 50%–80%. The analysis uses tracks reconstructed in the ITS and TPC with $1 < p_T < 8$ GeV/$c$ and in a fiducial region of $|\eta| < 0.8$. The track selection is described in Refs. [26,27]. The efficiency and purity of the primary charged-particle selection are estimated from a Monte Carlo (MC) simulation using the HIJING 1.383 event generator [28] (for Pb-Pb) and the PYTHIA 6.4 event generator [29] with the tune Perugia-0 [30] (for $pp$) with particle transport through the detector using GEANT3 [31]. The combined efficiency and acceptance for the track reconstruction in $|\eta| < 0.8$ is about 82%–85% at $p_T = 1$ GeV/$c$, and decreases to about 76%–80% at $p_T = 8$ GeV/$c$ depending on collision system, data sample, and event centrality. The contamination originating from secondary particles from weak decays and interactions in the detector material decreases from 2.5%–4.5% to 0.5%–1% in the $p_T$ range from 1 to 8 GeV/$c$. The contribution from fake tracks is negligible. From these quantities a correction factor is computed as a function of $\eta$, $p_T$, $z_{\text{vtx}}$, and event centrality, which is applied as a weight for each trigger particle and particle pair in the analysis.

The correlation between two charged particles (denoted trigger and associated particle) is measured as a function of $\Delta \phi$ (defined within $-\pi/2$ and $3\pi/2$) and $\Delta \eta$ [32]. The correlation is expressed in terms of the associated yield per trigger particle for intervals of $p_{T,\text{trig}}$ and $p_{T,\text{assoc}}$, measured as

$$\frac{1}{N_{\text{trig}}} \frac{d^2N_{\text{assoc}}}{d\Delta \eta d\Delta \phi} = \frac{S(\Delta \eta, \Delta \phi)}{B(\Delta \eta, \Delta \phi)},$$

where $N_{\text{trig}}$ is the total number of trigger particles in the centrality class and the $p_{T,\text{trig}}$ interval, ranging from 0.18 to 36 per event. The signal distribution $S(\Delta \eta, \Delta \phi)$ is $1/N_{\text{trig}}d^2N_{\text{same}}/d\Delta \eta d\Delta \phi$ is the associated yield per trigger particle for particle pairs from the same event. The background distribution $B(\Delta \eta, \Delta \phi) = \alpha d^2N_{\text{mixed}}/d\Delta \eta d\Delta \phi$ accounts for the acceptance and efficiency of pair reconstruction. It is constructed by correlating the trigger particles in one event with the associated particles from other events. The background distribution is scaled by a factor $\alpha$ which is chosen such that $B(0,0)$ is unity for pairs where both particles travel in approximately the same direction (i.e., $\Delta \phi \approx 0$, $\Delta \eta \approx 0$), and thus the efficiency and acceptance for the two particles are identical by construction.

A selection on the opening angle of the particle pairs is applied to both signal and background to avoid a bias due to the reduced efficiency for pairs with small opening angles. Furthermore, correlations induced by secondary particles from long-lived neutral-particle decays ($K_0^0$ and $\Lambda$) and $\gamma$ conversions are suppressed by rejecting pairs in the corresponding invariant mass region. A correction is performed for a mild $\Delta \eta$ dependence of the structures in the two-particle correlation, which is due to a minor dependence of particle production and anisotropic flow on pseudorapidity. For further details on the analysis procedure, see Ref. [33].

In order to characterize the near-side peak shape, a simultaneous fit of the peak, the combinatorial background, and the long-range correlation background stemming from collective effects is performed. This exploits that in two-particle correlations the near-side peak is centered around $\Delta \phi = 0$, $\Delta \eta = 0$, while long-range correlation structures
are mostly independent of $\Delta \eta$ [34]. The away-side peak, however, is elongated in $\Delta \eta$; therefore, this strategy cannot be applied to studying the away side. The fit function used is a combination of a constant, a generalized two-dimensional Gaussian function and $\cos(n \Delta \phi)$ terms for $n = 2, 3, 4$.

$$F(\Delta \phi, \Delta \eta) = C_1 + \sum_{n=2}^{4} 2V_{n \Delta} \cos(n \Delta \phi)$$

$$+ C_2 \frac{G_{\Delta \phi, w_{\Delta \eta}}(\Delta \phi)G_{\Delta \eta, w_{\Delta \eta}}(\Delta \eta)}{\Delta \phi \Delta \eta},$$

$$G_{\gamma, w}(x) = \frac{\gamma_x}{2w_x \Gamma(1/\gamma_x)} \exp \left[ -\left( \frac{|x|}{w_x} \right)^{\gamma_x} \right].$$

Thus, in Pb-Pb collisions, the background is characterized by four parameters ($C_1$, $V_{n \Delta}$), where $V_{n \Delta}$ are the Fourier components of the long-range correlations [35], and it should be noted that the inclusion of orders higher than four does not significantly change the fit results. In $pp$ collisions, the background consists effectively only of the pedestal $C_1$. The peak magnitude is characterized by $C_2$, and the shape, which is the focus of the present analysis, by four parameters ($\gamma_{\Delta \phi}, w_{\Delta \phi}, \gamma_{\Delta \eta}, w_{\Delta \eta}$). The aim of using this fit function is to allow for a compact description of the data rather than attempting to give a physical meaning to each parameter. Therefore, the variance of $G$ is calculated, which reduces the description of the peak shape to two parameters ($\sigma_{\Delta \phi}$ and $\sigma_{\Delta \eta}$). To describe the evolution of the peak shape from peripheral to central collisions the ratios of the widths in the central bin (0%–10%) and the peripheral bin (50%–80%), denoted by $\sigma_{\Delta \phi}^{CP}$ and $\sigma_{\Delta \eta}^{CP}$, are also calculated.

In the data a depletion around $\Delta \phi = 0$, $\Delta \eta = 0$ is observed at low $p_T$, however, the fit function does not include such a depletion. To avoid a bias on the extracted peak width, some bins in the central region are excluded from the fit. The size of the excluded region varies with $p_T$ and collision centrality (from no exclusion to 0.3). Thus, by definition, the peak width describes the shape of the peak outside of the central region. The depletion in the central region is quantified below by computing the difference between the fit and the per-trigger yield within the exclusion region.

In Pb-Pb collisions, the obtained $\chi^2/n_{\text{dof}}$ values of the fits are in the range 1.0–2.5; most are around 1.5. In the highest two $p_T$ bins (i.e., in $3 < p_{T,\text{assoc}} < 8$ GeV/$c$ and $4 < p_{T,\text{trig}} < 8$ GeV/$c$) the values increase up to about 2.5 showing that at high $p_T$ the peak shape starts to depart from the generalized Gaussian description. In $pp$ collisions, the $\chi^2/n_{\text{dof}}$ values are in the range 1.3–2.0.

Systematic uncertainties connected to the measurement are determined by modifying the event and track selections. In addition, uncertainties related to the cut on pairs with small opening angles and neutral-particle decays, as well as the sensitivity to the pseudorapidity range are considered.

The difference in the extracted parameters is studied as a function of $p_T$, centrality, and collision system, but these dependencies are rather weak and in most cases one uncertainty value can be quoted for each type of systematic uncertainty. Finally, the different sources of systematic uncertainties are added in quadrature. The extracted peak widths are rather insensitive to changes in the selections (total uncertainty of about 2%–4.5%), while the near-side depletion yield is more sensitive (about 24%–45% uncertainty). The contribution from resonance decays was studied by performing the analysis separately for like and unlike sign pairs, and a significant influence on the results presented below was not found.

Figure 1 shows the near-side peak in $1 < p_{T,\text{trig}} < 2$ GeV/$c$ and $1 < p_{T,\text{assoc}} < 2$ GeV/$c$ for the 10% most central collisions. In addition to the two-dimensional representation, projections are shown where the background estimated with Eq. (2) has been subtracted. The near-side peak is asymmetric, i.e., wider in $\Delta \eta$ than in $\Delta \phi$. It is also broader than in peripheral Pb-Pb and $pp$ collisions, where it is mostly symmetric in $\Delta \phi$ and $\Delta \eta$ (not shown, see Ref. [33]). Furthermore, a depletion around $\Delta \phi = 0$, $\Delta \eta = 0$ develops which will be discussed in more detail below. Also at higher $p_T$, the near-side peak is broader in central collisions than in peripheral or $pp$ collisions. This broadening is less pronounced at high $p_T$ than at low $p_T$, and the asymmetry between $\Delta \phi$ and $\Delta \eta$ disappears at the two highest $p_T$ bins; see Ref. [33].

The extracted shape parameters $\sigma_{\Delta \phi}$ and $\sigma_{\Delta \eta}$ are presented in Fig. 2. In $pp$ collisions, the $\sigma$ values range from about 0.14 to about 0.43 showing a $p_T$ dependence qualitatively expected due to the boost of the evolving parton shower; at larger $p_T$ the peak is narrower. In the $\Delta \phi$ direction (left panel) the values obtained in $pp$ collisions are consistent with those in peripheral Pb-Pb collisions. The peak width increases towards central events, which is most pronounced in the lowest $p_T$ bin (20% increase). In the higher $p_T$ bins no significant width increase can be observed. In the $\Delta \eta$ direction (right panel) a much larger broadening is found towards central collisions. Already in
peripheral collisions the width is larger than in \( pp \) collisions. From peripheral to central collisions the width increases further up to \( \sigma_{\Delta\eta} = 0.67 \) in the lowest \( p_T \) bin and the largest relative increase of about 85% is observed for \( 2 < p_T,\text{trig} < 3 \text{ GeV/c} \) and \( 2 < p_T,\text{assoc} < 3 \text{ GeV/c} \). For all but the two largest \( p_T \) bins a significant broadening can be observed. This increase is quantified for all \( p_T \) bins in Fig. 3 by \( \sigma_{\Delta\eta}^{\text{CP}} \) and \( \sigma_{\Delta\phi}^{\text{CP}} \). The increase is quantified with respect to peripheral Pb-Pb instead of \( pp \) collisions to facilitate the MC comparisons discussed below.

In \( pp \) collisions, the peak shows circular symmetry in the \( \Delta\eta-\Delta\phi \) plane for all \( p_T \) values. In Pb-Pb collisions, the peak becomes asymmetric towards central collisions for all but the two highest \( p_T \) bins. The magnitude of this asymmetry depends on \( p_T \) and is largest with about 70% (\( \sigma_{\Delta\eta} > \sigma_{\Delta\phi} \)) in the range \( 2 < p_T,\text{assoc} < 3 \text{ GeV/c} \) and \( 2 < p_T,\text{trig} < 3 \text{ GeV/c} \). These results are compatible with a similar study by the STAR Collaboration at \( \sqrt{s_{NN}} = 200 \text{ GeV} \) [12], which is detailed in the companion paper [33].

In Ref. [17] it was suggested that the interplay of longitudinal flow with a fragmenting high \( p_T \) parton can lead to the observed asymmetric peak shape. The authors argue that hard partons interact with a medium which shows collective behavior, contrary to the simpler picture where the parton propagates through an isotropic medium with respect to the parton direction. In their calculation the scattering centers are Lorentz boosted by applying a momentum shift depending on the collective component transverse to the parton-propagation direction. The calculation in Ref. [17] for Au-Au collisions at \( \sqrt{s_{NN}} = 200 \text{ GeV} \) predicts a 20% increase from peripheral to central events in the \( \Delta\phi \) direction and a 60% increase in the \( \Delta\eta \) direction, which is in good agreement with the measurements by the STAR Collaboration. Despite the different center of mass energy and collision systems, the calculation is in quantitative agreement with the results presented in this Letter as well.

In order to study further the possibility that an interplay of flow and jets can cause the observation, the data is compared to results from a multiphase transport model (AMPT) [36,37]. Two mechanisms in AMPT produce collective effects: partonic and hadronic rescattering. Before partonic rescattering, the initially produced strings may be broken into smaller pieces by the so-called string melting. Three different AMPT settings are considered which have either string melting [configuration (a)] or hadronic rescattering (b) or both activated (c) [38].

The peak widths are extracted from particle-level AMPT simulations in the same way as for the data. None of the AMPT settings provides an accurate description of the measured absolute widths. Further discussion and the corresponding figure can be found in Ref. [33]. In order to provide, nevertheless, a meaningful comparison of the relative increase, \( \sigma_{\Delta\eta}^{\text{CP}} \) and \( \sigma_{\Delta\phi}^{\text{CP}} \) from the models are shown together with the data in Fig. 3. In the \( \Delta\phi \) direction, the setting with string melting deactivated and hadronic rescattering active follows the trend of the data closest. The two other settings show a more uniform distribution across \( p_T \) and only differ in the two lowest \( p_T \) bins. In the \( \Delta\eta \) direction, the setting with string melting deactivated and hadronic rescattering active quite remarkably follows the trend of the data including the large increase for intermediate \( p_T \). The two other settings show qualitatively a similar trend but miss the data quantitatively.

The presented results have focused up to now on the overall shape of the near-side peak. In addition to the broadening, a distinct feature is observed, a depletion around \( \Delta\phi = 0, \Delta\eta = 0 \) (see Fig. 1). An extensive set of studies was carried out to exclude that this depletion could arise from a detector effect. A similar structure is found in AMPT simulations with hadronic rescattering regardless of the string melting setting [33].

In order to quantify this depletion, the difference is computed between the fit (where the depletion region has been excluded, see above) and the per-trigger yield for each \( p_T \) bin. This is normalized by the total peak yield and it is referred to as depletion yield in the following. The region, where effects are expected from the limited two-track reconstruction efficiency (\( |\Delta\phi| < 0.04 \) and \( |\Delta\eta| < 0.05 \), which corresponds to 0.5%–6% of the integrated region), is excluded from this calculation. Figure 4 presents the depletion yield as a function of centrality for the \( p_T \) bins, where it is different from 0. It can be seen that \( (2.2 \pm 0.5)\% \)
and by the AMPT configuration with the largest relative increase of the peak width (Fig. 3) is best described the two configurations (b) and (c) with depletion, and the lowest depletion observed in the AMPT events is present only in the configuration (a) without the depletion has the smallest expansion velocity without hadronic rescattering the depletion yield is con-

The coefficient of the yield is missing in the lowest centrality and with the elliptic flow coefficient \( \beta \) is, as well as to an interplay of partons traversing the longitudinally expanding medium [17]. Thus, a possible scenario is that the presented observations are caused by the interplay of the jet with the collective expansion.

The ALICE Collaboration would like to thank all its engineers and technicians for their invaluable contributions to the construction of the experiment and the CERN accelerator teams for the outstanding performance of the LHC complex. The ALICE Collaboration gratefully acknowledges the resources and support provided by all Grid centers and the Worldwide LHC Computing Grid (WLCG) collaboration. The ALICE Collaboration acknowledges the following funding agencies for their

FIG. 4. Missing yield in the depletion region relative to the overall peak yield extracted from the fit. The arrows indicate the upper limit in case the uncertainty bands touch 0. For compari-

In summary, we have presented a detailed characteriza-

tion of the flow-subtracted near-side peak in two-particle correlations in Pb-Pb collisions at \( \sqrt{s_{NN}} = 2.76 \) TeV together with a measurement in \( pp \) collisions at the same energy. The near-side peak shows several untypical characteristics in Pb-Pb collisions: the peak gets broader and more asymmetric from peripheral towards central collisions over a wide \( p_T \) range, and an unexpected depletion develops in central collisions at low \( p_T \). The broadening is present both in the \( \Delta \eta \) and the \( \Delta \phi \) directions, but it is significantly stronger in the \( \Delta \eta \) direction, leading to the asymmetric shape of the peak. The near-side peak also shows a characteristic \( p_T \) dependence in both Pb-Pb and \( pp \) collisions.

AMPT simulations show also an asymmetric broadening, and the depletion is present when hadronic rescattering is included. The AMPT configuration with hadronic rescattering and without string melting reproduces quantitatively the relative peak broadening as well as the size of the depletion, underlining the importance of the hadronic phase in heavy-ion collisions. The extraction of the radial-

flow expansion velocity suggests that the stronger the radial flow, the stronger the observed effects are. In addition, earlier theoretical and phenomenological work connected the longitudinal broadening of the near-side jetlike peak to strong longitudinal flow in AMPT [41], as well as to an interplay of partons traversing the longitudinally expanding medium [17]. Thus, a possible scenario is that the presented observations are caused by the interplay of the jet with the collective expansion.

The ALICE Collaboration would like to thank all its engineers and technicians for their invaluable contributions to the construction of the experiment and the CERN accelerator teams for the outstanding performance of the LHC complex. The ALICE Collaboration gratefully acknowledges the resources and support provided by all Grid centers and the Worldwide LHC Computing Grid (WLCG) collaboration. The ALICE Collaboration acknowledges the following funding agencies for their
support in building and running the ALICE detector: A. I. Alikhanyan National Science Laboratory (Yerevan Physics Institute) Foundation (ANSL), State Committee of Science and World Federation of Scientists (WFS), Armenia; Austrian Academy of Sciences and Nationalstiftung für Forschung, Technologie und Entwicklung, Austria; Conselho Nacional de Desenvolvimento Científico e Tecnológico (CNPq), Financiadora de Estudos e Projetos (Finep), and Fundação de Amparo à Pesquisa do Estado de São Paulo (FAPESP), Brazil; Ministry of Education of China (MOE of China), Ministry of Science & Technology of China (MOST of China), and National Natural Science Foundation of China (NSFC), China; Ministry of Science, Education and Sport and Croatian Science Foundation, Croatia; Centro de Investigaciones Energéticas, Medioambientales y Tecnológicas (CIEMAT), Cuba; Ministry of Education, Youth and Sports of the Czech Republic, Czech Republic; Danish National Research Foundation (DNRF), The Carlsberg Foundation, and The Danish Council for Independent Research—Natural Sciences, Denmark; Helsinki Institute of Physics (HIP), Finland; Commissariat à l’Énergie Atomique (CEA) and Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), and Centre National de la Recherche Scientifique (CNRS), France; Bundesministerium für Bildung, Wissenschaft, Forschung und Technologie (BMBF), and GSI Helmholtzzentrum für Schwerionenforschung GmbH, Germany; Ministry of Education, Research and Religious Affairs, Greece; National Research, Development and Innovation Office, Hungary; Department of Atomic Energy Government of India (DAE), India; Indonesian Institute of Science, Indonesia; Centro Fermi—Museo Storico della Fisica e Centro Studi e Ricerche Enrico Fermi, and Istituto Nazionale di Fisica Nucleare (INFN), Italy; Institute for Innovative Science and Technology, Nagasaki Institute of Applied Science (IIST), Japan Society for the Promotion of Science (JSPS) KAKENHI and Japanese Ministry of Education, Culture, Sports, Science and Technology (MEXT), Japan; Consejo Nacional de Ciencia (CONACYT) y Tecnología, through Fondo de Cooperación Internacional en Ciencia y Tecnología (FONCYCYT), and Dirección General de Asuntos del Personal Académico (DGAPA), Mexico; Nationaal instituut voor subatomaire fysica (Nikhef), Netherlands; The Research Council of Norway, Norway; Commission on Science and Technology for Sustainable Development in the South (COMSATS), Pakistan; Pontificia Universidad Católica del Perú, Peru; Ministry of Science and Higher Education and National Science Centre, Poland; Ministry of Education and Scientific Research, Institute of Atomic Physics and Romanian National Agency for Science, Technology and Innovation, Romania; Joint Institute for Nuclear Research (JINR), Ministry of Education and Science of the Russian Federation and National Research Centre Kurchatov Institute, Russia; Ministry of Education, Science, Research and Sport of the Slovak Republic, Slovakia; National Research Foundation of South Africa, South Africa; Korea Institute of Science and Technology Information and National Research Foundation of Korea (NRF), South Korea; Centro de Investigaciones Energéticas, Medioambientales y Tecnológicas (CIEMAT) and Ministerio de Ciencia e Innovacion, Spain; Knut & Alice Wallenberg Foundation (KAW) and Swedish Research Council (VR), Sweden; European Organization for Nuclear Research, Switzerland; National Science and Technology Development Agency (NSDTA), Office of the Higher Education Commission under NRU project of Thailand and Suranaree University of Technology (SUT), Thailand; Turkish Atomic Energy Agency (TAEK), Turkey; National Academy of Sciences of Ukraine, Ukraine; Science and Technology Facilities Council (STFC), United Kingdom; National Science Foundation of the U.S. (NSF) and United States Department of Energy, Office of Nuclear Physics (DOE NP), United States.


(ALICE Collaboration)

1A.I. Alikhanyan National Science Laboratory (Yerevan Physics Institute) Foundation, Yerevan, Armenia
2Benemérita Universidad Autónoma de Puebla, Puebla, Mexico
3Bogolyubov Institute for Theoretical Physics, Kiev, Ukraine
4Bose Institute, Department of Physics and Centre for Astroparticle Physics and Space Science (CAPSS), Kolkata, India
5Budker Institute for Nuclear Physics, Novosibirsk, Russia
6California Polytechnic State University, San Luis Obispo, California, USA
7Central China Normal University, Wuhan, China
8Centre de Calcul de l’IN2P3, Villeurbanne, Lyon, France
9Centro de Aplicaciones Tecnológicas y Desarrollo Nuclear (CEADEN), Havana, Cuba
10Centro de Investigaciones Energéticas Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain
11Centro de Investigación y de Estudios Avanzados (CINVESTAV), Mexico City and Mérida, Mexico
12Centro Fermi - Museo Storico della Fisica e Centro Studi e Ricerche “Enrico Fermi”, Rome, Italy
13Chicago State University, Chicago, Illinois, USA
14China Institute of Atomic Energy, Beijing, China
15Commissionariat à l’Energie Atomique, IRFU, Saclay, France
16COMSATS Institute of Information Technology (CIIT), Islamabad, Pakistan

†Deceased.
‡Also at Georgia State University, Atlanta, Georgia, USA.
§Also at Department of Applied Physics, Aligarh Muslim University, Aligarh, India.
∥Also at M.V. Lomonosov Moscow State University, D.V. Skobeltsyn Institute of Nuclear, Physics, Moscow, Russia.