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gA-driven shapes of electron spectra of forbidden β decays in the nuclear shell model

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The evolution of the shape of the electron spectra of 16 forbidden β− decays as a function of gA was studied using the nuclear shell model in appropriate single-particle model spaces with established, well-tested nuclear Hamiltonians. The β spectra of 96Nb(6+ → 96Mo(4+) and 98Tc(6+) → 98Ru(4+)) were found to depend strongly on gA, which makes them excellent candidates for the determination of the effective value of gA with the spectrum-shape method (SSM). A strong gA dependence is also seen in the spectrum of 96Zr(0+) → 96Nb(6+). This decay could be used for determining the quenching of gA in sixth-forbidden decays in the future, when the measurement of the spectrum becomes experimentally feasible. The calculated shell-model electron spectra of the ground-state-to-ground-state decays of 89Rb, 99Tc, and 137Cs and the decay of 137Cs to the isomeric 11/2− state in 137Ba were found to be in excellent agreement with the spectra previously calculated using the microscopic quasiparticle-phonon model. This is further evidence of the robust nature of the SSM observed in the previous studies.

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I. INTRODUCTION

At the nuclear level β decay can be considered as a mutual interaction of the hadronic and leptonic currents mediated by a massive vector boson W± [1]. The leptonic and hadronic currents can be expressed as a mixture of both vector and axial-vector components [2–4]. The weak vector and axial-vector coupling constants gV and gA enter the theory when the hadronic current is renormalized at the nucleon level [5]. The conserved vector-current hypothesis (CVC) and partially conserved axial-vector-current hypothesis (PCAC) yield the free-nucleon values gV = 1.00 and gA = 1.27 [6] but inside nuclear matter the value of gA is affected by many-nucleon correlations and a quenched value might be needed to reproduce experimental data [1]. Precise information on the effective value of gA is crucial when predicting half-lives of neutrinoless double beta decays since the half-lives are proportional to the fourth power of gA [7,8].

The effective value of gA has earlier been probed using a half-life comparison method, in which the predicted and experimental values are compared for different values of gA. This has been done for Gamow-Teller and first-forbidden decays using nuclear matrix elements calculated with the proton-neutron quasiparticle random-phase approximation (pnQRPA) [9–13]. The half-life comparison method predicts that the value of gA is quenched significantly in these transitions. Also in the old studies [14–17] of first-forbidden β decays, strong quenchings were predicted. In a recent shell-model study [18] of r-process waiting-point nuclei, a strong quenching of gA was confirmed. The half-life comparison method could be used to find out if similar quenching of gA is seen in highly forbidden unique beta decays by using the half-lives calculated in Ref. [19], once experimental data become available.

In Ref. [20] the spectrum-shape method (SSM) was introduced as a complementary way to study the effective value of the weak coupling constants. In the SSM the shapes of computed and experimental (normalized) electron spectra are compared in order to find the ratio gA/gV for nonunique forbidden beta β− decays, for which the shape factors depend on gA/gV in a very nontrivial way. In Ref. [21] the SSM was applied to the fourth-forbidden ground-state-to-ground-state transition of 113Cd using the microscopic quasiparticle-phonon model (MQPM) [22,23], the nuclear shell model (NSM) [24–26], the microscopic interacting boson-fermion model (IBFM-2) [27,28], and experimental spectrum of Ref. [29]. The closest match between the theoretical and experimental spectra was found when gA/gV ≈ 0.92 for all three nuclear models. The half-life comparison method, on the other hand, gave very different results when different nuclear models were used. In Ref. [30] it was noticed that the shape of the spectrum was largely unaffected by modest changes in the MQPM wave function, even though the predicted half-life was affected significantly. The observations of these previous studies suggest that the SSM is very robust, quite insensitive to the details of the adopted nuclear mean field and Hamiltonian. It is thus a reliable tool for probing the ratio gA/gV.

The gA-driven evolution of the shapes of several experimentally measurable forbidden nonunique decay spectra in medium-heavy odd-A nuclei was studied using the MQPM in Ref. [30]. Only four good candidates for the application of the spectrum-shape method have been found thus far. These are the ground-state-to-ground-state decays of 87Rb, 99Tc, 113Cd, and 115In [20,30]. In the present paper we set out to find potential new candidates for the application of the SSM in light to medium-heavy nuclei using the nuclear shell model. We also compare the NSM and MQPM β spectra of 87Rb, 99Tc, and 137Cs to see if the model independence observed in Ref. [21] applies for these decays as well. The decays of 137Cs to the 3/2+ ground state and the 11/2− isomeric state of 137Ba are particularly interesting for comparing the differences in the MQPM and NSM spectra, since the MQPM calculations predict a gA-independent spectrum shape for these decays.

This article is organized as follows. In Sec. II we give the theoretical background behind the β spectrum shape. In
Sec. III we describe the application of the nuclear shell model to the spectrum shape of forbidden $\beta^-$ decays. In Sec. IV we present our results and in Sec. V we draw the conclusions.

II. THEORETICAL FORMALISM

We begin the description of the $\beta^-$ decay by making the so-called impulse approximation, in which at the exact moment of decay the decaying nucleon only interacts via the weak interaction and the strong interaction with the other $A-1$ nucleons can be ignored. Since the vector boson $W^-$ has a large mass and thus propagates only a short distance, the flow lines of the nucleons, i.e., the hadronic current, and the flow lines of the emitted leptons, i.e., the leptonic current, can be considered to interact at a pointlike weak-interaction vertex with an effective coupling constant $G_F$, the Fermi constant. The parity nonconserving nature of the weak interaction is reflected in the fact that the hadronic current can be written at the quark level (up $u$ and down $d$ quarks) as a mixture of vector and axial-vector components as [2–4]

$$J_{\mu}^H = \bar{u}(x)\gamma^\mu (1 - \gamma_5) d(x), \quad (1)$$

where $\gamma^\mu$ are the usual Dirac matrices and $\gamma_5 = i\gamma^0 \gamma^1 \gamma^2 \gamma^3$. Renormalization effects of strong interactions must be taken into account when moving from the quark level to the hadron level. The hadronic current at the nucleon level (neutron $n$ and proton $p$) can then be written as [5,6]

$$J_{\mu}^H = \bar{p}(x)\gamma^\mu (g_V - g_A \gamma_5) n(x), \quad (2)$$

where $g_V$ and $g_A$ are the weak vector and axial-vector coupling constants, respectively. The conserved vector-current hypothesis of the standard model (CVC) gives the free-nucleon value $g_V = 1.0$ for the weak vector coupling constant and the partially conserved axial-vector-current hypothesis (PCAC) gives the free-nucleon value $g_A = 1.27$ for the weak axial-vector coupling constant. Inside the nuclear matter the value of $g_A$ is affected by many-nucleon correlations and so the free nucleon value might not be the one to use in practical calculations [6]. The present paper is a step towards solving the problem of what value to use for $g_A$ inside finite nuclei.

In the impulse approximation the probability of the electron being emitted with kinetic energy between $W_e$ and $W_e + dW_e$ is

$$P(W_e)dW_e = \frac{G_F}{(\hbar c)^6} \frac{1}{2\pi^3 \hbar} C(W_e)$$

$$\times p_e c W_e (W_0 - W_e)^2 F_0(Z, W_e) dW_e, \quad (3)$$

where $p_e$ is the momentum of the electron, $Z$ is the proton number, $F_0(Z, W_e)$ is the Fermi function, and $W_0$ is the end-point energy of the $\beta$ spectrum. The nuclear structure information is buried in the shape factor $C(w_e)$.

The half-life of a $\beta$ decay can be written as

$$t_{1/2} = \frac{\kappa}{C}, \quad (4)$$

where $C$ is the integrated shape factor and the constant $\kappa$ has the value [31]

$$\kappa = \frac{2\pi^3 \hbar^7 \ln 2}{m_e^4 c^4 (G_F \cos \theta_C)^2} = 6147 \text{ s}, \quad (5)$$

$\theta_C$ being the Cabibbo angle. In order to simplify the formalism it is usual to introduce unitless kinematic quantities $w_e = W_e/m_e c^2$, $w_0 = W_0/m_e c^2$, and $p = p_e c/(m_e c^2) = \sqrt{w_e^2 - 1}$. With the unitless quantities the integrated shape factor can be expressed as

$$C = \int_0^{w_0} C(w_e) p w_e (w_0 - w_e)^2 F_0(Z, w_e) d w_e. \quad (6)$$

The shape factor $C(w_e)$ of Eq. (6) contains complicated combinations of both (universal) kinematic factors and NMEs. In this paper we use the nuclear shell model to calculate the one-body transition densities related to the NMEs of the shape factor. The details of the shape factor and the constitution of its NMEs can be found from Refs. [32,33]. As in the previous SSM studies [20,21,30] we go beyond the earlier studies [32,33] and take into account the next-to-leading-order

FIG. 1. Normalized electron spectra for the second-forbidden nonunique $\beta^-$ decays of $^{36}$Cl [panel (a)] and $^{60}$Fe [panel (c)], and the fourth-forbidden nonunique decay of $^{50}$V [panel (b)]. The value $g_V = 1.00$ was adopted and the color coding represents the value of $g_A$. 

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terms of the shape function. The details of the next-to-leading order shape factors are discussed in [21].

The shape factor \( C(w_r) \) can be decomposed into vector, axial-vector, and mixed vector-axial-vector parts. In this decomposition the shape factor is

\[
C(w_r) = g_V^2 C_V(w_r) + g_A^2 C_A(w_r) + g_V g_A C_{VA}(w_r). \tag{7}
\]

Integrating Eq. (7) over the electron kinetic energy, we get an analogous expression for the integrated shape factor

\[
\tilde{C} = g_V^2 \tilde{C}_V + g_A^2 \tilde{C}_A + g_V g_A \tilde{C}_{VA}, \tag{8}
\]

where the factors \( \tilde{C}_i \) in Eq. (8) are just constants, independent of the electron energy.

III. ADOPTED MODEL SPACES AND NUCLEAR HAMILTONIANS

The electron spectra of 16 forbidden \( \beta^- \) transitions were calculated for different values of the coupling constant \( g_A \) using the nuclear matrix elements produced by the nuclear shell model. The calculations were done using the shell model code NuShellX@MSU [34], with appropriate model spaces and well established Hamiltonians chosen for each studied decay individually. No additional truncations, beyond those used in the original works, were introduced. The calculations were run on a desktop computer with a 3.3 GHz processor, so the computational burden had to be taken into account when choosing the model space.

The adopted model spaces and Hamiltonians are as follows. \( sd \) shell: for the calculation of the one-body transition densities related to the decay \( ^{36}\text{Cl}(2^{+}) \to ^{36}\text{Ar}(0^{+}) \) we adopted the 0–1s model space and the phenomenological USDB interaction [35]. \( pf \) shell: as in the earlier shell-model studies regarding the half-lives of the transitions \( ^{40}\text{Ca}(0^{+}) \to ^{40}\text{Ca}(2^{+}) \) we adopted the 0d–1s model space and the phenomenological interaction [40,41]. \( 0f \) – \( 0p \) model space (see endnote on Ref. [28] of [42]) was used for the calculation of the \( ^{86}\text{Rb}(3/2^-) \to ^{86}\text{Sr}(9/2^+ \) transition. For the decays of the \( A = 94–99 \) nuclei, the effective interaction of Gloczener [43] in the model space spanning the \( \pi 1p_{1/2}, \pi 0g_{9/2}, \nu 2s_{1/2}, \) and \( 2d_{5/2} \) was used. \( sdgh1_{1/2} \) shell: for the decay transitions \( ^{126}\text{Sn}(0^{+}) \to ^{126}\text{Sb}(2^+ \) and \( ^{137}\text{Cs}(7/2^+ \to ^{137}\text{Ba}(11/2^+ \) the Sn100pn interaction [44] in the \( 0g_{7/2}2s1d0h_{11/2} \) model space was adopted.

IV. RESULTS AND DISCUSSION

Below we present our results: the electron spectra of 16 forbidden \( \beta^- \) decays (Figs. 1–5) and their integrated shape factors (Table I). The electron spectra are discussed in Sec. IV A and the integrated shape factors in Sec. IV B.

A. Electron spectra and the effective value of \( g_A \)

The studied spectra fall into three groups.

GROUP 1. The first group, for which the spectra can be found in Figs. 1 and 2, consists of six transitions for which the SSM analysis has not been done prior to this study. These decays are potentially interesting for the practical application of the spectrum-shape method, since most of them are second-forbidden nonunique transitions with branching ratios above 90%. The only exception is the undetected fourth-forbidden \( \beta^- \) decay branch of \( ^{50}\text{V} \), for which the branching ratio has
FIG. 2. Same as Fig. 1 but for the second-forbidden nonunique decays of $^{94}$Nb [panel (a)], $^{98}$Tc [panel (b)], and $^{126}$Sn [panel (c)].

been predicted to be $\approx 2\%$ by a recent shell-model calculation [36]. An earlier study [30] of the $g_A$-driven evolution of the electron spectra of odd-$A$ nuclei demonstrated that the second and especially the fourth-forbidden nonunique $\beta$-decay spectra depend strongly on the effective value of the weak axial-vector coupling constant.

**GROUP 2.** The second group consists of the fourth-to-sixth-forbidden decays $^{48}$Ca($0^+\rightarrow 4^+-, 5^+, 6^+$) and $^{96}$Zr($0^+\rightarrow 6^+$) $\rightarrow ^{96}$Nb($4^+, 5^+, 6^+$), for which the beta spectra are presented in Figs. 3 and 4. The decays of $^{48}$Ca and $^{96}$Zr appear to be dominated by the two-neutrino double beta decay mode [36,45,46], so the study of these decays is somewhat challenging experimentally. However, they can shed light on whether the spectra of highly forbidden beta decays are $g_A$ dependent or not.

**GROUP 3.** The third group consists of four decays: $^{137}$Cs($7/2^+\rightarrow 137$Ba($11/2^-, 3/2^+$), $^{90}$Tc($9/2^+\rightarrow 9^+$Ru($5/2^+$), and $^{87}$Rb($3/2^-\rightarrow 87$Sn($9/2^+$). The electron spectra of these decays was studied previously in [30] using the MQPM. The spectra of the $^{137}$Cs decays were found to be independent of the value of $g_A$, while the dependence was significant for $^{87}$Rb and $^{99}$Tc. The shell-model and MQPM spectra of these decays are compared in order to see if the previously observed model independence of $^{113}$Cd and $^{115}$In spectra [21] holds also for these decays.

The spectra presented in Figs. 1 and 2 (GROUP 1) fall into two subgroups. The spectra of $^{36}$Cl, $^{50}$V, and $^{60}$Fe [Fig. 1, panels (a), (b), and (c)], and $^{126}$Sn [Fig. 2, panel (c)] do not exhibit noticeable dependence on the value of the axial-vector coupling constant. This is very surprising since the decay of $^{50}$V is fourth forbidden, and all the seven studied fourth-forbidden decays in Ref. [30] depend heavily on $g_A$. The spectra of the second subgroup, $^{94}$Nb and $^{98}$Tc [Fig. 2, panels (a) and (b)], are significantly $g_A$ dependent. The $g_A$ driven evolution of the spectra is not only very similar for these two decays, but also nearly identical to the one of $^{99}$Tc, shown in Fig. 5, panel (c), even though the $Q$ value of these decays varies between 300 keV and 450 keV. A similar grouping phenomenon was seen for several decays in [30]: for example the shapes of the spectra of fourth-forbidden
ground-state-to-ground-state $\beta^-$ decays of $^{97}$Zr, $^{101}$Mo, $^{117}$Cd, and $^{119}$In have an almost identical $g_A$ evolution.

Decays of GROUP 2, $^{48}$Ca($0^+$) → $^{48}$Sc($4^+, 5^+, 6^+$) and $^{96}$Zr($0^+$) → $^{96}$Nb($4^+, 5^+, 6^+$), were studied in order to see if there appears some interesting systematic behavior. Since the shape of the spectra of fourth-forbidden nonunique decays of odd-$A$ nuclei depends on $g_A$ in a significant way, it is interesting to see if this holds for the even-$A$ nuclei. Since the spectrum on $^{50}$V is largely independent of the values of the weak coupling constants, the dependence is not strong for at least some fourth-forbidden decays. The shape of the electron spectra of $^{48}$Ca($0^+$) → $^{48}$Sc($4^+, 5^+, 6^+$) and $^{96}$Zr($0^+$) → $^{96}$Nb($4^+, 5^+, 6^+$) transitions is independent of $g_A$, as shown in Fig. 3, panels (a), (b), and (c) and in Fig. 4, panels (a) and (b), revealing a second exception, $^{48}$Ca($0^+$) → $^{48}$Sc($4^+$), to the $g_A$-dependent fourth-forbidden nonunique decays. The sixth-forbidden decay $^{96}$Zr($0^+$) → $^{96}$Nb($6^+$) [Fig. 4, panel (c)], on the other hand, exhibits a strong $g_A$ dependence. The evolution is very different from the one of the $^{94}$Nb and $^{99}$Tc [Fig. 2, panels (a) and (b)], and $^{99}$Tc [Fig. 5, panel (c)] studied in the present paper and the fourth-forbidden ones studied in [30]. The branching ratio is, however, so small that the measurement of this spectrum is not currently within reach of experiments. It is somewhat striking that, from the 38 different forbidden decays studied in the present paper and in Ref. [30], strong $g_A$ dependence is seen only in even-forbidden decays.

In Fig. 5 the shell-model and MQPM spectra of the experimentally interesting decays found in Ref. [30] are shown for comparison. For $^{87}$Rb [panels (a) and (b)], which is the best candidate for the determination of the quenching of $g_A$ in third-forbidden decays, the MQPM and shell-model spectra are practically identical. For $^{99}$Tc [panels (c) and (d)] the evolution predicted by the two models is also similar, but in the range $g_A = 0.9–1.1$ there is some difference in the shape, leading to a potential small uncertainty in the value of $g_A$ when comparing to the experimental spectrum, hopefully available in the (near) future. The two studied decays of $^{137}$Cs [panels (e) and (f)] were found to be $g_A$ independent when using the MQPM model in Ref. [30]. The shell-model spectra calculated for these decays were not only also $g_A$ independent, but overlap with the MQPM spectra perfectly. Measurements of these spectra would by very desirable, since they would shed light on the quality of the calculated spectra, independent of the chosen nuclear model.

B. Integrated shape factor $\tilde{C}$

The integrated shape factors $\tilde{C}$ and their decomposition to the components of Eq. (8) are presented in Table I. For all the studied decays the sign of the vector and axial-vector components is positive and the sign of the mixed vector-axial-vector component varies. For the second-forbidden nonunique decays of $^{94}$Nb and $^{98, 99}$Tc which have very similar spectrum-shape evolution, the vector- and axial-vector components are roughly equal and the vector-axial-vector component is roughly twice as large, but with a negative sign. The resulting total integrated shape factor is two magnitudes smaller than its components. For the sixth-forbidden decay $^{96}$Zr($0^+$) → $^{96}$Nb($6^+$) the axial-vector component is 30% larger than the vector component and the vector-axial-vector component is again roughly the sum of the two other components and negative in sign.

For $^{87}$Rb the decomposition is similar to the case of the MQPM: the vector component $\tilde{C}_V$ is the largest, the axial-vector component $\tilde{C}_A$ is about 20% of $\tilde{C}_V$, and the vector-axial-vector component $\tilde{C}_{VA}$ is the smallest and has a negative sign. However, the components and the total integrated shape factor are approximately 30% larger for the MQPM. In the case of $^{99}$Tc the difference is much larger and the MQPM results are roughly twice the shell model ones. For the two studied decays of $^{137}$Cs there is no $g_A$ dependence. The decay to the isomeric $11/2^+$ state is unique and thus (practically) $g_A$ independent. On the other hand, the decay to the $3/2^+$ ground state is nonunique but $g_A$ independent since the three components of $\tilde{C}$ have the same sign and no interference between the components occurs. For these two decays the normalized shell model and MQPM electron spectra overlap perfectly but the difference in $\tilde{C}$ is considerable. Again, the relative fractional decomposition is almost identical for the two nuclear models, but the absolute
FIG. 5. Presently calculated shell-model spectra and the MQPM spectra, first published in Ref. [30], for the ground-state-to-ground-state decays of $^{87}$Rb [panels (a) and (b)], $^{99}$Tc [panels (c) and (d)], and $^{137}$Cs [panel (f)], and the decay of $^{137}$Cs to the isomeric $11/2^-$ state in $^{137}$Ba [panel (e)]. The color coding in the spectra of $^{87}$Rb and $^{99}$Tc refers to the value of $g_A$. The shape of the electron spectra of the two $^{137}$Cs decays is $g_A$ independent, and the color coding in these decays refers to the adopted nuclear model. The free-nucleon value $g_V = 1.00$ was used in all calculations.

The electron spectrum of $^{137}$Cs is $40$-fold the shell model one for the decay to the $11/2^-$ state, and $8$-fold for the decay to the $3/2^+$ state. Due to the $\tilde{C}^{-1}$ dependence of the half-life, the half-lives predicted by the two models do not agree with each other at all. This is a strong indication that the spectrum-shape method is a more robust tool for determining the effective value of $g_A$ than just a simple half-life comparison.

V. CONCLUSIONS

The evolution of the electron-spectrum shapes of $16$ forbidden $\beta^-$ decays, driven by the value of the axial-vector coupling constant $g_A$, were studied using nuclear matrix elements derived from the nuclear shell model. Established and well-tested nuclear Hamiltonians were used in these investigations. In the $\beta$-decay shape factors the usually omitted next-to-leading-order terms were taken into account. The main objective of the study was to find experimentally detectable transitions for which the shape of the electron spectra depends sensitively on the value of the weak coupling constants. Comparing the calculated and measured spectra the effective value of $g_A$ can be extracted.

The electron spectra of $^{94}$Nb($6^+$) $\rightarrow$ $^{94}$Mo($4^+$) and $^{98}$Tc($6^+$) $\rightarrow$ $^{98}$Ru($4^+$) depend significantly on the effective value of $g_A$ and these transitions are thus excellent new candidates for the spectrum-shape method (SSM). Both of these decays have almost a $100\%$ branching ratio, and have been experimentally observed. The shell-model spectra of the previously found candidates, the ground-state-to-ground-state decays of $^{87}$Rb and $^{99}$Tc, agree well with the microscopic quasiparticle-phonon model predicted spectra of a previous study, which makes the experimental measurement of these spectra very desirable. The $g_A$-independent shell model and MQPM spectrum shapes of the
$^{137}\text{Cs}(7/2^+) \rightarrow ^{137}\text{Ba}(11/2^-, 3/2^+)$ transitions overlap perfectly. The measurement of these spectra would shed light on the quality of the calculated spectra.

The transitions for which a heavy $g_A$ dependence has been recorded thus far are all even forbidden. Unlike in the case of the fourth-forbidden transitions in odd-$A$ nuclei studied in Ref. [30], which all showed a strong $g_A$ dependence, such a dependence is not seen in the corresponding transitions in the studied even-$A$ nuclei. Thus far all the found transitions with a very strongly $g_A$-dependent spectrum shape are in the mass region $A = 94$–119. Most of the experimentally accessible forbidden nonunique transitions in odd-$A$ nuclei with $A \leq 169$ and second or higher-forbidden nonunique transitions in even-$A$ nuclei with $A \leq 98$ have now been studied. The potential new candidates for the spectrum-shape method are heavier, and the study of these decays requires the adoption of some other nuclear model due to the large computational burden and the other problems with shell-model calculations for heavy deformed nuclei.

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