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On a topology optimization problem governed by two-dimensional Helmholtz equation

Jaroslav Haslinger & Raino A. E. Mäkinen

Abstract *The paper deals with a class of shape/topology optimization problems governed by the Helmholtz equation in 2D. To guarantee the existence of minimizers, the relaxation is necessary. Two numerical methods for solving such problems are proposed and theoretically justified: a direct discretization of the relaxed formulation and a level set parametrization of shapes by means of radial basis functions. Numerical experiments are given.*

Keywords Helmholtz equation, topology optimization, level set method, radial basis functions

1 Introduction

Helmholtz equation is used in modelling e.g. polarized electromagnetic waves in two dimensions or acoustic waves in two and three dimensions. Applications involving shape and topology optimization governed by the Helmholtz or Maxwell equation include among others: stealth technology [15], noise reduction [1], enhancing performance of acoustic devices (horn-lens combinations) [24], anti-reflection surfaces [9], and optimal distribution of conductivity minimizing the dissipative energy [11]. Further, in combination with the homogenization approach it can be used to improve of

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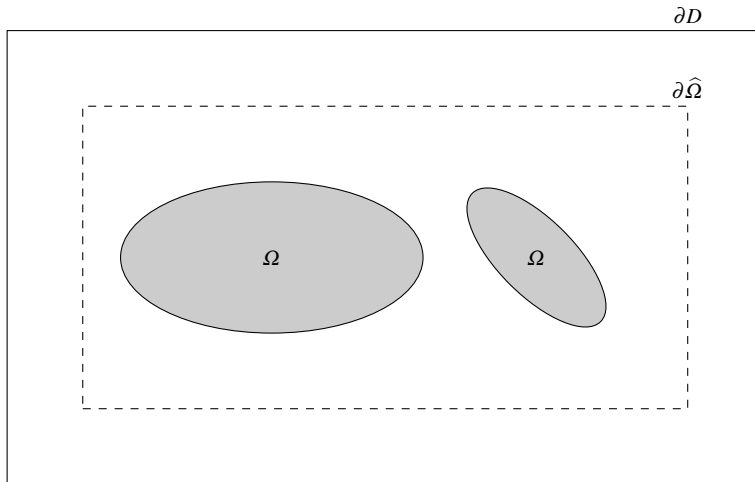


Fig. 1 Schematic representation of the computational geometry.

metamaterial properties in different wave propagation problems (e.g. electromagnetic waveguides, band gap structures using photonic crystals) [12], [19], [17], [7], [14], [21], and reconstruction of electromagnetic inclusions by boundary measurements [4].

In this paper we restrict ourselves to topology optimization problems governed by the scalar time-harmonic wave equation, i.e. the Helmholtz equation. Theoretical studies of the topic are rare. Usually they are proposed by the inverse problems community. For example Mark Kac asked in article [13] “Can one hear the shape of a drum?”. Recently, Bao et al. [3], have proved some useful results applicable also in studies of certain topology optimization problems.

We formulate a model topology optimization problem that is motivated by the above mentioned papers dealing with electromagnetic waveguides and band gap structures. We assume that there exists an obstacle in the free space that has different (but constant) relative permittivity compared to that of the free space. When the incoming wave travels between media of different relative permittivity it may both reflect and refract off a material interface. Consequently, in some areas the wave field is amplified or damped. This phenomenon can be utilized to prevent the incoming wave from entering some regions or propagating in certain directions. Our model problem, however, does not directly correspond to any concrete engineering design problem as we want to keep the formulation compatible with the one presented in [3].

As we have mentioned, a numerical realization will be based on an optimal control approach. The state problem will be given by the two-dimensional Helmholtz equation in which the relative permittivity ε_r plays the role of the control variable and has the following form: $\varepsilon_r = \varepsilon_{r,0} + (\varepsilon_{r,1} - \varepsilon_{r,0})\chi_\Omega$, where $\varepsilon_{r,0}, \varepsilon_{r,1}$ is the constant relative permittivity of the free space, and the scatterer, respectively, and χ_Ω stands for the characteristic function of a measurable set Ω which represents the scatterer. The choice of a cost functional J depends on particular optimization goals. Our aim

is to find the distribution of $\varepsilon_{r,0}$ and $\varepsilon_{r,1}$ in such a way that J attains its minimum. It is well-known that this class of optimization problems has no solution, in general, i.e. no minimizer of J over the set the admissible ε_r introduced above exists. For this reason an extension of this set is necessary. Since ε_r appears in the lower order term of the Helmholtz equation, the respective extension is represented by all measurable functions ranging between $\varepsilon_{r,0}$ and $\varepsilon_{r,1}$. A discretization of this new relaxed formulation can be directly used for solving the problem. Besides this, we propose another approach in which the parametrization of the shape of the scatter is done by a level set method involving radial basis functions. The advantage of this approach is the fact that shape/topology optimization can be transformed into a parametric optimization problem.

The paper is organized as follows. In Section 2 we introduce the physical setting of the state problem. In Section 3 we present a class of shape optimization problems governed by the Helmholtz equation and derive their relaxed form. Section 4 is devoted to a discretization and convergence analysis of both approaches mentioned above. In Section 5 implementation aspects are discussed. Finally, numerical results of two model examples are shown in Section 6.

Throughout of the paper we use the following notation: if $Q \subset \mathbb{R}^s$ is a domain then $H^k(Q)$, $k \geq 0$ integer, denotes the space of all functions which are together with their generalized derivative up to order k square integrable in Q ($L^2(Q) = H^0(Q)$). The norm, seminorm in $H^k(Q)$ will be denoted by $\|\cdot\|_{k,Q}$, and $|\cdot|_{k,Q}$, respectively. $L^\infty(Q)$ stands for the space of all bounded measurable functions in Q . Finally, $\Re f$, $\Im f$, stands for the real, and imaginary part of a complex valued function f , respectively and \bar{f} for its conjugate.

2 The state problem

Let us consider electromagnetic wave propagation in a dielectric, non-magnetic three-dimensional infinite medium where no charges and currents exist. Classical Maxwell's equations can be reduced to the vector wave equation for the electric field $E = (E_1, E_2, E_3)$ as follows:

$$\mu_0 \varepsilon_0 \varepsilon_r \frac{\partial^2 E}{\partial t^2} - \nabla^2 E = 0, \quad (1)$$

where ε_r is the relative permittivity of the material and ε_0 and μ_0 are the permittivity and permeability of the free space, respectively.

If we assume the time-harmonic transverse electric (TE) mode, the x_3 -component of the electric field is of the form

$$E_3(x_1, x_2, x_3, t) = \Re \left\{ u_{tot}(x_1, x_2) e^{-i\omega t} \right\}, \quad (2)$$

where u_{tot} is the complex amplitude of the wave and ω is the angular frequency.

From the above assumptions it follows that (1) reduces to the scalar Helmholtz equation

$$-\Delta u_{tot} - k_0^2 \varepsilon_r u_{tot} = 0 \quad \text{in } \mathbb{R}^2, \quad (3)$$

where the wave number $k_0 = \omega \sqrt{\varepsilon_0 \mu_0}$.

Let $\varepsilon_{r,0}$ be the relative permittivity of the infinite medium. A scatterer having relative permittivity $\varepsilon_{r,1} > \varepsilon_{r,0}$ is represented by a possibly multiply connected set Ω . Next we assume that the scatterer is contained in a domain $\widehat{\Omega}$.

Let us introduce a computational domain $D =]-a, a[\times]-b, b[$, $a, b > 0$ containing $\widehat{\Omega}$ as shown in Figure 1. The relative permittivity can be now represented in \mathbb{R}^2 as

$$\varepsilon_r(x) = \varepsilon_{r,0} + (\varepsilon_{r,1} - \varepsilon_{r,0})\chi_\Omega(x), \quad (4)$$

where χ_Ω is the characteristic function of Ω . Without loss of generality, we may assume that $\varepsilon_{r,0} = 1$ so that $\varepsilon_r = 1 + q_1\chi_\Omega =: 1 + q$ with $q_1 = \varepsilon_{r,1} - \varepsilon_{r,0} > 0$ and $q \in L^\infty(\widehat{\Omega})$. For purely technical reasons and to simplify notations we will assume that $q_1 = 1$ in Sections 3 and 4.

An incident plane wave $u_0 = e^{ik_0d \cdot x}$ propagating into the direction $d \in \mathbb{R}^2$, $\|d\| = 1$ is scattered from the obstacle. The total electric field u_{tot} can be split into the incident field u_0 and the reflected field u :

$$u_{tot} = u_0 + u. \quad (5)$$

As $-\Delta u_0 - k_0^2 u_0 = 0$ it follows that u satisfies

$$-\Delta u - k_0^2(1+q)u = k_0^2 q u_0 \quad \text{in } \mathbb{R}^2. \quad (6)$$

The scattered field u should not be reflected back from the artificial boundary ∂D . Therefore, as usual, we impose an absorbing boundary condition as an approximation to the (physically correct) Sommerfeld radiation condition

$$\sqrt{r} \left(\frac{\partial u}{\partial r} - ik_0 u \right) \xrightarrow{r \rightarrow \infty} 0, \quad r = |x|. \quad (7)$$

In the theoretical analysis of the problem we employ the first-order approximation of (7) as the boundary condition on the artificial boundary ∂D :

$$\frac{\partial u}{\partial n} - ik_0 u = 0 \quad \text{on } \partial D. \quad (8)$$

To give a weak formulation of (6) and (8), we introduce the bilinear form $a : H^1(D) \times H^1(D) \rightarrow \mathbb{C}$

$$a(u, v) = (\nabla u, \nabla v)_{0,D} - k_0^2((1+q)u, v)_{0,D} - ik_0(u, v)_{0,\partial D}$$

and the linear functional $b : H^1(D) \rightarrow \mathbb{C}$

$$b(v) = k_0^2(qu_0, v)_{0,D}.$$

Here $H^1(D)$ stands for the Sobolev space of complex valued functions, $(z, w)_D = \int_D z \bar{w} dx$, and $(z, w)_{0,\partial D} = \int_{\partial D} z \bar{w} ds$.

The *weak formulation* of the state problem reads:

$$\text{Find } u \in H^1(D) : \quad a(u, v) = b(v) \quad \forall v \in H^1(D). \quad (9)$$

3 Setting of the optimization problem

The two dimensional Helmholtz equation introduced in the previous section will be used as the state problem and the relative permeability as the control variable in a class of optimization problems.

Recall the classical formulation of the state problem: Find a scattered field $u : D \rightarrow \mathbb{C}$ solving the following boundary value problem:

$$\begin{cases} -\Delta u - k_0^2(1 + \chi_\Omega)u = k_0^2\chi_\Omega u_0 & \text{in } D \\ \frac{\partial u}{\partial n} - ik_0u = 0 & \text{on } \partial D, \end{cases} \quad (10)$$

where χ_Ω is the characteristic function of a measurable set $\Omega \subset \widehat{\Omega}$. The system of all such sets will be denoted by \mathcal{O} and it will be identified with

$$\mathcal{U} = \{q \in L^\infty(D) \mid q = \chi_\Omega \text{ in } D, \Omega \in \mathcal{O}\}. \quad (11)$$

The weak formulation of (10) reads as follows: given $q \in \mathcal{U}$,

$$\begin{cases} \text{Find } u := u(q) \in H^1(D) \text{ such that} \\ (\nabla u, \nabla v)_{0,D} - k_0^2((1+q)u, v)_{0,D} - ik_0(u, v)_{0,\partial D} = k_0^2(qu_0, v)_{0,D} \quad \forall v \in H^1(D). \end{cases} \quad (\mathcal{P}(q))$$

Let $J : H^1(D) \rightarrow \mathbb{R}$ be a functional and $\mathcal{J} : L^\infty(D) \rightarrow \mathbb{R}$ be defined by

$$\mathcal{J}(q) = J(u(q)) + \alpha \left(\int_{\widehat{\Omega}} q dx - \gamma \right)^+, \quad q \in \mathcal{U}, \quad (12)$$

where $u(q)$ is a solution to $(\mathcal{P}(q))$, $q \in \mathcal{U}$ and $\alpha \geq 0, \gamma > 0$ are given positive parameters.

\mathcal{J} will play the role of the objective function in the following problem:

$$\inf_{q \in \mathcal{U}} \mathcal{J}(q). \quad (\mathbb{P})$$

Remark 1 If $\alpha > 0$, then the last term in (12) can be interpreted as a penalty functional related to the constraint $\text{meas}\Omega \leq \gamma$.

As we shall see, problem (\mathbb{P}) has no solution, in general. For this reason we shall introduce its relaxed form. The set \mathcal{U} will be extended to the larger set $\mathcal{U}^\#$, where

$$\mathcal{U}^\# = \left\{ q \in L^\infty(D) \mid 0 \leq q \leq 1 \text{ a.e. in } \widehat{\Omega}, q \equiv 0 \text{ in } D \setminus \widehat{\Omega} \right\} \quad (13)$$

and the state problem $(\mathcal{P}(q))$ will be defined also for $q \in \mathcal{U}^\#$.

The existence and uniqueness of a solution to $(\mathcal{P}(q))$, $q \in \mathcal{U}^\#$ has been established in [3]. In addition, the estimate

$$\|u\|_{H^1(D)} \leq C \|q\|_{L^\infty(D)} \|u_0\|_{L^2(D)} \quad (14)$$

holds, where the constant $C > 0$ depends only on k_0 and D .

The relaxed form of (\mathbb{P}) is given by

$$\inf_{q \in \mathcal{U}^\#} \mathcal{J}(q). \quad (\mathbb{P}^\#)$$

Our aim will be to show that $(\mathbb{P}^\#)$ has a solution, i.e. a minimizer of \mathcal{J} on $\mathcal{U}^\#$ exists and, in addition, $\min_{q \in \mathcal{U}^\#} \mathcal{J}(q) = \inf_{q \in \mathcal{U}} \mathcal{J}(q)$. To this end we shall need the following continuity assumption on J :

$$y_n \xrightarrow[n \rightarrow \infty]{} y \text{ in } H^1(D) \implies J(y_n) \xrightarrow[n \rightarrow \infty]{} J(y). \quad (15)$$

We start with an auxiliary stability result.

Proposition 1 *Let $q_n \xrightarrow[n \rightarrow \infty]{} q$ weakly * in $L^\infty(D)$, $q_n, q \in \mathcal{U}^\#$ and $u_n := u(q_n)$ be the solution to $(\mathcal{P}(q_n))$, $n = 1, 2, \dots$. Then*

$$u_n \xrightarrow[n \rightarrow \infty]{} u \text{ in } H^1(D)$$

and $u := u(q)$ solves $(\mathcal{P}(q))$.

Proof: The sequence $\{u_n\}$ is bounded in $H^1(D)$ as follows from (14). Thus, one can find a subsequence $\{u_{n_j}\}$ such that

$$u_{n_j} \xrightarrow[j \rightarrow \infty]{} u \text{ in } H^1(D) \quad (16)$$

and, hence

$$u_{n_j} \xrightarrow[j \rightarrow \infty]{} u \text{ in } L^2(D). \quad (17)$$

Passing to the limit with $j \rightarrow \infty$ in $(\mathcal{P}(q_{n_j}))$ and using (16), (17) we see that u solves $(\mathcal{P}(q))$. Since $(\mathcal{P}(q))$ has a unique solution, (16) and (17) hold for the whole sequence $\{u_n\}$. To prove strong convergence we proceed as follows:

$$\begin{aligned} |u_n|_{1,D}^2 &= k_0^2((1+q_n)u_n, u_n)_{0,D} + ik_0(u_n, u_n)_{0,\partial D} + k_0^2(q_n u_0, u_n)_{0,D} \\ &\xrightarrow[n \rightarrow \infty]{} k_0^2((1+q)u, u)_{0,D} + ik_0(u, u)_{0,\partial D} + k_0^2(q u_0, u)_{0,D} = |u|_{1,D}^2 \end{aligned}$$

making use of (16), (17) and the definition of $(\mathcal{P}(q_n))$, $(\mathcal{P}(q))$. \square

Theorem 1 *Problem $(\mathbb{P}^\#)$ has a solution.*

Proof: Let $\{q_n\}$, $q_n \in \mathcal{U}^\#$ be a minimizing sequence:

$$l = \inf_{q \in \mathcal{U}^\#} \mathcal{J}(q) = \lim_{n \rightarrow \infty} \mathcal{J}(q_n). \quad (18)$$

There exists a subsequence $\{q_{n_j}\}$ and $q^* \in \mathcal{U}^\#$ such that

$$q_{n_j} \xrightarrow[j \rightarrow \infty]{} q^* \text{ weakly * in } L^\infty(D) \quad (19)$$

and

$$u(q_{n_j}) \xrightarrow[j \rightarrow \infty]{} u(q^*) \text{ in } H^1(D) \quad (20)$$

by Proposition 1. From (15), (19), and (20) we see that $l = \mathcal{J}(q^*)$. \square

If $q^* \in \mathcal{U}^\#$ is a solution of $(\mathbb{P}^\#)$ and $u^* \in H^1(D)$ the respective solution of $(\mathcal{P}(q^*))$ then (q^*, u^*) will be termed an *optimal pair* of $(\mathbb{P}^\#)$.

Now we introduce an auxiliary problem with a piecewise constant approximation of $\mathcal{U}^\#$. Let $\{\mathcal{T}_\kappa\}$, $\kappa \rightarrow 0+$ be a family of partitions of $\widehat{\Omega}$. Each \mathcal{T}_κ consists of a finite number of mutually disjoint subdomains $G_i \subset \widehat{\Omega}$, $i = 1, \dots, n(\kappa)$:

$$\widehat{\Omega} = \bigcup_{i=1}^{n(\kappa)} \overline{G_i}, \quad \text{diam } G_i \leq \kappa \quad \forall G_i \in \mathcal{T}_\kappa.$$

With any \mathcal{T}_κ we associate the set

$$\mathcal{U}_\kappa^\# = \{q_\kappa \in L^\infty(D) \mid q_\kappa|_{G_i} \in P_0(G_i) \quad \forall G_i \in \mathcal{T}_\kappa\} \cap \mathcal{U}^\#$$

and define the problem

$$\inf_{q_\kappa \in \mathcal{U}_\kappa^\#} \mathcal{J}(q_\kappa). \quad (\mathbb{P}_\kappa^\#)$$

We easily obtain the following existence result.

Theorem 2 *Problem $(\mathbb{P}_\kappa^\#)$ has a solution for any $\kappa > 0$.*

Proof: is parallel to the one of Theorem 1. \square

Next we establish a relation between solutions to $(\mathbb{P}^\#)$ and $(\mathbb{P}_\kappa^\#)$ when $\kappa \rightarrow 0+$.

Theorem 3 *Let $\{(q_\kappa^*, u_\kappa^*)\}$ be a sequence of optimal pairs of $(\mathbb{P}_\kappa^\#)$, $\kappa \rightarrow 0+$. Then there exist: a subsequence $\{(q_{\kappa_j}^*, u_{\kappa_j}^*)\}$ and a pair $(q^*, u^*) \in \mathcal{U}^\# \times H^1(D)$ such that*

$$\begin{cases} q_{\kappa_j}^* \rightharpoonup q^* \text{ weakly } * \text{ in } L^\infty(D) \\ u_{\kappa_j}^* \rightarrow u^* \text{ in } H^1(D). \end{cases} \quad (21)$$

Moreover, (q^, u^*) is an optimal pair of $(\mathbb{P}^\#)$. Any accumulation point of $\{(q_\kappa^*, u_\kappa^*)\}$ in the sense of (21) has this property.*

Proof: The existence of a subsequence satisfying (21) follows from the definition of $\mathcal{U}^\#$ and Proposition 1. To verify that (q^*, u^*) is an optimal pair of $(\mathbb{P}^\#)$ we use the fact that the system $\{\mathcal{U}_\kappa^\#\}$, $\kappa \rightarrow 0+$ is dense in $\mathcal{U}^\#$ in the weak * topology: for any $\bar{q} \in \mathcal{U}^\#$ given there exists a sequence $\{\bar{q}_\kappa\}$, $\bar{q}_\kappa \in \mathcal{U}_\kappa^\#$ such that

$$\bar{q}_\kappa \rightharpoonup \bar{q} \text{ weakly } * \text{ in } L^\infty(D) \quad (22)$$

and

$$u(\bar{q}_\kappa) \rightarrow u(\bar{q}) \text{ in } H^1(D) \quad (23)$$

by Proposition 1. From the definition of $(\mathbb{P}_{\kappa_j}^\#)$, where κ_j is the filter of indices from (21) it follows that

$$\mathcal{J}(q_{\kappa_j}^*) \leq \mathcal{J}(\bar{q}_{\kappa_j}) \quad \forall j.$$

Letting $j \rightarrow \infty$ and using (15), (21), (22), and (23) we get

$$\mathcal{J}(q^*) \leq \mathcal{J}(\bar{q}) \quad \forall \bar{q} \in \mathcal{U}^\#.$$

□

Since

$$\min_{q \in \mathcal{U}^\#} \mathcal{J}(q) \leq \min_{q_\kappa \in \mathcal{U}_\kappa^\#} \mathcal{J}(q_\kappa) \quad \forall \kappa > 0,$$

it follows from (21) that

$$\min_{q \in \mathcal{U}^\#} \mathcal{J}(q) = \liminf_{\kappa \rightarrow 0^+} \left(\min_{q_\kappa \in \mathcal{U}_\kappa^\#} \mathcal{J}(q_\kappa) \right)$$

or briefly

$$(\mathbb{P}^\#) = \liminf_{\kappa \rightarrow 0^+} (\mathbb{P}_\kappa^\#). \quad (24)$$

Any $q_\kappa \in \mathcal{U}_\kappa^\#$, $\kappa > 0$ fixed, being piecewise constant on \mathcal{T}_κ can be attained as the weak * limit of locally periodic, rapidly oscillating functions. We briefly describe the construction of such functions.

Let $Y =]0, 1[\times]0, 1[$ be the unit periodic cell and denote by \mathcal{O}_Y the system of all measurable sets $\omega \subset Y$. The characteristic function χ_ω of $\omega \in \mathcal{O}_Y$ will be periodically extended from Y on \mathbb{R}^2 . For any $\varepsilon > 0$ we define the scaled characteristic function $\chi_\omega^\varepsilon(x) := \chi_\omega(x/\varepsilon)$, $x \in \mathbb{R}^2$ and the set

$$\mathcal{U}_\kappa^\varepsilon = \left\{ q_\kappa^\varepsilon \in L^\infty(D) \mid q_\kappa^\varepsilon|_{G_i} = \chi_{\omega_i}^\varepsilon|_{G_i}, \omega_i \in \mathcal{O}_Y \forall G_i \in \mathcal{T}_\kappa \text{ and } q_\kappa^\varepsilon \equiv 0 \text{ in } D \setminus \overline{\Omega} \right\},$$

i.e. $q_\kappa^\varepsilon \in \mathcal{U}_\kappa^\varepsilon$ if and only if there exist sets $\omega_i \in \mathcal{O}_Y$ such that $q_\kappa^\varepsilon|_{G_i}$ is the scaled characteristic function of ω_i on G_i , $i = 1, \dots, n(\kappa)$.

Let $q_\kappa \in \mathcal{U}_\kappa^\#$ be given and denote $q_\kappa^i := q_\kappa|_{G_i}$, $i = 1, \dots, n(\kappa)$. It is well known (see [6]) that if $\omega_i \in \mathcal{O}_Y$ is such that $\text{meas } \omega_i = q_\kappa^i$ then

$$\chi_{\omega_i}^\varepsilon|_{G_i} \xrightarrow{\varepsilon \rightarrow 0^+} q_\kappa^i \text{ weakly * in } L^\infty(G_i) \quad \forall G_i \in \mathcal{T}_\kappa.$$

Consequently,

$$q_\kappa^\varepsilon \xrightarrow{\varepsilon \rightarrow 0^+} q_\kappa \text{ weakly * in } L^\infty(D),$$

where $q_\kappa^\varepsilon|_{G_i} = \chi_{\omega_i}^\varepsilon|_{G_i}$, $G_i \in \mathcal{T}_\kappa$.

For any $\kappa, \varepsilon > 0$ fixed, define the problem

$$\inf_{q_\kappa^\varepsilon \in \mathcal{U}_\kappa^\varepsilon} \mathcal{J}(q_\kappa^\varepsilon). \quad (\mathbb{P}_\kappa^\varepsilon)$$

Proposition 2 For any $\kappa > 0$ it holds:

$$(\mathbb{P}_\kappa^\#) = \liminf_{\varepsilon \rightarrow 0^+} (\mathbb{P}_\kappa^\varepsilon).$$

Proof: Clearly

$$\min_{q_\kappa \in \mathcal{U}_\kappa^\#} \mathcal{J}(q_\kappa) \leq \inf_{q_\kappa^\varepsilon \in \mathcal{U}_\kappa^\varepsilon} \mathcal{J}(q_\kappa^\varepsilon) \quad \forall \varepsilon > 0. \quad (25)$$

Let q_κ^* be a solution to $(\mathbb{P}_\kappa^\#)$. Using the approach described above, one can find a sequence $\{\bar{q}_\kappa^\varepsilon\}, \bar{q}_\kappa^\varepsilon \in \mathcal{U}_\kappa^\varepsilon$ such that

$$\bar{q}_\kappa^\varepsilon \rightharpoonup_{\varepsilon \rightarrow 0^+} q_\kappa^* \text{ weakly } * \text{ in } L^\infty(D)$$

and

$$u(\bar{q}_\kappa^\varepsilon) \rightarrow_{\varepsilon \rightarrow 0^+} u(q_\kappa^*) \text{ in } H^1(D).$$

This and (15) entails the equality in (25).

Now we are ready to prove the main result of this section.

Theorem 4 *It holds:*

$$\min_{q \in \mathcal{U}^\#} \mathcal{J}(q) = \inf_{q \in \mathcal{U}} \mathcal{J}(q). \quad (26)$$

Proof: From (24) and Proposition 2 we have

$$(\mathbb{P}^\#) = \liminf_{\substack{\kappa \rightarrow 0^+ \\ \varepsilon \rightarrow 0^+}} (\mathbb{P}_\kappa^\varepsilon).$$

On the other hand

$$\inf_{q \in \mathcal{U}} \mathcal{J}(q) \leq \inf_{q_\kappa^\varepsilon \in \mathcal{U}_\kappa^\varepsilon} \mathcal{J}(q_\kappa^\varepsilon) \quad \forall \varepsilon, \kappa > 0$$

so that

$$\inf_{q \in \mathcal{U}} \mathcal{J}(q) \leq \liminf_{\substack{\kappa \rightarrow 0^+ \\ \varepsilon \rightarrow 0^+}} \mathcal{J}(q_\kappa^\varepsilon) = \min_{q \in \mathcal{U}^\#} \mathcal{J}(q).$$

Since the opposite inequality is automatically satisfied we proved (26). \square

4 Discretization of $(\mathbb{P}^\#)$. Convergence analysis

This section deals with an approximation of $(\mathbb{P}^\#)$. We present two ways: *i*) a direct discretization of $(\mathbb{P}^\#)$ *ii*) discretization based on the level set approach parametrized by radial basis functions.

4.1 Direct discretization of $(\mathbb{P}^\#)$

In Section 3 we introduced auxiliary problem $(\mathbb{P}_\kappa^\#)$ which is already a partial discretization of $(\mathbb{P}^\#)$. The admissible set $\mathcal{U}^\#$ in $(\mathbb{P}^\#)$ was replaced by $\mathcal{U}_\kappa^\#$ which consists of piecewise constant functions on \mathcal{T}_κ . In what follows we shall study the full discretization of $(\mathbb{P}^\#)$.

Let $\{\mathcal{T}_h\}$, $h \rightarrow 0+$ be a regular system of triangulations of \bar{D} which satisfies the standard assumptions on the mutual position of triangles $T \in \mathcal{T}_h$. Although \mathcal{T}_h , \mathcal{T}_κ are two different partitions of \bar{D} and $\bar{\Omega}$, respectively, they will not be entirely independent. Next we shall suppose that each $G_i \in \mathcal{T}_\kappa$ is *polygonal* and $\mathcal{T}_h|_{G_i}$ defines a standard triangulation of \bar{G}_i , $i = 1, \dots, n(\kappa)$. Moreover we shall suppose that the norm κ of \mathcal{T}_κ is a decreasing function of h and $\kappa(h) \rightarrow 0+ \iff h \rightarrow 0+$. With any \mathcal{T}_h we associate the space of piecewise affine complex valued functions V_h :

$$V_h = \left\{ v_h \in C(\bar{D}) \mid v_h|_T \in P_1(T) \quad \forall T \in \mathcal{T}_h \right\}.$$

Let $\kappa, h > 0$ be fixed. The fully discrete problem $(\mathbb{P}_{\kappa h}^\#)$ reads as follows:

$$\begin{cases} \text{Find } q_\kappa^* \in \mathcal{U}_\kappa^\# \text{ such that} \\ \mathcal{J}_h(q_\kappa^*) \leq \mathcal{J}_h(q_\kappa) \quad \forall q_\kappa \in \mathcal{U}_\kappa^\#, \end{cases} \quad (\mathbb{P}_{\kappa h}^\#)$$

where

$$\mathcal{J}_h(q_\kappa) := J(u_h(q_\kappa)) + \alpha \left(\int_{\bar{\Omega}} q_\kappa dx - \gamma \right)^+$$

and $u_h := u_h(q_\kappa) \in V_h$ is a solution of the Galerkin approximation of $(\mathcal{P}(q_\kappa))$:

$$\begin{aligned} (\nabla u_h, \nabla v_h)_{0,D} - k_0^2((1 + q_\kappa)u_h, v_h)_{0,D} - ik_0(u_h, v_h)_{0,\partial D} \\ = k_0^2(q_\kappa, v_h)_{0,D} \quad \forall v_h \in V_h. \end{aligned} \quad (\mathcal{P}_h(q_\kappa))$$

Using classical compactness arguments, one can easily prove that $(\mathbb{P}_{\kappa h}^\#)$ has a solution for any $h, \kappa > 0$.

In what follows, we shall study the relation between solutions to $(\mathbb{P}^\#)$ and $(\mathbb{P}_{\kappa h}^\#)$ when $h \rightarrow 0+$. To this end, we shall need the following stability result on the solutions $u_h(q_\kappa)$ of $(\mathcal{P}_h(q_\kappa))$, $q_\kappa \in \mathcal{U}_\kappa^\#$: there exist constants $c > 0$ and $h_0 := h_0(k_0) > 0$ such that

$$\|u_h(q_\kappa)\|_{1,D} \leq c \quad \forall q_\kappa \in \mathcal{U}_\kappa^\# \quad \forall h \leq h_0 \quad \forall \kappa \leq \kappa(h_0). \quad (27)$$

Comments on the satisfaction of (27) will be done later on.

We start with the following auxiliary result which is parallel of Proposition 1.

Proposition 3 *Let (27) be satisfied and $\{q_\kappa\}$, $q_\kappa \in \mathcal{U}_\kappa^\#$, be a sequence such that*

$$q_\kappa \xrightarrow{\kappa \rightarrow 0+} q \quad \text{weakly } * \text{ in } L^\infty(D) \quad (28)$$

and $u_h := u_h(q_\kappa)$ be the solution to $(\mathcal{P}_h(q_\kappa))$, $h \rightarrow 0+$. Then

$$u_h \xrightarrow{h \rightarrow 0+} u \quad \text{in } H^1(D)$$

and $u := u(q)$ is a solution to $(\mathcal{P}(q))$.

Proof: In view of (27), one can find a subsequence $\{u_{h_j}\}$ such that

$$u_{h_j} \xrightarrow{j \rightarrow \infty} u \quad \text{in } H^1(D). \quad (29)$$

Let $\tilde{v} \in H^1(D)$ be arbitrary but fixed. Since the system $\{V_h\}, h \rightarrow 0+$ is dense in $H^1(D)$, there exists a sequence $\{\tilde{v}_h\}, \tilde{v}_h \in V_h$ such that

$$\tilde{v}_h \xrightarrow{h \rightarrow 0+} \tilde{v} \quad \text{in } H^1(D). \quad (30)$$

The definition of $(\mathcal{P}_{h_j}(q_{\kappa_j}))$, where $\kappa_j = \kappa(h_j)$ with $\{h_j\}$ as in (29) yields:

$$(\nabla u_{h_j}, \nabla \tilde{v}_{h_j})_{0,D} - k_0^2((1 + q_{\kappa_j})u_{h_j}, \tilde{v}_{h_j})_{0,D} - ik_0(u_{h_j}, \tilde{v}_{h_j})_{0,\partial D} = k_0^2(q_{\kappa_j}u_0, \tilde{v}_{h_j})_{0,D}.$$

Letting $j \rightarrow \infty$, using (28), (29), (30) and the fact that $\tilde{v} \in H^1(D)$ is arbitrary we see that u solves $(\mathcal{P}(q))$ and (29) holds for the whole sequence. Strong convergence of $\{u_{h_j}\}$ to u can be established exactly as in Proposition 1. \square

The next theorem is parallel of Theorem 3.

Theorem 5 *Let (27) be satisfied and $\{(q_{\kappa_j}^*, u_{h_j}^*)\}$ be a sequence of optimal pairs of $(\mathbb{P}_{h\kappa}^\#)$, $h \rightarrow 0+$. Then there exist: a subsequence $\{(q_{\kappa_j}^*, u_{h_j}^*)\}$ and a pair $(q^*, u^*) \in \mathcal{U}^\# \times H^1(D)$ such that*

$$\begin{cases} q_{\kappa_j}^* \xrightarrow{j \rightarrow \infty} q^* & \text{weakly } * \text{ in } L^\infty(D) \\ u_{h_j}^* \xrightarrow{j \rightarrow \infty} u^* & \text{in } H^1(D). \end{cases} \quad (31)$$

Moreover (q^*, u^*) is an optimal pair of $(\mathbb{P}^\#)$. Any accumulation point of $\{(q_{\kappa_j}^*, u_{h_j}^*)\}$ in the sense of (31) has this property.

Proof: The existence of a subsequence $\{(q_{\kappa_j}^*, u_{h_j}^*)\}$ satisfying (31) is obvious. Let $\tilde{q} \in \mathcal{U}^\#$ be an arbitrary element. Then there exists a sequence $\{\tilde{q}_\kappa\}, \tilde{q}_\kappa \in \mathcal{U}_\kappa^\#$ such that

$$\tilde{q}_\kappa \xrightarrow{\kappa \rightarrow 0+} \tilde{q} \quad \text{weakly } * \text{ in } L^\infty(D)$$

and at the same time

$$u_h(\tilde{q}_\kappa) \xrightarrow{h \rightarrow 0+} u(\tilde{q}) \quad \text{in } H^1(D)$$

by Proposition 3. From the definition of $(\mathbb{P}_{\kappa_j, h_j}^\#)$ it follows:

$$\mathcal{J}_{h_j}(q_{\kappa_j}^*) \leq \mathcal{J}_{h_j}(\tilde{q}_{\kappa_j})$$

where h_j, κ_j is a filter of indices for which (31) holds. Passing to the limit with $j \rightarrow \infty$ and using the continuity assumption (15) we obtain

$$\mathcal{J}(q^*) \leq \mathcal{J}(\tilde{q}) \quad \forall \tilde{q} \in \mathcal{U}^\#.$$

\square

Remark 2 In what follows, we justify the validity of (27). Let $f \in L^2(Q)$ and $G \in H^1(Q)$ be given. Denote $g := G|_{\partial Q}$ and suppose that f, g satisfy the balance condition. It has been shown in [10] that every solution of the Neumann problem

$$\begin{cases} -\Delta u = f & \text{in } Q \\ \frac{\partial u}{\partial n} = g & \text{on } \partial Q \end{cases}$$

belongs to $H^2(Q)$ provided that Q is a convex polygonal domain and, in addition, there exists a constant $C > 0$ which does not depend on f, G and such that

$$\|u\|_{2,Q} \leq C(\|f\|_{0,Q} + \|G\|_{1,Q}).$$

From this and (6), one can easily show that the unique solution $u(q)$ to $(\mathcal{P}(q))$ belongs to $H^2(D)$ for any $q \in \mathcal{U}^\#$ and there exists a constant $C > 0$ which depends only on k_0 and D such that

$$\|u(q)\|_{2,D} \leq C \quad \forall q \in \mathcal{U}^\#. \quad (32)$$

Since the bilinear form $a(\cdot, \cdot)$ appearing in $(\mathcal{P}(q))$ satisfies the Gårding inequality, there exist: $h_0 > 0$ and a constant $\tilde{C} > 0$ which depend only on k_0 and D such that $(\mathcal{P}_h(q_\kappa))$ has a unique solution for any $h \leq h_0$ and any $q_\kappa \in \mathcal{U}_\kappa^\#, \kappa \leq \kappa(h_0)$. Moreover

$$\|u(q_\kappa) - u_h(q_\kappa)\|_{1,D} \leq \tilde{C} \inf_{v_h \in V_h} \|u(q_\kappa) - v_h\|_{1,D}$$

holds for any $h \leq h_0$ and any $q_\kappa \in \mathcal{U}_\kappa^\#, \kappa \leq \kappa(h_0)$ (see [18], [20]). From this and (32) we obtain (27).

4.2 Discretization by the level set method parametrized by radial basis functions

In this section we describe and theoretically justify the partial discretization of $(\mathbb{P}^\#)$ using a level set approximation of $\mathcal{U}^\#$ but still keeping the continuous setting of the state problem $(\mathcal{P}(q))$, $q \in \mathcal{U}^\#$. Level-set techniques have been gained attention during the past decade. See [23] for a review on the topic. We have chosen radial basis function (RBF) approach which was already used by the authors in the paper [22].

First we introduce the following notation. For any square Q we denote by $\mathcal{H}(Q)$ the set of all positive rationals δ defined as follows: $\delta \in \mathcal{H}(Q)$ if and only if there exists a partition $\mathcal{T}_\delta(Q)$ of \bar{Q} into squares of size δ whose interiors are mutually disjoint. *In the rest of the paper we shall assume that $\widehat{\Omega}$ is a square.*

Let $\kappa \in \mathcal{H}(\widehat{\Omega})$. For any $\Delta \in]0, 1[$ denote

$$\mathcal{U}_{\kappa\Delta}^\# = \left\{ q \in L^\infty(D) \mid 0 \leq q \leq 1 - \Delta \text{ a.e. in } \widehat{\Omega}, q|_{G_i} \in P_0(G_i) \quad \forall G_i \in \mathcal{T}_\kappa(\widehat{\Omega}), q \equiv 0 \text{ in } D \setminus \widehat{\Omega} \right\}$$

and define the perturbed problem

$$\min_{q \in \mathcal{U}_{\kappa\Delta}^\#} \mathcal{J}(q). \quad (\mathbb{P}_{\kappa\Delta}^\#)$$

Problem $(\mathbb{P}_{\kappa\Delta}^\#)$ has a solution for any $\kappa \in \mathcal{H}(\widehat{\Omega})$, $\Delta \in]0, 1[$ and

$$\min_{q \in \mathcal{U}_{\kappa}^\#} \mathcal{J}(q) \leq \min_{q \in \mathcal{U}_{\kappa\Delta}^\#} \mathcal{J}(q) \quad \forall \Delta \in]0, 1[.$$

Hence,

$$(\mathbb{P}_{\kappa}^\#) \leq \liminf_{\Delta \rightarrow 0+} (\mathbb{P}_{\kappa\Delta}^\#). \quad (33)$$

It is readily seen that the equality holds in (33). Indeed, if $q^* \in \mathcal{U}_{\kappa}^\#$ is a solution to $(\mathbb{P}_{\kappa}^\#)$ and $\{\Delta_n\}$ is a sequence such that $\Delta_n \searrow 0$ then $q_{\Delta_n}^* := q^* - \Delta_n \xrightarrow[n \rightarrow \infty]{} q^*$ in $L^\infty(D)$, and $u(q_{\Delta_n}^*) \xrightarrow[n \rightarrow \infty]{} u(q^*)$ in $H^1(D)$. Since $q_{\Delta_n}^* \in \mathcal{U}_{\kappa\Delta_n}^\#$, the opposite inequality holds in (33). Consequently

$$(\mathbb{P}^\#) = \liminf_{\substack{\kappa \rightarrow 0+ \\ \Delta \rightarrow 0+}} (\mathbb{P}_{\kappa\Delta}^\#) \quad (34)$$

making use of (15) and (24).

Next we shall approximate functions from $\mathcal{U}_{\kappa\Delta}^\#$, $\kappa \in \mathcal{H}(\widehat{\Omega})$, $\Delta \in]0, 1[$ by sequences of scaled characteristic functions of a special class of sets which are periodically distributed in each $G_i \in \mathcal{T}_{\kappa}(\widehat{\Omega})$.

To this end we construct another family of partitions of $\widehat{\Omega}$. Let $G_i \in \mathcal{T}_{\kappa}(\widehat{\Omega})$ be a square and $0 < \varepsilon \ll \kappa$ be a scaling factor such that $\varepsilon \in \mathcal{H}(G_i)$, $i = 1, \dots, n(\kappa)$, or equivalently, each $G_i \in \mathcal{T}_{\kappa}(\widehat{\Omega})$ is the union of small squares G_{ij} of size ε , $j = 1, \dots, m(\varepsilon)$ and moreover, $m(\varepsilon)$ is the *same* for all $G_i \in \mathcal{T}_{\kappa}(\widehat{\Omega})$. The system $\{G_{ij}\}$, $i = 1, \dots, n(\kappa)$, $j = 1, \dots, m(\varepsilon)$ defines a refined partition of $\mathcal{T}_{\kappa}(\widehat{\Omega})$ (see Figure 3).

Here and in what follows we shall suppose that $\kappa \in \mathcal{H}(\widehat{\Omega})$ and $\varepsilon \in \mathcal{H}(G_i) \forall G_i \in \mathcal{T}_{\kappa}(\widehat{\Omega})$. For any $m \in \mathbb{N}$, $m \geq 1$ we construct a regular grid of points $\{C_{kl}\}_{k,l=1}^m$, $C_{kl} = (\frac{k}{m+1}, \frac{l}{m+1})$, lying in the unit periodic cell $Y =]0, 1[\times]0, 1[$. Denote by $B_{kl,r}$ the disc with the center at C_{kl} and of radius r and set $\omega_{m,r} = \bigcup_{k,l=1}^m B_{kl,r}$. It is easy to see that for any $\Delta \in]0, 1[$ one can find $\bar{m} \in \mathbb{N}$, $\bar{r} > 0$ which depends solely on Δ and such that

$$\omega_{\bar{m},\bar{r}} \subset Y \quad (35)$$

$$\text{meas } \omega_{\bar{m},\bar{r}} = 1 - \Delta. \quad (36)$$

Clearly, for any $\rho \in [0, 1 - \Delta]$ there exists $r \in [0, \bar{r}]$ such that $\text{meas } \omega_{\bar{m},r} = \rho$ and $\omega_{\bar{m},r} \subset Y$. To emphasize the role of $\omega_{\bar{m},r}$, we denote by Y_r , εY_r the periodic cell containing $\omega_{\bar{m},r}$, and its scaled form, respectively (see Figure 2). With any $G_i \in \mathcal{T}_{\kappa}(\widehat{\Omega})$ we associate the parameter r_i which may take any value between 0 and \bar{r} . The scaled cell εY_{r_i} will be “stamped” onto all $G_{ij} \in \mathcal{T}_{\kappa}(G_i)$ (see Figure 3). The characteristic function of all inclusions in G_i generated by $\varepsilon \omega_{\bar{m},r_i} \subset \varepsilon Y_{r_i}$ (black part of G_i) will be denoted by $\chi_{\bar{m},r_i}^\varepsilon$ in what follows.

For any $\kappa, \varepsilon > 0$ and $\Delta \in]0, 1[$ we introduce the set

$$\mathcal{U}_{\kappa\Delta}^\varepsilon = \left\{ q \in L^\infty(D) \mid q|_{G_i} = \chi_{\bar{m},r_i}^\varepsilon, r_i \in [0, \bar{r}] \forall G_i \in \mathcal{T}_{\kappa}(\widehat{\Omega}), q \equiv 0 \text{ in } D \setminus \widehat{\Omega} \right\}. \quad (37)$$

We prove the following result:

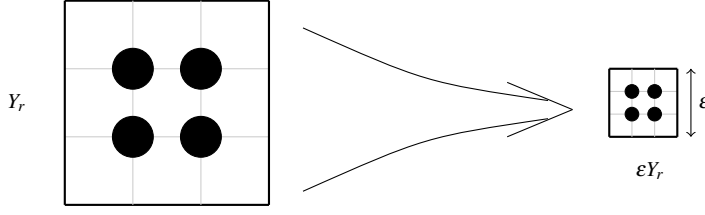


Fig. 2 Periodic cell Y_r and its scaling.

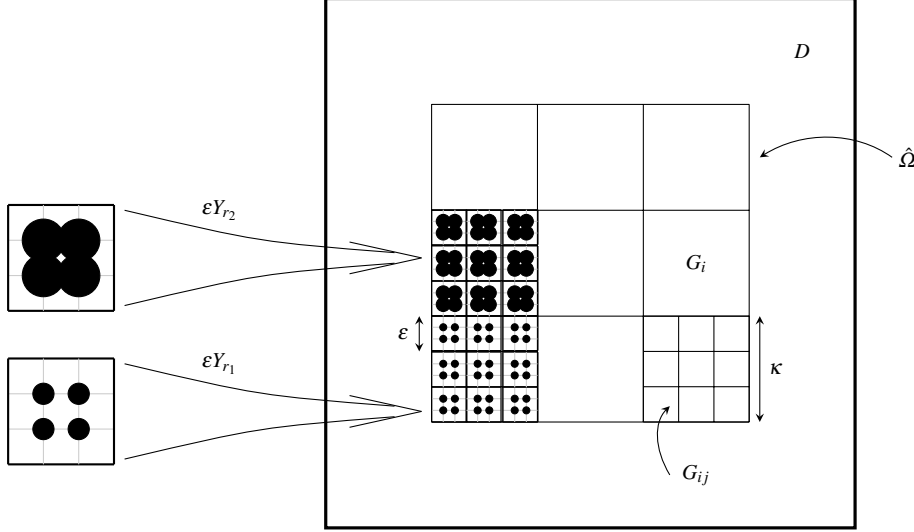


Fig. 3 Partition of $\widehat{\Omega}$ into $G_i \in \mathcal{T}_\kappa(\widehat{\Omega})$ and $G_{ij} \in \mathcal{T}_\varepsilon(G_i)$.

Proposition 4 Let $\kappa > 0$ and $\Delta \in]0, 1[$ be given. Then for any $q \in \mathcal{U}_{\kappa\Delta}^\#$ there exists a sequence $\{q_\varepsilon\}$, $q_\varepsilon \in \mathcal{U}_{\kappa\Delta}^\varepsilon$ such that

$$q_\varepsilon \xrightarrow{\varepsilon \rightarrow 0+} q \text{ weakly * in } L^\infty(D). \quad (38)$$

Proof: For $\kappa > 0$ and $\Delta \in]0, 1[$ given there exist $\bar{m} \in \mathbb{N}$ and $\bar{r} > 0$ for which (35) and (36) are satisfied. If $q \in \mathcal{U}_{\kappa\Delta}^\#$ then $q^i \in [0, 1 - \Delta]$ for any $G_i \in \mathcal{T}_\kappa(\widehat{\Omega})$, where $q^i := q|_{G_i}$ and one can find $\rho_i \in [0, \bar{r}]$ such that

$$\text{meas } \omega_{\bar{m}, \rho_i} = q^i \quad \forall i = 1, \dots, n(\kappa).$$

The sequence of the resulting scaled characteristic functions $\{\chi_{\bar{m}, \rho_i}^\varepsilon\}$ satisfies

$$\chi_{\bar{m}, \rho_i}^\varepsilon|_{G_i} \xrightarrow{\varepsilon \rightarrow 0+} q^i \text{ weakly * in } L^\infty(G_i) \quad (39)$$

for any $G_i \in \mathcal{T}_\kappa(\widehat{\Omega})$. Let $q_\varepsilon \in \mathcal{U}_{\kappa\Delta}^\varepsilon$ be defined by

$$q_\varepsilon|_{G_i} = \chi_{\bar{m}, \rho_i}^\varepsilon|_{G_i} \quad G_i \in \mathcal{T}_\kappa(\widehat{\Omega}), \quad \forall \varepsilon > 0.$$

Then (38) is satisfied in virtue of (39). \square

Recall that (38) and Proposition 1 entail

$$u(q_\varepsilon) \xrightarrow{\varepsilon \rightarrow 0+} u(q) \quad \text{in } H^1(D), \quad (40)$$

where $u(q_\varepsilon)$, $u(q)$ is a solution to $(\mathcal{P}(q_\varepsilon))$, and $(\mathcal{P}(q))$, respectively, $q_\varepsilon \in \mathcal{U}_{\kappa\Delta}^\varepsilon$, $q \in \mathcal{U}_{\kappa\Delta}^\#$. For any $\kappa, \varepsilon > 0$ and $\Delta \in]0, 1[$ given, define the problem

$$\inf_{q \in \mathcal{U}_{\kappa\Delta}^\varepsilon} \mathcal{J}(q). \quad (\mathbb{P}_{\kappa\Delta}^\varepsilon)$$

It is easy to show that $(\mathbb{P}_{\kappa\Delta}^\varepsilon)$ has a solution for any $\kappa, \varepsilon > 0$ and $\Delta \in]0, 1[$. Since

$$\min_{q \in \mathcal{U}_{\kappa\Delta}^\#} \mathcal{J}(q) \leq \inf_{q \in \mathcal{U}_{\kappa\Delta}^\varepsilon} \mathcal{J}(q) \quad \forall \varepsilon > 0, \Delta \in]0, 1[,$$

it follows from this, (38), and (40) that

$$(\mathbb{P}_{\kappa\Delta}^\#) = \liminf_{\varepsilon \rightarrow 0+} (\mathbb{P}_{\kappa\Delta}^\varepsilon)$$

which together with (34) yields:

$$(\mathbb{P}^\#) = \liminf_{\substack{\kappa \rightarrow 0+ \\ \Delta \rightarrow 0+ \\ \varepsilon \rightarrow 0+}} (\mathbb{P}_{\kappa\Delta}^\varepsilon). \quad (41)$$

Remark 3 From (41) it follows that for any $\eta > 0$ there exist $\kappa_\eta, \varepsilon_\eta > 0$ and $\Delta_\eta \in]0, 1[$ such that

$$l \leq \mathcal{J}(q^*) \leq l + \eta, \quad l = \min_{q \in \mathcal{U}^\#} \mathcal{J}(q)$$

holds for any optimal pair (q^*, u^*) of $(\mathbb{P}_{\kappa_\eta \Delta_\eta}^{\varepsilon_\eta})$.

Problem $(\mathbb{P}_{\kappa\Delta}^\varepsilon)$, $\kappa, \varepsilon > 0$ and $\Delta \in]0, 1[$ can be easily formulated as a parameter identification problem. Indeed, let $\Phi : [0, \infty[\rightarrow \mathbb{R}$ be a function which is positive in $[0, 1[$ and $\Phi \equiv 0$ elsewhere. Define

$$\begin{cases} \Psi_{kl,r}(x) = \Phi\left(\frac{\|x - C_{kl}\|}{r}\right), & x \in Y, \quad i, j = 1, \dots, m, \quad r > 0, \\ \Psi_{kl,0} \equiv 0, \end{cases} \quad (42)$$

where $\{C_{kl}\}_{k,l=1}^m$ is the regular grid of points in Y constructed above.

Let $\Delta \in]0, 1[$ be given and $\bar{m} \in \mathbb{N}$ and $\bar{r} > 0$ be such that (35) and (36) are satisfied. Since $\text{supp } \Psi_{kl,r} = \overline{B_{kl,r}}$ then $\omega_{\bar{m},r} = \text{supp } \Psi_{\bar{m},r}$, where $\Psi_{\bar{m},r} : Y \rightarrow \mathbb{R}$,

$$\Psi_{\bar{m},r} = \sum_{k,l=1}^{\bar{m}} \Psi_{kl,r}, \quad r \in [0, \bar{r}]. \quad (43)$$

If H denotes the Heaviside function then $H(\Psi_{\bar{m},r})$ is the characteristic function of $\omega_{\bar{m},r} \subset Y$, $r \in [0, \bar{r}]$. Next, $\Psi_{\bar{m},r}$ will be periodically extended from Y to the whole \mathbb{R}^2

(keeping the same notation) and then scaled: $\Psi_{\bar{m},r}^\varepsilon(x) = \Psi_{\bar{m},r}(\frac{x}{\varepsilon})$, $x \in \mathbb{R}^2$, $\varepsilon > 0$. It is easy to see that after an appropriate translation of the grid εZ , where $Z = (k, l)$, k, l integer, the restriction $H(\Psi_{\bar{m},r}^\varepsilon)|_{G_i}$ is the scaled characteristic function $\chi_{\bar{m},r}^\varepsilon$ from (37).

Let $\kappa, \varepsilon > 0$ and $\Delta \in]0, 1[$ be fixed. Denote by \mathcal{U}_{ad} the following subset of $\mathbb{R}^{n(\kappa)}$:

$$\mathcal{U}_{ad} = \{r = (r_1, \dots, r_{n(\kappa)}) \in \mathbb{R}^{n(\kappa)} \mid r_i \in [0, \bar{r}] \quad \forall i = 1, \dots, n(\kappa)\}$$

and define

$$Q(r)|_{G_i} := H(\Psi_{\bar{m},r_i}^\varepsilon|_{G_i}) \quad \forall r_i \in [0, \bar{r}] \quad \forall i = 1, \dots, n(\kappa). \quad (44)$$

Clearly, any function $q \in \mathcal{U}_{\kappa\Delta}^\varepsilon$ can be identified with a vector $r \in \mathcal{U}_{ad}$ and optimization problem $(\mathbb{P}_{\kappa\Delta}^\varepsilon)$ can be equivalently written in the following form:

$$\begin{cases} \text{Find } r^* \in \mathcal{U}_{ad} \text{ such that} \\ \mathcal{J}(r^*) \leq \mathcal{J}(r) \quad \forall r \in \mathcal{U}_{ad}, \end{cases} \quad (\mathbb{P}_r)$$

where

$$\mathcal{J}(r) := J(u(r)) + \alpha \left(\int_{\hat{\Omega}} Q(r) dx - \gamma \right)^+$$

and $u(r) \in H^1(D)$ solves the state problem

$$\begin{aligned} (\nabla u(r), \nabla v)_{0,D} - k_0^2((1 + Q(r))u(r), v)_{0,D} - ik_0(u(r), v)_{0,\partial D} \\ = k_0^2(Q(r)u_0, v)_{0,D} \quad \forall v \in H^1(D). \end{aligned} \quad (\mathcal{P}(r))$$

Remark 4 Instead of (43), one can use a more general definition of $\Psi_{\bar{m},r}$, namely

$$\Psi_{\bar{m},r,\alpha} = \sum_{k,l=1}^{\bar{m}} \alpha_{kl} \Psi_{kl,r} \quad \text{in } Y, \quad (45)$$

where $r \in [0, \bar{r}]$, $\alpha_{kl} \in [\alpha_{\min}, \alpha_{\max}] \forall k, l$, $\alpha_{\min} < \alpha_{\max}$ given and $\alpha \in \mathbb{R}^{m^2}$ is the vector of α_{kl} , $k, l = 1, \dots, \bar{m}$. The appropriate choice of the coefficients α_{kl} gives us more flexibility in building $\text{supp } \Psi_{\bar{m},r,\alpha} \subset Y$. Indeed, if $\alpha_{kl} = 0$ for some k, l then the respective $\text{supp } \Psi_{k,l,r}$ contributing to $\omega_{\bar{m},r}$ is missing. If for example $\Delta \in]1 - \frac{\pi}{4}, 1[$ then instead of m^2 small discs of radius r at C_{kl} the union of which satisfies (36) one can use only one disc of radius $r \in [0, \frac{1}{2}]$ and center $(\frac{1}{2}, \frac{1}{2}) \in Y$ to fulfill (36). Analogously (44), let Q be a function of α and r :

$$Q(\alpha, r)|_{G_i} = H(\Psi_{\bar{m},r_i,\alpha_i}^\varepsilon|_{G_i}) \quad \forall G_i \in \mathcal{T}_\kappa(\hat{\Omega}),$$

where $r \in \mathcal{U}_{ad}$ and $\alpha = (\alpha^1, \dots, \alpha^{n(\kappa)}) \in (\mathbb{R}^{m^2})^{n(\kappa)}$. In this case, $(\mathbb{P}_{\kappa\Delta}^\varepsilon)$ can be formulated as a minimization problem for the control variables α and r .

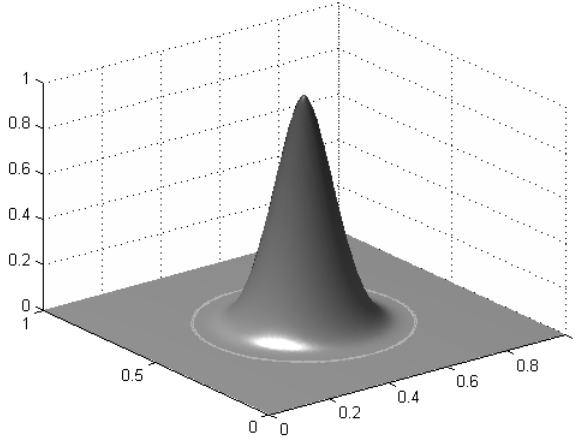


Fig. 4 Graph of the radial basis function $\Psi_{i_j, r_{i_j}}$ defined by (47).

5 Numerical realization

In this section we describe numerical realization of the radial basis function approach. As in computations one is not able to use a high number of radial basis functions, we use just one regular grid of points $\{C_{ij}\}$ uniformly distributed in $\widehat{\Omega}$ (see Figure 5) and define the level set function as the linear (affine) combination of RBF basis functions as follows:

$$\Psi_{\alpha, r} = \Psi_0 + \sum_{i, j=1}^m \alpha_{i_j} \Psi_{i_j, r_{i_j}}, \quad (46)$$

where $\text{supp} \Psi_{i_j, r_{i_j}}$ is the disc with center at C_{i_j} and of radius r_{i_j} . Thus, the design variables of the parametrized optimization problem are represented by the vector $\beta := (\alpha, r) \in \mathbb{R}^N$, $N := 2m^2$. The constant shift Ψ_0 is optional and its role will be established later. The function Φ in (42) used for constructing $\Psi_{i_j, r_{i_j}}$ is chosen as ([25]):

$$\Phi(t) = [(1-t)^+]^4 (4t+1), \quad t \in \mathbb{R}. \quad (47)$$

The respective radial basis function $\Psi_{i_j, r_{i_j}}$ is depicted in Figure 4.

The function q in (11) can be now written with the aid of the level set function $\Psi_\beta := \Psi_{\alpha, r}$ and the Heaviside step function as follows:

$$q = H(\Psi_\beta). \quad (48)$$

The weak formulation of the state problem (10) is discretized by piecewise linear finite elements as discussed in Section 4.1. By q_β we denote any of the finite dimensional parametrizations of q discussed in the previous section. The finite element

approximation of (10) then reads as follows:

$$\begin{cases} \text{Find } u_h(q_\beta) \in V_h \text{ such that} \\ a(q_\beta; u_h(q_\beta), v_h) = b(q_\beta; v_h) \quad \forall v_h \in V_h, \end{cases} \quad (49)$$

where

$$\begin{aligned} a(q_\beta; z, w) &= (\nabla z, \nabla w)_{0,D} - k_0^2((1 + q_\beta)z, w)_{0,D} - ik_0(z, w)_{0,\partial D} \\ b(q_\beta; w) &= k_0^2(q_\beta u_0, w)_{0,D}. \end{aligned}$$

Remark 5 In order to obtain physically realistic results, the simplest first-order boundary condition (8) on ∂D is inappropriate unless ∂D is very far from $\partial \widehat{\Omega}$. Therefore, the second-order approximation of (7) proposed by Engquist and Majda [8] and Bamberg et al. [2] was implemented. Instead of $a(q_\beta; z, w)$ we use its modified form

$$\tilde{a}(q_\beta; z, w) := a(q_\beta; z, w) + \frac{i}{2k_0} \left(\frac{\partial z}{\partial s}, \frac{\partial w}{\partial s} \right)_{0,\partial D} + \frac{3}{4} \sum_{x \in CP} z(x) \overline{w(x)}, \quad z, w \in V_h,$$

where s is the unit tangent vector on ∂D and CP denotes the set of four corners of ∂D . The finite element approximation of (10) used in computations reads as follows:

$$\begin{cases} \text{Find } u_h(q_\beta) \in V_h \text{ such that} \\ \tilde{a}(q_\beta; u_h(q_\beta), v_h) = b(q_\beta; v_h) \quad \forall v_h \in V_h. \end{cases} \quad (50)$$

A structured and uniform triangulation \mathcal{T}_h of \overline{D} for constructing V_h is shown in Figure 5.

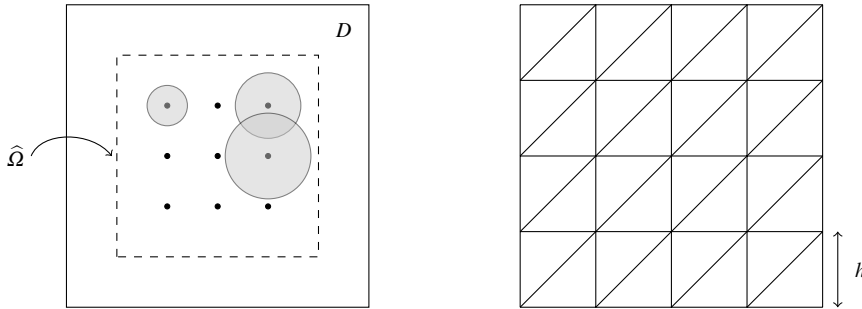


Fig. 5 Positions of $\{C_{ij}\}$ and supports of a few radial basis functions (left). The topology of the triangular finite element mesh \mathcal{T}_h (right).

Let $h > 0$ and $m \in \mathbb{N}$ be fixed. The approximate solution $u_h \in V_h$ is represented as a linear combination $u_h = \sum_i c_i \varphi_i$, where $c_i \in \mathbb{C}$ and $\{\varphi_i\}$ are the standard real-valued Courant basis functions. The matrix form of (50) reads:

$$A(\beta)u(\beta) = b(\beta), \quad (51)$$

where the matrix $A(\beta)$ and the right-hand side $b(\beta)$ can be decomposed as follows:

$$A(\beta) = K - k_0^2 M(\beta) - ik_0 C + \frac{i}{2k_0} D + \frac{3}{4} G, \quad b(\beta) = k_0^2 f(\beta).$$

The entries of the above matrices and of the vector $f(\beta)$ are given by

$$\begin{aligned} K_{ij} &= \int_D \nabla \varphi_i \cdot \nabla \varphi_j dx, & M(\beta)_{ij} &= \int_D (1 + q_\beta) \varphi_i \varphi_j dx \\ C_{ij} &= \int_{\partial D} \varphi_i \varphi_j ds, & D_{ij} &= \int_{\partial D} \frac{\partial \varphi_i}{\partial s} \frac{\partial \varphi_j}{\partial s} ds \\ G_{ij} &= \sum_{x \in CP} \varphi_i(x) \varphi_j(x), & f(\beta)_i &= \int_D u_0 q_\beta \varphi_i dx. \end{aligned}$$

Matrices K , C , D , and G are evaluated exactly. Their sum is a pentadiagonal sparse matrix. To compute the design dependent ‘‘mass’’ matrix $M(\beta)$ and the excitation vector $f(\beta)$, we utilize numerical integration. Let $\mathcal{V}(T)$ be the set of nodal numbers of vertices associated with the triangle T . Moreover, let x_T, x_l denote the center of gravity and the l :th nodal point of T , respectively. The following integration formulae are used:

$$\begin{aligned} M(\beta)_{ij} &\approx \sum_{T \in \mathcal{T}_h} \left[(1 + q_\beta(x_T)) \frac{1}{3} \text{meas}(T) \sum_{k \in \mathcal{V}(T)} \varphi_i(x_k) \varphi_j(x_k) \right] \\ f(\beta)_i &\approx \sum_{T \in \mathcal{T}_h} \left[u_0(x_T) q_\beta(x_T) \frac{1}{3} \text{meas}(T) \sum_{k \in \mathcal{V}(T)} \varphi_i(x_k) \right]. \end{aligned}$$

Note that the former yields a diagonal approximation of $M(\beta)$.

Finally, we arrive at the following fully discrete nonlinear programming problem:

$$\begin{cases} \text{Find } \beta^* \in \mathcal{A}_{\text{ad}} \text{ such that} \\ \mathcal{J}(\beta^*) \leq \mathcal{J}(\beta) \quad \forall \beta \in \mathcal{A}_{\text{ad}}, \end{cases} \quad (\mathbf{P})$$

where $\mathcal{J}(\beta) := J(u_h(q_\beta)) + \alpha \left(\int_{\hat{\Omega}} H(\Psi_\beta) dx - \gamma \right)^+$ and

$$\mathcal{A}_{\text{ad}} = \left\{ \beta = (\alpha, r) \in \mathbb{R}^N \mid \alpha_{\min} \leq \alpha_{ij} \leq \alpha_{\max}, 0 \leq r_{ij} \leq r_{ij, \max} \quad \forall i, j \right\}.$$

is the set of admissible design parameters, where $\alpha_{\min}, \alpha_{\max}$ and $r_{ij, \max} \quad \forall i, j$ are given.

The discretized cost function \mathcal{J} in (\mathbf{P}) is not continuously differentiable due to the discontinuous Heaviside function H . If one wishes to use descent-type methods, then smoothing of H is necessary. Therefore, H will be replaced by a smooth function H_s defined by

$$H_s(x) = \begin{cases} 0 & \text{if } x < -s \\ 1 & \text{if } x > s \\ \frac{3}{4} \left(\frac{x}{s} - \frac{x^3}{3s^3} \right) + \frac{1}{2} & \text{otherwise,} \end{cases} \quad (52)$$

where $s > 0$ is given. As $H_s(0) = \frac{1}{2}$, it is necessary to add a negative shift Ψ_0 to the level set function Ψ_β to ensure that $H_s(\Psi_\beta(x)) = 0$ for $x \notin \hat{\Omega}$. Alternatively, one can omit the shift and modify the definition of H_s .

To compute the gradient of \mathcal{J} , the standard adjoint variable approach will be used:

$$\frac{\partial \mathcal{J}(\beta)}{\partial \beta_k} = \Re \left[p^* \left(\frac{\partial b(\beta)}{\partial \beta_k} - \frac{\partial A(\beta)}{\partial \beta_k} u(\beta) \right) \right], \quad k = 1, \dots, N. \quad (53)$$

Here p solves the adjoint problem

$$[A(\beta)]^* p = \nabla_u \mathcal{J}(u(\beta)), \quad (54)$$

where \star stands for the conjugate transpose of a matrix and the complex gradient ∇_u is defined by

$$\nabla_u \mathcal{J}(u(\beta)) := \nabla_{\Re u} \mathcal{J}(u(\beta)) + i \nabla_{\Im u} \mathcal{J}(u(\beta)).$$

The partial derivatives of $M(\beta)$ and $f(\beta)$ in (53) are easy to compute. From (48) we have:

$$\begin{aligned} \frac{\partial H_s(\Psi_\beta)}{\partial \alpha_{ij}} &= H'_s(\Psi_\beta) \Psi_{ij,r_{ij}} \quad \forall i, j \\ \frac{\partial H_s(\Psi_\beta)}{\partial r_{ij}} &= H'_s(\Psi_\beta) \frac{\partial \Psi_{ij,r_{ij}}}{\partial r_{ij}} \quad \forall i, j. \end{aligned}$$

6 Numerical experiments

This section presents and compare numerical results of two model examples which are computed using the relaxation formulation and the radial basis approach. The basic geometry in both problems is the same, namely $D =]-1, 1[\times]-1, 1[$, $\widehat{\Omega} =]-\frac{7}{10}, \frac{7}{10}[\times]-\frac{7}{10}, \frac{7}{10}[$, and $d = (0, 1)$.

In the relaxed formulation, the coefficient $q \in \mathcal{U}^\#$ is represented by

$$q = q_1 \sum_{i,j=1}^m \alpha_{ij} \chi_{G_{ij}}, \quad \alpha_{ij} \in [0, 1] \quad \forall i, j,$$

where G_{ij} is the system of m^2 small squares which define the partition of $\widehat{\Omega}$ and $q_1 = \varepsilon_{r,1} - \varepsilon_{r,0}$ (see (4)).

The level set function used in computations by the radial basis approach has the following form:

$$\Psi_r = \Psi_0 + \sum_{i,j=1}^m r_{ij} \Psi_{ij,r_{ij}}, \quad (55)$$

where $\Psi_{ij,r_{ij}}$ is defined by (42). Unlike (46) where the coefficients of the linear combination α_{ij} are independent of r_{ij} we set $\alpha_{ij} = r_{ij}$ in (55). The motivation for this choice of α_{ij} is the following: the maximum of $\Psi_{ij,r_{ij}}$ is always equal to one for any $r_{ij} > 0$. Therefore, when $r_{ij} \rightarrow 0$ the basis function $\Psi_{ij,r_{ij}}$ becomes an increasingly sharp peak resulting in numerical difficulties. The maximal radius $r_{ij,\max}$ of $\text{supp} \Psi_{ij,r_{ij}}$ in the definition of \mathcal{A}_{ad} is $1.4/(m+1)$ for all C_{ij} which are closest to $\partial \widehat{\Omega}$ and $2.8/(m+1)$ for the rest.

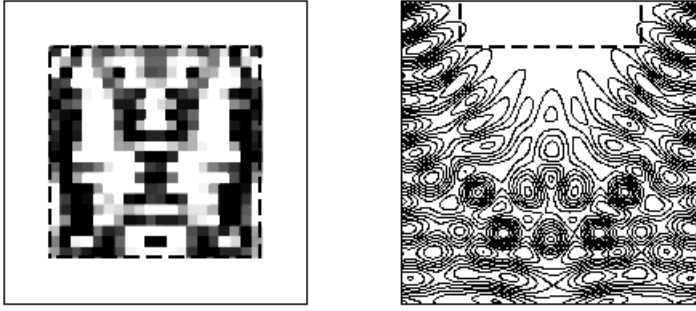


Fig. 6 Left: Optimized design. Right: Total wave $\Re(u+u_0)$. The location of D_0 is shown by a dashed line.

The smoothing parameter s in (52) is chosen by using the following heuristic rule: $s = s_0 h$, where $s_0 = 4$. We do not employ more sophisticated adaptive techniques to determine s used e.g. in [22] as they would increase the nonlinearity of the optimization problem and the heuristic choice works well enough. The shift in (55) is chosen as $\Psi_0 = -(s_0 + 0.1)h$.

The state and adjoint solver as well as the cost function evaluation were implemented using MATLAB [16]. Minimization was carried out by `fmincon` with `'sqp'` option from the MATLAB Optimization Toolbox. All computations were done on a desktop computer.

Example 1 This example deals with a simplified problem of designing the so-called bandgap structures. The aim is to design the scatterer Ω in such a way that the wave having the wave number k_0 is prevented to enter a specified subdomain $D_0 \subset D$. We use the simplest choice of the cost function, namely

$$J(u) = \frac{1}{2} \int_{D_0} |u + u_0|^2 dx,$$

i.e. the absolute value of the total wave u_{tot} (see (5)) is minimized in D_0 .

The physical parameters are $k_0 = 6\pi$ and $q_1 = 0.75$. We choose $D_0 =]-0.6, 0.6[\times]0.7, 1[$, and the mesh size $h = 2/300$. The problem was again solved by both the relaxed and the radial basis function formulations. We used grid of 20×20 design variables in both cases.

The initial guess for the relaxed formulation was $\alpha_{ij} = 0.5 \forall i, j$. After 121 optimization iterations (122 function evaluations) the cost function was reduced to the value 3.40×10^{-4} . The optimized design and the corresponding total wave are shown in Figure 6.

The optimized design has a significant amount of grey cells. In topology optimization during optimization process, usually some filtering or penalizing techniques are used to force the design to the upper or lower bound [5]. However, sophisticated filtering techniques are out of scope of this paper. Therefore, we only threshold the optimized design variables α_{ij}^* to either 0 or 1, i.e. we set $\hat{\alpha}_{ij}^* = H(\alpha_{ij}^* - \frac{1}{2})$. The

thresholded optimized design (the cost function value 3.76×10^{-3}) and the corresponding total wave are shown in Figure 7.

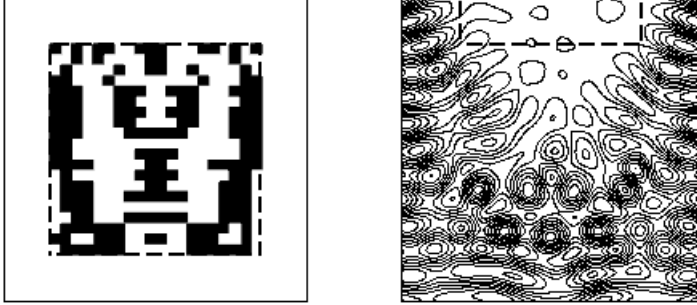


Fig. 7 Left: Thresholded design. Right: Corresponding total wave $\Re(u + u_0)$

The initial guess for the radial basis function formulation was $r_{ij} = 0.05\forall i, j$. After 47 optimization iterations (50 function evaluations) the cost function decreased to the value 5.95×10^{-5} . The optimized design and the corresponding total wave are shown in Figure 8.

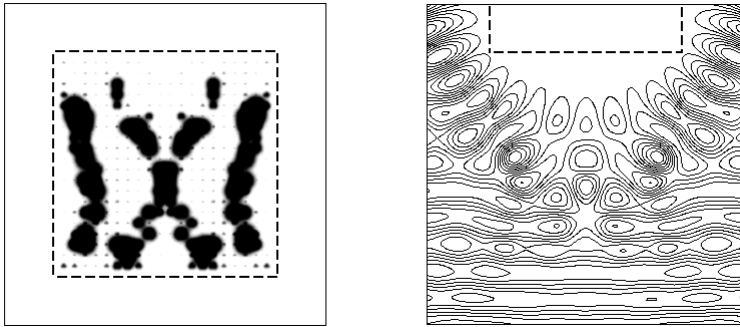


Fig. 8 Left: Optimized design. Right: Corresponding total wave $\Re(u + u_0)$.

The optimized design contains many tiny discs which are obviously due to numerical noise and having no physical meaning. Therefore we repeated the optimization by taking as the initial guess the final solution of the previous run but setting those r_{ij} smaller than 0.08 to zero. After 38 optimization iterations (44 function evaluations) the cost function was reduced to 6.07×10^{-5} . The results are shown in Figure 9. The cost is only slightly higher but the obtained design is now almost free from tiny discs.

Comparing the optimized designs obtained by the relaxed and radial basis function formulations, we see that they differ significantly on details. However, both de-

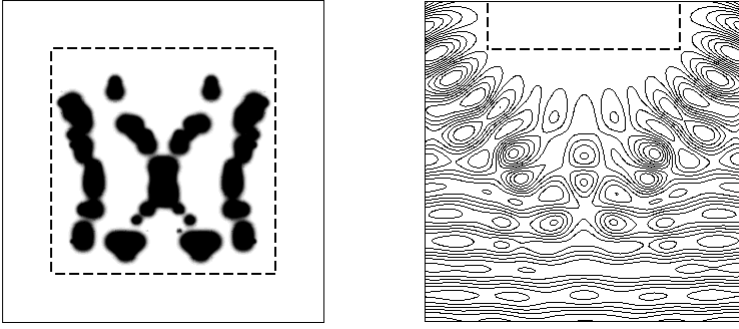


Fig. 9 Left: Optimized design. Right: Total wave $\Re(u + u_0)$.

signs essentially meet the design target, i.e., they prevent wave entering into D_0 . It is obvious that the optimization problem has many global optima and/or has many local minima with slightly higher cost function value than the global minimum.

Finally, we recomputed the example (20×20 radial basis function discretization, initial guess $r_{ij} = 0.05 \forall i, j$) but with the regularized cost function $\tilde{J}(u) = J(u) + \alpha \int_D H_s(\Psi)$, $\alpha = 10^{-2}$. After 117 optimization iterations (123 function evaluations) the cost function was reduced to 3.84×10^{-3} . The optimized design and the corresponding total wave are shown in Figure 10.

The optimized design is almost free of isolated tiny disks. However, it is not symmetric any more.

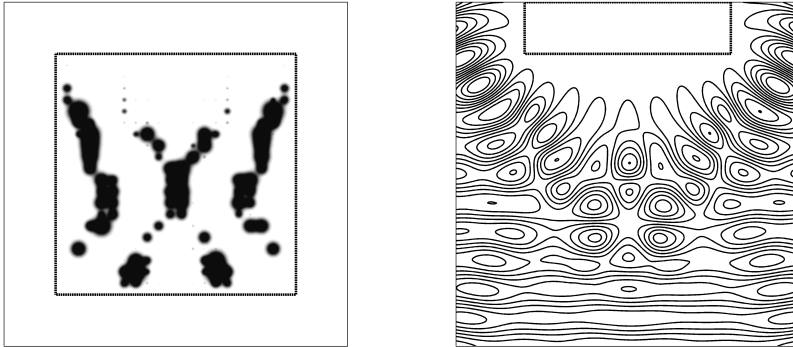


Fig. 10 Left: Optimized design. Right: Total wave $\Re(u + u_0)$.

Example 2 In this example, the basic setting of the problem is the same as in Example 1 except the area constraint $\text{meas}(\Omega) = \gamma$ added as a penalty term into the cost function:

$$\tilde{J}(u) = J(u) + \alpha \int_D [(H_s(\Psi) - \gamma)^+]^2 dx$$

We choose the target area as $\gamma = \frac{3}{10}$ and the penalty parameter as $\alpha = 10$. The initial guess for the relaxed formulation (30×30 basis functions) was $\alpha_{ij} = 0.25\forall i, j$. After 250 optimization iterations (253 function evaluations) the cost function was reduced to the value 4.87×10^{-3} . The optimized design and the corresponding total wave are shown in Figure 11. Convergence of the cost functional with respect to optimization iterations is shown in Figure 12.

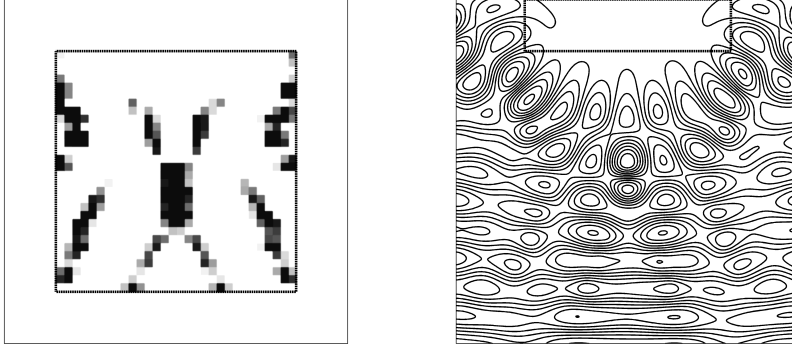


Fig. 11 Left: Optimized design. Right: Total wave $\Re(u + u_0)$.

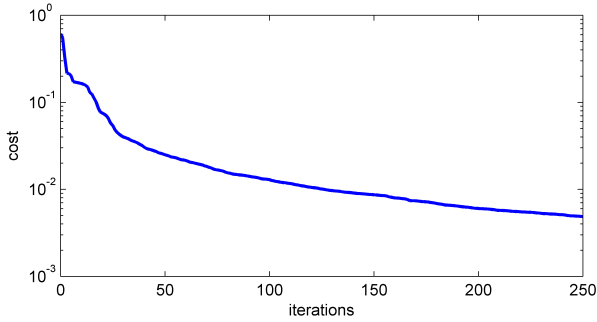


Fig. 12 Convergence of the cost functional with respect to the number of iterations

The initial guess for the radial basis function formulation (20×20 basis functions) was $r_{ij} = 0.05\forall i, j$. After 162 optimization iterations (183 function evaluations) the cost function decreased to the value 2.14×10^{-3} . The optimized design and the corresponding total wave are shown in Figure 13. Convergence of the cost functional with respect to optimization iterations is shown in Figure 14.

As can be noticed, the additional area constraint also serves as a good way to avoid oscillating designs.

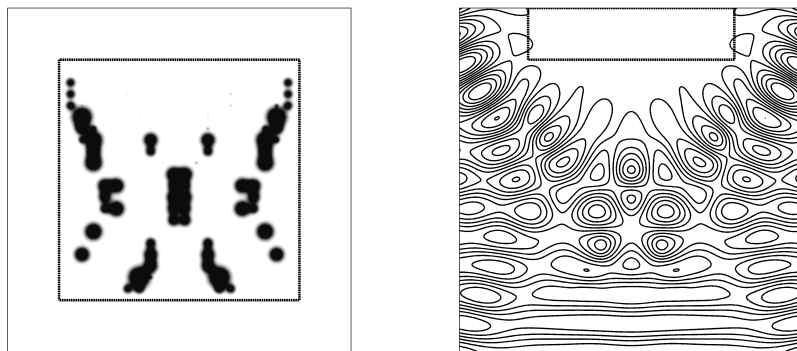


Fig. 13 Left: Optimized design. Right: Total wave $\Re(u + u_0)$.

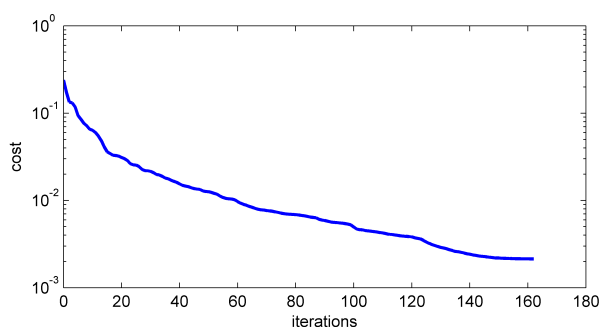


Fig. 14 Convergence of the cost functional with respect to the number of iterations

7 Conclusions

This paper deals with a class of topology optimization problems governed by the two dimensional Helmholtz equation and the permittivity as the control variable. Its main contribution consists in the rigorous convergence analysis using both, the classical relaxed formulation and a level set/radial basis function approach.

The numerical tests confirmed that this type of problems is difficult from the implementation point of view, since many designs produce very similar solutions to the state problem. Also the discretization of the control variable using radial basis functions having variable radii seems to increase the nonlinearity of the optimization problem. However, the presence of a penalty or regularization functional helped to stabilize the computations.

Computations which are based on the relaxed formulation with a penalization or filtering techniques are currently dominant in practical applications. On the other hand it is worthwhile to study alternative techniques, among them level set methods which attained much interest in the past years. The advantage of level set approaches is the fact that they work on the “safe side”, i.e. the designs are automatically required to have values $\varepsilon_{r,0}$ and $\varepsilon_{r,1}$ except in a very small region near the material boundaries

due to the smoothing of the Heaviside function. Also smooth boundaries of simple geometrical shapes are better represented by the radial basis function approach.

Variants of the level set/ radial basis functions approach presented in this paper, in particular the case when a single fixed radius defines the basis functions is the object of future studies.

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