STRUCTURAL STUDIES OF DIELECTRIC POLYMER NANOCOMPOSITES

BY

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Academic Dissertation for the Degree of Doctor of Philosophy

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ABSTRACT

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A constant need for the development of new and superior materials is always present; for example, a better electrical insulator would enable more efficient use of electrical power. Polymeric nanocomposites, i.e. nanodielectrics, are thought to have unique electrical properties. The basic chemical constitution of a material alone fails to provide an understanding of how desired properties of a material originate or predict the long term behavior of a material. We need to define the structure behind the functionality of a material. To do that, the structure of the material must be studied on several scales. This research was part of the Finnish Funding Agency for Technology and Innovation (TEKES) consortium projects NANOCOM and NANOPOWER. The general objective of these projects was to create truly new theoretical, experimental and practical knowledge of novel polymer nanocomposites to be further developed and finally used both in electrical and electronics insulation technology, as well as in other fields of technology.

Raman imaging was found to be a good tool for studying the structure of the materials: it provides information on the chemical species as well as dispersion of the filler. Alongside traditional confocal Raman imaging, which gives detailed information even at the submicron scale, coarse Raman imaging was used allowing the screening of a large area from a sample which is important for quality control of the composites to be used in industrial scale.

A new synthetic approach was used to afford well-dispersed silica particles with electroactive core functionalization in epoxy in order to study the effect of the charge layer at the interface of nanoparticle and polymer matrix. If the achieved distribution of particles was dense enough, the dielectric breakdown strength (DBS) increased considerably. The demonstrated increase in DBS and permittivity of the material leads to an increase of up to 125% in theoretical capacitive energy storage capability, and this is promising for future applications. These changes in properties are achieved with only 2 wt-% filler loading. Dielectric losses in the frequency range critical to the planned application stayed at the level of unfilled epoxy. This work is the first step towards new type functionalization of filler that offers good dispersion of nanoparticles without using harsh mixing conditions. These preliminary results indicate that the charge layer in the nanoparticle core, alongside a sufficiently dense enough distribution of particles at nanoscale, could be one way to improve dielectric properties of polymer materials.
Keywords: nanodielectrics, polymer nanocomposite, dispersion, transmission electron microscopy, Raman microscopy
PREFACE

The work presented in this thesis has been carried out at the Department of Chemistry, Nanoscience Center at the University of Jyväskylä during the years 2008-2013. For the last manuscript the research is done during 10 months ASLA-Fulbright Pre-Doctoral Research Fellows visit at Rensselaer Polytechnic Institute (RPI), Troy, NY 2012-2013.

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First I would like to thank my supervisor Prof. Mika Pettersson for giving me opportunity to step in to this challenge, and also for being such a great role model. There were so many people that provided help and the work that they have done, has kept this research going forward. No amount of gratitude could ever be too much for XPD study and support I have got from Adj. Prof. Manu Lahtinen. I wish to express my gratitude to Dr. Viivi Nuottajärvi for tomography analysis; you had such a deep interest to do it. MSc. Tuomas Turpeinen was an indispensable help in particle analysis; you were there always to help me in the hour of need. I could not have done this without your help. TEM lab technicians: Paavo Niutanen and Petri Papponen are greatly appreciated. MSc. Susanna Ahonen stepped in to the project during time of my maternity leave and did such a great job at it. MSc. Jaakko Koivisto is remembered for all the help he has given me over the years.

My collaborators were the key to accomplish any results: most the dielectric analysis of the studied materials and theoretical calculations relying on the experimental results are conducted in Tampere University of Technology (TUT). The dielectric materials were manufactured at the Technical Research Centre of Finland (VTT), so, I was fortuned to have a strong interdisciplinary collaboration network providing support. Especially I want to thank Dr. Markus Takala for the introduction to the field of electrical testing and also being so good company. MSc. Hannes Ranta is also sincerely acknowledged for the dielectrical testing. I am grateful for MSc. Tommi Kortelainen to have the patience and wisdom to help me understand the great mysteries on computational chemistry; it was truly a pleasure to work with you. I would like especially want to thank late Assoc. Prof. Kari Kannus for encouraging me to try to go on exchange visit to RPI and of course Prof. emer. J. Keith Nelson and Prof. Linda S. Schadler for giving me opportunity to work there; you were great. I had opportunity to try so many new things! It was a life changing experience for me and my family, and we will never forget it. I will cherish the hours in the basement of Materials Research Center (MRC) with Prof. emer. Robert MacCrone and Prof. emer. J. Keith Nelson and the EPR setup. I got also a chance to work alongside with Timothy Krentz: thank you for all the times you helped me to start all over again after everything seemed gone wrong. Thousands of thanks go to the whole Bullpen at MRC for camping gear and all the support I got during my
visit. I want to acknowledge Dr. Robert Smith for guiding me in pulsed electroacoustic analysis, Michael Goodman for all the time you spent teaching us voltage endurance testing, Michael Topka for the tests we did with “click-reaction”. I sincerely want to thank Michael Bell from University of South Carolina for providing me “long and short brushes”. I will also remember the valuable discussions over the phone with Assoc. Prof. Henrik Hillborg.

I appreciate my colleagues at Physical Chemistry and Nanoscience Center for the nice atmosphere; especially university teachers and Dr. Saara Kaski and Dr. Tiina Kiviniemi, MSc. Päivi Ruokola, Dr. Riikka Reitzer have been a real support to me during these years.

I am truly grateful for my friends and family. The support from my parents; Kari and Tuija has been solid. I want to thank my children Senja and Angus for being born and my husband Ian for love and support.

Jyväskylä, November 2013
Suvi Virtanen
LIST OF ORIGINAL PUBLICATIONS

The main results in this thesis have been reported in the following publications and they are herein referred by their Roman numerals (I-IV).


Author’s contribution

In Publication I the Author has either done or coordinated the structural analysis and written the corresponding part of the paper. In publication II the Author has done or coordinated the experimental part and has written the paper. In publication III the Author has done all the structural characterization and written the paper. In publication IV the Author has done the nanoparticle surface modification syntheses, structural characterization and participated in composite processing and electrical testing. The Author has also written the first draft of the manuscript.
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<thead>
<tr>
<th>Abbreviation</th>
<th>Description</th>
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<tr>
<td>AC</td>
<td>alternating current</td>
</tr>
<tr>
<td>DBS</td>
<td>dielectric breakdown strength</td>
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<tr>
<td>DC</td>
<td>direct current</td>
</tr>
<tr>
<td>DMF</td>
<td>dimethylformamide</td>
</tr>
<tr>
<td>CXμT</td>
<td>computerized x-ray micro-tomography</td>
</tr>
<tr>
<td>NNI</td>
<td>nearest neighbor index</td>
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<tr>
<td>MEK</td>
<td>methyl ethyl ketone</td>
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<td>Om-POSS</td>
<td>octamethyl polyhedral oligomeric silsesquioxane</td>
</tr>
<tr>
<td>PMDETA</td>
<td>pentamethyldiethylenetriamine</td>
</tr>
<tr>
<td>PGMA</td>
<td>polyglycidylmethacrylate</td>
</tr>
<tr>
<td>POSS</td>
<td>polyhedral oligomeric silsesquioxane</td>
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<tr>
<td>PD</td>
<td>partial discharge</td>
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<tr>
<td>PP</td>
<td>polypropylene</td>
</tr>
<tr>
<td>SEI</td>
<td>special energy input</td>
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<tr>
<td>TEM</td>
<td>transmission electron microscopy</td>
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<tr>
<td>TGA</td>
<td>thermogravimetric analysis</td>
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<td>THF</td>
<td>tetrahydrofuran</td>
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1 INTRODUCTION

Considering our present lifestyle, electricity and its distribution are crucial to human society. Electrical insulation materials, i.e. dielectric materials, are a vital part of the components used in the utilization of electrical energy. The shift from ceramic electric insulating materials and oil-paper insulations to polymeric materials has been a major change in the field of high voltage insulation technology during the past three decades. Today, polymers are widely used in high voltage equipment as they can withstand high voltages. They have high breakdown strength, but their charge storage ability is low and they are not as stable over long periods of time as ceramics are. There has been an attempt to combine ceramic materials into polymers to increase their permittivity or gain better thermal and mechanical properties. However, the breakdown strength of the composites is reduced at high filler concentrations, which are needed to improve permittivity and any potential increase in energy density is lost. Polymeric also age and that is not easily predicted in applications. At the beginning of the century, as nanoparticles came onto the market, there opened a possibility to try adding only small amount of ceramic particles, ranging in size from 1 to 100 nm, into polymer. The idea of these so called nanocomposites was to produce a large interfacial area between the matrix polymer and nanofillers, as the interface can act as a trapping site for electrons and prevent breakdown, or ramp up efficiency of the charge storage ability of dielectric material in an electric field.

The main objective of this work was to improve dielectric properties of polymers used in high voltage capacitors by adding inorganic filler into polymers. Simultaneously the structure of the composites was characterized to understand the features in the structure that determine the functionality of the material. Two types of organic polymers were used as a matrix for different types of inorganic fillers, ranging from molecules to colloidal structures. One matrix is thermoplastic; polypropylene, and it is already used in high voltage capacitors. The other is thermosetting copolymer; epoxy that can be cured at rather low temperatures and it is relatively easy to handle at laboratory scale. As it is a good electric insulator, it can be used as a model system with an en-
semble known as bimodal-polymer-brush-grafted particles and focus the study on the charge layer at the interface between nanoparticles and polymers to see if it has an effect on the dielectric properties, as has been hypothesized. In this composite, silica nanoparticles are grafted with bimodal ligands. A short ligand, oligothiophene or ferrocene, controls the electrical properties and a long epoxy compatible ligand, polyglycidylmethacrylate (PGMA), ensures optimal dispersion.

Various characterization methods were used in an attempt to correlate dielectric properties with the multiscale structure of the nanocomposites. The used methods were transmission electron microscopy (TEM), Raman microscopy, optical microscopy, x-ray powder diffraction (XPD) and computerized x-ray micro-tomography (CXμt).
2 A DIELECTRIC IN ELECTRIC FIELD

Dielectric material does not contain free electrons and becomes polarized when it is placed under an electric field. Our present view of electronic transport in polymeric insulators is based on the energy band theory of crystalline solids. Although polymers are amorphous or microcrystalline and have no long-range order that is characteristic of a single crystal, the atoms are arranged like that of a corresponding crystal. The short-range order is expected to cause some of the features in the electronic properties of the crystalline state: the existence of conduction and valence bands separated by a gap and localized states caused by imperfections. There are also extended states present, i.e. an electron in one of these states is not localized but is free to move through the solid with reasonably good mobility.\textsuperscript{11} (see Figure 2.1)

\begin{figure}  
\centering  
\includegraphics[width=\textwidth]{density_of_states.png}  
\caption{Density of electron states as a function of electron energy, and the location of these states through the solid. This image is from reference 11. Copyright \textcopyright1975 IEEE}  
\end{figure}
When an electric field is applied to an insulating material, three basic processes can take place: the dipoles tend to rotate, ions migrate and a space charge can be injected at the interfaces, depending on parameters such as temperature or an applied electric field.\textsuperscript{12} (see Figure 2.2) Polarization in a dielectric material can be classified into electronic, orientational, atomic or ionic and interfacial polarization. The prevailing temperature and frequency of the applied field will determine what process is triggered. At low frequencies, the polarizability of a material has contributions from all the high

\begin{figure}[h]
\centering
\includegraphics[width=\textwidth]{figure2_2}
\caption{Development of a charge distribution $p(z)$ in a dielectric material subjected to an electric field (a) associated with dipole orientation, (b) by ion migration, and (c) by charge transfer at the interfaces. This image is from reference 12. Copyright © 1986 IEEE.}
\end{figure}
frequency mechanisms that can maintain the speed of an alternating electric field, which results in an increase in the dielectric constant.\textsuperscript{13} (see Figure 2.3) The complex dielectric constant $\varepsilon$, that can be divided into real and imaginary parts $\varepsilon'$ and $\varepsilon''$, respectively, describes the amount of interaction that the material has with the electric field, also known as electric permittivity. The imaginary part is associated with the amount of energy that is absorbed by the material and is called dielectric loss. Permittivity or dielectric constant are terms used for the real part and are used for all frequencies of the electromagnetic spectrum; for visible light, the term refractive index is commonly used.\textsuperscript{14} There is also time required for the polarization to reach its maximum, like a viscous lag. This is called dielectric relaxation and is present when an electric field is applied or removed and must be taken into account when considering frequency variation of permittivity.\textsuperscript{15}

Figure 2.3 Schematic diagram of $\varepsilon'$ and $\varepsilon''$ as a function of frequency. This image is from reference 13. Copyright © 1937 Alcatel-Lucent.
The effects of polarization phenomena can be discussed using electromagnetic theory. One way dielectric polarization at a macroscopic scale is easily illustrated is to place the material between two parallel conducting plates, having area $A$ distance $d$ apart, and apply a constant potential difference between them; this system is called a capacitor. The capacitance of this system can be measured and it is the property of a capacitor to store charge under an electric field. If edge effects are neglected, the electric field is inversely proportional to the distance between the plates. A charge stored can be represented as an increase of the field compared to the situation without the dielectric material. This system of charges is neutral and possesses a dipole moment:

$$\mu = AE\varepsilon_0(\varepsilon - 1)d$$  \hspace{1cm} (3.1)

where $\varepsilon_0$ is the electric permittivity of a vacuum.

Induced dielectric polarization can be represented as the dipole moment per unit volume:

$$P = \frac{\mu}{Ad} = E(\varepsilon - 1)\varepsilon_0$$  \hspace{1cm} (3.2)

The term $(\varepsilon - 1) = \chi$ is called susceptibility.\textsuperscript{16}

The polarization charges near the interface of the two media make a potential well in which a charge carrier can be trapped (see Figure 2.4). It is important to note that from this electrostatic point of view, electrons and holes both can be trapped in the same region of the medium of lowest susceptibility.\textsuperscript{17}
For each material under an electric field there is a certain voltage when current starts to increase suddenly, this is known as the breakdown voltage and material is permanently damaged. There is no uniform theory that could explain this phenomenon, and it can be considered to be a sum of different mechanisms. Most current theories relate breakdown to a charge accumulation and displacement of one type or another. Dielectric breakdown is always a thermal process, one way or the other; however, breakdown mechanisms are often still categorized as either an electronic or a thermal breakdown. Breakdown happens rather quickly, usually in ~10 ns in a solid, and is reminiscent of an electron avalanche process that is a well-accepted mechanism for breakdown in gases. Using an electronic breakdown approach, an intrinsic breakdown voltage can be calculated for each material, but solid material cannot survive these kinds of field magnitudes. There is always gas or impurities inside the solid that can cause high local fields in reasonably low voltages, and erosion due to partial discharges can occur, which will eventually erode the whole material. This can be detected from thicker samples as destructed patterns called electrical trees. The geometry of a sample will have an effect on how the breakdown proceeds, or at least how it will be categorized. The electrostatic attraction force between the electrodes can result in an electromechanical breakdown. When it is identified that a material deteriorates chemically due to environmental effects, the mechanism is called electrochemical breakdown. Thermal breakdown is caused by heat generation due to polarization losses and it happens when the material fails to dissipate heat and there is enough thermal energy to release electrons to a conduction band.
3 DIELECTRIC NANOCOMPOSITES

Dielectric nanocomposites are dielectric polymeric materials that are doped with small amounts of inorganic particles, with a size of 1 to 100 nm, to produce a large interfacial area between the matrix polymer and the nanofillers and, hence, to tailor the dielectric properties. The idea is to use very low amounts of the filler to avoid any percolation. Percolation threshold is a limit of the formation of long-range connectivity in random systems. Below the threshold, particles are not connected to each other. The aim is, for example, to increase the material’s dielectric breakdown strength, or the dielectric constant in combination with low loss. It is suggested that nanoparticles could increase breakdown strength by hindering the electrical tree growth within the material.

There is considerable interest in using these new materials in capacitor applications. Currently, there are three major models describing the effect brought on by interfaces. Lewis’s model concentrates on electrical properties of interfaces in a way that the particle introduces the interface region to the polymer based on overlapping Gouy-Chapman layers. Louis Georges Gouy and David Chapman introduced a diffuse model of the electrical double layer, in which the electric potential decreases exponentially away from the surface to the fluid bulk. Tsagaropoulos’ model tries to explain the chemical and physical features of interfaces based on differences seen on glass transition temperatures and assumes that interfacial regions are formed as the polymer bonds, loosely or tightly, to the particle. Tanaka’s multicore model is the most recent and tries to combine these two preceding models.

There are experimental results that suggest that it is not enough just to control the dispersion of particles. In order to improve the charge storage ability of the material, the relative polarity of the particle surface ought to be controlled as well. The added interfacial control, directly bonding the particle to the polymer matrix, can prevent conductive percolation across particle surfaces and increase dielectric breakdown strength. Furthermore, significant reduction in leakage currents and dielectric losses as well as improvement in dielec-
tric breakdown strength has resulted when electropositive, or electron-withdrawing functional groups, were located at the particle surface.\textsuperscript{38}

There are several scales that must be considered when trying to elucidate the relationship of structure of nanocomposite to its function: molecular, intermediate and micro-scale. (see Figure 3.1) The local interactions, molecular conformations and the properties of the matrix material need to be considered within the interphase regions. Intriguingly, what still is unanswered is how far interphase regions extend into the matrix and what the properties of these regions are. Also, it should be taken into account whether the interactions that are predicted to occur between the matrix and the nanofiller tend to promote dispersion, or are they even preventing it to happen. Many questions still remain without clarification.\textsuperscript{39}

\begin{figure}[h]
\centering
\includegraphics[width=0.5\textwidth]{figure3_1.png}
\caption{Schematic diagram indicating the range of different dimensional levels that need to be considered when attempting to characterize a nanocomposite. This image is from reference 39. Copyright ©2008 IEEE}
\end{figure}

Nanoparticles can be used in both cross-linked thermosetting polymers and in semi-crystalline thermoplastics. Potentially, when using thermoplastics that are recyclable, it is possible to achieve more applicable and economical insulation structures.\textsuperscript{40} The heterogeneous nature of the composition of these polymer nanocomposites brings challenges and requires the development of processing methods\textsuperscript{10} and characterization; in practice, it is difficult to assign morphology to an increase in DBS.\textsuperscript{41} Also, relatively high filler loadings, such as 10 wt-% TiO\textsubscript{2} particles in epoxy, are reported to be optimal in order to gain good dielectric breakdown strength.\textsuperscript{42} To develop these new materials for industrial use, a fundamental understanding of dielectric properties at the nanoscale level is of great importance. There are many ways to mix particles to the matrix polymer. Nanoparticles are usually coated with polar or non-polar organic molecules to make their surface more compatible with the matrix. It was found that dry mixing of these coated particles is problematic; once nanoparticles are dried, they inevitably agglomerate and become very difficult to disperse as they have a very high surface to volume ratio. The harsh mixing condi-
tions needed to redisperse the nanoparticles may degrade the properties of the matrix polymer. \(43-50\)

In the following chapters, the components of composites used in this study are presented. The used nanofillers range from molecular to colloidal, and surface treatments used are either surfactant or covalent modification.

### 3.1 Melt compounding: polypropylene (PP)

Polypropylene is a mildly polar hydrocarbon and identical macromolecules of it do not exist.\(^51\) It is semi-crystalline consisting of lamellaes that form birefringent spherulites that grow until they touch another spherulite. Spherulites are aggregates of crystals with a radial fibrillar structure.\(^{52-54}\) (see Figure 3.2) Three different crystalline phases are found to exist in PP depending on the crystallization conditions used: \(\alpha\) (monoclinic), \(\beta\) (hexagonal), \(\gamma\) (orthorhombic).\(^{55, 56}\) Polymer crystallization has been studied since the 1960s, and yet no molecular theory exists. It is controversial whether preordering occurs before the onset of crystallization or not.\(^57\) Partially crystalline polymers have both a glass transition state and a melting point. Below glass transition temperature, molecular motion is restricted to the rotation of side groups or their parts.\(^16\)

Thermoplastic polymer composites are processed physically using extrusion. Using this approach, good dispersion of the filler can be achieved, as polymer gains increased mobility through the input of thermal energy, and the nanofillers are mechanically dispersed and mixed under the influence of shear forces. When the effect of processing conditions was studied, it was noticed that a specific energy input (SEI) is related to the shear intensity, which correlates closely to the degree of mixing.\(^{58}\) Small crystal size and low defect amounts are related to improved DBS of PP.\(^{59}\) There is still a need for resolving the connection between the morphology of PP and its DBS.\(^{60}\)

![Figure 3.2](image-url) **Figure 3.2** Stages in the formation of a spherulite from a stack of lamellae; a side view and TEM micrograph of spherulites; a top view.
3.1.1 Molecular filler: octamethyl polyhedral silsesquioxane (Om-POSS)

Cage-like silsesquioxanes are usually called polyhedral oligosilsesquioxanes or polyhedral oligomeric silsesquioxanes (POSS). Interest in these molecules has been increasing since the early 1990s, and attempts have been made to try to understand the structure-property relationship of the molecules blended into polymers. This class of highly symmetric molecules usually features a nanoscopic size, approximately 1.5 nm in diameter when the vertex (R) groups are included. These molecules have a hybrid inorganic core/organic shell structure (see Figure 3.3) They can be regarded as the smallest possible silica particles. POSS-based materials, including fire retardants, biomaterials, dielectric materials, organic light-emitting diodes, lithography resists, catalysts, membrane fuel cells and battery membranes have recently been introduced.

Octamethyl POSS (Om-POSS) is a molecule that has methyl vertex groups. It is a non-reactive, crystalline solid. Physical mixing is believed to be an efficient method to disperse POSS in the polymer matrix mediated by van der Waals or hydrogen-bonding interactions. However, proof of molecular level dispersion has not been reported using this mixing approach; instead, there is evidence for strong particle-particle interactions. Additionally, it has been reported that in PP, om-POSS acts as a nucleating agent and is mixed mainly as microcrystals. It must be taken into account when processing om-POSS that it sublimes at elevated temperatures (>200°C) under a nitrogen atmosphere. Phase changes inside the polymer is a probable cause for observations of the average crystal size of om-POSS decreasing as it is mixed into the polymer matrix. The om-POSS as a filler in PP has the ability to increase the dielectric breakdown strength of PP and epoxy.

![Figure 3.3](image_url) Chemical structure of om-POSS molecular nanofiller. This image is from reference 71. Copyright © 2007 Wiley Periodicals, Inc.
3.1.2 Crystalline filler: surfactant coated nano-CaCO$_3$

Synthetic, so-called precipitated calcium carbonate is commonly produced by a recarbonizing process in which natural calcium carbonate is decomposed to calcium oxide and carbon dioxide. Calcium oxide with water forms calcium hydroxide that reacts with carbon dioxide and as a result pure, synthetic calcium carbonate precipitates. It is commonly used in large mixture ratios to reinforce polypropylene. The improved toughness is related to the $\beta$-nucleating effect of CaCO$_3$ and the increase in crystallinity of PP. It is well excepted that improvements in mechanical properties are related to good dispersion of the filler. The $\beta$-nucleation efficiency and matrix interactions are dependent on the surface treatment of the particles; it has a crucial role in determining the properties of the composite. When considering the composite processing, aggregated nano-CaCO$_3$ causes anisotropic polymer chain orientation in extrusion and introduces porosity into PP films under biaxial stretching. This is detrimental to the application. Dielectric breakdown strength of neat polypropylene is reported to decrease if it contains 0.1-1.0 nm diameter voids. CaCO$_3$ hardly affects the flammability of PP, but it is found to increase its thermal conduction.

There are only a few studies of the dielectric properties of CaCO$_3$ polymer composites and the increase of permittivity and losses of composites are explained by the moisture absorbed by CaCO$_3$. The loading of 3 wt-% of the surfactant modified CaCO$_3$ in PP has not shown any increase of permittivity. For CaCO$_3$ filled PP composites, a raised AC breakdown strength has been reported. Dielectric spectroscopy measurements have shown that the smaller particle size raises the permittivity compared to pure continuous CaCO$_3$ at low frequencies.

3.1.3 Amorphous filler: surface modified fumed silica

Fumed silica, also known as pyrogenic silica as it is produced in a flame, can have a very small particle size (2 nm) and if properly dispersed would provide a very high surface/volume ratio and large interfacial volume. It is also a cheap alternative as a filler for possible up-scaled application. Silica is an excellent electrical insulator. Amorphous silica has been studied to see whether it can act as an electron trap under electron bombardment. Fumed silica decreases the size of spherulites and the crystallinity of PP in nanocomposites. Silica has been mixed into polymers together with other fillers. Even considerable improvements in voltage endurance of polymers filled with fumed silica have been reported. For example, 12.5 wt-% vinylsilane-treated fumed silica has improved the voltage endurance of cross-linked polyethylene, a composite useful for cable applications. For this particular material, the degree of crystallinity was claimed not to be a predominant factor in increasing breakdown strength and the highest voltage endurance was achieved for composites believed to have a covalent bonding between the matrix and the filler. Below 10 wt-% fumed silica in biaxially oriented blown PP films have shown
breakdown strength, energy density, corona resistance and mechanical properties such as impact strength, tensile strength and ductility comparable to or superior to blown films of the un-filled polymer composition. There are, however, no indications that fumed silica would have considerably improved the permittivity of the polymer composites.

### 3.2 In Situ polymerization: Epoxy

Epoxy is the most abundant thermosetting polymer. It is highly cross linked, glassy and brittle at room temperature. The reaction between an epoxy resin and a hardener is an irreversible exothermic poly-addition, i.e. no by-products are formed, and the epoxy plastic cannot be decomposed into epoxy resin and hardener. There are more than 50 different substances that fulfill the definition for an epoxy resin. The epoxies are versatile as there are several hundred different hardeners available. Thermosetting composites with epoxy are made by in situ polymerization (see Figure 3.4) If the nanofiller can be mixed with the resin in a non-dried form, it is better for particle dispersion; however, particles that solvate to solution are not necessarily compatible with the matrix, and the solvent might also deteriorate the matrix polymer or introduce undesirable ions to the composite.

Figure 3.4 Typical curing, i.e. cross-linking reactions, of epoxy resin using an aliphatic diamine curing agent below 150°C. This image is from reference 111. Copyright © 1999 John Wiley & Sons, Inc.

#### 3.2.1 Colloidal filler with electro active surface: bimodal-polymer-brush-grafted silica

A polymer brush is simply a polymer chain tethered to a surface. They have been studied for over 30 years. There are two methods used to produce these brushes: the “grafting to” approach, where chains are polymerized and subsequently attached to the surface, and a “grafting from” approach, used where polymerization takes place from a site on the surface of a particle. The
latter has the advantage of achieving higher graft densities since previously
grafted chains do not inhibit the attachment of additional chains. “Grafting to”,
used in this study, is quick and easier to scale up and provide some flexibility in
the chemistry of the attached molecule.\textsuperscript{112} The limitation is in the lack of control
over graft density, but good dispersion of particles has been modeled using this
approach,\textsuperscript{113} and a phase diagram has been experimentally validated to predict
the dispersion for bimodal-polymer-brush-grafted nanoparticles.\textsuperscript{114} “Grafting to”
can be done using so called “click” chemistry,\textsuperscript{115, 116} where the idea is that it has
a high thermodynamic driving force. This type of reaction proceeds rapidly to
completion and also tends to be highly selective for a single product. The anal-
ogy would be a spring loaded for a single trajectory.\textsuperscript{117} This approach has been
reported for functionalization of silica\textsuperscript{118-122} and also used as a way to make so
called “matrix free” silica polymer composite by using alkyne and azide modi-
fied polymer brushes on silica.\textsuperscript{123} PGMA has been used for dispersing high
loadings of TiO\textsubscript{2} to epoxy with a high refractive index and transparency.\textsuperscript{124} Elec-
tro-active alkyne terminated oligoaniline has been previously attached to poly-
mers on silica particle surfaces retaining its electro-activity.\textsuperscript{125}
4 EXPERIMENTAL METHODS

4.1 Nanoparticle surface coating and composite processing

The filler materials used in this study were om-POSS, CaCO$_3$ and SiO$_2$. Om-POSS was used without any surface functionalization with an average crystal size of 170 nm. CaCO$_3$ was precipitated as 60 nm particles with stearic acid as a surface coating. In PP, SiO$_2$ was 12 nm fumed silica with a hydrophobic surface. In epoxy systems, SiO$_2$ was added as a solution of 15 nm colloidal particles with bimodal-grafted-brushes tethered on the surface. This ensured a high density of short chains with tailored dielectric behavior, and a low density of long chains that are compatible with the matrix. The PP composites were mixed using extrusion. First PP and the filler were premixed in dry form and subsequently melt mixed. The epoxy composites were mixed in shear mixer, degassed and cured in silicone molds for 10 h in room temperature followed by 10 h in 60°C and finally 4 h in 100°C.

4.1.1 Bimodal-polymer-brush-grafted silica

Silica particles were modified using copper(I)-catalyzed Huisgen 1,3-dipolar cycloaddition of azides and terminal alkynes, [3+2] cycloaddition (CuAAC) reaction.$^{126}$

A suspension (40 ml) of 30 wt-% colloidal silica particles was added by funnel into a dried 250 ml two-necked, round-bottom flask with 3-(Chloropropyl)-trimethoxysilane (1.18 g, 6.0 mmol) and dried THF (40 ml, 72 h over a molecular sieve). The reaction mixture was refluxed in 100°C oil bath under an argon balloon overnight and then cooled to room temperature. The reaction mixture was precipitated into a large amount of hexanes (500 ml). The particles were recovered by centrifugation at 3000 rpm for 15 min. The particles were then dissolved in 20 ml of acetone and precipitated in 200 ml of hexanes. The functionalized particles were dispersed directly into 50 ml of DMF for subsequent use. An aliquot of the functionalized silica nanoparticles was dried and
subjected to thermal gravimetric analysis to determine the amount of silane agent anchored onto the particles.\textsuperscript{127}

3-(Chloropropyl)trimethoxysilane functionalized silica particles (<0.59 g (3 mmol)) and sodium azide (<0.3 g, 5.52 mmol) were added into 100 ml round bottom flask in DMF and refluxed in 100°C oil bath for 4 h. After the reaction, excess sodium azide was washed with deionized water three times. The particles were dissolved in THF to form a clear solution. An aliquot of TGA was taken (weight loss due to linker was 5-6%) and concentration was determined (mg/ml).

0.04 g ethynylferrocene or 0.08 g alkyne terminated oligothiophene [1 equiv] and 0.2 g alkyne terminated PGMA [1:10 equiv], (see Figure 4.1) 0.8 g functionalized particles; azide [1 equiv], and PMDETA (40 µL [0.5 equiv]) was added in 40 ml of THF. The mixture was degassed by bubbling argon gas for 5 min to get rid of oxygen before adding CuBr (approx. 14 mg [0.5 equiv]). CuBr was purified with glacial acetic acid and washed in ethanol before use, bubbled with argon gas for additional 5 min, and then stirred for 24 h. Particles were precipitated by methanol and centrifugation (4000 x g 10 min) and then mixed into bisphenol-A based epoxy resin and solvent residue was evaporated before adding the aliphatic amine based hardener.

\begin{figure}[h]
\centering
\includegraphics[width=\textwidth]{figure.png}
\caption{Chemical structures of entities tethered on the silica surface by click reaction: a) alkyne terminated oligothiophene b) ethynylferrocene c) alkyne terminated PGMA}
\end{figure}
Figure 4.1 Surface modification process of silica to supply bimodal brush grafted core functionalized nanoparticles. Colors indicate the reactive parts and product of used grafting to “click” reaction.

Functionalized silica was analyzed with TGA, infrared spectroscopy (IR) and UV-vis spectroscopy after each processing step. (see Figure 4.1) Attachment of each linker to particles was detected by IR: PGMA can be confirmed from a C=O vibration at 1733 cm$^{-1}$ and a characteristic peak of an azide group is seen around 2100 cm$^{-1}$. The completion of “click” reaction was verified by the disappearance of the azide peak. Attachment of electroactive molecule was verified by UV-vis spectroscopy: oligothiophene at 360 nm and ferrocene at 440 nm. The grafting density $\sigma$ was calculated from the weight loss ratio determined by TGA:

$$\frac{z}{(1 - z)} = \frac{w_{\text{polymer}}}{w_{\text{silica}}} = \frac{w_{\text{polymer}}}{\rho N_\varnothing \frac{4}{3} \pi r^3}$$

(4.1)
The grafting density $\sigma$ is simply number of entities grafted onto silica particles: the number of grafted chains ($w_{\text{polymer}} N_A / M_n$) is divided by the surface area of colloidal silica ($4\pi r^2 N$). silica (chains/nm$^2$):

$$\sigma = \frac{w_{\text{polymer}} N_A / M_n}{4\pi r^2 N} = \frac{r \rho z \times 10^{-21} N_A}{3(1-z) M_n}$$

(4.2)

where $w$, $N_A$, $N$ and $z$ are the weight, Avogadro’s number, the number of particles and weight loss of polymer chains respectively. Residue weight (1-$z$) is the fraction of silica. Supposing that the SiO$_2$ particles are spherical with a radius $r$ of 7.5 nm and density $\rho$ of 2.2 g/cm$^3$ for amorphous quartz, and the molecular weight of the polymer $M_n$ was 15000 g/mol; the grafting density of PGMA was estimated to be between 0.02-0.05 polymer chains/nm$^2$. When both a short and long brush was attached simultaneously, the estimation of achieved grafting density could be determined since oligothiophene decomposes right after 250°C and PGMA decomposes after 400°C.

### 4.2 Structural characterization

There are two scales to be considered, micro and nano, if the dispersion of fillers in original structure of PP composites was to be illustrated. Mostly, the structural characterization was done using two techniques: Raman microscopy and TEM. Two kinds of samples were used: cast film samples that have an approximate thickness of 500 µm and biaxially oriented films with a thickness of around 20 µm, called oriented samples.

#### 4.2.1 Raman imaging

For the Raman analysis of cast film samples, a dispersive micro-Raman spectrometer (Bruker Senterra R200-785) equipped with a linearly polarized diode laser (785 nm) in back-scattering geometry was used.

For confocal imaging of om-POSS composite, the pinhole used was 25 µm with a 100x objective (NA 0.90). The laser power was kept low to avoid over-heating the sample (25 mW). At each point, the measurement time was 10 s. The area of interest in the sample was scanned with a motorized stage of 0.1 µm accuracy in two or three dimensions using 0.5 µm steps.

For coarse imaging of CaCO$_3$ composites, the laser power was 100 mW and a slit width of 50 x 1000 µm was used with a 50 x objective in order to get a good signal from larger area. This way, it could be determined whether the dis-
persion of nano-CaCO$_3$ is homogeneous over large areas (~200 µm). The area of interest in the sample was scanned with a motorized stage with 0.1 µm accuracy in two dimensions, using a 3 µm step size. The spectra of composites were measured from a 210 x 210 µm$^2$ area. Four areas of each sample were then scanned and visualized as an intensity contour plot of a selected characteristic Raman peak of CaCO$_3$ and PP. The PP signal was used as a reference to check that the overall signal level was stable over the measured area. The analysis was done for a PP composite sample series containing 1.8 to 8.1 wt% nano-CaCO$_3$. Altogether, there were samples of eight different nano-CaCO$_3$ concentrations. Pure PP, processed similarly as the composites, was used as a reference. Both cast and oriented film samples were characterized.

4.2.2 Transmission electron microscopy

TEM analysis can be used to analyze both micro- and nanoscale particles. This method is the preferred method for nanoscale dispersion studies in cast films. The PP composites are quite soft materials and, therefore, it is challenging to obtain good quality, ultrathin sections for TEM imaging. Especially, the contrast of silica nanoparticles with respect to the PP matrix is also a limiting factor for quantitative TEM analysis. On the other hand, the microcrystallinity of PP can be used as an internal standard to ensure that the structure studied is really in original form. The quality of thin sections has been the limiting factor in acquiring good quality images. For TEM analysis, ultrathin sections of the cast film samples or epoxy blocks were obtained using a Diatome 35° diamond knife at room temperature with a Leica Reichert Ultracut ultramicrotome. The 50-90 nm thick sections were collected on a 400 mesh copper grid and imaged with JEOL-JEM-1200EX or 1400 electron microscope with 2.5 - 75 k magnification. Statistical analysis was performed using images that were converted to a binary format. For PP composites, the particle area and center of mass were determined using the ImageJ® program. Microparticles were excluded from this analysis. The data was further processed to emphasize the distribution of the filler material in the matrix media; the relative volume of the agglomerates was estimated assuming that they are spherical objects. (see Figure 4.2)
Dispersion quantification was performed for PP composites using two methods with numerical descriptor: the quadrat method (skewness), which measures the uniformity of the spacing of particles and agglomerates, and the 1st nearest neighbor index (NNI), which indicates how well individual particles are separated from one another. These two methods are all relatively insensitive to volume fraction of filler and primary particle size since clusters are treated as point objects and they can be considered to give a value of the quality of dispersion.

The quadrat method divides the sample into cells, counts the number of particles in each cell and reports the skewness of the particle distribution:

$$\frac{N}{(N-1)(N-2)} \sum_{i=1}^{N} \frac{(x_i - \bar{x})^3}{\sigma}$$  \hspace{1cm} (4.2)

where $x_i$ denotes the number of particles in the $i$th quadrat, $\bar{x}$ is the mean of $x_i$ ($i = 1, 2, \ldots, N$), and $\sigma$ is the standard deviation of $x_i$. Skewness evaluates the asymmetry of distribution. When counting particles in quadrats, the size of the quadrat was selected to be two times the average diameter of agglomerates, as is commonly used.\(^{130}\) When a reasonable size is selected for the quadrats, the number of particles tends to be the same in each quadrat when the particles are well mixed, resulting in a skewness that is close to zero. Aggregates will result in a large number of empty or less populated quadrats and skewness will be then non-zero. The higher the deviation is from zero, the poorer the dispersion.

In the NNI analysis, each agglomerate or particle is treated as one body, and the mean distance to the closest agglomerate or particle $\langle L_n \rangle$ is compared to the distance expected for an equilibrium distribution $\langle R_n \rangle$ of particles having the same diameter and volume fraction.\(^{131}\) NNI is simply the average of the actual closest neighbor distances divided by the randomly distributed case:
NNI \( >1 \) indicates regularity and \( <1 \) indicates clustering. In previous studies of nanodielectrics NNI values of \( >0.7 \) and a low absolute value of skewness (\( <1.5 \)) are related to improved DBS.

For epoxy composites, using MATLAB\textsuperscript{®}, free-space length \( L_f \) was used to quantify dispersion from binary images. The free-space length is defined as the width of the largest randomly placed square for which the most probable number of intersecting particles is zero. This parameter is well-suited for comparing different loadings as the measurement is a physical quantity; the distance between particles. But it is not relative to random dispersion as for example NNI is. For the same filler loading, the further the dispersion is from uniform, the longer the \( L_f \) becomes (see Figure 4.3). The code was retrieved and used according to the literature.

![Figure 4.3](image_url) Varying dispersions of 30 nm particles at 6% loading: a) uniform dispersion \( L_f = 73 \text{ nm} \); b) random dispersion \( L_f = 106 \text{ nm} \); c) clustered dispersion \( L_f = 220 \text{ nm} \); d) agglomerated dispersion \( L_f = 460 \text{ nm} \). Open squares have a length equal to the free-space length which increases as the dispersion worsens. This image is from reference 131.

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4.2.3 Other methods

As there were high density differences between filler particles and the matrix, computerized X-ray micro-tomography (CXμT) could be used to study the micron-scale 3D structure of composites. The resolution of the technique is 0.9 μm. CXμT techniques, based on synchrotron radiation sources, have been used in probing the internal 3D structure of materials for decades. The method has become more feasible as table-top tomographic scanners utilizing x-ray tubes become available. Although such scanning techniques were initially developed for the medical field at a macro-level, they are increasingly being adopted in various research sectors on microstructure analysis.\textsuperscript{135, 136} In CXμT, a 3D image of the sample is calculated from the transmission or reflection data collected by illuminating the object from many different directions. The scanner uses a continuous incoherent x-ray beam to produce 2D shadowgraphs from multiple angles. Reconstruction of the shadowgraphs produces cross-sectional images of the absorption coefficients of the x-ray beam.

For describing the average crystallite size of the constituents, powder x-ray diffraction was used. The average crystallite sizes of PP and fillers were estimated using the simple Scherrer method that gives the average crystal size of the crystallites in direction normal to the lattice planes. The method is based on the broadening of diffraction peak profiles as a function of decreasing size of the diffracting crystals.

4.3 Dielectric properties

4.3.1 Dielectric breakdown strength (DBS)

The DC breakdown strength is a very important feature and, therefore, was measured from all the composite samples. It was measured for PP from biaxially oriented films and for epoxy composites it was measured from recessed samples to gain quasi-uniform electric field.\textsuperscript{137} In oriented films, PP sample breakdown results were fitted with a 2-parameter Weibull distribution, whose cumulative distribution function is given by

\[ F(x) = 1 - \exp \left( - \left( \frac{x}{\alpha} \right)^\beta \right) \]

where \( F(x) \) is the probability of breakdown at a certain electric field strength \( x \) (kV/mm). The scale parameter \( \alpha \) is related to the 63.2% probability of breakdown at field strength \( \alpha \), and the shape parameter \( \beta \) describes the shape of the distribution. The higher the value of \( \beta \), the narrower the spread of indi-
Individual breakdown strength results are. The $\beta$ is also known as shape factor. The shape of the distribution and deviation between parallel samples for each composite can be visualized by percentiles, corresponding to the probability of breakdown at the respective electric field magnitude.\textsuperscript{138, 139}

When recessed epoxy samples were analyzed, 3-parameter Weibull distribution was found to fit the results better.

$$F(x) = 1 - \exp \left\{ - \left( \frac{x - t}{\alpha} \right)^\beta \right\}$$  \hspace{1cm} (4.5)

This three parameter Weibull is commonly used in failure analysis.\textsuperscript{140} There is an obvious threshold value, $t$, required for breakdown to happen.

\textbf{4.3.2 Dielectric spectroscopy}

Dielectric permittivity and loss factor of the materials were determined by measuring the complex impedance of the samples. The real and imaginary parts of the complex permittivity ($\varepsilon'$ and $\varepsilon''$) and relative complex permittivity ($\varepsilon_r$) were calculated from the measured parallel capacitance and resistance:

$$\varepsilon' = \left( \frac{C_p - C_e}{C_0} \right); C_0 = \frac{\varepsilon_0 A}{d}$$  \hspace{1cm} (4.6)

$$\varepsilon'' = \varepsilon' \tan \delta = \frac{\varepsilon'}{2\pi R_p C_p f}$$  \hspace{1cm} (4.7)

where $C_p$ and $R_p$, respectively, are the measured parallel capacitance and resistance at frequency $f$, and $C_e$ is the error capacitance due to small electric field distortions at the electrode edges.\textsuperscript{141} $C_0$ is the vacuum (geometric) capacitance defined by electrode area $A$, $\varepsilon_0$ is the permittivity of a vacuum, $d$ is the average thickness of the sample, and $\tan \delta$ is the dielectric loss tangent.
5. RESULTS AND DISCUSSION

Studies of nanodielectrics rarely concentrate on structure, as it is challenging to determine properly. Vast parts of the literature also lack understanding about the interfacial effect. Results published seem to vary tremendously, and all findings must be restricted to each composite in question. There seems only a little dependence of the improvement of properties on nanoparticle concentration. This can be explained by difficulty of mixing particles with high surface energy evenly into the polymer matrix, as particles will form agglomerates and the interfacial area decreases. It is also difficult to study one particular property of dielectric material as the results can be dependent on sample preparation, rather than properties of the material itself. Also, the precursors are commercially made and their whole content is often not revealed entirely. For possible future applications, it is still important to understand the detailed composition of the material in order to avoid possible drawbacks caused by microparticles formed from possibly agglomerated nanoparticles. According to structural characterization done in this work; it seems that reason why poorly dispersed nanoparticles still improve the dielectric properties of the polymer is the increase in the stability of the composite caused by structural changes that the filler imposes on the composition and structure of the polymer. But what is also found out is when the filler is truly dispersed at nanoscale, it seems that the increased polarizability introduced by nanofiller, plays a major role in improving the dielectric properties of the material.

5.1 PP composites

This section briefly describes the main results of the work presented in Publications I-III. The main motivation for these studies was to perform structural characterization of various composites and correlate that information to dielectric properties in order to gain insight into the structure-function relation-
ship. Challenging for the structure property relations of these materials were that the DBS level of reference material was quite low compared to values for different industrial PP capacitor films indicating that there might be some other effects than caused by fillers, that are also important for gaining preferred material properties.

5.1.1 Molecular filler: om-POSS

The molecular filler om-POSS was found to be dispersed in PP mainly as agglomerates formed from nanocrystallites (see Figure 5.1) From CXμT, the concentration of om-POSS at the microscale was estimated at 2.3 wt-%, as the added amount of om-POSS was 3 wt-%, so only a small portion from the added filler was properly dispersed as nanoparticles. Microagglomerates were found to be elongated and evenly distributed in the matrix. (see Figure 5.1b) The diameter distribution extends from about 1 μm to 16 μm and the most probable diameter is approximately 6 μm.

Crystal agglomerates of om-POSS were formed when the PP melt was cooled, as the average crystallite size determined from powder x-ray diffraction (XPD) was unexpectedly diminished during processing, from 170 nm to 40 nm, apparently due to om-POSS sublimation-condensation phase transition during melt mixing at ~200°C. TEM reveals that there is also the polymer seen inside these crystal structures. (see Figure 5.2a) The XPD patterns are presented in Figure 5.2b. Based on the XPD pattern of PP, the crystalline zones in a sample are mainly in an α-crystal phase. However, at least one weak diffraction peak at ~16.0° 2θ is observed only in PP sample, which indicates the presence of a low fraction of the β-phase which is missing from composite, so the addition of om-POSS will create differences in the microstructure of PP.
Figure 5.1 Dispersion in 3 wt-% om-POSS PP composite: a) TEM image and b) CXμT visualization of the particle phase. Physical dimensions of the sample are (1.760 x 0.813 x 0.339) mm. (CXμT courtesy of Viivi Nuottajärvi)

Figure 5.2 a) TEM images of an om-POSS particle in composite sample. It can be seen how the polymer material has intruded the om-POSS microcrystal. b) XPD patterns of PP, om-POSS-PP composite and om-POSS powder samples. Diffraction peaks used for crystallite size evaluation are indicated by arrows. The average crystallite size of om-POSS in the composite is 40 nm; showing that in the fabrication process a dramatic reduction of the crystallite size takes place. (XPD patterns courtesy of Manu Lahtinen.)
Figure 5.3 Raman intensity contour plots of an om-POSS particle in a composite generated by integrating intensity of the peak at 157 cm\(^{-1}\) (left image) and at 206 cm\(^{-1}\) (right image). The crosses situated in the intense regions of the contour plots mark the place where the spectra below are collected. The difference in the images reflects the relative difference of the peak intensities in different regions of the agglomerate. This image is from Publication II. Copyright © 2011 NORD-IS

In the confocal Raman studies, it was noticed that relative intensities of the characteristic peaks vary depending on the measurement spot (see Figure 5.3). If different crystals were experimentally measured with changing the polarization of laser and the analyzer; phenomena was further emphasized, and the peak intensities varied tremendously when polarization was changed. (see Figure 5.4a) But, if the measurement volume was increased, the relative peak intensities approached constant value. The theoretical DFT calculations were done by Abinit-code by changing the crystal orientation, in a linearly polarized electric field, and they showed qualitatively similar changes in relative peak intensities.\[\text{III}\] (see Figure 5.4b) A plausible explanation of this effect is the orientation of the crystallites in agglomerates. The tensorial nature of Raman scattering leads to a dependence of the signal intensity on the relative orientation between the electric field of the incoming electromagnetic field and the molecule. Within the confocal volume, there must be orientation in the ensemble of individual crystallites. This effect is not visible when the measurement volume is large enough, compared to the different orientations, to cancel the effect of each other. These findings suggest that Raman imaging can be used to study molecular orientation at quite high resolution. They also indicate that intensity variations in Raman maps\[\text{72}\] may originate partially from orientation effects, in addition to concentration variations.
This material, 3 wt-% of om-POSS, had somewhat increased AC and DC breakdown strength; simultaneously, permittivity was slightly higher than reference PP.\textsuperscript{72} This might be due to that om-POSS effects on the microstructure of PP or also that it might locally increase the thermal stability of PP under electric stress when high temperatures are present. It can act as a synergist in intumescent, a swelling of material, where a barrier is formed that acts as physical and thermal protection against a rapid increase of temperature by limiting heat and mass transfer between the gas and the condensed phase.\textsuperscript{62} The amount of filler dispersed at nanoscale was very low in this composite despite all the effort to achieve good dispersion.

### 5.1.2 Crystalline filler: nano-CaCO$_3$

The PP with crystalline filler CaCO$_3$ was analyzed with coarse Raman imaging in addition to TEM analysis, and composites were found to contain well-dispersed nanoparticles and steadily increasing amounts of microparticles as the filler concentration was increased. (see Figure 5.5) There was no obvious correlation between NNI and skewness to DBS, but the parameters evaluating
the quality of dispersion were both under the limits previously suggested for optimal dielectric properties. These numerical dispersion assessment methods are all relatively insensitive to volume fraction and primary particle size since they consider clusters as point objects. The evolution of these parameters with concentration is complex and no straight correlation can be expected. The maximum increase in DBS was achieved with 1.8 wt-% filler loading. The DBS decreased as the microparticles became more abundant, but the effect of the filler on increasing the permittivity of the material was subtle. (see Figure 5.6) Overall the addition of nano-CaCO₃ to PP did not much change the dielectric properties.

Figure 5.5 Raman contour maps (210 x 210 µm²) of cast samples, generated by integrating the most intense characteristic peak of CaCO₃: a) pure PP b) 1.8 wt-% c) 2.7 wt-% d) 3.6 wt-% e) 4.6 wt-% f) 5.6 wt-% g) 6.4 wt-% h) 7.4 wt-% and i) 8.1 wt-% nano-CaCO₃/PP composite sample. The overall level of the CaCO₃ signal rises as the concentration of nano-CaCO₃ increases in the composites. Shown below are TEM images from the same composites, showing gradually increasing concentration of the filler dispersed at nanoscale.
5.1.3 Amorphous filler: fumed silica

Both AC and DC DBS of the composite increased substantially for 5 wt-% silica composite and the $\beta$-value, also known as the shape factor, is higher than with reference PP. The higher the value of $\beta$, the narrower the spread of individual results; this is an indication of a homogeneous material. Silica was found to disperse in PP as fractal agglomerates, around 100 nm in size. (see Figure 5.7) XPD and Raman studies found that this composite also contains CaCO$_3$, a contamination from an extruder. 3D-characterization of this composite by $\mu$XT revealed CaCO$_3$ microcrystal loading to be 2.3 wt-% in this composite. But CaCO$_3$ alone is not responsible of the increased DBS, so it must be effect caused by silica filler.
There was dramatic enhancement in partial discharge (PD) endurance of the composite compared to neat PP. Optical microscopy shows fractal structure formed in the composite due to PD stress; these might indicate that silica has some role in restricting the degradation of the composite as not as prominent structures were found in neat PP reference at any stress voltage used. (see Figure 5.8)

**Figure 5.8** Optical microscope images of (i) PP and (ii) 5 wt-% silica PP composite: a) unstressed b) 2 kV c) 3 kV d) 4 kV stressed samples. See the fractal structures formed in composite under electrical stress.
5.2 Epoxy Composite

This section describes the results from Publication IV. The motivation was to get fundamental understanding of the role of the nanoparticle interface in controlling the material properties. It focuses on synthesis and dielectric properties of polymeric nanocomposites where silica nanoparticles are grafted with bimodal ligands: a short ligand; oligothiophene or ferrocene to control the electrical properties and a long epoxy compatible ligand (PGMA) that ensures optimal dispersion.

5.2.1 Colloidal filler: bimodal-polymer-brush-grafted silica

A new synthetic approach was used to provide well-dispersed silica particles with electroactive core functionalization in epoxy in order to study the effect of the charge layer at the interface of the nanoparticle and polymer matrix. (see Figure 5.9) Dispersion of the filler using this approach is so effective that most of the materials are optically transparent, (see Figure 5.10) especially composites with bimodal silica with thiophene as its functionality, had outstanding dispersion and electrical properties. Significant increases in the DBS were represented by the 63.2-% ($\alpha$) parameter from the Weibull distribution plus the threshold parameter. 2 wt-% silica grafted with PGMA and oligothiophene and 4 wt-% silica grafted with PGMA and ferrocene provided an increase in the 63.2% average DBS $\alpha$ of greater than 40%. The free-space length, $L_f$, and DBS $\alpha$ 63.2% average with shape factor are shown in Table 5.1. The increase in DBS is well correlated to a decrease in $L_f$. For example, a comparison between the bare silica and the PGMA modified silica shows a significant improvement in DBS corresponding to a large decrease in $L_f$. It seems that distribution of electroactive nanofiller must be dense enough to give the greatest improvement in DBS. If dense enough distribution of particles was achieved, then the DBS increased considerably. (see Figure 5.11 and table 5.1)

It is obvious that the electroactivity at the silica particle core is largely responsible for the improvements to the DBS. Even if similar good dispersion is achieved with PGMA modified silica filler, the small improvement in DBS indicates that the silica filler itself is not alone responsible for the improvements seen in the modified filler composites. Another indication is that even if the particle dispersion is not optimal, as in silica-PGMA-ferrocene composites, improvement is still seen in DBS.

Promising for future applications is that permittivity increased 25% with only 2 wt-% filler loading as dielectric loss stayed at the level of reference epoxy. If there were also microaggregates in the material, as in 4 wt-% silica-PGMA-ferrocene, low frequency losses increased. (see Figure 5.12)
Figure 5.9 Schematic model of bimodal polymer grafted core functionalized colloidal silica particle.

Figure 5.10 Visual transparencies of some epoxy composite blocks with different silica contents/surface modifications. From left: 1 wt-% bare silica, 2 wt-% silica-PGMA, 2 wt-% silica-PGMA-oligothiophene and 4 wt-% silica-PGMA-ferrocene in epoxy.
Figure 5.11 TEM images of epoxy composites with: a) 1 wt-% bare silica b) 2 wt-% silica-PGMA c) 2 wt-% silica-PGMA-oligothiophene d) 1 wt-% silica-PGMA-ferrocene e) 2 wt-% silica-PGMA-ferrocene f) 4 wt-% silica-PGMA-ferrocene

Table 5.1 The free-space length ($L_f$) between filler particles and 63.2- % ($\alpha$) DBS and shape factor ($\beta$) from measured samples.

<table>
<thead>
<tr>
<th>Sample</th>
<th>$L_f$ (nm)</th>
<th>$\alpha$ (kV/mm)</th>
<th>$\beta$</th>
</tr>
</thead>
<tbody>
<tr>
<td>neat epoxy</td>
<td>NA</td>
<td>185</td>
<td>1.4</td>
</tr>
<tr>
<td>1 wt-% bare SiO$_2$</td>
<td>900</td>
<td>182</td>
<td>1.6</td>
</tr>
<tr>
<td>2 wt-% SiO$_2$-PGMA</td>
<td>250</td>
<td>213</td>
<td>1.4</td>
</tr>
<tr>
<td>1 wt-% SiO$_2$-PGMA-ferro</td>
<td>600</td>
<td>236</td>
<td>1.5</td>
</tr>
<tr>
<td>2 wt-% SiO$_2$-PGMA-ferro</td>
<td>550</td>
<td>235</td>
<td>1.4</td>
</tr>
<tr>
<td>4 wt-% SiO$_2$-PGMA-ferro</td>
<td>300</td>
<td>265</td>
<td>1.8</td>
</tr>
<tr>
<td>2 wt-% SiO$_2$-PGMA-thio</td>
<td>250</td>
<td>266</td>
<td>1.9</td>
</tr>
</tbody>
</table>
Figure 5.12  a) Dielectric permittivity  b) Dielectric losses as a function of frequency for neat epoxy, bare silica filled composite and for composites that have an electroactive layer and epoxy compatible polymer layer on the core of the filler particle core.
6 SUMMARY AND CONCLUSIONS

Studies of multiscale structure and dielectric properties of PP doped with om-POSS, nano-CaCO$_3$ and fumed silica were carried out. All these composites showed improvement in their dielectric properties, especially DBS was increased. In all of these composites only part of the filler was found to be dispersed at nanoscale. When composite material consisting of a molecular filler om-POSS and PP was studied by confocal Raman micro-spectroscopy, it was found that the relative intensities of different Raman peaks of om-POSS varied depending on the measurement spot. Multiscale characterization yielded a detailed picture of the samples consisting of homogeneously distributed micron-sized elongated om-POSS agglomerates. The variation in the Raman spectra of these can be explained by the orientation of crystallites, showing that, within the confocal volume, there is orientation in the ensemble of crystallites forming large agglomerates. The results are important considering the interpretation of Raman images in composite materials and imply a possibility for orientational imaging. The structure of the composite with crystalline filler CaCO$_3$ was determined using TEM and coarse Raman imaging; which allow screening of larger areas from a sample. Dielectric measurements indicate that DC breakdown strength increases slightly with 1.8 wt-% of CaCO$_3$, but then decreases as a function of concentration, largely owing to an increase of low-field breakdowns likely caused by the large microparticles present in the material. Thus, the positive effect of nanodispersion is masked by the negative effect caused by microparticles. Raman microscopy is shown to be a particularly informative technique to reveal structural information at the nano- and microscale. The amorphous silica composite showed improved the DBS and especially high PD endurance. Silica was dispersed as fractal agglomerates in the composite. All the results demonstrate the importance of multiscale structural analysis of composite samples when studying the effect of structure on dielectric properties.

A new synthetic approach was used to afford well-dispersed silica particles with electroactive core functionalization in epoxy in order to be able to study the effects of the charge layer at the interface of the nanoparticle and polymer matrix. The long, epoxy compatible PGMA brush ensured good disper-
sion and allowed study of the effect of the electroactive layer on the breakdown strength. The DBS and permittivity increased considerably in the bimodal brush particle composites. Still more studies are needed to get more detailed picture of why the electroactive layer at the nanoparticle core improves dielectric breakdown strength of the polymer composite. However, this can only be done if the dispersion of particles in the composite at the nanoscale is possible. Bimodal ligand engineering seems to be a good approach to achieve that.
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