KNEE JOINT LOADING IN THE OPEN AND SQUARE STANCE TENNIS FOREHANDS

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ABSTRACT

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The aim of this study was to calculate and compare the resultant forces and moments in the knee joint of the open and square stance tennis forehands. An additional aim was to incorporate the ground reaction forces to shed light on the role of the lower body in aiding pelvic rotation in these two different forehands.

The two types of forehand strokes hit by 7 right-handed university team level tennis players were filmed and the ground reaction forces on each foot were measured simultaneously. A link segment model of the lower limbs was used along with the inverse dynamics approach to calculate the resultant joint forces and moments in both knees during both the forward swing and follow-through phases of the shot. The ground reaction forces and resultant knee joint forces and moments of the open stance (OS) strokes were compared with those obtained for the square stance (SS) forehands.

Although the average swing times, peak racket velocities and peak transverse plane pelvic velocities were similar, significant kinetic differences (P<0.05) were recorded in the knee joints between the strokes. The SS created greater abduction moments in the left knee during the forward swing and follow-through (0.7 and 0.95 Nm/kg) than the OS (0.2 and 0.32 Nm/kg), respectively. The largest external rotation moments were found in the SS left knee during the follow-through (0.31 Nm/kg) and in the OS right knee prior to impact (0.24 Nm/kg). Despite this, significantly (P<0.05) greater ground reaction torque (T_Z) was generated by the right leg in the OS stroke before impact (0.27 Nm/kg), than by the left leg in the SS during the follow-through (0.09 Nm/kg).

The data did not support the hypothesis that the peak knee external rotation moments would be greater in the square stance. Analysis of the ground reaction forces together with the knee kinetic and pelvic motion data suggests that the square stance utilises the forward step as an aid to pelvic rotation, whereas; the open stance relies more on the extension, abduction and external rotation of the right leg to aid pelvic rotation.

Keywords: tennis forehand, inverse dynamics, knee joint loading, free moment.

ABBREVIATIONS AND DEFINITIONS

OS = Open stance

SS = Square stance

CS = Closed stance

GRF = Ground reaction force

GRFV = Ground reaction force vector

Fz = Vertical component of the GRF

Tz = Vertical torque (free moment)

CoP = Centre of pressure

CoM = Centre of mass

JCS = Joint coordinate system

LCS = Local coordinate system

GCS = Global coordinate system

ASIS = Anterior superior iliac spine

Prox = Proximal

Dist = Distal

Med = Medial

Lat = Lateral

Ant = Anterior

- Pos = Posterior
- Flex = Flexion
- Ext = Extension
- Abd = Abduction
- Add = Adduction

Int Rot = Internal rotation

Ext Rot = External rotation

In the literature, the terms moment and torque are often used interchangeably. In this study, the term moment is used when referring to the resultant joint moment, whereas the term torque is used to describe the moment applied to the foot about the vertical axis from a force plate.

A vector is donated in **boldface** type and scalars are printed in *italic*. The axes denoting the global coordinate system are in capitals whereas the axes of the local coordinate system are in lower case.

The following descriptions of tennis technique all refer to right-handed players.

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1 INTRODUCTION

Tennis is a popular sport played throughout the world. It is estimated about 75 million people play tennis regularly (Pluim et al. 2007). Despite this there is very little published data concerning three dimensional tennis biomechanics and almost nothing related to lower limb kinetics of the forehand.

After the serve, the forehand is the most important shot in a player's arsenal. Rotation of both the upper body and lower body has been described as a significant source of power in the forehand stroke (Groppel 1992, 41; Bollettieri 2001, 118; Elliott 2002, 11; Knudson 2006, 75). The energy is transferred upward from the legs to the pelvis, through the trunk to the arm and then to the racket (Perry et al. 2004). In the kinetic chain of the lower body, the knee joint is regarded as the "critical middle link" in the proximal transfer of force (Whiting & Zernicke 1998, 157). The rotation of the pelvis and trunk involves torsional forces in the lower body, not only during the forward swing but also during the follow-through in which this rotational energy is being dissipated. Research on the lower limb kinetics of the closed stance (CS) forehand has shown that a leg drive was essential to create high axial hip rotational torques to aid trunk rotation (Iino & Kojima 2001; 2003). Only sagittal plane knee moments were described in these studies.

However, most forehands are now hit with an open stance (OS) (Roetert et al. 2009). To date, no data on the lower body kinetics or ground reaction forces of the OS forehand or square stance (SS) forehands have been published. Also, no data concerning the lower body forces found during the follow-through phase of neither the OS nor the SS forehand have been published. Therefore, the main aim of this study is to calculate and compare the forces and moments in the knee joint in both the OS and SS forehand during both the forward swing and follow-through phases. The secondary aim is to incorporate the ground reaction force data to gain information about the role of the lower body in rotating the pelvis in the OS and SS forehands.

2 BIOMECHANICS OF THE TENNIS FOREHAND

The following is a description of the biomechanics of the tennis forehand that is most relevant to this study. This description is based on four major temporal phases of the stroke: preparation, backswing, swing to impact and follow-through.

2.1 Preparation

The two biomechanical issues relating to the preparation for a tennis stroke are readiness and stance.

Readiness. A good ready position or "split step" is important as it enables the player to move quickly to intercept the opponent's shot (Groppel 1992, 41). A split step involves timing a small hop that coincides with the opponent striking the ball. This hop involves knee flexion followed by extension. By using a split step the player benefits from the stretch-shortening cycle (Komi & Nicol 2000, 87) which has been quoted as enhancing the speed of movement by approximately 15-20% (Knudson & Elliott 2004, 154).

Stance. As much of the force used to hit the ball is transferred up through the body to the racket arm (Groppel 1992, 79), the way the player positions their feet will effect how this force is generated and transferred. Stances fall on a continuum between open and closed (Knudson 2006, 81). The positions of the feet for the three main stances are illustrated in figure 2.1.

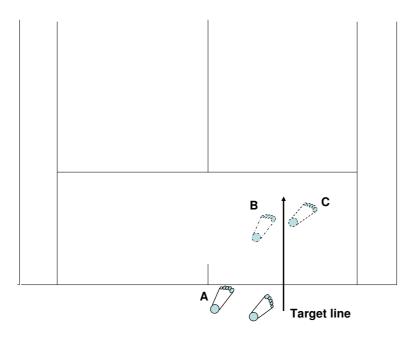


FIGURE 2.1. Position of the feet for the three main type of stance. Open stance (A), square stance (B) and closed stance (C). The right foot is in the same position for the three stances depicted. Modified from Knudson (2006, 82).

2.1.1 The open stance forehand

To execute an open stance (OS) forehand, the player faces the court with their pelvis roughly parallel to the net (hence the name "open") and rotates the shoulders away from the court during the backswing and towards the court during the forward swing (Crespo & Reid 2003, 24). The OS forehand is likely to utilise both linear and angular momentum, although the amount of linear momentum is likely to be less than in the square stance (SS) forehand (Knudson 2006, 84). An example of an OS forehand is shown in figure 2.1 (a). This is the type of forehand that is used the most in tennis (Roetert et al. 2009).

2.1.2 The square stance forehand

The square stance (SS) forehand has been described as the "traditional" forehand as opposed to the "modern" or OS forehand (Crespo & Higueras 2001, 149). To execute a SS forehand, the player pivots on the rear leg and takes a step toward the ball thereby placing their pelvis and feet perpendicular to the intended direction of the shot (hence the name "square"). This stepping action transfers their weight from the rear foot onto

the leading foot just prior to contact with the ball (Bahamonde & Knudson 2003). An example of the SS forehand is shown in figure 2.2.

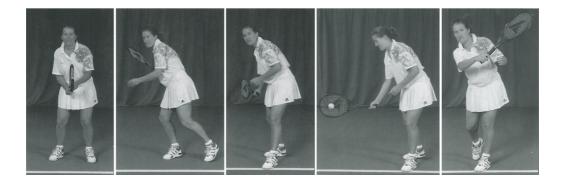


FIGURE 2.2. The SS forehand. Notice that the weight has been transferred to the leading leg before contact in the SS forehand

According to Groppel (1992, 10) and Knudson (2006, 84), the weight transfer has been highlighted as an important source of power in this type of stroke because the forward step generates an amount of linear momentum that is converted into angular momentum during the forward swing.

2.1.3 Ground reaction forces in the forehand

Published data from the closed and SS forehands has shown that the weight transfer occurs before impact and that the body first accelerates and then decelerates prior to impact. An example of the weight transfer and forward and braking impulses occurring before impact in the sideways on type forehands is shown in figure 2.3.

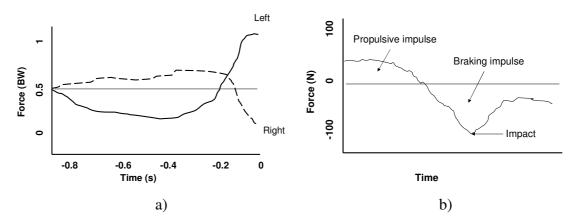


FIGURE 2.3. a) The vertical ground reaction forces showing the transfer of weight from the right leg to the left before impact during the closed stance forehand. Ball impact occurs at time = 0 seconds. Adapted from Iino and Kojima (2001). b) Anterior/posterior force plate graph that shows the body first accelerates and then decelerates just prior to impact. Adapted from Van Gheluwe and Hebbelinck (1986).

No ground reaction torque values have been published for tennis shots, although they have been used to study the kinetics at the shoe-surface interface in golf. In golf the term range has been used to quantify the change in direction of the ground reaction torque. Range of the torque was defined as from peak clockwise T_Z to peak counter-clockwise T_Z torque (Worsfold et al. 2008).

2.1.4 Differences between the open stance and square stance forehands

A comparison of the main features of the OS and SS forehand strokes is given in table 2.1. Although the position of the feet is different between the strokes, the motions of the trunk, arm and racket have been shown to be similar (Crespo & Higueras 2001, 150).

	Open stance (OS)	Square stance (SS)
Preparation time	Shorter	Longer
Initial footwork	Slight step to the sideline	Step forward
Hip position	Facing the net	Facing the sideline
Optimal hitting zone	Between shoulder and	Between knee and waist
	waist height	height
Amount of pelvic rotation	More	Less
Recovery time	Shorter	Longer

TABLE 2.1. Comparison of the OS and SS forehands. (Crespo & Higueras 2001, 151; Saviano1997.

Since modern tennis is played at a high speed on hard surfaces that give a higher bouncing ball (Saviano 1997), the OS is the shot of choice for three main reasons: a) less preparation and recovery time is required, b) it enables powerful shots to be hit off high bouncing balls and c) modern light-weight rackets allow for powerful strokes even with an abbreviated backswing.

There have been concerns about the injury risk of the modern OS forehands compared to the traditional forehand. Trunk muscle activity (Knudson & Blackwell 2000) and the loads in the racket arm (Bahamonde & Knudson 2003) have been studied with the conclusion that there was no significant injury risk of the OS compared to the SS.

2.2 Backswing

The backswing is essential for power generation. In the backswing phase the player typically flexes the knees while the shoulders rotate backwards more than the pelvis creating "a coiling effect" which pre-stretches the muscles of the chest/shoulder and trunk (Elliott 2003, 35). The difference between the alignment of the upper part of the trunk and the pelvis is known as the "separation angle" (figure 2.4). The pre-stretching of the muscles around the pelvis during the backswing is sometimes called "hip loading" in coaching texts (Bollettieri 2001, 114).

The backswing also helps to generate power by increasing the distance and time over which the speed of the racket can be increased during the forward stroke (Elliott 2003, 33).

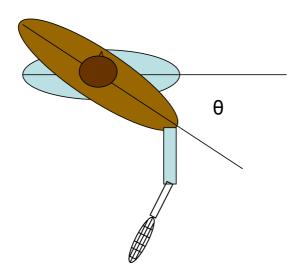


FIGURE 2.4. Separation angle. The angle between the alignment of the shoulders and the pelvis (Elliott 2003, 35).

2.3 Forward stroke

The purpose of the forward stroke is to generate racket head speed and accurate racket movement to intercept the ball. The racket head speed at impact is the major contributor to the ball speed of the shot (Elliot et al. 1997). The biomechanical issues are recovery of the elastic energy in the stretch-shortening cycle and coordination of the body segments necessary to hit the ball.

In the forward swing the energy stored in the muscle tendon unit during the backswing is converted to the kinetic energy of the segment. In addition, the initial rotation of the trunk creates a stretch of the shoulder muscles which enhances its ability to create a forceful forward swing of the arm (Elliott 2003, 35).

The coordinated activation of the body segments in order to hit the ball is commonly referred to in coaching as the kinetic chain (Kibler & van der Meer 2001, 100). The chain consists of the lower body, pelvis, trunk, arm, hand and finally the racket. The

goal of a kinetic chain is to place the racket in the optimum position at the optimum velocity to strike the ball. By using the kinetic chain, the player can transfer the large forces created by the ground reaction forces up through the body to the racket (Perry et al. 2004; Knudson 2006, 8). In the forehand, the hips, trunk, shoulder, elbow, wrist and the racket have been shown to reach peak velocity in sequence (Landlinger et al. 2010). This type of sequential coordination conforms to the "summation of speed principle" (Marshall & Elliott, 2000). The one notable exception in the proximal-to-distal sequencing of the arm segments is internal rotation of the upper arm which reaches its peak angular velocity just before (Elliott 2002, 5) or even just after impact (Landlinger et al. 2010).

2.3.1 The role of the lower body in rotating the pelvis.

Research on the closed stance (CS) forehand has highlighted the roles of the hips in rotating the pelvis. The right hip abduction and external rotation moments plus right hip and knee extension moments all contribute to the generation of pelvic torque in the transverse plane (Iino & Kojima 2001; 2003). The reason these hip moments contribute to pelvic torque can be explained by viewing the projections of these vectors onto the superior/inferior axis of the pelvis (figure 2.5).

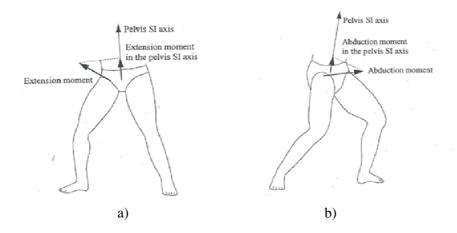


FIGURE 2.5. The effect of the attitude of the right hip joint on the contributions of a) the hip extension moment and b) the abduction moment to the pelvic torque (Iino and Kojima 2001).

Thus for the right hip to rotate the pelvis using extension and abduction moments, the hip must be in a flexed and abducted position. The abducted position aids the role of

the hip extension moment to rotate the pelvis (figure 2.5a) while hip flexion aids the role of the hip abduction moment to rotate the pelvis (figure 2.5b). The projection of the right hip external rotation vector will be parallel to the pelvis superior/inferior axis when the lower body is in the anatomical position. This means that the external rotation moment of the hip can contribute the most when the thigh is neither flexed nor abducted (Iino & Kojima 2001).

2.4 Follow-through

After impact the body can retain as much as 88% of the kinetic energy of the stroke (Knudson & Elliott 2004, 172). Therefore, the main purpose of the follow-through is to safely dissipate this kinetic energy. The main biomechanical principle is impulsemomentum (Knudson 2006, 41). In accordance with the impulse momentum principle (McGinnis 2005, 91), the greater the deceleration time the lower the peak force and the smaller the load on the musculoskeletal system. During the follow- through phase the antagonist muscles are likely to be active and so will be eccentrically contracting in order to slow down the moving body segments (Bahamonde & Knudson 2003, 58). During the deceleration phase the antagonistic muscles may have to tolerate larger loads than the agonist muscles during the acceleration phase (Reid et al. 2003, 125). Therefore, to decrease the risk of injury, a long follow-through is recommended with the racket and arm typically wrapping around the body (Knudson & Elliott 2004, 172).

3 THREE DIMENSIONAL JOINT KINETICS

The direct measurement of the internal forces acting on the joint components such as the joint surface, tendons and ligaments requires very invasive techniques (Robertson et al. 2004, 99). Studies have been performed to directly measure tendon forces in humans using buckle transducers (Gregory at al. 1991) and fibre-optic transducers (Komi et al. 1996; Finni et al. 1998). A more common non-invasive method is to create a model and obtain the resultant joint reaction forces and moments using the inverse dynamics approach.

3.1 Inverse dynamics

The inverse dynamics approach involves obtaining information about segment kinematics, anthropometric measures and external forces and uses this data to calculate the resultant internal joint reaction forces and moments responsible for the movement under focus (Enoka 2008, 124). A descriptive overview of the inverse approach is shown in figure 3.1. The key parts of the inverse dynamics approach are the free body diagram and the equations of motion.

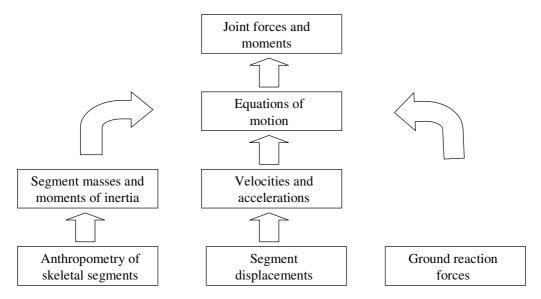


FIGURE 3.1. Flowchart for the inverse dynamics approach (Adapted from Vaughan et al. (1999, 4).

3.1.1 Free body diagram

A free body diagram contains the segments under focus that have been separated from their adjacent segments and all the external forces and moments acting on these segments (Hamill & Knutzen 2003, 350). The free body diagram of the six lower body segments is shown in figure 3.2.

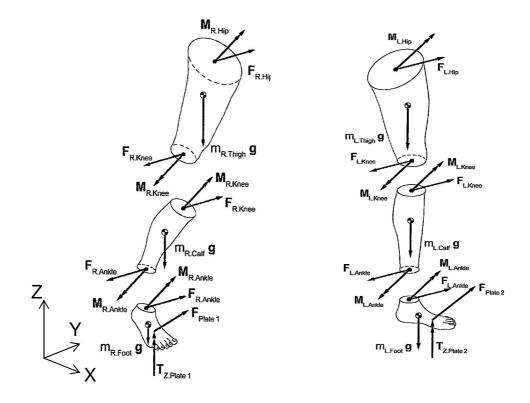


FIGURE 3.2. Free body diagram for the lower body segments. The diagram shows the resultant external forces acting on each segment. In accordance with Newton's 3^{rd} law of motion, the forces and moments at the knee and ankle joints are equal in magnitude but opposite in direction, depending on the segment concerned. \mathbf{M} = joint moment, \mathbf{F} = joint force, $\mathbf{T}_{\mathbf{Z}}$ = ground reaction torque, m = segment mass and \mathbf{g} = acceleration due to gravity (Vaughan et al. (1999, 102-106).

Segment masses and moments of inertia. The segment masses and moments of inertia can be estimated by taking anthropometric measurements such as foot breadth and calf circumference and inputting these values into specific regression equations (Winter 2004, 63; Vaughan et al 1999, 85; Zatsiorsky 2002, 575-582).

Centre of mass. Tracking markers must be filmed in order to record segment displacements and to calculate the location of the centre of mass (CoM) of a segment. The location of the centre of mass (CoM) of a segment can be estimated from a ratio of the length of the segment (Winter 2004, 63; Vaughan et al 1999, 85; Zatsiorsky 2002, 575-582).

Velocities and accelerations. Once the location of the centre of mass is known, the segment velocities and accelerations can be computed. Velocity is calculated by taking the first derivative of the displacement of the CoM with respect to time (Zatsiorsky 1998, 149) and acceleration is calculated by taking the second derivative of the displacement of the CoM with respect to time (Zatsiorsky 1998, 196).

3.1.2 Ground reaction forces

In order to perform inverse dynamic calculations of the lower body a force plate must be used to provide the ground reaction force (GRF), the centre of pressure (CoP) and the ground reaction torque (T_Z).

The ground reaction force is the reaction force exerted on the performer by the supporting surface (Zatsiorsky 2002, 627). This is in accordance with Newton's third law, which states: that for every (action) force that is applied by the body, a (reaction) force is exerted by the second body on the first that is equal in magnitude and opposite in direction. GRFs are typically measured with a force plate. Therefore, the GRF recorded by the force plate is in magnitude but opposite in direction to the force applied on it by the subject's foot.

As the GRF is the algebraic sum of all mass-acceleration products of all the body segments, the GRF reflects the acceleration of the total body centre of mass (Miller 1990, 204).

For three dimensional (3D) movements, the GRF can be decomposed into three orthogonal forces: force in the X direction (F_X), force in the Y direction (F_Y), and the vertical force (F_Z). If X-axis of the force plate points in the direction of travel, a

positive F_X will be a propulsive force while a negative F_X will be a braking force. Examples of propulsive and braking forces are shown in figure 2.3b. The same principle applies to motion in the Y direction.

In addition, the location of the resultant GRF in the X, Y plane is required. This is known as the centre of pressure (Zatsiorsky 2002, 46). The location of the CoP is critically important for calculating flexion/extension and abduction/adduction joint moments as it determines the length of the moment arm used in calculating the moments acting on the ankle (Winter 2005, 102).

The final value that the force plate must provide is the ground reaction torque, which is the torque applied to the foot about the vertical Z axis (Vaughan et al. 1999, 36). This torque acts on a plane parallel to the surface of the force plate (figure 3.3) and is sometimes called the free moment (Miller 1990, 220). The ground reaction torque is especially important for calculating internal/external rotation moments of the lower body during movements that involve twisting (Grimshaw et al. 2006, 332). The location of (T_z) is the coordinates of the centre of pressure.

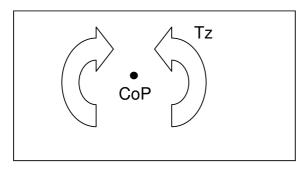


FIGURE 3.3. Depiction of the ground reaction torque when looking down onto the force plate (Grimshaw et al. 2006, 332).

3.1.3 Equations of motion: Forces

A variety of methods can be used to obtain the equations of motion in dynamic musculoskeletal modelling, such as Lagrange (Andrews 1995, 154), the Kane method (Yamaguchi 2004, 173) and Newton-Euler (Vaughan et al. 1992, 83). The advantage of the Newton-Euler equations is that they enable the joint reaction forces to be calculated.

In the Newton-Euler approach, the joint reaction forces are calculated by applying the linear form of Newton's second law, and the joint moments are calculated by applying Euler's equations of motion to the angular form of Newton's second law (Siegler & Liu 1997). The joint reaction forces are calculated first and are then used as inputs to calculate the joint moments.

The forces acting on each segment are represented by Newton's second law (Zatsiorsky, 2002, 368) which can be written in vector form as:

$$\Sigma \mathbf{F} = m \mathbf{a}_{\rm CoM} \tag{1}$$

where $\Sigma \mathbf{F}$ is the sum of all the external forces acting on the segment, *m* is the mass of the segment and \mathbf{a}_{COM} is the acceleration of the centre of mass of the segment in the global coordinate system.

To calculate the forces in the three orthogonal directions of the global coordinate system (X,Y and Z), equation (1) can be expressed in scalar form as:

$$\Sigma F_{x} = ma_{x}$$

$$\Sigma F_{y} = ma_{y}$$

$$\Sigma F_{z} = ma_{z}$$
(2)

where F_z is the compressive force and F_x and F_y are shear forces (Enoka 2001, 120).

The resultant joint force is the vector sum of all the forces acting across the joint. This includes bone, ligament and muscular forces (Vaughan et al. 1999, 40). It is an abstract quantity that is defined by Nigg (2007, 527) as:

$$F = \sum_{i=1}^{m} F_i + \sum_{j=1}^{l} F_j + \sum_{k=1}^{n} F_k + \sum_{p=1}^{s} F_p$$
(3)

(muscle) (lig) (bone) (other)

where:

- \mathbf{F} = resultant intersegmental joint force
- \mathbf{F}_{i} = forces transmitted by muscles
- \mathbf{F}_{j} = forces transmitted by ligaments
- $\mathbf{F}_{\mathbf{k}}$ = bone-to-bone contact forces
- $\mathbf{F}_{\mathbf{p}}$ = forces transmitted by soft tissue, etc.

3.1.4 Equations of motion: Moments

In a similar way to the forces, the resultant joint moment is the net moment produced by all the intersegmental forces acting across the joint. The resultant joint moment includes the resultant muscle, bone and ligament moments and is defined by Nigg (2007, 527) as:

$$\mathbf{M} = \sum_{i=1}^{m} \mathbf{M}_{i} + \sum_{j=1}^{l} \mathbf{M}_{j} + \sum_{k=1}^{n} \mathbf{M}_{k} + \sum_{p=1}^{s} \mathbf{M}_{p}$$
(4)
(muscle) (lig) (bone) (other)

where:

M = resultant intersegmental joint moment

- M_i = moments transmitted by muscles
- M_j = moments transmitted by ligaments
- M_k = bone-to-bone contact moments
- M_p = moments transmitted by soft tissue, etc.

The resultant joint moment only shows the net effect of the agonist and antagonist muscles. That is why it is also known as the net joint moment (Nigg 2007, 527).

The equation for the net moment acting on a segment is based on the principle of angular momentum (Andrews 1995, 153):

$$\Sigma \mathbf{M}_{\rm CoM} = \mathbf{H}_{\rm CoM} \tag{5}$$

where $\Sigma \mathbf{M}_{COM}$ is the sum of all the external moments acting on the segment and \mathbf{H}_{COM} is the time-derivative of the segment's angular momentum about its CoM. The units for

moment and angular momentum are Nm and kgm²/s respectively (Robertson et al. 2004, 241).

In order to use equation (5), this equation must first be expressed in scalar form with respect to the principal axes of the segment.

The time derivative of a vector in a rotating coordinate system is defined as:

$$(\mathbf{P}_{\text{CoM}})_{\text{G}} = (\mathbf{P}_{\text{CoM}})_{\text{L}} + (\boldsymbol{\omega} \times \mathbf{P})$$
(6)

where $(\mathbf{P}_{COM})_{G}$ and $(\mathbf{P}_{COM})_{L}$ are the time derivative of the vector \mathbf{P} viewed from the fixed GCS and rotating LCS, respectively, and $\boldsymbol{\omega}$ is the angular velocity of the LCS viewed from the GCS (Zatsiorsky 1998, 205). This is illustrated in figure (3.4).

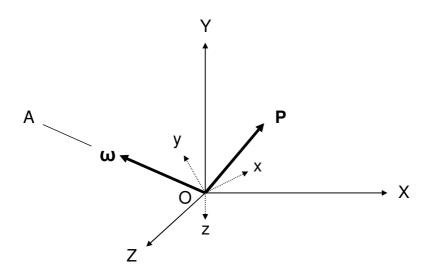


FIGURE 3.4. Vector **P** represented in fixed and rotating reference frames. X,Y,Z and x,y,z represent the GCS and LCS respectively, **P** is a vector and $\boldsymbol{\omega}$ is the angular velocity vector of the rotating LCS. The two reference frames are centred at O. The local coordinate system x,y,z rotates about the axis OA. Modified from Beer & Johnston (2004, 972).

If the LCS is translating, the rate of change of the vector \mathbf{P} is the same with respect to the fixed and the translating frame (Beer & Johnston 2004, 971). Therefore, only the rotation of the LCS needs to be described.

If **P** was to represent the position of a particle then the expression ($\boldsymbol{\omega} \times \mathbf{P}$) in equation (6) would give the instantaneous tangential velocity of the particle situated at the tip of **P** (Beer & Johnston 2004, 972).

Since angular momentum (**H**) is also a vector, **H** can be substituted into equation (6) in place of **P**. This gives:

$$(\mathbf{\dot{H}}_{CoM})_{G} = (\mathbf{\dot{H}}_{CoM})_{L} + (\boldsymbol{\omega} \times \mathbf{H}_{CoM})$$
(7)

where \mathbf{H}_{G} and \mathbf{H}_{L} are the time derivative of the angular momentum viewed from the GCS and LCS, respectively, $\boldsymbol{\omega}$ is the angular velocity of the LCS viewed from the GCS and **H** is the angular momentum about its CoM (Zatsiorsky 2002, 392).

Substituting equation (7) into equation (5) gives (Beer & Johnston 2004, 1166):

$$\Sigma \mathbf{M}_{\text{CoM}} = (\mathbf{H}_{\text{CoM}})_{\text{L}} + (\boldsymbol{\omega} \times \mathbf{H}_{\text{CoM}})$$
(8)

Angular momentum is defined as:

$$\mathbf{H}_{\rm CoM} = [\mathbf{I}] \boldsymbol{\omega} \tag{9}$$

where \mathbf{H}_{COM} is the angular momentum of the body, $\boldsymbol{\omega}$ the angular velocity and [I] the inertia tensor (Goldstein 1980, 192). The inertia tensor is a 3x3 symmetrical matrix of the form:

$$[\mathbf{I}] = \begin{bmatrix} I_{xx} & -I_{xy} & -I_{xz} \\ -I_{xy} & I_{yy} & -I_{yz} \\ -I_{xz} & -I_{yz} & I_{zz} \end{bmatrix}$$
(10)

where the diagonal elements are the moments of inertia and the off-diagonal elements are the products of inertia (Seliktar & Bo 1995, 227). The subscripts for the moments of inertia represent the axes and the subscripts for the products of inertia represent the planes of rotation respectively (Zatsiorsky 2002, 276).

The moment of inertia tensor will change with each change in orientation of the LCS. However, if the LCS is located at the CoM of the segment with the axes aligned with the principal axes, the products of inertia become zero (Zatsiorsky 2002, 281) and the matrix becomes:

$$[\mathbf{I}] = \begin{bmatrix} I_{xx} & 0 & 0\\ 0 & I_{yy} & 0\\ 0 & 0 & I_{zz} \end{bmatrix}$$
(11)

Applying equation (8) the components of the angular momentum **H** can then be written:

$$\mathbf{H}_{\text{CoM}} = I_x \,\omega_x \,\mathbf{i} + I_y \,\omega_y \,\mathbf{j} + I_z \,\omega_z \,\mathbf{k} \tag{12}$$

where I_x , I_y and I_z represent the moments of inertia about the principal axes. ω_x , ω_y , and ω_z are the angular velocities and **i**, **j** and **k** are the unit vectors along the respective axis of the LCS.

Substituting for \mathbf{H}_{CoM} from equation (12) into Equation (8) and expanding the cross product, 3 scalar equations are obtained:

$$\Sigma M_{x} = I_{x} \overleftrightarrow{\omega}_{x} + (I_{y} - I_{z}) \, \omega_{y} \, \omega_{z}$$

$$\Sigma M_{y} = I_{y} \overleftrightarrow{\omega}_{y} + (I_{z} - Ix_{x}) \, \omega_{z} \, \omega_{x}$$

$$\Sigma M_{z} = I_{z} \overleftrightarrow{\omega}_{z} + (I_{x} - I_{y}) \, \omega_{x} \, \omega_{y}$$
(13)

where I_x , I_y , and I_z represent the moments of inertia, ω_x , ω_y , and ω_z the angular velocities and ω_x , ω_y , and ω_z the angular accelerations about the respective principal axes. Equation (13) is known as Euler's equations of motion (Goldstein 1980, 205).

The units for angular velocity and the angular acceleration are rad/s and rad/s^2 respectively (Robertson et al. 2004, 240).

3.1.5 Equations of motion: Angular kinematics

In order to obtain the angular velocity and angular acceleration of a segment, the orientation of the segment with respect to time must be obtained. There are several methods to calculate orientation of a segment such as direction cosines or the helical technique (Zatsiorsky 1998, 25) but the most common is with the use of Euler angles (Robertson et al. 2004, 52).

Euler angles are three angles that completely describe the orientation of one coordinate system relative to another. These three angles can be described either geometrically or by using three separate transformation matrices to rotate the local coordinate system from the GCS to its actual orientation (Zatsiorsky 1998, 46). These rotations are performed in a specific sequence (Goldstein 1980, 145). There are a total of 12 different rotation sequences (Winter 2005, 183), an example of one sequence is illustrated in figure 3.5. This sequence is described by Greenwood (2006, 142) as the classical Euler angles.

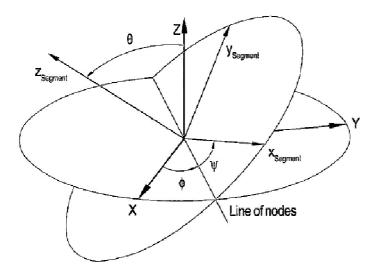


FIGURE 3.5. The three Euler angles (φ , θ , ψ) used to define the orientation of a rotating coordinate system (x, y, z) relative to the a fixed coordinate system (X, Y, Z). The sequence of rotation is: φ about the Z axis; θ about the line of nodes; and ψ about the z_{segment} axis. The line of nodes represents the intersection of the x,y and X,Y planes (Vaughan et al. 1999, 98: Goldstein 1980, 145).

For the rotation sequence illustrated in figure 3.5, the equation for the angular velocity

about the local axes of the segment is given by (Goldstein 1980, 176):

$$\omega_{segment.x} = \dot{\phi}\sin\theta\sin\psi + \dot{\theta}\cos\psi$$

$$\omega_{segment.y} = \dot{\phi}\sin\theta\cos\psi - \dot{\theta}\sin\psi \qquad (14)$$

$$\omega_{segment.z} = \dot{\phi}\cos\theta + \dot{\psi}$$

where the dot above the Euler angles represents the time derivative of the angle. The derivation of the angular velocity equation (equation (13)) is given in the Appendix.

The segment angular accelerations can be calculated by taking the derivative of equation (14) which gives the following (Vaughan et al 1999, 97):

$$\dot{\omega}_{segment.x} = \ddot{\phi}\sin\theta\sin\psi + \dot{\phi}\dot{\theta}\cos\theta\sin\psi + \dot{\phi}\dot{\psi}\sin\theta\cos\psi + \ddot{\theta}\cos\psi - \dot{\theta}\dot{\psi}\sin\psi$$
$$\dot{\omega}_{segment.y} = \ddot{\phi}\sin\theta\cos\psi + \dot{\phi}\dot{\theta}\cos\theta\cos\psi - \dot{\phi}\dot{\psi}\sin\theta\sin\psi - \ddot{\theta}\sin\psi - \dot{\theta}\dot{\psi}\cos\psi$$
(15)
$$\dot{\omega}_{segment.z} = \ddot{\phi}\cos\theta - \dot{\phi}\dot{\theta}\sin\theta + \ddot{\psi}$$

where the double dot above the Euler angles represents the second time derivative of the angle.

The derivation of the segment angular velocity and angular acceleration equations from the Euler angles (equations (14) and (15)) is given in the Appendix.

The major disadvantage of the Euler angles is that errors in the calculation of the angles will occur when the axes of the rotating system and the fixed system coincide. This is known as gimbal lock (Zatsiorsky 1998, 52).

3.1.6 Limitations of the inverse dynamics model

Simplifications are a necessary part of the modelling process. For example, it is assumed that each segment is non-deformable and has a fixed mass in the form of a point mass located at its centre of mass (CoM) instead of mass unevenly distributed along the whole segment. The location of CoM is assumed to remain fixed relative to

the segment during the movement. In reality, the movement of muscles and other soft tissues would influence the CoM position. The mass moment of inertia of each segment about its CoM is also assumed to remain constant. Also, all the joints are considered to be frictionless ball and socket joints, free to move in all three planes (Kirtley 2006, 124; Robertson et al. 2004, 116; Winter 2005, 87).

Inverse dynamic calculations for joint moments can only be performed when the CoP can accurately be assessed to be under the foot of the performer. Due to limitations of the force plates, when the vertical GRF is small (e.g. during the beginning and end of the stance phase) the CoP calculation becomes unstable and will give erroneous CoP locations (Robertson et al. 2004, 92).

3.2 Estimating muscle activity with the GRF vector

The ground reaction force vector (GRFV) can be used to give a simple indication of which group of muscles are likely to be dominant when the foot is in contact with the force plate (Kirtley 2006, 118). For example, for the ankle joint during quiet standing (Figure 3.6) the ground reaction force is anterior to the ankle and would create an external moment that would cause a dorsiflexor movement of the foot. To maintain equilibrium, the Achilles tendon would need to produce an equal but opposite (internal) moment (Newton's third law). Therefore, the active muscle is likely to be on the opposite side of the joint to the GRF (Kirtley 2006, 77).

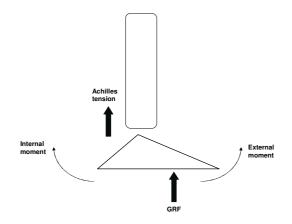


FIGURE 3.6. Balanced equilibrium between external and internal moments at the ankle joint (Kirtley 2006, 77).

The same principle can be applied to the knee joint. For example, during a squatting posture shown in figure 3.7, the GRFV passes posterior to the knee joint. This creates an external moment that would cause the knee to flex. Therefore, the quadriceps muscles are likely to be active to produce an internal extensor moment to maintain equilibrium.

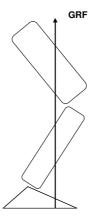


FIGURE 3.7. Diagram to show the position of the GRFV during squatting. The GRFV passes anterior to the ankle joint and posterior to the knee joint.

Use of the GRFV to estimate muscle activity is also called the floor reaction force vector approach (FRFV) in the literature (Simonsen et al. 1997). The FRFV approach should not be used to calculate joint moments as it produces large errors in calculating moments in joints that are superior to the ankle (Winter 2005, 105; Simonsen et al. 1997). However, it is used in clinical environments to provide instant qualitative interpretation (Kirtley 2006, 120).

4 PURPOSE OF THE STUDY

The SS forehand has traditionally been taught as the main way of hitting the forehand (Bollettieri 2001, 87). Although the OS forehand has previously been described as poor technique and one that does not optimally utilise the kinematic chain of the lower body (Bahamonde & Knudson 2003), it is the most common type of forehand in tennis today (Roetert et al. 2009).

The advantages and disadvantages between the OS and SS forehands is a subject of much debate amongst the coaching profession. One of the aspects has been knee joint loading. In particular the loading of the left knee in the SS stroke has been stated as being greater than that would be found in either knee in the OS (Langevad 2003).

However, to date, no data on the lower body kinetics of the OS and SS forehands have been published. As a result of this, a scientific comparison of knee joint loading in the OS and SS forehands has not been possible.

Therefore, the primary goal of this study was to calculate and compare the resultant forces and moments of the knee joint in both the OS and SS forehand during both the forward swing and follow-through phases. The hypothesis of this study was that the greatest transverse plane moments in the knee joint would be found in the left knee of the SS forehand during the follow-through phase.

An additional goal was to shed light on the role of the lower extremities in aiding pelvic rotation, the latter being a critical process in the force generation of the forehand.

5 METHOD

5.1 Subjects

Seven male, right-handed university team level players (mean \pm standard deviation, age 33 \pm 11.5 years; mass 77.7 \pm 14.2 kg) participated in the study. They all provided informed consent. Each subject had 15 (2.5 cm diameter) spherical reflective markers attached to the lower body as listed in table 5.1. In addition, 2 cm long strips of reflective tape were placed on the tip and both sides of the subjects' rackets.

TABLE 5.1. Names of the marker positions

Number	Name
l	Right metatarsel head II
2	Right heel
3	Right lateral malleolus
4	Right tibial tubercle
5	Right femoral epicondyle
6	Right greater trochanter
7	Right anterior superior iliac spine (ASIS)
8	Left metatarsel head II
9	Left heel
10	Left lateral malleolus
11	Left tibial tubercle
12	Left femoral epicondyle
13	Left greater trochanter
14	Left anterior superior iliac spine (ASIS)
15	Sacrum

5.2 Experimental procedure and measurements

The experimental setup is shown in figure 5.1. The filming took place on an acrylic indoor tennis court. Each subject performed OS and SS forehands with each foot on a separate force plate (figure 5.2). The players were allowed to use their own rackets during the data collecting process. As the surface of the force plates were 16.5 cm above the court surface, a badminton net was fixed to the tennis net so it was 16.5 cm above the normal height of the net. The tennis ball was fed via a tennis ball machine (Lobster

Elite III, Lobster, California, USA) at approximately 20 m/s. The subjects were instructed to hit each ball with rally-paced ball speed and medium topspin over the badminton net so that the ball would land into the 2.5 m by 2.5 m target area.

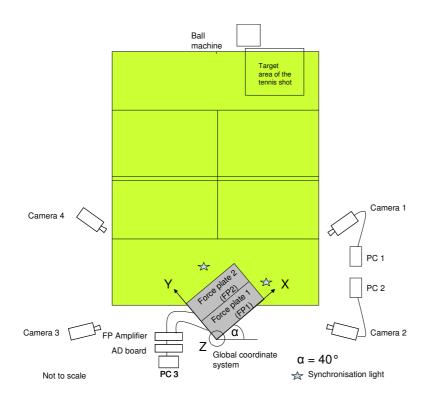


FIGURE 5.1. The experimental setup showing the positions of the force plates, global coordinate system, cameras and ball machine relative to the court.

Only the shot the subject decided represented a normal version of their forehand and where the ball landed in the target zone was saved for later analysis. Altogether, three OS and three SS forehands were saved for each subject. Four cameras were used. Two cameras (Sony HDR-HC3 1080i, Sony, Japan) recorded at 200 frames per second (fps) directly onto MiniDV tape whilst the other two cameras, (Memrecam fxK4 NAC) recorded at 1000 fps onto laptop computers (PCs 1 and 2).



FIGURE 5.2. The picture shows a subject with one foot on each force plate, the location of the markers and the global origin (X,Y,Z) located at the back right corner of force plate 1.

The outputs from the two force plates (Z13216, Kistler Instruments, Winterhur, Switzerland) were amplified (Kistler 9861A) then converted to a digital signal (CED DAC, CED Ltd. Oxford, UK) and recorded at 400Hz on a separate laptop (PC 3) running Signal 2.16 program (CED Ltd). A pair of timing lights was used to help synchronise the force plate data with the camera data. The lights were programmed to flash twice during each shot, one short and one long flash. This was recorded along with the force plate data on Signal 2.16.

An 8 rod, 25 point calibration frame (Peak Performance Technologies, Eaglewood, CO, USA) of dimensions 2.2 m x 1.6 m x 1.9 m was filmed prior to and after each filming session. The frame was large enough and was positioned on the force plates so that it encompassed the space used by the players during their shots.

The subjects were weighed and lower body measurements were taken according to figure 5.3.

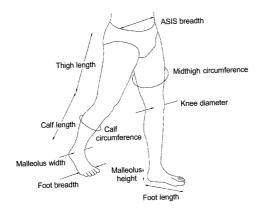


FIGURE 5.3. The anthropometric measurements needed to calculate the masses and moments of inertia of the six lower body segments (Vaughan et al 1999, 17). Since the subjects wore tennis shoes, the foot model was modified to include the mass of the shoe.

5.3 Analytical procedures

The 25 points of the calibration frame were digitised using Vicon Motus 9.0 (Vicon Motion Systems, Oxford, UK). The origin of the calibration frame was translated to the back right corner of force plate 1 (as shown on figures 5.1 and 5.2) and the coordinate system rotated using the Vicon Motus program.

One OS and one SS shot that was considered to be the closest to a text-book forehand for each subject, was chosen for analysis. The markers were digitised using Vicon Motus. This program uses the direct linear transformation method (Vicon User Manual) to obtain the 3D coordinates of the markers using the information from the calibration points. Since the filming was carried out by two pairs of different cameras filming at different frame rates, the images from all cameras had to be processed to the same camera speed (200 Hz). A short computer program was written for Avisynth 2.0 (free open-source software) to deinterlace the images from the Sony cameras. The images from the two Memrecam cameras were post-processed by a technician at the Finnish Research Institute for Olympic Sports (KIHU) and recorded onto a digital video disc (DVD).

The force plate data was imported into a specifically designed Excel (Microsoft Corporation) template for processing. The force plate data was converted from volts into Newtons and then the six GRF quantities were calculated (F_X , F_Y , F_Z , T_Z , CoP_X and CoP_Y). The equations for computing the six quantities of the GRF are shown in table 5.2 (ISB, Kistler).

The scale factors were; FP1 37.3 mV/kg in the Z direction, FP2 35.8 mV/kg in the Z direction, and 77.0 mV/kg in the X and Y directions for both plates. The force plate data was recorded as voltages, therefore to convert it into Newtons; each sampled value was multiplied by 1000, divided by the relevant scale factor and multiplied by 9.81 m/s².

The final conversion involved translating and rotating the Kistler coordinate system to the global coordinate system (GCS) used for the inverse dynamic computations. This ensured the origin of the CoP was at the GCS origin and not at the centre of each plate.

Parameter	Calculation	Description*
Fx	fx12 + fx34	Medial-lateral force ¹
Fy	fy14 + fy23	Anterior force ¹
Fz	fz1+fz2+fz3+fz4	Vertical force
Mx	b*(fz1+fz2-fz3-fz4)	Plate moment about X-axis ³
My	a*(-fz1+fz2+fz3-fz4)	Plate moment about Y-axis ³
Mz	b*(-fx12+fx34)+a*(fy14-fy23)	Plate moment about Z-axis ³
Mx'	Mx+Fy*azO	Plate moment about top plate surface ²
My'	My-Fx*azO	Plate moment about top plate surface ²
CoP x	-My'/Fz	X-Coordinate of the centre of pressure ²
CoP y	Mx'/Fz	Y-Coordinate of the centre of pressure ²
Tz	Mz-Fy*CoPx+Fx*CoPy	Vertical torque

Table 5.2. The equations used to calculate the ground reaction forces, CoP and vertical torque (ISB Kistler).

* Note, these parameters are in the Kistler coordinate system, not the GCS used for the inverse dynamics calculations. 2 azO = top plane offset (-0.08m). 3 sensor offset, a = 0.21m and b = 0.35m

Analysis of the obtained force-time graphs showed that the force plate data suffered from varying amounts of drift, even though the reset button on the Kistler amplifier was regularly pressed. This was rectified by finding when on the force-time curve the foot was off the plate, and applying a correction factor for each of the 8 FP outputs. The correction factor gave zero Newtons in all three directions (X,Y,Z) while the foot was off the plate and a value of approximately $\frac{1}{2}$ body weight during quite standing. This enabled the force plates to be recalibrated after the data was collected. This was performed on both force plates, for each subject and for every shot.

Further analysis of the CoP data revealed that the CoP data was not accurate enough for the right leg of two subjects during the SS shot and so their right leg data was removed from the analysis.

Certain calculations such as those shown in figures 6.2 and 6.3 in chapter 6, required the positive axis for the medial/lateral ground reaction forces (GRFs) to be parallel to the baseline pointing to the left and the positive axis for the anterior/posterior GRFs to be perpendicular to the baseline in the direction of the net. The following formulae were used to rotate the coordinate system of the force plates (ISB Kistler):

$$Fx_{lcs} = (FX * \cos 50) + (FY * \sin 50)$$
(16)

$$Fy_{lcs} = (FY * \cos 50) - (FX * \sin 50)$$
(17)

where Fx_{lcs} = anterior/posterior force, Fy_{lcs} = medial/lateral force. FX = force in the global X direction, FY = force in the global Y direction as shown in figure (5.1) and figure (5.2).

Various computer programs were written using Matlab (The Mathworks, Massachusetts, USA) in order to calculate the net lower body joint forces and torques. The detailed mathematics of these calculations can be found in vector algebraic form in Appendix B of Vaughan et al (1992) and Vaughan et al (1999).

Each lower limb was modelled as three rigid link segments of the foot, calf and thigh. The free body diagram is shown in figure 3.2. In terms of the moments of inertia, the calf and thigh were modelled as cylinders and the foot was modelled as a right pyramid. The position of the centre of mass (CoM) of each segment relative to its proximal end and the inertial properties of each segment were estimated using the coefficients and regression equations of Vaughan et al (1999, 85).

The anatomical joint angles were determined according the convention adopted in the joint coordinate system (JCS) of Grood and Suntay (1983). This enables the joint angles to have a functional, anatomical meaning. Joint angles are defined as a rotation of the distal segment relative to the proximal segment. An example of how the JCS is defined for the knee is illustrated in figure 5.4. The rotations are defined as follows:

- Flexion and extension (plus dorsiflexion and plantar flexion) take place about the mediolateral axis of the proximal segment (i.e., the **k**_{thigh} axis)
- Internal and external rotation take place about the longitudinal axis of the distal segment (i.e., the **i**_{calf} axis).
- Abduction and adduction take place about a floating axis that is orthogonal to both the flexion/extension and internal/external rotation axes.

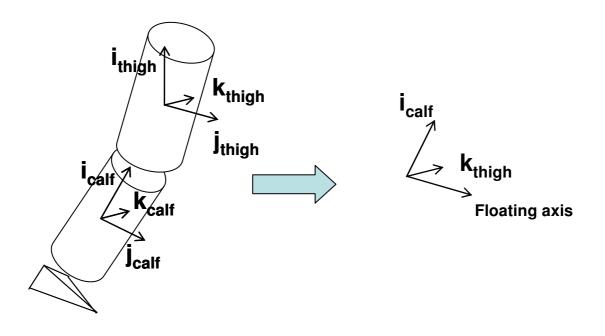


FIGURE 5.4. The JCS for the right knee. The axes are i_{calf} , k_{thigh} and the floating axis. The floating axis is the cross product of these two axes. (Based on diagram in Robertson et al. 2004, 50).

The JCS gives a geometric description of the Cardan angles (Zatsiorsky 1998, 100) and is the International Society of Biomechanics standard for reporting kinematic data (Livesay et al. 1997).

The input for calculating segment linear acceleration was the CoM displacement data, whereas the input for calculating segment angular velocity and angular acceleration were the Euler angles. The order of rotation of the Euler angles was defined as: Z, line of nodes, z (see figure 3.5 in chapter 3).

The equations given in the books by Vaughan et al (1992) and Vaughan et al (1999) for calculating the Euler angles were used initially. However, polarity errors of the angular velocities were found. These problems have also been identified by other researchers (Ariel; Shen) using the books Dynamics of Human Gait by Vaughan et al (1992) and (1999). Therefore, the following equations were used instead (Wikipedia):

 $\varphi = \operatorname{atan2} \left(\mathbf{k} \mathbf{x}, -\mathbf{k} \mathbf{y} \right) \tag{18}$

$$\theta = \cos^{-1} (\mathbf{k}z) \tag{19}$$

$$\Psi = \operatorname{atan2}\left(\mathbf{I}\mathbf{Z}, \mathbf{J}\mathbf{Z}\right) \tag{20}$$

where: **i**, **j**, **k** are the unit vectors of the segment in the global coordinate system and x, y, z represent the x, y or z component of that unit vector, and atan2 is a standard trigonomic computer function. An additional algorithm was used to ensure the polarity of the angle was consistent.

The coordinate data for the markers, centre of mass (CoM) of each segment and the segment Euler angles were smoothed using a 4th order, zero phase lag Butterworth low-pass digital filter (Winter 2005, 45). The cutoff frequency for each coordinate (X, Y and Z) for each segment was determined individually by residual analysis (Winter 2005, 49).

The segment linear acceleration and the time derivatives of the Euler angles were calculated using finite difference calculus (Robertson et al 2004, 21). The following formulae were used:

$$\frac{dx_n}{dt} = \dot{x}_n = \frac{x_{n+1} - x_{n-1}}{(2\Delta t)}$$
(21)

$$\frac{d^2 x_n}{dt^2} = \ddot{x}_n = \frac{x_{n+1} - 2x_n + x_{n-1}}{(\Delta t)^2}$$
(22)

where x is an input data point, n refers to the nth sample and Δt is the time between adjacent frames. Velocity was calculated using equation (21) and acceleration using equation (22).

The segment angular velocities and angular acceleration were calculated from equations (14) and (15) respectively. The segment angular velocities and angular accelerations were then used along with the segment inertia parameters and the moments created from the joint reaction forces, as inputs into equation (13) for calculating the net moment acting on a segment.

To aid in comparison of the results, the net forces and moments were normalised and the stroke was divided into two temporal phases. The net joint reaction forces were normalised by divided by body weight and multiplied by 100%, and the net joint moments were normalised by dividing my body mass. This is a standard way of normalising gait analysis data (Kirtley 2005, 127; Allard et al. 1997, 315). The temporal phases of the stroke were termed forward swing and follow-through. These were defined as:

- Forward swing: from the beginning of trunk left rotation to impact.
- Follow-through: from ball impact to the point were the trunk had stopped rotating.

The force plate readings (particularly the centre of pressure) became unreliable when the vertical force was below approximately 10% body weight. The amount varied depending on the subject and the shot. This was most noticeable for the right leg both during and after impact with the ball when the vertical components were at their lowest as the weight was now over the front foot. The individual subjects' resultant moments were not analysed when the CoP could not be accurately assessed to be under the subject's foot. That is why no kinetic data has been recorded in tables 6.4 to 6.7 and figure 6.8 for the right leg during the follow-through phase.

5.4 Expression of joint forces and moments

The resultant joint reaction forces and joint moments are expressed in terms of the axes used in describing the JCS. These are defined as follows (Vaughan et al 1999, 41):

Forces

- A mediolateral force takes place along the mediolateral axis of the proximal segment.
- A proximal/distal force takes place along the longitudinal axis of the distal segment.
- An anterior/posterior force takes place along a floating axis that is perpendicular to the mediolateral and longitudinal axes.

Moments

- A flexion/extension moment takes place about the mediolateral axis of the proximal segment.
- An internal/external rotation moment takes place about the longitudinal axis of the distal segment.
- An abduction/adduction moment takes place about a floating axis that is perpendicular to the mediolateral and longitudinal axes.

In should be noted the axes used in the JCS are defined relative to the body segments and therefore, are not necessarily in the same directions as the force plate axes. Also, the JCS is not an orthogonal system (Zatsiorsky 1998, 101) the resultant forces and moments at each joint may not be exactly at right angles to each other.

5.5 Statistical Methods

Statistical analysis was performed using Matlab. The data was first tested for normality using the Kolmogorov-Smirnov (K-S) test. The data failed this test and so significance testing was performed using the non-parametric Wilcoxon matched pairs test instead of the *t* test. The statistical significance was set at a level of $p \le 0.05$.

Only when there was sufficient reliable data from all seven subjects were the results statistically analysed. This means that no statistical analysis was performed on the right knee for the SS shot and for the right knees in the follow-through phase of both strokes.

6 **RESULTS**

6.1 Stroke parameters

In order to assess whether the biomechanical parameters of the OS and SS strokes were easily comparable with each other, racket speed, ball velocity and normalised ball height at impact were analysed.

The mean resultant racket tip speed just before ball impact for the OS and SS strokes were 28.45 ± 2.27 m/s and 29.25 ± 1.5 m/s, respectively. These were not significantly different (n.s.). The mean resultant ball speed just prior to impact in the OS and SS strokes were 7.54 ± 0.22 and 7.76 ± 0.24 m/s, respectively (n.s.). The average forward swing times of the OS and SS strokes were 0.22 ± 0.03 and 0.24 ± 0.04 s, respectively (n.s.). The mean ratio of ball impact height to the height of the right hip at impact for the OS and SS strokes were 1.10 ± 0.13 and 0.96 ± 0.15 , respectively (n.s.).

6.2 Ground reaction forces

As the GRF reflects the summation of the mass acceleration products of all body segments, some of the fundamental differences between the OS and SS techniques can be seen in the recorded GRFs.

6.2.1 Weight transfer

Typical examples of the vertical ground reaction forces (Fz) as a function of time in the OS and SS strokes are shown in figures 6.1a and 6.1b. These figures clearly show the transfer of body weight from the right leg to the left leg. This was observed for all participants in both the OS and SS strokes. The difference was in the timing of the weight transfer. For all subjects except one, the weight transfer occurred after contact in the OS forehand. In the SS forehand, the weight transfer occurred before contact for all subjects.

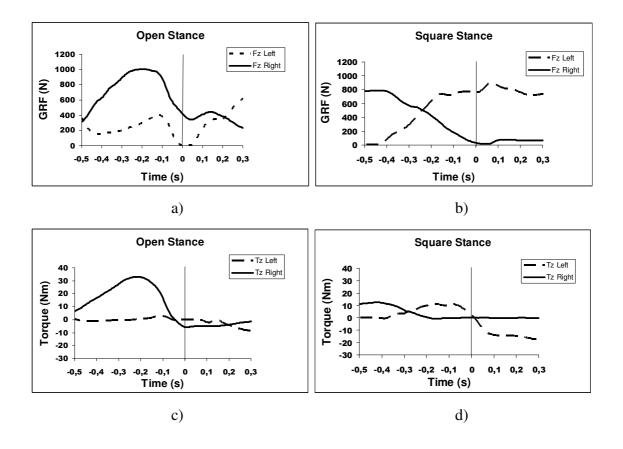


Figure 6.1. The vertical ground reaction forces (Fz), figures a and b, and vertical torques (Tz), figures c and d for one representative subject. Positive values reflect a counter-clockwise torque. Negative values reflect a clockwise torque. Time zero = ball contact.

6.2.2 Vertical torque

The vertical torque (T_Z) applied to the feet about the vertical Z axis as a function of time in the OS and SS strokes is shown in figures 6.1c and 6.1d. Maximum values for the vertical torque were recorded before contact in the right leg in the OS forehand and after contact in the left leg in the SS forehand. The values for the mean peak vertical torque are shown in table 6.1. The mean peak vertical torque in the OS right leg prior to impact, 0.27 ± 0.09 Nm/kg was significantly greater than the SS left leg follow-through, 0.17 ± 0.07 Nm/kg (p=0.016). The total torque generated (i.e. the range) by the right leg in the OS was nearly significantly greater than the left leg in the SS (p=0.078).

Nm/kg ± (standard dev.)				
CW Torque	CCW Torque	Range		
0.07 (0.04)	0.11 (0.06)	0.18 (0.08)		
0.05 (0.03)	0.27 (0.09)	0.32 (0.11)		
	CW Torque 0.07 (0.04)	CW Torque CCW Torque 0.07 (0.04) 0.11 (0.06)		

Square Stance (SS)

Foot	CW Torque	CCW Torque	Range
Left	0.17 (0.07)	0.09 (0.03)	0.26 (0.08)
Right	$0.13 (0.05)^5$	$0.10 (0.09)^5$	$0.23 (0.07)^5$

 $()^n$ = number of subjects (if n < 7).

6.2.3 Medial/lateral and anterior/posterior GRFs

The GRFs in the anterior/posterior and medial/lateral directions in both types of shot for the same subject are shown in figures 6.2 and 6.3. The anterior/posterior direction (x) is at right angles to the baseline pointing towards the net. The medial/lateral direction (y) is parallel to the baseline pointing to the left.

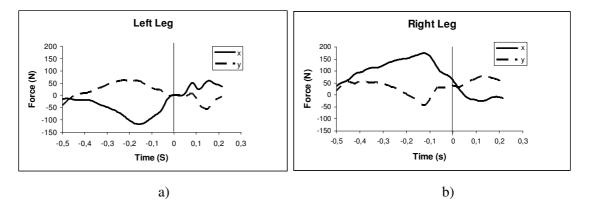


FIGURE 6.2 OS anterior/posterior and medial/lateral GRFs for one representative subject. Time zero = ball contact.

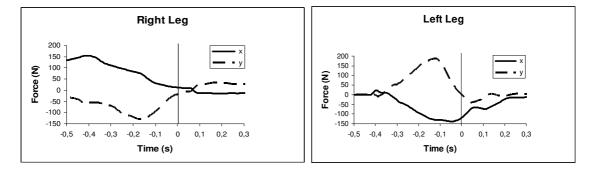


FIGURE 6.3 SS anterior/posterior and medial/lateral GRFs for one representative subject. x = towards the net. y = along the baseline. Time zero = ball contact.

Figures 6.2 and 6.3 show that the anterior/posterior and medial/lateral ground reaction forces peak before impact for both shots. The effects of these forces on the feet are illustrated in sections 7.3 and 7.4 in the Discussion.

6.3 Body movements based on marker data

6.3.1 Comparison of pelvic motion

The graphs in figures 6.4 and 6.5 illustrate the motion of a point midway between the two ASIS markers for both types of forehands. Motion in the transverse plane is shown in figure 6.4.

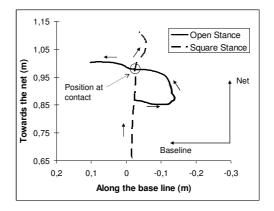


FIGURE 6.4. Transverse view of the motion of the midpoint between the ASIS markers during one complete OS and SS stroke for one representative player.

The mean horizontal translation of the midpoint towards the net, for OS and SS in the 0.5 s prior to impact were 11.4 ± 6.5 cm and 29.2 ± 5.8 cm, respectively. This was significantly different (p=0.016).

Pelvic angular velocity. There was no significant difference between the average peak transverse plane angular velocity of the pelvis in the OS and SS strokes during the forward swing phase $(5.54 \pm 1.10 \text{ rad/s} \text{ and } 5.02 \pm 1.56 \text{ rad/s} \text{ respectively}).$

Vertical motion of the pelvis. The vertical motion (along the global Z-axis) of the midpoint between the two ASIS markers as function of time for the two types of forehands is shown in figure 6.5. The mean maximum vertical translation of the midpoint for OS and SS during the forward swing was 12.0 ± 3.0 cm and 6.9 ± 3.0 cm, respectively. This was significantly different (p=0.016).

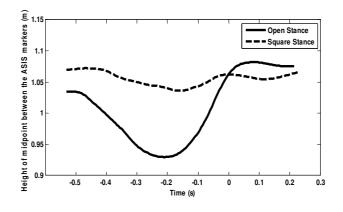


FIGURE 6.5. Illustration of the vertical motion of the midpoint between the ASIS markers. Data from two different players. Time zero = ball contact.

6.3.2 Knee angles

Maximum knee flexion occurred at the start of the forward swing phase which was followed by extension towards impart. There were no significant differences between the maximum knee flexion angles in the two strokes (table 6.2). Although, on average, the amount of right knee extension before impact was greater for the OS, there was also no significant difference in the amount of knee extension recorded from maximum knee flexion to ball impact (table 6.3).

Table 6.2 Maximum joint angles of the knees. Data is presented as mean (degrees) \pm (standard dev.)

Open Stance (OS)		Square Stance (SS)
Left Knee	61.9 (±13.7)	53.9 (±14.3)
Right Knee	47.4 (±14.9)	42.6 (±15.8)

Table 6.3. Amount of knee extension from maximum flexion to ball impact. Data is presented as mean (degrees) \pm (standard dev.)

Open Stance (OS		Square Stance (SS)	
Left Knee	27.3 (±4.1)	26.1 (±9.1)	
Right Knee	36.2 (±8.6)	22.0 (±10.3)	

6.3.3 Peak knee external rotation moment versus knee flexion angle

There was, however, a difference between the knee flexion angles when the peak external rotation moment occurred. The mean OS right knee and SS left knee flexion angles were 49.5 degrees (± 13.2) and 19 (± 4.2) degrees respectively, p=0.031. This indicates that the peak external rotation moments occur with the knee in a more extended position in the SS follow-through.

6.4 Forces and moments of the knee joint

6.4.1 Open stance

Typical examples of the resultant reaction forces and moments for the knee as a function of time in the OS stroke are shown in figure 6.6.

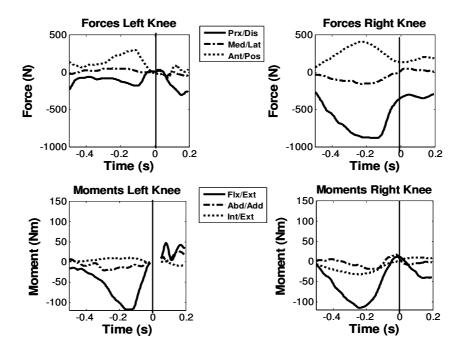


FIGURE 6.6. OS forces and moments for one representative subject. Positive values correspond to proximal, medial and anterior forces and flexion, abduction and internal rotation moments. Negative values correspond to distal, lateral and posterior forces and extension, adduction and external rotation moments. Time zero = ball contact. For the period when the foot was off the force plate, no moments were calculated.

Forces. Table 6.4 lists the average maximum normalised knee joint reaction forces recorded during the forward swing and follow-through phases. The dominant knee force was the distal (weight bearing) force. In the forward swing phase, the greatest distal forces were found in the right knee (p=0.014), whereas, in the follow-through, greater distal forces were found in the left knee. During the beginning of the forward swing phase the right knee distal, lateral and anterior forces all reached their peak. In the left knee the anterior and distal forces were of similar magnitude during the forward swing. The anterior force peaked during the forward swing phase; the distal force peaked during the forward swing phase; the distal force peaked during the follow-through.

TABLE 6.4. Average maximum knee forces recorded during the forward swing and followthrough phases for the OS forehand. Data is represented as maximum force (F) divided by body weight (BW).

FORWARD SWIN	G	F/BW (%) [± standard dev.]				
	Prox	Dist	Med	Lat	Ant	Pos
Left Knee	NSD	36.8 (±28.2)*	6.7 (±5.6)	NSD	39.6 (±13.5)	NSD
Right Knee	NSD	104.5 (±25.8)	NSD	14.4 (±8.8)	24.0 (±16.1)	NSD
FOLLOW-THROUGH						
	Prox	Dist	Med	Lat	Ant	Pos
Left Knee	NSD	71.2 (±21.7)	NSD	4.7 (±2.7)	14.5 (±7.3) ⁶	7.8 (±2.9) ⁵
Right Knee	NSD	NSD	NSD	NSD	NSD	NSD

NSD = Not sufficient data (data for less than 4 subjects), ⁿ Number of subjects (if not seven). * Significantly less than the value for the same knee in the SS stroke.

Moments. The largest knee moments (in decreasing order of magnitude) were flexion/extension, abduction and external rotation (table 6.5). The greatest moments were left knee extension moments which peaked during the forward swing for all subjects. During the forward swing phase six subjects displayed both extension followed by flexion moments in the right knee. One subject only showed a flexion moment for the right knee. The right knee abduction moments peaked during the follow-through. The major transverse plane moments were right knee external rotation which peaked at the start of the forward swing phase (figure 6.8b).

TABLE 6.5. Average maximum knee moments recorded during the forward swing and followthrough phases for the OS forehand.

FORWARD SWING			Nm/kg (± standard dev.)				
	Flex	Ext	Abd	Add	Int Rot	Ext Rot	
Left Knee	NSD	1.46 (±0.53)	0.2 (±0.17) [▲]	0.23 (±0.10) ⁶	0.14 (±0.07)	0.05 (±0.07)•	
Right Knee	0.55 (±0.29)	0.71 (±0.43) ⁶	0.33 (±0.14) ⁶	NSD	NSD	0.24 (±0.13)	
FOLLOW-TH	ROUGH						
	Flex	Ext	Abd	Add	Int Rot	Ext Rot	
Left Knee	0.54 (±0.26) ⁶	NSD	0.32 (±0.16)	NSD	NSD	0.10 (±0.09)	
Right Knee	NSD	NSD	NSD	NSD	NSD	NSD	

NSD = Not sufficient data (data for less than 4 subjects), ⁿ Number of subjects (if not seven).

[★]■[●]▲ Significantly less than the value for the same knee in the SS stroke

6.4.2 Square stance

Typical examples of the resultant reaction forces and moments for the knee as a function of time in the SS stroke are shown in figure 6.7.

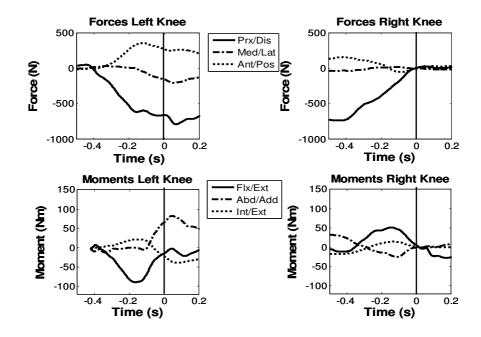


FIGURE 6.7. SS forces and moments for one representative subject. Positive values correspond to proximal, medial and anterior forces and flexion, abduction and internal rotation moments. Negative values correspond to distal, lateral and posterior forces and extension, adduction and external rotation moments. Time zero = ball contact.

Forces. Table 6.6 lists the average maximum normalised knee joint reaction forces recorded during the forward swing and follow-through phases. The major right leg force was distal which decreased during the forward swing phase reducing to close to zero during the follow-through. The right knee medial/lateral forces moved from lateral at the beginning of the forward swing to medial just before ball impact. Greater medial/lateral and anterior/posterior forces were found in the left knee. The left knee distal force reached its peak during the follow-through. The anterior forces peaked during the forward swing phase. Medial forces were present in the left knee at the beginning of the forward swing phase which turned to lateral forces just before ball impact for all subjects

TABLE 6.6. Average maximum knee forces recorded during the forward swing and followthrough phases for the SS forehand. Data is represented as maximum force (F) divided by body weight (BW).

FORWARD SWIN	G		F/BW (%) [± standard dev.]			
	Prox	Dist	Med	Lat	Ant	Pos
Left Knee	NSD	79.2 (±10.5)*	13.7 (±5.3)	10.8 (±3.5) ⁵	46.9 (±8.4)	NSD
Right Knee	NSD	98.9 (±14.6) ⁵	7.0 (±9.0) ⁴	12.2 (±13.2) ⁵	31.8 (±13.1) ⁵	9.0 (±6.6) ⁵
FOLLOW-THROU	GH					
	Prox	Dist	Med	Lat	Ant	Pos
Left Knee	NSD	89.6 (±7.4)	NSD	12.7 (±8.2)	18.5 (±12.3)	NSD
Right Knee	NSD	NSD	NSD	NSD	NSD	NSD

NSD = Not sufficient data (data for less than 4 subjects), ⁿ = number of subjects (if not seven).* Significantly greater than the value for the same knee in the OS stroke.

Moments. The largest knee moments (in decreasing order of magnitude) were flexion/extension, abduction/adduction and internal/external rotation (table 6.7). The greatest moments were left knee extension moments which peaked during the forward swing for all subjects. The frontal plane left knee moments changed from adduction at the very start of the forward swing to abduction moments which peaked during the follow-through phase. During the forward swing phase all subjects displayed a net internal rotation moment in the left knee which turned to an external rotation moment before impact and peaked during the follow-through (figure 6.8c).

TABLE 6.7. Average maximum knee moments recorded during the forward swing and followthrough phases for the SS forehand. Data is represented as peak moment divided by body mass.

FORWARD SWING			Nm/kg (± standard dev.)			
	Flex	Ext	Abd	Add	Int Rot	Ext Rot
Left Knee	NSD	1.31 (±0.20)	0.70 (±0.14) [▲]	0.23 (±0.28)	0.13 (±0.09)	0.22 (±0.07)•
Right Knee	NSD	NSD	NSD	NSD	NSD	NSD
FOLLOW-TH	Flex	Ext	Abd	Add	Int Rot	Ext Rot
Left Knee	0.58 (±0.21) ⁵	0.62 (±0.59) ⁴	0.95 (±0.26)	NSD	NSD	0.31 (±0.10)
Right Knee	NSD	NSD	NSD	NSD	NSD	NSD

NSD = Not sufficient data (data for less than 4 subjects), ⁿ = number of subjects (if not seven).

[★]■[●]▲ Significantly greater than the value for the same knee in the OS stroke

6.4.3 Internal/external rotation moments

Since the temporal phase patterns of internal and external rotation moments for all subjects were very similar, an ensemble average curve could be plotted. The average curve profiles of the transverse plane knee moments during the OS and SS forehands for all subjects are shown in figures 6.8a, 6.8b and 6.8c respectively. In the left knee for both the OS and SS strokes, all subjects had a similar pattern of internal rotation (positive peak) before ball impact and an external rotation moment (negative peak) after impact.

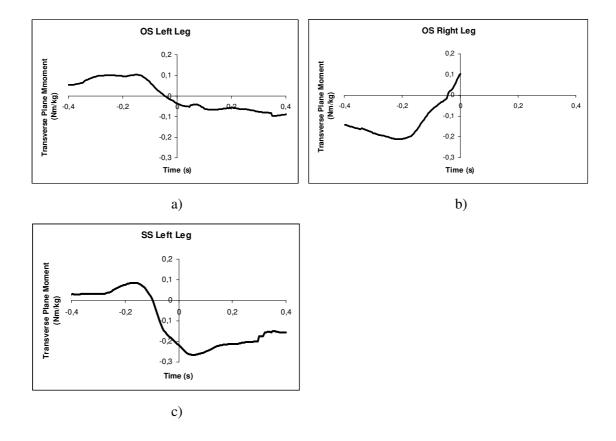


FIGURE 6.8. Left and right knee average transverse plane moment profile (7 subjects). a) OS left knee b) OS right knee and c) SS left knee. Positive values correspond to internal rotation moments. Negative values correspond to external rotation moments. Time zero = ball contact.

In the right knee for the OS stroke, all subjects had a similar pattern of external rotation (negative peak) before ball impact (figure 6.8b).

Since the location of the CoP was inconsistent for most subjects for the right force plate during the OS follow-through and SS forward swing and follow-through, it was not possible to produce a coherent average plot for the SS right knee. This was because most of the weight had been transferred to the left leg at the beginning of the forward swing phase.

6.4.4 Kinetic comparison of the open stance and square stance strokes

Forces. The left knee distal force was greater for the SS in the left leg during the forward swing phase (79.2 ±10.5 % BW) compared to the left leg in the OS forward swing phase (36.8 ±28.2 % BW), (p=0.016). No differences were found between the OS and SS left knee medial forces in the forward swing phase. No differences were found between the OS and SS left knee anterior forces in the forward swing phase. During the follow-through, there was a trend for the lateral forces in the SS left knee (12.7 ±8.2 % BW), to be greater than the OS left knee (4.7 ±2.7 % BW), which approached significance (p=0.063). No differences were found when comparing the left knee distal forces in the follow-through.

Moments. There was a greater abduction moment in the left knee in the SS forward swing phase (0.70 ±0.14 Nm/kg) compared with the left knee in the OS forward swing phase (0.20 ±0.17 Nm/kg), (p=0.016). There was also a greater abduction moment in the left knee in the SS follow-through stage (0.95 ±0.26 Nm/kg) compared to the left knee in the OS follow-through (0.32 ±0.16 Nm/kg), (p=0.016). The largest moments in the transverse plane were the external rotation moments of the left knee in the SS follow through. The external rotation moments in the left knee were significantly greater in the SS stroke than the OS stroke during both the forward swing (0.22 ±0.07 Nm/kg) compared with (0.05 ±0.07 Nm/kg), respectively (p=0.016) and the follow-through phases (0.31 ±0.10) compared to (0.10 ±0.09), respectively (p=0.031). This can be seen from the average curves in figure 6.8. However, there was no significant difference between the right knee OS external rotation moment in the forward swing phase (0.24 ±0.13 Nm/kg) when compared with the SS left knee external rotation moment of the follow-through (0.31 ±0.10 Nm/kg), (p=0.30).

7 DISCUSSION

The primary purpose of this study was to calculate and compare the reaction forces and moments of the knee joint in the OS and SS tennis forehand.

The main findings of this study were: (a) the largest abduction moments occurred in the left knee of the SS stroke during the follow-through. (b) The largest external rotation moments were found in the SS left knee during the follow-through and in the OS right knee prior to impact. (c) The peak OS external rotation moment occurred with the right knee in a more flexed position compared to the SS left knee peak external rotation moment. (d) The greatest peak ground reaction torque (T_Z) was generated by the right leg in the OS stroke before impact.

As there was no significant difference between the right knee OS external rotation moment in the forward swing phase when compared with the SS left knee external rotation moment of the follow-through, the research hypothesis was rejected.

It was expected that the greatest external rotation moments would be found in the left knee during the follow-through phase of the SS stroke. The fact that the OS generates a similar level of external rotation moment in the right knee in the forward swing phase was not expected initially. Although, considering that the OS does not use any initial forward stepping action before impact, it is reasonable that the right knee external rotation moment could have an important role in assisting pelvic rotation.

7.1 Stroke parameters

The obtained results show that in terms of angular velocity of the pelvis, racket head speed, ball speed, swing time and ratio of ball impact height; there were no significant differences between the two strokes. This means that the differences between the two strokes were likely to be due to the differences in the technique and not because of differences in for example, how fast the subjects swung the racket.

7.2 Comparison with the literature

In order to evaluate the validity of the kinetic and other data values obtained in this study, a comparison was made with the values found in the literature. As very little relevant tennis data exists, data from other sports with similar type of movement patterns was used.

Knee moments. A comparison of the knee moments found in this study and those reported in the literature for similar types of athletic movements is shown in table 7.1.

TABLE 7.1. Comparison between mean peak moments experienced by the knee during the tennis forehand, golf, walking, running, a running side step manoeuvre and baseball pitching.

Peak Knee Moment	Peak Knee Moment Nm/kg (± standard de	
Extension	Abduction	External Rotation
Moment	Moment	Moment
1.31 (0.20)	0.95 (0.26)	0.31 (0.10)
1.0 (0.26)		
0.71 (0.43)	0.33 (0.14)	0.24 (0.13)
1.26	0.83	0.36
		0.25
	0.54 (0.25)	
0.2	0.3	0.15
0.4	0.7	0.23
2.2	0.2	0.32
e) ^g	1.00	
e) ^g		0.95
	Extension Moment 1.31 (0.20) 1.0 (0.26) 0.71 (0.43) 1.26 0.2 0.4	Extension Moment Abduction Moment 1.31 (0.20) 0.95 (0.26) 1.0 (0.26) 0.33 (0.14) 1.26 0.83 0.54 (0.25) 0.2 0.4 0.7 2.2 0.2 e) ^g 1.00

a = Iino and Kojima 2003, b = Gatt et al. (1999, 20), c = Lynn and Noffal (2010), d = Allard (1997, 315), e = Kerrigan et al. 2009, f = Besier et al. (2001), and g = Huang and Lin (2011). Values in bold refer to the present study.

The average maximum extension moment found in this study in the SS left knee was slightly greater than that reported by Iino and Kojima (2003) for the closed stance forehand, although the mean resultant racket speeds described in their article and in this study were similar, 28.2 ± 3.3 and 29.2 ± 1.5 m/s respectively.

No published data for the abduction moment of the knee of the tennis forehand was found. However, when compared to data for similar types of movements in other sports, the values obtained for the maximum abduction moment in the leading knee of the SS in this study were similar to those reported for the leading knee in baseball pitching but greater than those reported for walking and running.

The peak external rotation moments found in the tennis shots are similar to those calculated for a golf drive and running sidestep manoeuvre, but are slightly greater than the peak external rotation moments reported in walking and running.

Weight transfer. The timing of the observed weight transfer for both the OS and SS strokes was consistent with the descriptions of proper technique found in coaching texts and previous research on the OS and SS forehand (Bahamonde & Knudson 2003). As noted in section 2.1.1 and 2.1.2.

Vertical torque. Force plate torque values for sports techniques are rarely reported in the literature (Lees & Lake 2008, 65; Miller 1990, 220). The mean peak force plate T_Z values for the right leg (Table 6.1) in the OS are similar to those reported for the leading leg in golf (0.27 Nm/kg) using metal spiked shoes (Worsfold et al. 2008).

Anterior/posterior ground reaction forces. The anterior/posterior ground reaction forces for the SS forehand shown in figure 6.3 agree with the propulsion followed by braking prior to impact forces illustrated by van Gheluwe and Hebbelinck (1986) and lino and Kojima (2001) shown in section 2.1.3.

Knee flexion angles. The maximum flexion angles of the left and right knees of 42.6° and 53.9° in the SS stroke were similar to the flexion angles of the left and right knees of 50.2° and 54.7° , respectively, reported by and Iino and Kojima (2003) for the CS forehand.

Conclusion. Comparison with the literature has shown that the range of the data values obtained in this study were at a reasonable level. Although, the mean peak values are similar to those found in golf and the external rotation moment was lower than that

found in the trail knee in baseball pitching, the injury risk of tennis compared to these sports is not so clear. In tennis, more forehands are likely to be hit during a normal match than one would execute the drive in golf or pitch a baseball during a normal game of these respective sports.

7.3 Open stance forces and moments

Since there is no published data concerning OS knee forces and moments it is not possible to directly compare the OS results with previous studies. Therefore, only an indirect comparison can be made with actions from other aspects of human motion.

The loading of the knee joint corresponds with the flexion and extension of the knee and the location of the player's weight. During the backswing phase the players first flexed their knees with the majority of the weight over the right leg for 6 of the 7 subjects. During the flexion the right force plate torque and right knee forces increased. At approximately 0.2 seconds before ball impact the knees had reached their peak flexion angle (table 6.2), the pelvis had reached its lowest point (figure 6.5) and the force plate torques their maximum values (table 6.1). It is at this stage when the forward swing phase commenced. The right knee started to extend and the right vertical ground reaction force (F_Z), ground reaction torque (T_Z), and right knee distal, lateral and anterior forces all reached their peak values as shown in figures 6.1a, 6.1c, 6.6, 6.8 and tables 6.4 and 6.5. This was the peak loading period of the knee for the OS stroke which was characterised by the right leg forcefully extending and externally rotating.

An example of the GRFVs during the peak loading period is shown in figure 7.1a. This figure shows the greater resultant GRFV that was generated by the right leg at the start of the forward swing. During this peak loading period the right femur was pushing (in decreasing order of magnitude) distally, anteriorally and laterally. The left femur was pushing down onto the left tibia in a similar way although with less force. An illustration of the directions of these forces for the knee joint is shown in figure 7.1b.

Whereas the right knee forces were all distal during the forward swing phase, the right knee flexion/extension moments changed from first extension and then to flexion for all

subjects except one. This was mirrored by the GRFV moving from behind the knee to in front of the knee. This change from extension to flexion moment could possibly be caused by the centre of mass of the player moving anteriorally during the forward swing. This forwards motion of the CoM is likely to cause the GRFV to move in front of the knee, hence the net flexion moment just prior to impact.

For the subject who transferred their weight before impact, the GRFV was in front of the right knee throughout the forward swing phase and so displayed a net extension moment throughout the forward swing. This highlights how movement patterns that can appear similar to the naked eye can be caused by different joint moments (Knudson & Elliott 2004, 165).

The left knee flexion/extension moments were extension for all subjects during the forward swing phase. The abduction/adduction and internal/extension moments were lower than those for the right knee.

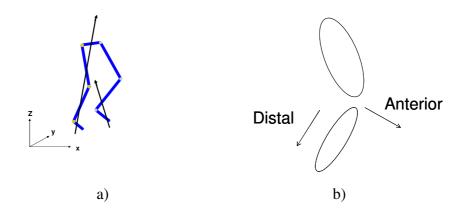


FIGURE 7.1. a) Stick figure displaying the left and right leg resultant GRFs during the right knee loading phase at the start of the forward swing for a representative subject, (t = -0.2s). b) A diagram to illustrate the directions of the distal and anterior forces acting on the right knee. The lateral knee force is perpendicular to the distal and anterior forces, pointing out of the page.

The anterior/posterior and medial/lateral components of the resultant GRFV shown in figure 7.1 are shown graphically in figure 7.2. At the start of the forward swing the right leg produced a propulsive force while the left leg produces a braking force. (These are the same forces that are illustrated with respect to time in figure 6.2). These transverse plane impulses resembled the couple found during the double support phase

of walking. During the double support phase of walking this type of couple induces body rotation around the vertical axis (Zatsiorsky 2002, 63). Therefore, in the tennis stroke, it is likely that these impulses play an important role in rotating the pelvis and form the starting point of the kinetic chain described in section 2.3.

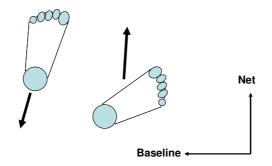


FIGURE 7.2. An illustration of the resultant x and y components of the ground reaction forces for the OS 0.2 s before impact for a representative subject. The ground reaction forces resemble a couple in the transverse plane.

For the player who had transferred the weight to the left leg before impact, the GRFs were approximately body weight at impact for the left leg and minimal for the right leg. The left knee forces reached their peak values and the right knee forces their minimal values during the forward swing for this subject.

At ball impact the pelvis had reached its maximum vertical position and the ground reaction forces are virtually zero for the left leg, and minimum values for the right leg for 6 of the 7 subjects. This indicated that the players had generated high levels of vertical impulse to nearly leave the surface of the force plate. The corresponding knee force and moment values were at their minimum values during ball impact.

During the follow-through phase most of the weight had been transferred to the left leg, and so the dominant left knee force was acting distally. The left knee abduction and external rotation moments both peaked during the follow-through as the body aimed to safely dissipate the rotation energy. Even though the greatest left knee moments during the follow-through were flexion moments, the values were not large and varied from player to player. This variation could be caused by how much the player leaned forward with their torso after hitting the ball as that would effect the position of the CoM and hence the degree to which the GRFV was in front of the left knee.

Despite the fact that the right leg was powerfully extending during the forward swing, higher extension moments were found in the left knee. This unexpected finding could be explained by analysing the forces and moments of the ankle joint which contribute to the resultant knee moment. The posterior reaction force of the left ankle joint appeared to be the major contributor to the left knee extension moment. That is, the left tibia is forcefully pushing forwards against the foot during the left knee extension before ball impact.

7.4 Square stance forces and moments

Similarly with the OS stroke, the loading of the knee joint in the SS forehand also corresponds with the location of the player's weight. During the forward step (approximately 0.5 to 0.4 seconds before ball contact) all the weight was over the right leg, the left F_Z was zero (figure 6.1b) and the left knee forces were minimal (figure 6.7). This forward stepping action is illustrated in figure 7.3. At this stage the right knee forces were predominately acting in the distal direction.

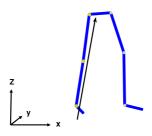


FIGURE 7.3. Diagram illustrating the stepping phase of the SS forehand, 0.5 s before impact for a representative subject.

When the left foot touches the force plate the left foot experienced an initial braking force (figure 6.3) and the left knee forces and moments rapidly increased as weight was being transferred to the left leg

Most subjects commenced the forward swing at the moment of weight transfer. Although, one subject started the forward swing just before and one subject just after weight transfer. At the start of the forward swing the right leg generated a forwards impulse and the left leg a braking impulse (figure 6.3) which are illustrated in figures 7.4 and 7.5. The ground reaction forces for the SS were of a similar pattern to those illustrated for the forward swing of the closed stance forehand illustrated by Iino and Kojima (2001).

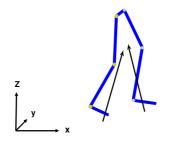


FIGURE 7.4. Weight transfer from right to left leg. Note the braking force from the left leg and the propulsive force from the right leg. The positions of the GRFVs were similar at the start of forward swing for all subjects.

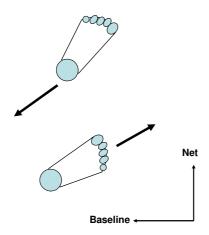


FIGURE 7.5. Transverse view of the braking (left foot) and propulsion (right foot) ground reaction forces shown in figure 7.4.

Although the position of the feet was different from the OS stroke these transverse plane ground reaction forces at the start of the forward swing phase also resemble the couple described in section 7.3 for the OS stroke. And therefore, these forces may also play an important role in rotating the pelvis.

At the start of the forward swing the GRFVs were anterior and medial to the right knee and posterior for the left knee (figure 7.4). This helps explain why the main right knee moments were flexion and abduction while the main left knee moments were extension for all subjects at the start of the forward swing (figure 6.7).

During the forward swing (approximately -0.2 to 0 seconds), the left knee distal and lateral forces and left knee abduction and external rotation moments all started to increase. During the same period, the right knee forces all decreased as the weight had been transferred to the left leg.

After ball impact and during the follow-through (0 to 0.3 seconds) virtually all the weight was over the left leg and so the right knee forces were almost zero (figure 6.7). Since virtually all the weight is over the left leg, the joints in left leg have to endure the forces and moments necessary to dissipate the rotation energy of the pelvis and upper body. This caused the distal and lateral forces, and abduction and external rotation moments to reach their peak (figures 6.7 and 6.8c and tables 6.6 and 6.7). The left knee abduction moment could be explained by the GRFV being medial to the left knee (figure 7.6).

FIGURE 7.6. 0.1s after impact. The GRFV is medial to the left knee, hence the abduction moment. At this stage of the follow-through, the left knee abduction and external rotation moments reach their peak. Note also that there is virtually no weight on the right leg.

7.5 Suggestion for the role of the lower extremities in generating pelvic rotation

Ino and Kojima (2001) have stated that the role of the legs in the forehand is to generate pelvic rotation. The OS and SS strokes appear to use the lower body in different ways to generate this pelvic motion.

In the OS stroke; the greatest vertical ground reaction force (F_Z), ground reaction torque (T_Z), distal, lateral and anterior forces, abduction and external rotation moments were found in the right leg during the forward swing phase. Therefore, in the OS stoke, the right leg seems to play the dominant role in aiding pelvic rotation by first flexing and then forcefully extending whilst pushing laterally and externally rotating. The high extension moments found in the left knee which contribute to the posterior ground reaction forces, suggest the left leg also plays a role in assisting pelvic rotation before impact.

Whereas, in the SS stroke; higher levels of ground reaction torque (T_z), anterior forces, extension, abduction and external rotation moments were found in the left leg during the forward swing phase. Therefore, the SS seems to use the stepping action to help generate pelvic rotation. As the player steps forward he is generating the linear momentum described in coaching texts such as (Groppel 1992, 77), but as the left foot hits the ground, the left leg experiences a breaking impulse that is unlikely to be directed through the CoM of the player. This is similar (although of lesser magnitude) to the breaking impulses seen in the planting of the front leg in the delivery stride of the javelin throw (Bartonietz 2000, 420; Bartlett 2007, 196). This combined with the propulsive force from the right leg (figure 7.4) probably helps to convert the linear motion into angular motion of the pelvis.

7.6 Error sources and limitations

The limitations of the applied model were noted earlier on section 3.1.6. Therefore, only the limitations of the measurements are discussed in this section.

The force plates suffered from drift during the data collection. This was corrected as best as possible as described in the Analytical Procedures (section 5.3). The location of the CoP and its location relative to the GCS is very important for the precision of the inverse dynamics calculations. Even with correctly calibrated force plates the accuracy of the CoP location is largely influenced by the magnitude of the vertical GRF (Robertson et al, 2004, 92). The inconsistent CoP locations which were mainly found for the right leg during the follow-through phase could be explained by the low vertical GRFs. This was because the player's weight had already shifted to the front (left) leg.

A limitation in the inverse dynamic equations programmed in Matlab meant that no joint moments could be calculated when the foot is not in contact with the force plate (e.g. left knee moments in figure 6.6). If the joint moments are required during the time the foot is off the force plate, the Matlab script could be modified accordingly.

The data from one subject had to be discarded because their joint angle values were unrealistic. Since joint angles are calculated from marker placement, this highlights how important not only correct marker placement is but also the markers must not move relative to their anatomical landmarks.

The knee moments were calculated up to the end of the follow-through phase. In a normal tennis match the player would quickly return to the ready position immediately after the follow-through. It is possible that during that recovery period the player would create larger external rotation moments than those found in this study. Due to the restrictions on the amount of high speed video data that could be captured, it was not possible to record the recovery step in this study.

The surface of the force plate was smooth aluminium. Tennis courts such as acrylic will have a higher coefficient of friction. Therefore, the vertical torque (T_Z) and

subsequent knee external rotation moments calculated from this study may have been larger if a representative surface could have been fixed onto the top of the force plate.

An additional limitation of this study was the size of the force plates and the size of the area that the data could be reliable captured. This did not allow a great deal of movement prior to hitting the ball. During match conditions the players are more likely to be running prior to stabilising themselves for the forward swing.

7.7 Coaching implications

On the basis of this study one cannot advocate one type of forehand over another. Although, coaches who train players with knee problems or those recovering from knee injury or knee surgery should be aware that the OS forehand loads the right knee joint before impact while the SS forehand loads the left knee after contact. Therefore, a player recovering from right knee injury may want to try to utilise more linear motion in their forehand. Similarly, a player recovering from a left knee injury may want to use a more open stance.

7.8 Suggestions for further study

The role of the lower body in rotating the pelvis in the OS forehand could be examined in greater detail. This has been performed on the CS forehand (Iino & Kojima 2001) and backhand (Akutagawa & Kojima, 2005), but not on the OS forehand. The results of this work suggest that the right leg may make a greater contribution to pelvic torque in the OS forehand compared to the levels reported by Iino & Kojima (2001) for the CS forehand.

Electromyographic techniques could be incorporated into an inverse dynamic study of the lower body. By incorporating the electromyography (EMG) data, an understanding of the muscle activation patterns and timing, and their role in generating the resultant joint moments and forces could be obtained.

In the OS stroke the player can transfer their weight onto the leading foot before or after

contact. In the SS stroke the weight transfer occurs before contact. So, it is possible that if the forward weight transfer occurs before contact in the OS stroke, the T_Z forces and external rotation moments in the right leg may be lower. Therefore, a study could be performed to investigate the differences between the loading of the lower body in the OS stroke for weight transfer before contact and weight transfer after contact.

The linear and angular momentum, and total body energy of the OS and SS forehands could be calculated. For example, if the total body mechanical energy can be calculated it could be used to show how much energy the supporting leg has to dissipate in the follow-through. Angular momentum has been calculated for the tennis backhand (Wang et al. 2010) and serve (Bahamonde, 2000), while power analysis has been used to analyse the golf swing (Nesbit, 2005), but to date no data is available for the tennis forehand. This is despite the fact that the conversion of linear to angular momentum is frequently described in coaching texts (Papas, 2008).

A more complex model of the knee joint could be used to estimate tendon, ligament and joint contact loads such as that used by Shelburne et al. (2005) or an EMG driven neuromuscular model used by Lloyd et al. (2005). Commercially available computer simulators such as LifeModeler[®] KneeSim (LifeModeler Inc. San Clemente, California, USA) could also be used. This could be combined with further investigations to highlight the potential for the tennis forehand to cause an overuse injury in the knee.

There are slight differences in the anatomy of the knee joint in men and women (Palastanga, 2002), and the loading of the knee joint has been shown to differ between men and women in other movements (Pollard et al, 2007). Therefore, a study to investigate the sex differences in the knee kinetics of the tennis strokes could be beneficial.

8 CONCLUSIONS

In this study, both OS and SS forehands displayed similar levels of pelvic angular velocity, swing times and racket head speeds. Therefore, although the two types of shot involved contrasting movement patterns of the lower body, the effect on the ball was the same and so both shots used the lower body in different ways to produce the same end result.

Although the peak loads are similar, the OS forehand creates greater loads on the right knee joint before impact while the SS forehand creates greater loads on the left knee during the follow-through. This loading occurs with the knee in a more flexed position in the OS forehand. The extension, abduction and torsional forces generated in the right leg assist in the generation of the angular motion of the pelvis in the OS forehand whereas, the checking action of the left leg assists in the conversion of linear to angular motion in the SS forehand.

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10.1 Definition of the Euler angles using the Z, line of nodes, z rotation sequence

To derive the equations for the angular velocity of the segment the in terms of the LCS (equation 14), three successive rotations of the LCS will be described. The rotations start with the axes of the LCS coinciding with the axes of the GCS and ends with the LCS back in its original orientation. The first rotation takes place about the global Z axis, the second about the line of nodes and the third about the z axis of the segment. The complete sequence is shown in figure 10.1. The line of nodes is defined as the intersection of the x[']_{segment}, y[']_{segment} and the X, Y planes (figure 10.1b). This can be expressed mathematically as the cross product of the Z and $z_{segment}$ axes (Goldstein 1980, 145; Jose & Saletan 1998, 526; Vaughan et al 1999, 98).

The sequence commences by rotating the LCS, by an angle φ counter-clockwise about the global Z axis (figure 10.1a), the resultant coordinate system is labelled $x''_{segment}$, $y''_{segment}$, $z''_{segment}$. Next, the $x''_{segment}$, $y''_{segment}$, $z''_{segment}$ axes are rotated counterclockwise about the line of nodes by an angle θ to produce another set of intermediate axes $x'_{segment}$, $y'_{segment}$, $z'_{segment}$ (figure 10.1b). Finally, the $x'_{segment}$, $y'_{segment}$, $z'_{segment}$ axes are rotated counter-clockwise about the $z'_{segment}$ by an angle ψ to produce the desired $x_{segment}$, $y_{segment}$, $z_{segment}$ axes (figure 10.1c). This is now the original orientation of the LCS.

These three rotations can be described numerically using transformation matrices. The three matrices are:

$$\mathbf{A} = \begin{vmatrix} \cos\phi & \sin\phi & 0 \\ -\sin\phi & \cos\phi & 0 \\ 0 & 0 & 1 \end{vmatrix}$$
(23)

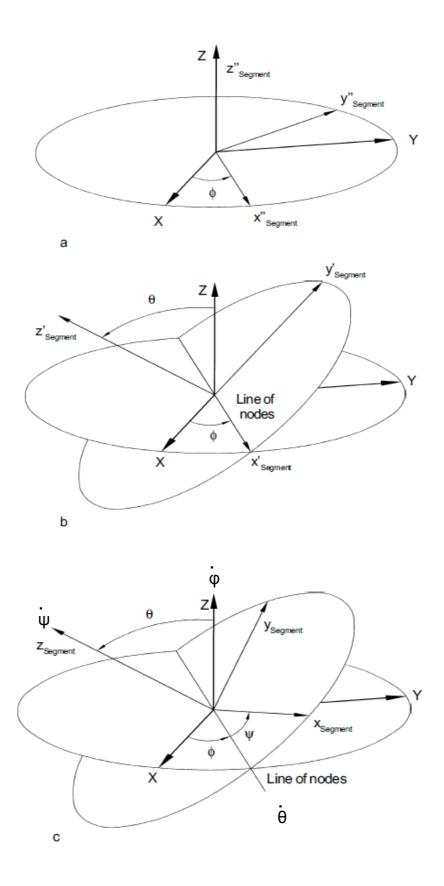


FIGURE 10.1. Definition of the 3 Euler angles ($\varphi_{segment}$, $\theta_{segment}$, $\psi_{segment}$). The rotations take place in the following order: (a) $\varphi_{segment}$ about the Z axis; (b) $\theta_{segment}$ about the line of nodes; and (c) $\psi_{segment}$ about the z segment axis (Vaughan et al. 1999, 98).

$$\mathbf{B} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos\theta & \sin\theta \\ 0 & -\sin\theta & \cos\theta \end{bmatrix}$$
(24)
$$\mathbf{C} = \begin{bmatrix} \cos\psi & \sin\psi & 0 \\ -\sin\psi & \cos\psi & 0 \\ 0 & 0 & 1 \end{bmatrix}$$
(25)

where **A**, **B** and **C** represent the rotation matrices for the first, second and third rotation respectively.

The complete sequence of rotation is described by taking the product of these 3 matrices

$$\mathbf{D} = \mathbf{CBA} \tag{26}$$

which gives the following transformation matrix:

 $\mathbf{D} = \begin{bmatrix} \cos\psi\cos\phi - \cos\theta\sin\phi\sin\psi & \cos\psi\sin\phi + \cos\theta\cos\phi\sin\psi & \sin\psi\sin\theta \\ -\sin\psi\cos\phi - \cos\theta\sin\phi\cos\psi & -\sin\psi\sin\phi + \cos\theta\cos\phi\cos\psi & \cos\psi\sin\theta \\ \sin\theta\sin\phi & -\sin\theta\sin\phi & \cos\theta \end{bmatrix}$ (27)

10.2 Angular velocity of the segment

If the vector $\boldsymbol{\omega}$ represents the resultant instantaneous angular velocity of the LCS, then $\boldsymbol{\omega}$ can be described as the sum of the angular velocity vectors pointing in the direction of the axes of rotation:

$$\boldsymbol{\omega} = \boldsymbol{\omega}_{\varphi} + \boldsymbol{\omega}_{\theta} + \boldsymbol{\omega}_{\psi} \tag{28}$$

where ω_{ϕ} , ω_{θ} , and ω_{ψ} represent the components of the angular velocity vector ω about the global Z axis, the line of nodes and the z_{segment} axes respectively.

However, for calculating the net moment of the segment, the components of ω along the segment axes (x_{segment}, y_{segment}, z_{segment}) need to be defined. The time derivatives of the Euler angles cannot be used directly, as that gives the components ω about 3 axes of rotation (as shown in figure 10.1c and described by equation (26)). Fortunately, the matrices in equations (22 to 25) can be used to find the components of the angular velocity vectors along each one of the axes of the LCS. This is achieved by finding the respective x_{segment}, y_{segment}, z_{segment} components of ω_{ϕ} , ω_{θ} , and ω_{ψ} and then simply adding them together (Goldstein 1980, 176).

Since ω_{φ} lies along the global Z axis, the complete rotation sequence (matrix **D**) needs to be used. This gives:

$$\mathbf{D}\begin{bmatrix}0\\0\\\vdots\\\phi\end{bmatrix} = \begin{array}{c} \dot{\phi}\sin\theta\sin\psi\\ \dot{\phi}\sin\theta\cos\psi\\ \dot{\phi}\cos\theta\end{array}$$
(29)

where $(\omega_{\varphi}) x_{\text{segment}} = \dot{\phi} \sin \theta \sin \psi$, $(\omega_{\varphi}) y_{\text{segment}} = \dot{\phi} \sin \theta \cos \psi$, $(\omega_{\varphi}) z_{\text{segment}} = \dot{\phi} \cos \theta$

 ω_{θ} lies along the line of nodes which coincides with the $x''_{segment}$ axis. The components of ω_{θ} are obtained by applying only the final transformation i.e. matrix **C**.

$$\mathbf{C} \begin{bmatrix} \dot{\boldsymbol{\theta}} \\ \boldsymbol{\theta} \\ \boldsymbol{0} \\ \boldsymbol{0} \end{bmatrix} = \begin{array}{c} \dot{\boldsymbol{\theta}} \cos \boldsymbol{\psi} \\ \dot{\boldsymbol{\theta}} \cos \boldsymbol{\psi} \\ \dot{\boldsymbol{\theta}} = - \boldsymbol{\theta} \sin \boldsymbol{\psi} \\ \boldsymbol{0} \end{array}$$
(30)

where $(\boldsymbol{\omega}_{\boldsymbol{\theta}}) \mathbf{x}_{\text{segment}} = \boldsymbol{\theta} \cos \boldsymbol{\psi}$, $(\boldsymbol{\omega}_{\boldsymbol{\theta}}) \mathbf{y}_{\text{segment}} = -\boldsymbol{\theta} \sin \boldsymbol{\psi}$, $(\boldsymbol{\omega}_{\boldsymbol{\theta}}) \mathbf{z}_{\text{segment}} = 0$

No transformation is necessary for the components of ω_{ψ} , becuase ω_{ψ} lies along the z_{segment} axis. This means:

$$(\boldsymbol{\omega}_{\boldsymbol{\psi}}) \mathbf{x}_{\text{segment}} = 0, \ (\boldsymbol{\omega}_{\boldsymbol{\psi}}) \mathbf{y}_{\text{segment}} = 0, \ (\boldsymbol{\omega}_{\boldsymbol{\psi}}) \mathbf{z}_{\text{segment}} = \boldsymbol{\psi}$$
 (31)

Adding their respective components gives equation (14), the angular velocity about the local axes of the segment (Goldstein 1980, 145; Jose & Saletan 1998, 526; Vaughan et al 1999, 98):

$$\omega_{segment.x} = \dot{\phi}\sin\theta\sin\psi + \dot{\theta}\cos\psi$$

$$\omega_{segment.y} = \dot{\phi}\sin\theta\cos\psi - \dot{\theta}\sin\psi \qquad (14)$$

$$\omega_{segment.z} = \dot{\phi}\cos\theta + \dot{\psi}$$

10.3 Angular acceleration of the segment

Since acceleration is defined as the second time derivative of displacement, the segment angular accelerations are calculated by taking the derivative of equation (14) which gives the following (Vaughan et al 1999, 97):

$$\dot{\omega}_{segment.x} = \ddot{\phi}\sin\theta\sin\psi + \dot{\phi}\dot{\theta}\cos\theta\sin\psi + \dot{\phi}\psi\sin\theta\cos\psi + \ddot{\theta}\cos\psi - \dot{\theta}\psi\sin\psi$$
$$\dot{\omega}_{segment.y} = \ddot{\phi}\sin\theta\cos\psi + \dot{\phi}\dot{\theta}\cos\theta\cos\psi - \dot{\phi}\dot{\psi}\sin\theta\sin\psi - \ddot{\theta}\sin\psi - \dot{\theta}\dot{\psi}\cos\psi$$
(15)
$$\dot{\omega}_{segment.z} = \ddot{\phi}\cos\theta - \dot{\phi}\dot{\theta}\sin\theta + \ddot{\psi}$$

10.4 Examples of selected raw data (unfiltered, before CoP corrections)

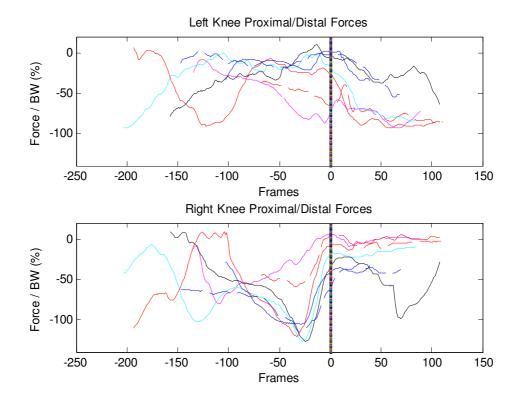


FIGURE 10.2. OS proximal/distal forces. Frame 0 = ball impact.

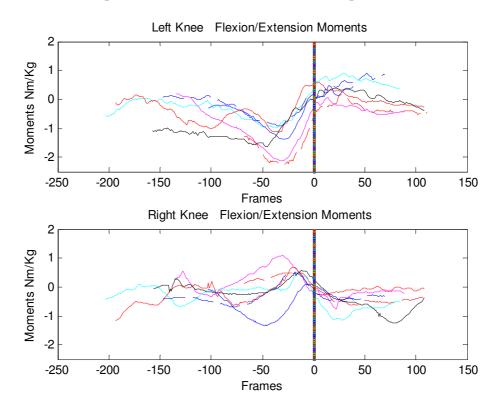


FIGURE 10.3. OS flexion/extension moments. Frame 0 = ball impact.

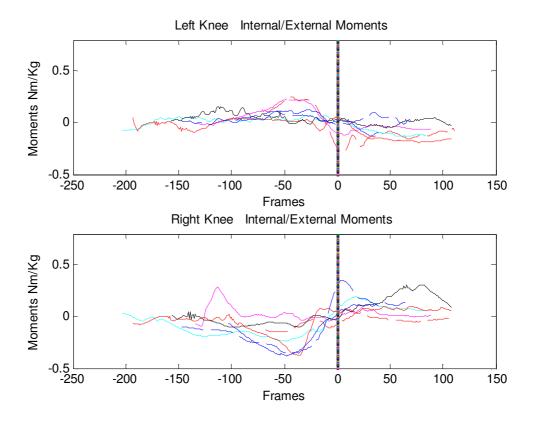


FIGURE 10.4. OS internal/external rotation moments. Frame 0 = ball impact.

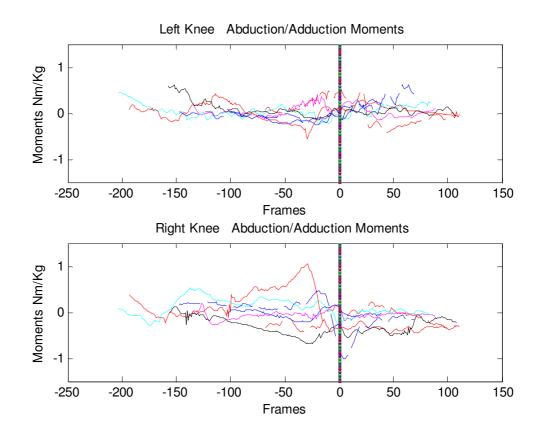


FIGURE 10.5. OS abduction/adduction moments. Frame 0 = ball impact.

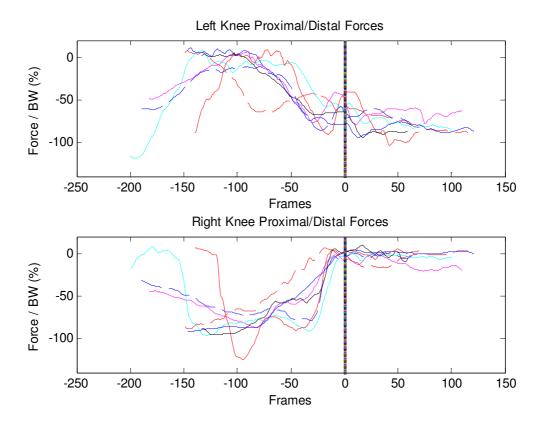


FIGURE 10.6. SS proximal and distal forces. Frame 0 = ball impact.

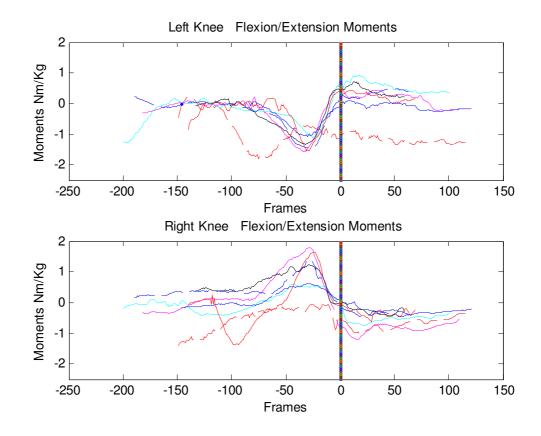


FIGURE 10.7. SS flexion/extension moments. Frame 0 = ball impact.

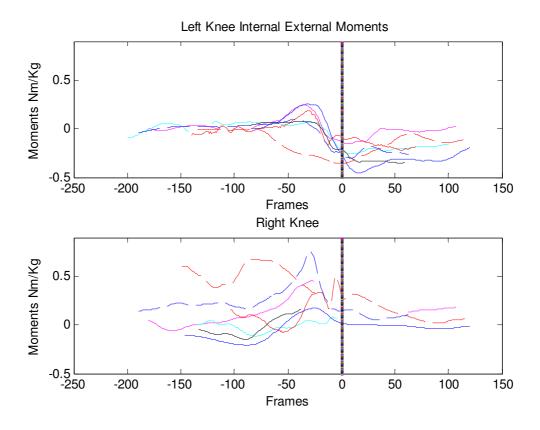


FIGURE 10.8. SS internal/external rotation moments. Frame 0 = ball impact

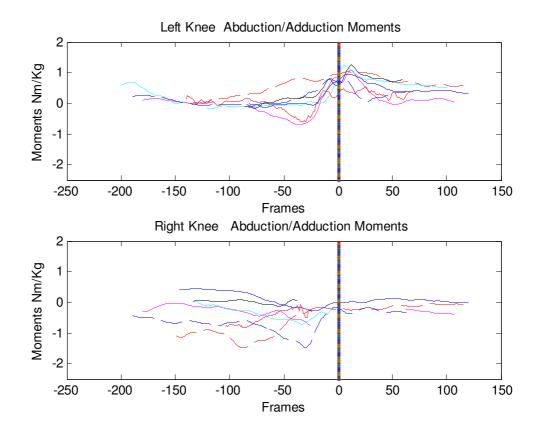


FIGURE 10.9. Square stance abduction/adduction moments. Frame 0 = ball impact.