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Title: Rigidity, counting and equidistribution of quaternionic Cartan chains

Year: 2022

Version: Published version

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## Please cite the original version:

Parkkonen, J., \& Paulin, F. (2022). Rigidity, counting and equidistribution of quaternionic Cartan chains. Annales Mathematiques Blaise Pascal, 28(1), 45-69. https://doi.org/10.5802/ambp. 399

## ANNALES MATHÉMATIQUES



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Volume 28, $\mathrm{n}^{0} 1$ (2021), p. 45-69.
<http://ambp.centre-mersenne.org/item?id=AMBP_2021 $\qquad$ 28_1_45_0>


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> Publication éditée par le laboratoire de mathématiques Blaise Pascal de l'université Clermont Auvergne, UMR 6620 du CNRS

> Clermont-Ferrand - France


MERSENNE

# Rigidity, counting and equidistribution of quaternionic Cartan chains 

Jouni Parkkonen<br>Frédéric Paulin


#### Abstract

In this paper, we prove an analog of Cartan's theorem, saying that the chain-preserving transformations of the boundary of the quaternionic hyperbolic spaces are projective transformations. We give a counting and equidistribution result for the orbits of arithmetic chains in the quaternionic Heisenberg group.


## Rigidité, comptage et équidistribution de chaînes de Cartan quaternioniennes

## Résumé

Dans ce papier, nous montrons un analogue d'un théorème de Cartan, disant que les transformations du bord des espaces hyperboliques quaternioniens qui préservent les chaînes sont des transformations projectives. Nous donnons un résultat de comptage et d'équidistribution pour les orbites de chaînes arithmétiques dans le groupe de Heisenberg quaternionien.

## 1. Introduction

The sphere at infinity $\partial_{\infty} X$ of a negatively curved symmetric space $X$ carries many rich structures, from the geometric, analytic and arithmetic points of view. When the sectional curvature is not constant, the possibilities are particularly rich, for instance with the Carnot-Carathéodory, sub-Riemannian or (hyper) CR structures (see for instance [4, $10,12,14,17]$ ), leading to strong rigidity properties, as Pansu's rigidity theorem for quasi-isometries [18]. Arithmetic subgroups of the isometry group of $X$ endow the sphere at infinity of $X$ with arithmetic structures, and problems of equidistribution of rational points or subvarieties in $\partial_{\infty} X$, as well as in other homogeneous manifolds, have been intensively studied (see for instance $[1,2,6,8,9,11,15,22]$ and many others).

In this paper, we study the quaternionic hyperbolic spaces $X$, whose extreme rigidity is exemplified by the Margulis-Gromov-Schoen theorem in [13], proving, contrarily to the real or complex case, the arithmeticity of lattices in the isometry group of $X$. As announced in [22], we prove a von Staudt-Cartan type of rigidity result for the family of all 3 -sphere chains in the sphere at infinity of $X$, and, analogously to the complex hyperbolic case treated in [20], an effective equidistribution result for the arithmetic

[^0]chains in orbits of arithmetic groups built using maximal orders in rational quaternion algebras.

More precisely, let $\mathbb{H}$ be Hamilton's quaternion algebra over $\mathbb{R}$, with $x \mapsto \bar{x}$ its conjugation, $\mathrm{n}: x \mapsto x \bar{x}$ its reduced norm, $\operatorname{tr}: x \mapsto x+\bar{x}$ its reduced trace. Let $q$ be the quaternionic Hermitian form on the right vector space $\mathbb{H}^{3}$ over $\mathbb{H}$ defined by

$$
q\left(z_{0}, z_{1}, z_{2}\right)=-\operatorname{tr}\left(\overline{z_{0}} z_{2}\right)+\mathrm{n}\left(z_{1}\right)
$$

and $\mathrm{PU}_{q}$ its projective unitary group. It is the isometry group of the quaternionic hyperbolic plane $\mathbf{H}_{\mathbb{H}}^{2}$, realised as the negative cone of $q$ in the right projective plane $\mathbb{P}_{\mathbf{r}}^{2}(\mathbb{H})$, and normalised to have maximal sectional curvature -1 . See Section 2 for a more complete description.

The boundary at infinity $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{2}$ of $\mathbf{H}_{\mathbb{H}}^{2}$ is the isotropic cone of $q$ in $\mathbb{P}_{\mathrm{r}}^{2}(\mathbb{H})$, and the intersections with $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{2}$ of the quaternionic projective lines meeting $\mathbf{H}_{\mathbb{H}}^{2}$ are called chains. We study them, giving their elementary properties and complete geometric descriptions in Section 3. Our first result is similar to Cartan's theorem (see [7, 10]) in the complex hyperbolic case. See Theorem 3.3 for a version in any dimension.

Theorem 1.1. A chain-preserving transformation from the boundary at infinity of the quaternionic hyperbolic plane to itself is a projective unitary transformation.

The boundary at infinity $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{2}$ of $\mathbf{H}_{\mathbb{H}}^{2}$, with the point $\infty=[1: 0: 0]$ removed, identifies by the map $\left(w_{0}, w\right) \mapsto\left[w_{0}: w: 1\right]$ with the quaternionic Heisenberg group

$$
\mathbb{H e i s}_{7}=\left\{\left(w_{0}, w\right) \in \mathbb{H} \times \mathbb{H}: \operatorname{tr} w_{0}=\mathrm{n}(w)\right\}
$$

with group law

$$
\begin{equation*}
\left(w_{0}, w\right)\left(w_{0}^{\prime}, w^{\prime}\right)=\left(w_{0}+w_{0}^{\prime}+\bar{w} w^{\prime}, w+w^{\prime}\right) \tag{1.1}
\end{equation*}
$$

We endow the metabelian simply connected real Lie group $\mathbb{H e i s}_{7}$ with its Cygan distance $d_{\mathrm{Cyg}}$, which is the unique left-invariant distance such that $d_{\mathrm{Cyg}}\left(\left(w_{0}, w\right),(0,0)\right)=$ $\left(4 \mathrm{n}\left(w_{0}\right)\right)^{\frac{1}{4}}$. The chains $C$ contained in $\mathbb{H}^{2} \mathrm{is}_{7}$ are ellipsoids, and have a natural center cen $(C)$ and radius (see Section 3).

Let $A$ be a definite $\left(A \otimes_{\mathbb{Q}} \mathbb{R}=\mathbb{H}\right)$ quaternion algebra over $\mathbb{Q}$, with discriminant $D_{A}$. Let $\mathscr{O}$ be a maximal order in $A$. We refer for instance to [25] for background on quaternion algebras and orders. The group $\mathrm{PU}_{q}(\mathscr{O})$ of elements of $\mathrm{PU}_{q}$ represented by matrices with coefficients in $\mathscr{O}$ is a (necessarily arithmetic) lattice in $\mathrm{PU}_{q}$. A chain $C_{0}$ is said to be arithmetic over $\mathscr{O}$ if the orbit of some point of $C_{0}$ under the stabiliser of $C_{0}$ in $\mathrm{PU}_{q}(\mathscr{O})$ is dense in $C_{0}$. The stabiliser $\mathrm{PU}_{q}(\mathscr{O})_{\infty}$ of $[1: 0: 0]$ in $\mathrm{PU}_{q}(\mathscr{O})$ preserves the diameters of the chains for $d_{\text {Cyg. }}$. The following result (see Theorem 4.2 for an explicit and more general version) is an asymptotic counting result of the arithmetic chains in an orbit under the arithmetic group $\mathrm{PU}_{q}(\mathscr{O})$ when their Cygan diameter tends to 0 .

Theorem 1.2. Let $C_{0}$ be an arithmetic chain in $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{2}$. There exists a constant $\kappa>0$ and an explicit constant $c>0$ such that, as $\epsilon \rightarrow 0$, the number of chains modulo $\mathrm{PU}_{q}(\mathscr{O})_{\infty}$ in the $\mathrm{PU}_{q}(\mathscr{O})$-orbit of $C_{0}$, with Cygan diameter at least $\epsilon$, is equal to $c \epsilon^{-10}\left(1+\mathrm{O}\left(\epsilon^{\kappa}\right)\right)$.

An arithmetic chain $C_{0}$ bounds in $\mathbf{H}_{\mathbb{H}}^{2}$ a homothetic copy of the real hyperbolic space of dimension 4 . We denote by $\operatorname{Covol}\left(C_{0}\right)$ the volume of the quotient of this real hyperbolic space, normalised to have sectional curvature -1 , by the stabiliser $\mathrm{PU}_{q}(\mathscr{O})_{C_{0}}$ of $C_{0}$ in $\mathrm{PU}_{q}(\mathscr{O})$, and by $m_{0}$ the order of the pointwise stabiliser of this real hyperbolic space in $\mathrm{PU}_{q}(\mathscr{O})$. We endow the real Lie group $\mathbb{H e i s}_{7}$ with its Haar measure Haar $_{\text {Heis }_{7}}$ normalised in such a way that the total mass of the induced measure on the quotient of $\mathbb{H e}{ }^{2}{ }_{7}$ by its (uniform) lattice $\mathbb{H e i s}_{7} \cap(\mathscr{O} \times \mathscr{O})$ is $\frac{D_{A}^{2}}{4}$ (see for instance [22, Lem. 8.4] for an explanation of this normalisation). Let $m_{A}=72$ if $D_{A}$ is even, and $m_{A}=1$ otherwise. Finally, we denote by $\Delta_{x}$ the unit Dirac mass at any point $x$. The following result proves that the centers of the arithmetic chains in an orbit under the arithmetic group $\mathrm{PU}_{q}(\mathscr{O})$ equidistribute in the quaternionic Heisenberg group.

Theorem 1.3. For the weak-star convergence of measures on $\mathbb{H e i s}_{7}$, we have

$$
\begin{aligned}
& \frac{m_{0} m_{A} \pi^{6} \prod_{p \mid D_{A}}(p-1)\left(p^{2}+1\right)\left(p^{3}-1\right)}{255152^{24} \operatorname{Covol}\left(C_{0}\right)} \epsilon^{10} \sum_{\substack{[g] \in \mathrm{PU}_{q}(\mathscr{O}) / \mathrm{PU}_{q}(\mathscr{O}) C_{0} \\
\epsilon \leq \operatorname{diam}_{d_{\mathrm{Cyg}}}\left(g C_{0}\right)<\infty}} \Delta_{\operatorname{cen}\left(g C_{0}\right)} \\
& \stackrel{*}{\sim} \text { Haar }_{\text {Heis }}{ }^{2} .
\end{aligned}
$$

We refer to Section 4 for a version with congruences and error terms, and a more developped study of explicit examples of arithmetic chains.

## Acknowledgements

The authors thank the snowy arctic conditions in Äkäslompolo in January 2020 which have provided an exceptional working environment. This research was supported by CNRS IEA BARP. The second author thanks the Laboratoire de Mathématiques Jean Leray at the Université de Nantes where this paper was completed.

## 2. Quaternionic hyperbolic spaces and Heisenberg groups

In this section, we briefly recall some background on the quaternionic hyperbolic spaces and quaternionic Heisenberg groups, as mostly contained in [22, $\S 3$ and $\S 6$ ], see also [16, 23] (with different choices of quaternionic Hermitian form and normalisation of the curvature).

Let $\mathbb{H}$ be Hamilton's quaternion algebra over $\mathbb{R}$, with $x \mapsto \bar{x}$ its conjugation, $\mathrm{n}: x \mapsto x \bar{x}$ its reduced norm, $\operatorname{tr}: x \mapsto x+\bar{x}$ its reduced trace and $\operatorname{Im}: x \mapsto \frac{1}{2}(x-\bar{x})$ its imaginary part map. We denote by $(1, i, j, k)$ the canonical basis of $\mathbb{H}$ as a real vector space, so that $\overline{x_{0}+x_{1} i+x_{2} j+x_{3} k}=x_{0}-x_{1} i-x_{2} j-x_{3} k$. Let

$$
\operatorname{Im} \mathbb{H}=\{x \in \mathbb{H}: \operatorname{tr} x=0\}=\mathbb{R} i+\mathbb{R} j+\mathbb{R} k
$$

be the $\mathbb{R}$-subspace of purely imaginary quaternions of $\mathbb{H}$. For all $w=\left(w_{1}, \ldots, w_{N}\right)$ and $w^{\prime}=\left(w_{1}^{\prime}, \ldots, w_{N}^{\prime}\right)$ in the right vector space $\mathbb{H}^{N}$ over $\mathbb{H}$, we denote by $\bar{w} \cdot w^{\prime}=\sum_{p=1}^{N} \overline{w_{p}} w_{p}^{\prime}$ their standard quaternionic Hermitian product, and we define $\mathrm{n}(w)=\bar{w} \cdot w=\sum_{p=1}^{N} \mathrm{n}\left(w_{p}\right)$. We endow $\mathbb{H}^{N}$ with the standard Euclidean structure $\left(w, w^{\prime}\right) \mapsto \frac{1}{2} \operatorname{tr}\left(\bar{w} \cdot w^{\prime}\right)$.

We fix $n \in \mathbb{N}-\{0,1\}$. On the right vector space $\mathbb{H} \times \mathbb{H}^{n-1} \times \mathbb{H}$ over $\mathbb{H}$ with coordinates $\left(z_{0}, z, z_{n}\right)$, let $q$ be the nondegenerate quaternionic Hermitian form

$$
\begin{equation*}
q\left(z_{0}, z, z_{n}\right)=-\operatorname{tr}\left(\overline{z_{0}} z_{n}\right)+\mathrm{n}(z) \tag{2.1}
\end{equation*}
$$

of Witt signature $(1, n)$, and let $\Phi: \mathbb{H}^{n+1} \times \mathbb{H}^{n+1} \rightarrow \mathbb{H}$, defined by

$$
\begin{equation*}
\Phi:\left(\left(z_{0}, z, z_{n}\right),\left(z_{0}^{\prime}, z^{\prime}, z_{n}^{\prime}\right)\right) \mapsto-\overline{z_{0}} z_{n}^{\prime}-\overline{z_{n}} z_{0}^{\prime}+\bar{z} \cdot z^{\prime} \tag{2.2}
\end{equation*}
$$

be the associated quaternionic sesquilinear form.
The Siegel domain model of the quaternionic hyperbolic $n$-space $\mathbf{H}_{\mathbb{H}}^{n}$ is

$$
\left\{\left(w_{0}, w\right) \in \mathbb{H} \times \mathbb{H}^{n-1}: \operatorname{tr} w_{0}-\mathrm{n}(w)>0\right\},
$$

endowed with the Riemannian metric

$$
\mathrm{d} s_{\mathbf{H}_{\mathbb{H}}^{n}}^{2}=\frac{1}{\left(\operatorname{tr} w_{0}-\mathrm{n}(w)\right)^{2}}\left(\mathrm{n}\left(\mathrm{~d} w_{0}-\overline{\mathrm{d} w} \cdot w\right)+\left(\operatorname{tr} w_{0}-\mathrm{n}(w)\right) \mathrm{n}(\mathrm{~d} w)\right)
$$

Its boundary at infinity is

$$
\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}=\left\{\left(w_{0}, w\right) \in \mathbb{H} \times \mathbb{H}^{n-1}: \operatorname{tr} w_{0}-\mathrm{n}(w)=0\right\} \cup\{\infty\} .
$$

A quaternionic geodesic line in $\mathbf{H}_{\mathbb{H}}^{n}$ is the image by an isometry of $\mathbf{H}_{\mathbb{H}}^{n}$ of the intersection of $\mathbf{H}_{\mathbb{H}}^{n}$ with the quaternionic line $\mathbb{H} \times\{0\}$. With our normalisation of the metric, a quaternionic geodesic line is a totally geodesic submanifold of real dimension 4 and constant sectional curvature -4 .

The closed horoballs in $\mathbf{H}_{\mathbb{H}}^{n}$ centred at $\infty \in \partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$ are the subsets

$$
\begin{equation*}
\mathscr{H}_{s}=\left\{\left(w_{0}, w\right) \in \mathbf{H}_{\mathbb{H}}^{n}: \operatorname{tr} w_{0}-\mathrm{n}(w) \geq s\right\}, \tag{2.3}
\end{equation*}
$$

and the horospheres centred at $\infty$ are their boundaries $\partial \mathscr{H}_{s}$, where $s$ ranges in $] 0,+\infty[$. Note that, for every $s \in] 0,1]$, we have

$$
\begin{equation*}
d\left(\partial \mathscr{H}_{1}, \partial \mathscr{H}_{s}\right)=-\frac{\ln s}{2} \tag{2.4}
\end{equation*}
$$

The Siegel domain $\mathbf{H}_{\mathbb{H}}^{n}$ embeds in the right quaternionic projective $n$-space $\mathbb{P}_{\mathrm{r}}^{n}(\mathbb{H})$ by the map (using homogeneous coordinates)

$$
\left(w_{0}, w\right) \mapsto\left[w_{0}: w: 1\right] .
$$

By this map, we identify $\mathbf{H}_{\mathbb{H}}^{n}$ with its image, which when endowed with the isometric Riemannian metric, is called the projective model of $\mathbf{H}_{\mathbb{H}}^{n}$. Note that this image is the negative cone of the quaternionic Hermitian form $q$ defined in Equation (2.1) : we have $\mathbf{H}_{\mathbb{H}}^{n}=\left\{\left[z_{0}: z: z_{n}\right] \in \mathbb{P}_{\mathrm{r}}^{n}(\mathbb{H}): q\left(z_{0}, z, z_{n}\right)<0\right\}$. This embedding extends continuously to the boundary at infinity, by mapping the point $\left(w_{0}, w\right) \in \partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}-\{\infty\}$ to $\left[w_{0}: w: 1\right]$ and $\infty$ to $[1: 0: 0]$, so that the image of $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$ is the isotropic cone of $q$ : we have $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}=\left\{\left[z_{0}: z: z_{n}\right] \in \mathbb{P}_{r}^{n}(\mathbb{H}): q\left(z_{0}, z, z_{n}\right)=0\right\}$. A projective point $\left[z_{0}: z: z_{n}\right] \in \mathbb{P}_{\mathrm{r}}^{n}(\mathbb{H})$ is positive if $q\left(z_{0}, z, z_{n}\right)>0$.

For every $N \in \mathbb{N}$, let $I_{N}$ be the identity $N \times N$ matrix. Let

$$
J=\left(\begin{array}{ccc}
0 & 0 & -1 \\
0 & I_{n-1} & 0 \\
-1 & 0 & 0
\end{array}\right)
$$

The conjugate-transpose matrix of a quaternionic matrix $X=\left(x_{p, p^{\prime}}\right)_{1 \leq p \leq r, 1 \leq p^{\prime} \leq s}$ in $\mathscr{M}_{r, s}(\mathbb{H})$ is $X^{*}=\left(x_{p, p^{\prime}}^{*}=\overline{x_{p^{\prime}, p}}\right)_{1 \leq p \leq s, 1 \leq p^{\prime} \leq r} \in \mathscr{M}_{s, r}(\mathbb{H})$. Let

$$
\mathrm{U}_{q}=\left\{g \in \mathrm{GL}_{n+1}(\mathbb{H}): q \circ g=q\right\}=\left\{g \in \mathrm{GL}_{n+1}(\mathbb{H}): g^{*} J g=J\right\}
$$

be the unitary group of $q$. Its left linear action on $\mathbb{H}^{n+1}$ induces a projective action on $\mathbb{P}_{\mathrm{r}}^{n}(\mathbb{H})$ with kernel its center, which is reduced to $\left\{ \pm I_{n+1}\right\}$. The projective unitary group

$$
\mathrm{PU}_{q}=\mathrm{U}_{q} /\left\{ \pm I_{n+1}\right\}
$$

of $q$ acts faithfully on $\mathbb{P}_{\mathrm{r}}^{n}(\mathbb{H})$, preserving $\mathbf{H}_{\mathbb{H}}^{n}$, and its restriction to $\mathbf{H}_{\mathbb{H}}^{n}$ is the full isometry group of $\mathbf{H}_{\mathbb{H}}^{n}$.

A matrix

$$
X=\left(\begin{array}{lll}
a & \gamma^{*} & b \\
\alpha & M & \beta \\
c & \delta^{*} & d
\end{array}\right) \in \operatorname{GL}_{n+1}(\mathbb{H})
$$

with $a, b, c, d \in \mathbb{H}, \alpha, \beta, \gamma, \delta \in \mathbb{H}^{n-1}$ (identified with their column matrices in $\mathscr{M}_{n-1,1}(\mathbb{H})$ ) and $M \in \mathscr{M}_{n-1, n-1}(\mathbb{H})$, belongs to $\mathrm{U}_{q}$ if and only if

$$
\left\{\begin{align*}
\bar{c} a-\alpha^{*} \alpha+\bar{a} c & =0  \tag{2.5}\\
\bar{d} b-\beta^{*} \beta+\bar{b} d & =0 \\
-\delta \gamma^{*}+M^{*} M-\gamma \delta^{*} & =I_{n-1} \\
\bar{d} a-\beta^{*} \alpha+\bar{b} c & =1 \\
\delta a-M^{*} \alpha+\gamma c & =0 \\
\delta b-M^{*} \beta+\gamma d & =0 .
\end{align*}\right.
$$

With $\operatorname{Sp}(n-1)=\left\{g \in \mathrm{GL}_{n+1}(\mathbb{H}): g^{*} g=I_{n-1}\right\}$, an easy computation shows that the block upper triangular subgroup of $\mathrm{U}_{q}$ is

$$
\mathrm{B}_{q}=\left\{\left(\begin{array}{ccc}
\mu r & \zeta^{*} & \frac{1}{2 r}(\mathrm{n}(\zeta)+u) \mu \\
0 & U & \frac{1}{r} U \zeta \mu \\
0 & 0 & \frac{\mu}{r}
\end{array}\right): \begin{array}{l}
\zeta \in \mathbb{H}^{n-1}, u \in \operatorname{Im} \mathbb{H}, \\
U \in \operatorname{Sp}(n-1), \mu \in \operatorname{Sp}(1), r>0
\end{array}\right\}
$$

Its image $\mathrm{PB}_{q}=\mathrm{B}_{q} /\left\{ \pm I_{n+1}\right\}$ in $\mathrm{PU}_{q}$ is equal to the stabiliser of $\infty$ in $\mathrm{PU}_{q}$.
The quaternionic Heisenberg group $\mathbb{H e i s}_{4 n-1}$ of dimension $4 n-1$ is the real Lie group structure on $\mathbb{H}^{n-1} \times \operatorname{Im} \mathbb{H}$ with law

$$
(\zeta, u)\left(\zeta^{\prime}, u^{\prime}\right)=\left(\zeta+\zeta^{\prime}, u+u^{\prime}+2 \operatorname{Im} \bar{\zeta} \cdot \zeta^{\prime}\right)
$$

and inverses $(\zeta, u)^{-1}=(-\zeta,-u)$. It identifies with the punctured boundary at infinity $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}-\{\infty\}$ by the map $(\zeta, u) \mapsto\left(w_{0}, w\right)$ where

$$
\begin{equation*}
\left(w_{0}, w\right)=\left(\frac{\mathrm{n}(\zeta)+u}{2}, \zeta\right) \text { hence }(\zeta, u)=\left(w, 2 \operatorname{Im} w_{0}\right) \tag{2.6}
\end{equation*}
$$

and with a subgroup of $\mathrm{PB}_{q} \subset \mathrm{PU}_{q}$, preserving every horoball $\mathscr{H}_{s}$ for $s>0$, by the map

$$
(\zeta, u) \mapsto \pm\left(\begin{array}{ccc}
1 & \zeta^{*} & \frac{\mathrm{n}(\zeta)+u}{2} \\
0 & I_{n-1} & \zeta \\
0 & 0 & 1
\end{array}\right)
$$

Equation (2.6) allows to recover the definition of $\mathbb{H e i s}_{7}$ given in the Introduction, for which the inverses are $\left(w_{0}, w\right)^{-1}=\left(-w_{0}+\mathrm{n}(w),-w\right)$.

For every $(\zeta, u) \in \mathbb{H e i s}_{4 n-1}$, the map $\left(\zeta^{\prime}, u^{\prime}\right) \mapsto(\zeta, u)\left(\zeta^{\prime}, u^{\prime}\right)$ is the Heisenberg translation by $(\zeta, u)$. For every $\zeta \in \mathbb{H}^{n-1}$, the Heisenberg translation by $(\zeta, 0)$ is called a horizontal (Heisenberg) translation. For every $u \in \operatorname{Im} \mathbb{H}$, the Heisenberg translation by $(0, u)$ is called a vertical (Heisenberg) translation. The canonical map $\Pi_{v}: \mathbb{H e i s}_{4 n-1} \rightarrow \mathbb{H}^{n-1}$ defined by $(\zeta, u) \mapsto \zeta$ is a real Lie group morphism, called the vertical projection, whose kernel is the center of $\mathbb{H e i s}_{4 n-1}$. For every $U \in \operatorname{Sp}(n-1)$,
the map $(\zeta, u) \mapsto(U \zeta, u)$ is the Heisenberg rotation by $U$. For every $\lambda>0$, the map $h_{\lambda}:(\zeta, u) \mapsto\left(\lambda \zeta, \lambda^{2} u\right)$ is the Heisenberg dilation by $\lambda$.

The Cygan distance $d_{\mathrm{Cyg}}$ on $\mathbb{H e i s}_{4 n-1}$ is the unique left-invariant distance on the real Lie group $\mathbb{H e i s}_{4 n-1}$ such that

$$
\begin{equation*}
d_{\mathrm{Cyg}}((\zeta, u),(0,0))=\left(\mathrm{n}(\zeta)^{2}+\mathrm{n}(u)\right)^{1 / 4} \tag{2.7}
\end{equation*}
$$

or equivalently $d_{\text {Cyg }}\left(\left(w_{0}, w\right),(0,0)\right)=\left(4 \mathrm{n}\left(w_{0}\right)\right)^{\frac{1}{4}}$ by Equation (2.6). We introduce (see $[19,20]$ in the complex case) the modified Cygan distance $d_{\text {Cyg }}^{\prime \prime}$, as the unique left-invariant map from $\mathbb{H e i s}_{4 n-1} \times \mathbb{H e i s}_{4 n-1}$ to $[0,+\infty[$ such that

$$
\begin{equation*}
d_{\mathrm{Cyg}}^{\prime \prime}((\zeta, u),(0,0))=\frac{\left(\mathrm{n}(\zeta)^{2}+\mathrm{n}(u)\right)^{1 / 2}}{\left(\left(\mathrm{n}(\zeta)^{2}+\mathrm{n}(u)\right)^{1 / 2}+\mathrm{n}(\zeta)\right)^{1 / 2}} \tag{2.8}
\end{equation*}
$$

or equivalently by Equation (2.6)

$$
d_{\mathrm{Cyg}}^{\prime \prime}\left(\left(w_{0}, w\right),(0,0)\right)=\frac{2 \mathrm{n}\left(w_{0}\right)^{1 / 2}}{\left(2 \mathrm{n}\left(w_{0}\right)^{1 / 2}+\mathrm{n}(w)\right)^{1 / 2}}
$$

Though not actually a distance, the map $d_{\text {Cyg }}^{\prime \prime}$ is symmetric and satisfies

$$
\frac{1}{\sqrt{2}} d_{\mathrm{Cyg}} \leq d_{\mathrm{Cyg}}^{\prime \prime} \leq d_{\mathrm{Cyg}} .
$$

For every nonempty bounded subset $E$ of $\mathbb{H e i s}_{4 n-1}$, we define the diameter of $E$ for this almost distance as

$$
\operatorname{diam}_{d_{\mathrm{Cyg}}^{\prime \prime}}(E)=\sup _{x, y \in E} d_{\mathrm{Cyg}}^{\prime \prime}(x, y)
$$

Note that the Cygan distance and the modified Cygan distance are invariant under Heisenberg translations and rotations, and that for every $\lambda>0$, the Heisenberg dilation $h_{\lambda}$ is a homothety of ratio $\lambda$ for both distances.

Lemma 2.1. For every geodesic line $] x, y\left[\right.$ in $\mathbf{H}_{\mathbb{H}}^{n}$ disjoint from the horoball $\mathscr{H}_{1}$, the distance in $\mathbf{H}_{\mathbb{H}}^{n}$ between $\mathscr{H}_{1}$ and $] x, y[$ is equal to

$$
d\left(\mathscr{H}_{1},\right] x, y[)=-\ln \left(\frac{1}{\sqrt{2}} d_{\mathrm{Cyg}}^{\prime \prime}(x, y)\right) .
$$

Proof. By the invariance under Heisenberg translations of $\mathscr{H}_{1}$, of the distance in $\mathbf{H}_{\mathbb{H}}^{n}$ and of the modified Cygan distance, we may assume that $x=\left(w_{0}, w\right) \in \partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}-\{\infty,(0,0)\}$ and $y=(0,0) \in \partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}-\{\infty\}$. By [22, Lem. 6.4], the geodesic line from $\left(w_{0}, w\right)$ to $(0,0)$ is, up to translation at the source, the map

$$
\gamma_{w_{0}, w}: t \mapsto\left(w_{0}\left(1+e^{2 t} w_{0}\right)^{-1}, w\left(1+e^{2 t} w_{0}\right)^{-1}\right)
$$

The point $\gamma_{w_{0}, w}(t)$ belongs to the horosphere $\mathscr{H}_{s(t)}$, where, since $\operatorname{tr} w_{0}=\mathrm{n}(w)$,

$$
s(t)=\operatorname{tr}\left(w_{0}\left(1+e^{2 t} w_{0}\right)^{-1}\right)-\mathrm{n}\left(w\left(1+e^{2 t} w_{0}\right)^{-1}\right)=\frac{2 e^{2 t} \mathrm{n}\left(w_{0}\right)}{\mathrm{n}\left(1+e^{2 t} w_{0}\right)} .
$$

Let $r=\mathrm{n}\left(w_{0}\right)^{1 / 2}$ be the norm of the vector $w_{0}$ and $\theta$ the angle between the vectors 1 and $w_{0}$ in the Euclidean space $\mathbb{H}$. Then the map

$$
t \mapsto s(t)=\frac{2 e^{2 t} r^{2}}{e^{4 t} r^{2}+2 r e^{2 t} \cos \theta+1}
$$

reaches its maximum at $e^{2 t}=\frac{1}{r}$. Since $\operatorname{tr} w_{0}=\mathrm{n}(w)$, the value of this maximum is

$$
s_{\max }=\frac{2 \mathrm{n}\left(w_{0}\right)^{1 / 2}}{2+\operatorname{tr}\left(w_{0} \mathrm{n}\left(w_{0}\right)^{-1 / 2}\right)}=\frac{2 \mathrm{n}\left(w_{0}\right)}{2 \mathrm{n}\left(w_{0}\right)^{1 / 2}+\mathrm{n}(w)}=\frac{1}{2} d_{\mathrm{Cyg}}^{\prime \prime}\left(\left(w_{0}, w\right),(0,0)\right)^{2}
$$

The result then follows from Equation (2.4).

## 3. Chains

In this section, we define the quaternionic Cartan chains and give their elementary geometric properties, see also [24]. In the complex case, the notion of chain is attributed to von Staudt by [7]. The exposition follows the one of [10] in the complex case. We fix $m \in\{1, \ldots, n-1\}$.

### 3.1. A vocabulary of chains

An $m$-chain $C$ in $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$ is the intersection with $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$ of a quaternionic projective space $L_{C}$ of dimension $m$ meeting $\mathbf{H}_{\mathbb{H}}^{n}$. Note that $C$ determines $L_{C}$ and conversely. A chain is a 1 -chain, and a hyperchain is an $(n-1)$-chain. An $m$-chain is vertical if it contains $\infty=[1: 0: 0]$, and finite otherwise.

If $P=\left[z_{0}: z: z_{n}\right] \in \mathbb{P}_{\mathrm{r}}^{n}(\mathbb{H})$, let

$$
P^{\perp}=\left\{\left[z_{0}^{\prime}: z^{\prime}: z_{n}^{\prime}\right] \in \mathbb{P}_{\mathrm{r}}^{n}(\mathbb{H}): \Phi\left(\left(z_{0}, z, z_{n}\right),\left(z_{0}^{\prime}, z^{\prime}, z_{n}^{\prime}\right)\right)=0\right\}
$$

be the orthogonal quaternionic projective subspace of $P$. The map $P \mapsto P^{\perp}$, from the set of positive projective points to the set of quaternionic projective hyperplanes in $\mathbb{P}_{\mathrm{r}}^{n}(\mathbb{H})$ meeting $\mathbf{H}_{\mathbb{H}}^{n}$, is a $\mathrm{PU}_{q}$-equivariant bijection. Therefore, the map

$$
P \mapsto C_{P}=P^{\perp} \cap \partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}
$$

is a $\mathrm{PU}_{q}$-equivariant bijection from the set of positive projective points to the set of hyperchains. The point $P$ is called the polar point of the hyperchain $C_{P}$, or of the
quaternionic projective hyperplane $P^{\perp}$. If $P=\left[z_{0}: z: z_{n}\right]$, we have

$$
\begin{equation*}
C_{P} \cap\left(\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}-\{\infty\}\right)=\left\{\left[w_{0}: w: 1\right]:-\left(\frac{\mathbf{n}(w)}{2}-\operatorname{Im} w_{0}\right) z_{n}+\bar{w} \cdot z-z_{0}=0\right\} . \tag{3.1}
\end{equation*}
$$

This hyperchain $C_{P}$ is hence vertical if and only if $z_{n}=0$, in which case $C_{P} \cap\left(\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}-\{\infty\}\right)$ is the preimage by the vertical projection $\Pi_{v}: \mathbb{H e i s}_{4 n-1} \rightarrow \mathbb{H}^{n-1}$ of the quaternionic affine hyperplane of $\mathbb{H}^{n-1}$ with equation $\bar{z} \cdot w=\overline{z_{0}}$ in the unknown $w$. Similarly a vertical chain is the preimage of a point of $\mathbb{H}^{n-1}$ by the vertical projection $\Pi_{v}$.

When $C=C_{P}$ is a finite hyperchain, that is, when $z_{n} \neq 0$, then $C$ is a codimension 4 ellipsoid in the Euclidean space $\mathbb{H}^{n-1} \times \operatorname{Im} \mathbb{H}$, whose vertical projection is the Euclidean sphere in $\mathbb{H}^{n-1}$ with real codimension 1 and equation $\mathrm{n}(w)-\operatorname{tr}\left(\bar{w} \cdot\left(z z_{n}^{-1}\right)\right)+\operatorname{tr}\left(z_{0} z_{n}^{-1}\right)=0$ in the unknown $w$, with center $z z_{n}^{-1}$ and radius

$$
R_{C}=\frac{q\left(z_{0}, z, z_{n}\right)^{1 / 2}}{\mathrm{n}\left(z_{n}\right)^{1 / 2}}
$$

This radius $R_{C}$ of the Euclidean sphere $\Pi_{v}(C)$ is called the radius of the finite hyperchain $C$. The map $\left.\Pi_{v}\right|_{C}$ from $C$ to $\Pi_{v}(C)$ is a homeomorphism. When $z=0$ and $z_{0} z_{n}^{-1} \in \mathbb{R}$, the hyperchain $C=C_{P}$ is contained in the horizontal subspace $\left\{(\zeta, u) \in \mathbb{H}^{n-1} \times \operatorname{Im} \mathbb{H}: u=0\right\}$ of $\mathbb{H e i s}_{4 n-1}$, by Equation (2.6).

Similarly, a finite chain is a 3-dimensional ellipsoid in the Euclidean space $\mathbb{H}^{n-1} \times \operatorname{Im} \mathbb{H}$, whose vertical projection is a Euclidean 3 -sphere in $\mathbb{H}^{n-1}$. In particular, any chain is homeomorphic to the 3 -sphere $\mathbb{S}^{3}$.

### 3.2. Transitivity properties of $\mathrm{PU}_{q}$ on chains

Through any two distinct projective points belonging to $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$ passes one and only one quaternionic projective line, and this projective line meets $\mathbf{H}_{\mathbb{H}}^{n}$. Hence through two distinct points of $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$ passes one and only one chain. By Witt's theorem, the group $\mathrm{PU}_{q}$ acts transitively on the set of quaternionic projective spaces $L$ of dimension $m$ meeting $\mathbf{H}_{\mathbb{H}}^{n}$, hence it acts transitively on the set of $m$-chains.

Note that two $m$-chains having the same vertical projection differ by a vertical Heisenberg translation, that the group generated by Heisenberg translations and rotations acts transitively on the set of vertical $m$-chains, and that $\mathrm{PB}_{q}$ (that contains the Heisenberg dilations, rotations and translations) acts transitively on the set of finite $m$-chains.

The next result gives the topological structure of a family of chains, called a fan in the complex hyperbolic case (see for instance [10, p. 131]).

Proposition 3.1. The union $F$ of all chains containing a given point $P \in \partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$ and passing through an m-chain $C$ of $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$ not containing $P$ is homeomorphic to the
topological quotient space $\left(\mathbb{S}^{3} \times \mathbb{S}^{4 m-1}\right) / \sim$ where $\sim$ is the equivalence relation generated by $\left(x_{0}, x\right) \sim\left(x_{0}, y\right)$ for all $x, y \in \mathbb{S}^{4 m-1}$, where $x_{0}$ is any fixed point in $\mathbb{S}^{3}$.

Proof. By the transitivity properties of $\mathrm{PU}_{q}$, we may assume that $P=\infty$. Hence $C$ is a finite chain, and by the transitivity properties of the Heisenberg translations, we may assume that $C$ is a Euclidean sphere of dimension $4 m-1$ contained in the horizontal space $\left\{(\zeta, u) \in \mathbb{H}^{n-1} \times \operatorname{Im} \mathbb{H}: u=0\right\}$. Thus $F=\bigcup_{(\zeta, u) \in C} \Pi_{v}^{-1}(\zeta, u)$ is clearly homeomorphic to the above quotient of $\mathbb{S}^{3} \times \mathbb{S}^{4 m-1}$.

### 3.3. Reflexions on chains

The chains are fixed point sets at infinity of natural isometries of $\mathbf{H}_{\mathbb{H}}^{n}$, that we now describe.
If $L$ is a proper quaternionic projective subspace of $\mathbb{P}_{\mathrm{r}}^{n}(\mathbb{H})$ meeting $\mathbf{H}_{\mathbb{H}}^{n}$, there exists a unique involution $\iota_{L}$ in $\mathrm{PU}_{q}$ with fixed point set $L$, called the reflexion on $L$. Note that the set of fixed points of $\iota_{L}$ in $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$ is the $m$-chain $L \cap \partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$, where $m$ is the quaternionic dimension of $L$, assuming that $m \neq 0$.

For instance, $C=\left\{\left[z_{0}: z_{1}: \cdots: z_{n}\right] \in \partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}: z_{m}=0, \ldots, z_{n-1}=0\right\} \cup\{\infty\}$ is a vertical $m$-chain, called the standard vertical $m$-chain and the reflexion $\iota_{L_{C}}$ is the map

$$
\left[z_{0}: z_{1}: \cdots: z_{n}\right] \mapsto\left[z_{0}: z_{1}: \cdots: z_{m-1}:-z_{m}: \cdots:-z_{n-1}: z_{n}\right]
$$

The vertical $m$-chains are the images of the standard vertical $m$-chain by the Heisenberg translations and Heisenberg rotations: they are the subspaces

$$
(E \times \operatorname{Im} \mathbb{H}) \cup\{\infty\}
$$

where $E$ is a quaternionic affine subspace of $\mathbb{H}^{n-1}$ with dimension $m-1$ (hence a point when $m=1$ ).

Lemma 3.2. Let L and $L^{\prime}$ be quaternionic projective subspaces of $\mathbb{P}_{r}^{n}(\mathbb{H})$ meeting $\mathbf{H}_{\mathbb{H}}^{n}$ such that one is not contained in the other, whose sum of dimensions is $n$. The following assertions are equivalent.
(1) The reflexions $\iota_{L}$ and $\iota_{L^{\prime}}$ commute.
(2) The reflexion $\iota_{L}$ preserves $L^{\prime}$.
(3) The reflexion $\iota_{L^{\prime}}$ preserves $L$.
(4) We have $\left(\iota_{L} \circ \iota_{L^{\prime}}\right)^{2}=\mathrm{id}$.
(5) The totally geodesic subspaces $L \cap \mathbf{H}_{\mathbb{H}}^{n}$ and $L^{\prime} \cap \mathbf{H}_{\mathbb{H}}^{n}$ intersect perpendicularly in the Riemannian manifold $\mathbf{H}_{\mathbb{H}}^{n}$.
(6) The subspace $L \cap \mathbf{H}_{\mathbb{H}}^{n}$ is a fiber of the orthogonal projection on $L^{\prime} \cap \mathbf{H}_{\mathbb{H}}^{n}$ in $\mathbf{H}_{\mathbb{H}}^{n}$.
(7) The subspace $L^{\prime} \cap \mathbf{H}_{\mathbb{H}}^{n}$ is a fiber of the orthogonal projection on $L \cap \mathbf{H}_{\mathbb{H}}^{n}$ in $\mathbf{H}_{\mathbb{H}}^{n}$.

Proof. The proof is similar to the one of [10, Lem. 4.3.1] in the complex hyperbolic case. Note that $L \cap \mathbf{H}_{\mathbb{H}}^{n}$, being the set of fixed points of the isometry $\iota_{L}$ of the negatively curved Riemannian manifold $\mathbf{H}_{\mathbb{H}}^{n}$, is indeed totally geodesic.

Two involutions commute if and only if their composition is an involution or the identity, hence Assertions (1) and (4) are equivalent. Since the centralizer of a projective transformation preserves its fixed point set, Assertion (1) implies Assertions (2) and (3). If Assertion (2) is satisfied, then $\iota_{L} \circ \iota_{L^{\prime}} \circ \iota_{L}^{-1}=\iota_{\iota_{L}\left(L^{\prime}\right)}=\iota_{L^{\prime}}$, so that Assertion (1) is satisfied. Similarly, Assertion (3) implies Assertion (1). Finally, the totally geodesic subspaces $L^{\prime} \cap \mathbf{H}_{\mathbb{H}}^{n}$ and $L \cap \mathbf{H}_{\mathbb{H}}^{n}$ in $\mathbf{H}_{\mathbb{H}}^{n}$

- either have disjoint closures in $\mathbf{H}_{\mathbb{H}}^{n} \cup \partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$, or
- are disjoint and have closures meeting in $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$, or
- meet in $\mathbf{H}_{\mathbb{H}}^{n}$.

In the first two cases, the composition $\iota_{L} \circ \iota_{L^{\prime}}$ has infinite order, and in the last case, $\iota_{L} \circ \iota_{L^{\prime}}$ can be an involution if and only if $L^{\prime} \cap \mathbf{H}_{\mathbb{H}}^{n}$ and $L \cap \mathbf{H}_{\mathbb{H}}^{n}$ are perpendicular.

An $m$-chain $C$ and an $(n-m)$-chain $C^{\prime}$ are orthogonal if neither of the corresponding quaternionic projective subspaces $L_{C}$ and $L_{C^{\prime}}$ contains the other and if they satisfy one of the equivalent assertions of Lemma 3.2. For instance, the hyperchains orthogonal to the standard vertical chain $(\{0\} \times \operatorname{Im} \mathbb{H}) \cup\{\infty\}$ are exactly the Euclidean spheres centered at $\left(0, u_{0}\right)$ in the horizontal subspace $\left\{(\zeta, u) \in \mathbb{H}^{n-1} \times \operatorname{Im} \mathbb{H}: u=u_{0}\right\}$ of $\mathbb{H e i s}_{4 n-1}$, for some $u_{0}$ in $\operatorname{Im} \mathbb{H}$.

### 3.4. Description of the center and radius of chains

We now define and study the centers of chains, whose equidistribution we will prove in Section 4.

The center of an $m$-chain $C$ is $\operatorname{cen}(C)=\iota_{L_{C}}(\infty)$. In particular, $\operatorname{cen}(C)=\infty$ if and only if $C$ is vertical. For every element $\gamma \in \mathrm{PB}_{q}$ (which fixes $\infty$ ), the reflexion on the $m$-chain $\gamma C$ is $\gamma \iota_{L_{C}} \gamma^{-1}$, so that the center of $\gamma C$ is

$$
\begin{equation*}
\operatorname{cen}(\gamma C)=\gamma \operatorname{cen}(C) \tag{3.2}
\end{equation*}
$$

When $P_{0}=\left[-\frac{1}{2}: 0: 1\right]$, the hyperchain $C_{P_{0}}$ with polar point $P_{0}$ is, by Equation (3.1), the sphere centered at $(0,0)$ with radius 1 in the horizontal codimension 3 Euclidean
subspace $\left\{(\zeta, u) \in \mathbb{H}^{n-1} \times \operatorname{Im} \mathbb{H}: u=0\right\}$ in $_{\mathbb{H e i s}_{4 n-1}}$. The reflexion on $L=L_{C_{P_{0}}}$ is the involutive map $\iota_{L}:\left(w_{0}, w\right) \mapsto\left(\frac{1}{4} w_{0}^{-1}, \frac{1}{2} w w_{0}^{-1}\right)$, induced by

$$
\pm\left(\begin{array}{ccc}
0 & 0 & 1 / 2 \\
0 & I_{n-1} & 0 \\
2 & 0 & 0
\end{array}\right) \in \mathrm{PU}_{q}
$$

Thus, $\operatorname{cen}\left(C_{P_{0}}\right)=\iota_{L}(\infty)=(0,0)$.
Let $P=\left[z_{0}: z: z_{n}\right]$ be a positive projective point with $z_{n} \neq 0$. An easy computation shows that the Heisenberg translation $\gamma$ by

$$
\left[\frac{\mathrm{n}(z)}{2 \mathrm{n}\left(z_{n}\right)}-\operatorname{Im}\left(z_{0} z_{n}^{-1}\right):-z z_{n}^{-1}: 1\right]
$$

maps $P$ to $\left[-\frac{R^{2}}{2}: 0: 1\right]$ where $R=R_{C_{P}}=\frac{q\left(z_{0}, z, z_{n}\right)^{1 / 2}}{\mathrm{n}\left(z_{n}\right)^{1 / 2}}$ is the radius of the finite hyperchain $C_{P}$, and the Heisenberg dilation

$$
h_{R}:\left(w_{0}, w\right) \mapsto\left(R^{2} w_{0}, R w\right)
$$

maps $P_{0}$ to $\left[-\frac{R^{2}}{2}: 0: 1\right]$. Hence the center of the finite hyperchain $C_{P}$ with polar point $P$ is, by Equation (3.2), equal to

$$
\operatorname{cen}\left(C_{P}\right)=\gamma^{-1} h_{R} \operatorname{cen}\left(C_{P_{0}}\right)=\gamma^{-1}(0,0)=\left[\frac{2 \operatorname{Im}\left(z_{0} \overline{z_{n}}\right)+\mathrm{n}(z)}{2 \mathrm{n}\left(z_{n}\right)}: z z_{n}^{-1}: 1\right],
$$

or cen $\left(C_{P}\right)=\left(z z_{n}^{-1}, 2 \operatorname{Im}\left(z_{0} z_{n}^{-1}\right)\right)$ in the $(\zeta, u)$-coordinates of $\mathbb{H e i s}_{4 n-1}$ by Equation (2.6). Thus, by Equation (3.1), if $C$ is a finite hyperchain in $\mathbb{H e i s}_{4 n-1}$ with center $\left(\zeta_{0}, u_{0}\right)$ and radius $r_{0}$, then

$$
C=\left\{(\zeta, u) \in \mathbb{H}^{n-1} \times \operatorname{Im} \mathbb{H}: \mathrm{n}\left(\zeta-\zeta_{0}\right)=r_{0}^{2} \text { and } u=u_{0}+2 \operatorname{Im}\left(\overline{\zeta_{0}} \zeta\right)\right\}
$$

In particular, a finite hyperchain is uniquely determined by its center and its radius, and the hyperchains contained in the horizontal Euclidean space $\left\{(\zeta, u) \in \mathbb{H}^{n-1} \times \operatorname{Im} \mathbb{H}: u=0\right\}$ are exactly the Euclidean spheres centered at $(0,0)$.

### 3.5. A von Staudt-Cartan rigidity theorem

The following theorem shows that the chain-preserving transformations of the boundary of the quaternionic hyperbolic spaces are projective transformations. This is a quaternionic version of the result of Cartan in the complex case (see for instance [10, Thm. 4.3.12]), close to von Staudt's fundamental theorem of real projective geometry.

Theorem 3.3. A bijection from $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$ to itself, mapping chains to chains, is (the restriction to $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$ of) an element of $\mathrm{PU}_{q}$.

Proof. Up to composing by an element of $\mathrm{PU}_{q}$, we may assume that $f$ fixes $\infty=[1: 0: 0]$. Hence $f$ preserves the set of vertical chains, which are the ones containing $\infty$. The set of vertical chains identifies with the horizontal space $\mathbb{H}^{n-1}$ of the quaternionic Heisenberg group by the vertical projection $\Pi_{v}$, which sends a vertical chain $C$ to the unique point of $\mathbb{H}^{n-1}$ whose preimage by $\Pi_{v}$ is $C$. Hence $f$ induces a bijection $\bar{f}$ from $\mathbb{H}^{n-1}$ to itself, which sends the vertical projections of the finite chains to the vertical projections of the finite chains.

The vertical projections of the finite chains are exactly all the Euclidean 3-spheres in $\mathbb{H}^{n-1}$. Given two distinct points $x, y$ in $\mathbb{H}^{n-1}$, the complement of the union of all the Euclidean 3 -spheres containing $x$ and $y$ is the real affine line containing $x$ and $y$, with $x$ and $y$ removed. Hence $\bar{f}$ is a bijection of $\mathbb{H}^{n-1}$ sending real affine lines to real affine lines. By the fundamental theorem of real affine geometry, this map is an affine transformation of $\mathbb{H}^{n-1}$. Since the affine transformations of $\mathbb{H}^{n-1}$ are vertical projections of elements of the stabiliser $\mathrm{PB}_{q}$ of $\infty$ in $\mathrm{PU}_{q}$, up to composing $f$ by an element of $\mathrm{PB}_{q}$, we may assume that $\bar{f}$ is the identity map of $\mathbb{H}^{n-1}$, and also that $f(0)=0$.

Let $x \in \partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}-\{\infty\}$, and let us prove that $f(x)=x$. First assume that $\Pi_{v}(x) \neq 0$. Then the unique chain $C_{x}$ passing through 0 and $x$ is a finite chain, and the vertical projections of $C_{x}$ and $f\left(C_{x}\right)$ coincide, since $\bar{f}=\mathrm{id}$. By the uniqueness of a chain with given vertical projection up to a vertical translation, since $f(0)=0$, we have $f\left(C_{x}\right)=C_{x}$. But if $f(x) \neq x$, then since $f(x)$ and $x$ have the same vertical projections, the chains $C_{x}$ and $f\left(C_{x}\right)$ through 0 would be different. Hence $f(x)=x$. This is in particular true for any given $x=x_{0} \neq 0$ in the horizontal space $\mathbb{H}^{n-1} \times\{0\}$. Replacing 0 by such an $x_{0}$ in the above argument allows to prove that $f(x)=x$ when $\Pi_{v}(x)=0$.

A similar proof shows that an injective map $f$ from $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$ to itself, such that any three points belong to a same chain if and only if their images by $f$ belong to a same chain, is the restriction of an element of $\mathrm{PU}_{q}$.

### 3.6. Relation with the hyper CR structure

In this subsection, we give a characterisation of the chains in terms of the natural hyper CR structure on $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$. We refer for example to [3] and [14] for background on hyperkähler manifolds and hyper CR manifolds, respectively.

We endow the manifold $\mathbb{P}_{\mathrm{r}}^{n}(\mathbb{H})$ with its natural hyperkähler structure, and we denote by $(\mathbb{I}, \mathbb{J}, \mathbb{K})$ the corresponding triple of almost complex structures. The boundary at infinity $W=\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$ is a smooth real hypersurface in the real manifold $\mathbb{P}_{\mathrm{r}}^{n}(\mathbb{H})$ of real dimension 4n, and $E=T W \cap \mathbb{I} T W \cap \mathbb{J} W \cap \mathbb{K} T W$ is a real codimension 3 subbundle of the real tangent bundle $\left.T \mathbb{P}_{\mathrm{r}}^{n}(\mathbb{H})\right|_{W}$, invariant under $\mathrm{PU}_{q}$, defining a hyper CR structure on $W$. When $x$ is the point $(0,0)$ in the $(\zeta, u)$-coordinates of $\mathbb{H e i s}_{4 n-1}=\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}-\{\infty\}$, then, identifying
$\mathbb{H}^{n-1} \times \operatorname{Im} \mathbb{H}$ with its real tangent space at $x$, the fiber $E_{x}$ of $E$ over $x$ is the horizontal subspace $\left\{(\zeta, u) \in \mathbb{H}^{n-1} \times \operatorname{Im} \mathbb{H}: u=0\right\}$.

A calibration of $E$ is a 1-form $\omega$ on $W$ with values in $\operatorname{Im} \mathbb{H}$ such that $E=\operatorname{ker} \omega$. Its Levi form is $\mathrm{d} \omega$. For instance, in the $(\zeta, u)$-coordinates of $\mathbb{H e i s}_{4 n-1}$, the form

$$
\omega=\mathrm{d} u-2 \operatorname{Im}(\bar{\zeta} \cdot \mathrm{~d} \zeta)
$$

is a calibration of $E$ (when restricted to $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}-\{\infty\}$ ). An easy computation shows that this calibration is invariant under Heisenberg translations and rotations: For every such transformation $\gamma$, we have $\gamma^{*} \omega=\omega$. The fact that $\omega$ is indeed a calibration follows by invariance since $\operatorname{ker} \mathrm{d} u=\left\{(\zeta, u) \in \mathbb{H}^{n-1} \times \operatorname{Im} \mathbb{H}: u=0\right\}$. This calibration $\omega$ is scaled by the Heisenberg dilations as follows : for every $\lambda>0$, we have $\left(h_{\lambda}\right)^{*} \omega=\lambda^{2} \omega$.

In the following result, we denote by $v=v_{1} i+v_{2} j+v_{3} k$ the standard coordinate in $\operatorname{Im} \mathbb{H}$, and by $\mathrm{d} v$ the tautological $(\operatorname{Im} \mathbb{H})$-valued 1 -form on $\operatorname{Im} \mathbb{H}$, so that for every $x \in \operatorname{Im} \mathbb{H}$, the map $\mathrm{d} v_{x}: T_{x} \operatorname{Im} \mathbb{H}=\operatorname{Im} \mathbb{H} \rightarrow \operatorname{Im} \mathbb{H}$ is the identity map. We denote by $\omega_{1}, \omega_{2}, \omega_{3}$ the standard coordinates of the calibration $\omega$, so that

$$
\omega=\omega_{1} i+\omega_{2} j+\omega_{3} k
$$

Given a chain $C$ in $\partial \mathbf{H}_{\mathbb{H}}^{n}$, let $\mu=\mu_{C}$ be the (Borel positive) measure on $\mathbb{H e i s}_{4 n-1}$ with support $C \cap \mathbb{H e i s}_{4 n-1}$ associated with the volume form $\omega_{1} \wedge \omega_{2} \wedge \omega_{3}$ on $C$. For instance, if $C=\left\{(\zeta, u) \in \mathbb{H}^{n-1} \times \operatorname{Im} \mathbb{H}: \zeta=0\right\} \cup\{\infty\}$ is the standard vertical chain, then $\left.\omega\right|_{C}=\left.\mathrm{d} u\right|_{C}$, so that $\mu_{C}$ is the (infinite) measure

$$
\mu_{C}=\mathrm{d} u_{1} \mathrm{~d} u_{2} \mathrm{~d} u_{3},
$$

whose restriction to the Euclidean space $C-\{\infty\}=\{0\} \times \operatorname{Im} \mathbb{H}$ is the standard Lebesgue measure.

Given a nonzero measure $\mu$ with compact support on a finite dimensional real affine space $V$, the barycenter (or centroid) of $\mu$ is the point $\operatorname{bar}(\mu)$ of $V$ defined by

$$
\operatorname{bar}(\mu)=\frac{1}{\mu(V)} \int_{x \in V} x \mathrm{~d} \mu(x) .
$$

For instance, when $\mu$ is supported on a finite set $S$, then $\operatorname{bar}(\mu)$ is the usual affine barycenter of the weighted family of points $\left\{\left(s, \frac{\mu(\{s\})}{\mu(S)}\right)\right\}_{s \in S}$.

We denote the open ball of center 0 and radius $r$ in the Euclidean space $\operatorname{Im} \mathbb{H}$ by $B(r)$. Recall that the radius of a finite chain $C$ is denoted by $R_{C}$.
Proposition 3.4. Let $C$ be a chain in $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$ and $c \in C$.
(1) If C is a finite chain, then the center of the chain $C$ is equal to the barycenter of the measure $\mu_{C}$ :

$$
\operatorname{cen}(C)=\operatorname{bar}\left(\mu_{C}\right)
$$

(2) If C is a vertical chain, there is a diffeomorphism $\tau=\tau_{C}: \operatorname{Im} \mathbb{H} \rightarrow C-\{\infty\}$ such that $\tau^{*} \omega=\mathrm{d} v$, unique up to postcomposition by a vertical Heisenberg translation. For every Heisenberg translation or rotation $\gamma$, we have $\tau_{\gamma C}=\gamma \circ \tau_{C}$.
(3) If $C$ is a finite chain, there exists a smooth diffeomorphism $\tau=\tau_{C, c}$ from $B\left(2 \pi R_{C}^{2}\right)$ to $C-\{c\}$, admitting a continuous extension to $\partial B\left(2 \pi R_{C}^{2}\right)$ sending this sphere to $c$, such that $\tau^{*} \omega=\mathrm{d} v$. This mapping is unique up to postcomposition by a Heisenberg rotation preserving $C$ and $c$, and $2 \pi R_{C}^{2}$ is the unique radius for which such a mapping exists.

For every Heisenberg translation or rotation $\gamma$, we have $\tau_{\gamma C, \gamma c}=\gamma \circ \tau_{C, c}$.
Proof. (1). Note that $\mathbb{H e i s}_{4 n-1}=\mathbb{H}^{n-1} \times \operatorname{Im} \mathbb{H}$ has a natural structure of a real affine space, and that the elements of $\mathrm{PB}_{q}$ act by affine transformations on $\mathbb{H e i s}_{4 n-1}$. This can be seen for instance by saying that $\mathbb{H e i s}_{4 n-1}$, identified with the boundary of the projective model of $\mathbf{H}_{\mathbb{H}}^{n}$ minus $\{\infty\}$, is a $\mathrm{PB}_{q}$-invariant affine subspace of the affine chart of the quaternionic projective space defined by the quaternionic projective hyperplane $\left\{\left[z_{0}: z: z_{n}\right] \in \mathbb{P}_{\mathrm{r}}^{n}(\mathbb{H}): z_{n}=0\right\}$, and that the quaternionic projective transformations preserving this hyperplane act by affine transformations on the associated affine chart. Another way is to check, by an easy computation, that the Heisenberg translations, rotations and dilations preserve the barycenters in the real affine space $\mathbb{H}^{n-1} \times \operatorname{Im} \mathbb{H}$ : For instance, for all $\left(\zeta_{0}, u_{0}\right),(\zeta, u),\left(\zeta^{\prime}, u^{\prime}\right) \in \mathbb{H e i s}_{4 n-1}$ and $t \in[0,1]$, we have

$$
\left(\zeta_{0}, u_{0}\right) \cdot\left(t(\zeta, u)+(1-t)\left(\zeta^{\prime}, u^{\prime}\right)\right)=t\left(\zeta_{0}, u_{0}\right) \cdot(\zeta, u)+(1-t)\left(\zeta_{0}, u_{0}\right) \cdot\left(\zeta^{\prime}, u^{\prime}\right)
$$

In particular, the barycenters of measures $\mu$ with compact support on $\mathbb{H e i s}_{4 n-1}$ are equivariant under the Heisenberg translations, rotations and dilations: For every such transformation $\gamma$, we have

$$
\begin{equation*}
\operatorname{bar}\left(\gamma_{*} \mu\right)=\gamma \operatorname{bar}(\mu) . \tag{3.3}
\end{equation*}
$$

In order to prove Assertion (1), by Equations (3.2) and (3.3), and by the transitivity properties of the Heisenberg translations and dilations on chains, we may assume that $n=2$ and that $C$ is a Euclidean sphere with center $(0,0)$ and radius 1 in the horizontal subspace $\left\{(\zeta, u) \in \mathbb{H}^{n-1} \times \mathbb{H}: u=0\right\}$. Since the $\operatorname{Im} \mathbb{H}$-valued 1 -form $\left.\omega\right|_{C}$ is invariant under the Heisenberg rotations, the volume form $\omega_{1} \wedge \omega_{2} \wedge \omega_{3}$ on $C$ is invariant under the Heisenberg rotations. Since the only measure on $C$ invariant under the Heisenberg rotations is, up to a scalar multiple, the Lebesgue measure on the Euclidean sphere $C$, the measure $\mu_{C}$ is a multiple of the Lebesgue measure on $C$. This can also be proved by a direct computation: On the Euclidean sphere $C$, with $\zeta=\zeta_{0}+\zeta_{1} i+\zeta_{2} j+\zeta_{3} k$, we have

$$
\omega_{1} \wedge \omega_{2} \wedge \omega_{3}=-8 \sum_{i=0}^{3}(-1)^{i} \zeta_{i} \mathrm{~d} \zeta_{0} \wedge \cdots \wedge \widehat{\mathrm{~d} \zeta}{ }_{i} \wedge \cdots \wedge \mathrm{~d} \zeta_{3}
$$

Since the barycenter of this measure is exactly the origin $(0,0)$, which is the center of the finite chain $C$, this proves Assertion (1).
(2). First assume that $C$ is the standard vertical chain

$$
C_{\infty}=\left\{(\zeta, u) \in \mathbb{H}^{n-1} \times \operatorname{Im} \mathbb{H}: \zeta=0\right\} \cup\{\infty\}
$$

Let $\tau=\tau_{C_{\infty}}: v \mapsto(0, v)$. Then $\tau$ is a diffeomorphism from $\operatorname{Im} \mathbb{H}$ onto $C_{\infty}-\{\infty\}$, such that $\tau^{*}(\mathrm{~d} u-2 \operatorname{Im}(\bar{\zeta} \mathrm{~d} \zeta))=\mathrm{d} v$. For every vertical Heisenberg translation $\gamma$, the map $\gamma \circ \tau$ is also a diffeomorphism from $\operatorname{Im} \mathbb{H}$ onto $C_{\infty}-\{\infty\}$, and since $\omega$ is invariant under the Heisenberg translations, we also have $(\gamma \circ \tau)^{*} \omega=\mathrm{d} v$.

If $\sigma: \operatorname{Im} \mathbb{H} \rightarrow C_{\infty}-\{\infty\}$ is another diffeomorphism such that $\sigma^{*} \omega=\mathrm{d} v$, then for every $v \in \operatorname{Im} \mathbb{H}$, we have $\sigma^{\prime}(v)-\tau^{\prime}(v) \in T C_{\infty} \cap \operatorname{ker} \omega=\{0\}$, thus the maps $\sigma$ and $\tau$ differ by an element of the vector subspace $C_{\infty}$. Therefore there exists a vertical Heisenberg translation $\gamma$ such that $\sigma=\gamma \circ \tau$.

Now, if $C$ is another vertical chain, there exists a composition $\gamma$ of Heisenberg translations and rotations such that $C=\gamma C_{\infty}$. Defining $\tau_{C}=\gamma \circ \tau_{C_{\infty}}$ gives a diffeomorphism from $\operatorname{Im} \mathbb{H}$ onto $C-\{\infty\}$ such that $\tau_{C}{ }^{*} \omega=\mathrm{d} v$, by the invariance of $\omega$ under the Heisenberg translations and rotations. This proves Assertion (2).
(3). First assume that $C$ is the Euclidean 3-sphere

$$
\left\{(\zeta, u) \in \mathbb{H}^{n-1} \times \operatorname{Im} \mathbb{H}: \mathrm{n}\left(\zeta_{1}\right)=R^{2} \text { and } u=\zeta_{2}=\cdots=\zeta_{n-1}=0\right\}
$$

and that $c=\left(\zeta_{c}=(-R, 0, \ldots, 0), u_{c}=0\right)$. Note that $R$ is the radius of the finite chain $C$. By the properties of the exponential map of the Lie group of unit quaternions, whose tangent space at the identity element 1 is $\operatorname{Im} \mathbb{H}$, the smooth map

$$
\tau=\tau_{C, c}: v \mapsto\left(\zeta=\left(R e^{-v /\left(2 R^{2}\right)}, 0, \ldots, 0\right), u=0\right)
$$

from $\operatorname{Im} \mathbb{H}$ to $C$ is a diffeomorphism from $B\left(2 \pi R^{2}\right)$ onto $C-\{c\}$. It extends continuously (and even smoothly) to the sphere $\partial B\left(2 \pi R^{2}\right)$, mapping this sphere to $c$. Considering $\zeta$ as a function of $v$, we have $\mathrm{d} \zeta=\left(-\frac{1}{2 R} e^{-v /\left(2 R^{2}\right)} \mathrm{d} v, 0, \ldots, 0\right)$. Hence, since $v$ and $\mathrm{d} v$ are purely imaginary quaternions, we have

$$
\tau^{*} \omega=-2 \operatorname{Im}(\bar{\zeta} \cdot \mathrm{~d} \zeta)=-2 \operatorname{Im}\left(\left(R e^{-\bar{v} /\left(2 R^{2}\right)}\right)\left(-\frac{1}{2 R} e^{-v /\left(2 R^{2}\right)} \mathrm{d} v\right)\right)=\mathrm{d} v
$$

The uniqueness of $\tau$ up to postcomposition by a Heisenberg rotation preserving $C$ and $c$, and the extension to the other chains, follow as previously from the fact that the chains are transverse to the quaternionic contact structure on $\mathbb{H e i s}_{4 n-1}$ and by invariance of the calibration $\omega$ under the Heisenberg translations and rotations.

## 4. Counting and equidistribution of arithmetic chains in hyperspherical geometry

In this section, we prove (generalised versions of) Theorems 1.2 and 1.3 of the introduction. We start by recalling a general statement, coming from a special case of the main results of [21], that has been made explicit in [22].

Let $\Gamma$ be a lattice in $\mathrm{PU}_{q}$. Let $D^{-}$and $D^{+}$be nonempty proper closed convex subsets of $\mathbf{H}_{\mathbb{H}}^{n}$, with stabilisers $\Gamma_{D^{-}}$and $\Gamma_{D^{+}}$in $\Gamma$ respectively, such that the families $\left(\gamma D^{-}\right)_{\gamma \in \Gamma / \Gamma_{D^{-}}}$ and $\left(\gamma D^{+}\right)_{\gamma \in \Gamma / \Gamma_{D^{+}}}$are locally finite in $\mathbf{H}_{\mathbb{H}}^{n}$. For all $\gamma, \gamma^{\prime}$ in $\Gamma$, the convex sets $\gamma D^{-}$and $\gamma^{\prime} D^{+}$have a common perpendicular if and only if their closures $\overline{\gamma D^{-}}$and $\overline{\gamma^{\prime} D^{+}}$in $\mathbf{H}_{\mathbb{H}}^{n} \cup \partial_{\infty} \mathbf{H}_{\mathbb{H}}^{n}$ do not intersect. We denote by $\alpha_{\gamma, \gamma^{\prime}}$ this common perpendicular, starting from $\gamma D^{-}$at time $t=0$, and by $\ell\left(\alpha_{\gamma, \gamma^{\prime}}\right)$ its length. The multiplicity of $\alpha_{\gamma, \gamma^{\prime}}$ is

$$
m_{\gamma, \gamma^{\prime}}=\frac{1}{\operatorname{card}\left(\gamma \Gamma_{D^{-}} \gamma^{-1} \cap \gamma^{\prime} \Gamma_{D^{+}} \gamma^{\prime-1}\right)},
$$

which equals 1 for all $\gamma, \gamma^{\prime} \in \Gamma$ when $\Gamma$ acts freely on $T^{1} \mathbf{H}_{\mathbb{H}}^{n}$ (for instance when $\Gamma$ is torsion-free). For all $s>0$ and $x \in \partial D^{-}$, let

$$
m_{s}(x)=\sum_{\gamma \in \Gamma / \Gamma_{D^{+}}: \overline{D^{-}} \cap \overline{\gamma D^{+}}=\emptyset, \alpha_{e, \gamma}(0)=x, \ell\left(\alpha_{e, \gamma}\right) \leq s} m_{e, \gamma}
$$

be the multiplicity of $x$ as the origin of common perpendiculars with length at most $s$ from $D^{-}$to the elements of the $\Gamma$-orbit of $D^{+}$.

For every $s>0$, let

$$
\mathscr{N}_{D^{-}, D^{+}}(s)=\sum_{\left(\gamma, \gamma^{\prime}\right) \in \Gamma \backslash\left(\left(\Gamma / \Gamma_{D^{-}}\right) \times\left(\Gamma / \Gamma_{D^{+}}\right)\right): \overline{\gamma D^{-}} \cap \overline{\gamma^{\prime} D^{+}}=\emptyset, \ell\left(\alpha_{\gamma, \gamma^{\prime}}\right) \leq s} m_{\gamma, \gamma^{\prime}},
$$

where $\Gamma$ acts diagonally on $\Gamma \times \Gamma$. When $\Gamma$ has no torsion, $\mathscr{N}_{D^{-}, D^{+}}(s)$ is the number (with multiplicities coming from the fact that $\Gamma_{D^{ \pm}} \backslash D^{ \pm}$is not assumed to be embedded in $\left.\Gamma \backslash \mathbf{H}_{\mathbb{H}}^{n}\right)$ of the common perpendiculars of length at most $s$ between the images of $D^{-}$and $D^{+}$in $\Gamma \backslash \mathbf{H}_{\mathbb{H}}^{n}$.

The following statement is a special case of [22, Thm. 8•1]. We denote by $\Delta_{x}$ the unit Dirac mass at a point $x$.

Theorem 4.1. Let $D^{-}$be a horoball in $\mathbf{H}_{\mathbb{H}}^{n}$ centred at a parabolic fixed point of $\Gamma$ and let $D^{+}$be a quaternionic geodesic line in $\mathbf{H}_{\mathbb{H}}^{n}$ such that $\Gamma_{D^{+}} \backslash D^{+}$has finite volume. Let $m^{+}$be the order of the pointwise stabiliser of $D^{+}$in $\Gamma$ and let

$$
c\left(D^{-}, D^{+}\right)=\frac{2(n-1)(2 n-1)}{\pi^{2} m^{+}} \frac{\operatorname{Vol}\left(\Gamma_{D^{-}} \backslash D^{-}\right) \operatorname{Vol}\left(\Gamma_{D^{+}} \backslash D^{+}\right)}{\operatorname{Vol}\left(\Gamma \backslash \mathbf{H}_{\mathbb{H}}^{n}\right)} .
$$

There exists $\kappa>0$ such that, as $s \rightarrow+\infty$,

$$
\mathscr{N}_{D^{-}, D^{+}}(s)=c\left(D^{-}, D^{+}\right) e^{(4 n+2) s}\left(1+\mathrm{O}\left(e^{-\kappa s}\right)\right)
$$

Furthermore, the origins of the common perpendiculars from $D^{-}$to the images of $D^{+}$ under the elements of $\Gamma$ equidistribute in $\partial D^{-}$to the induced Riemannian measure: As $s \rightarrow+\infty$, we have

$$
\begin{equation*}
\frac{2(2 n+1) \operatorname{Vol}\left(\Gamma_{D^{-}} \backslash D^{-}\right)}{c\left(D^{-}, D^{+}\right)} e^{-(4 n+2) s} \sum_{x \in \partial D^{-}} m_{s}(x) \Delta_{x} \stackrel{*}{\rightharpoonup} \operatorname{vol}_{\partial D^{-}} \tag{4.1}
\end{equation*}
$$

For smooth functions $\psi$ with compact support on $\partial D^{-}$, there is an error term in the equidistribution claim of Theorem 4.1 when the measures on both sides are evaluated on $\psi$, of the form $\mathrm{O}\left(e^{-\kappa s}\|\psi\|_{\ell}\right)$ where $\kappa>0$ and $\|\psi\|_{\ell}$ is the Sobolev norm of $\psi$ for some $\ell \in \mathbb{N}$.

From now on, we assume that $n=2$. Let $A, D_{A}, m_{A}$ and $\mathscr{O}$ be as in the Introduction. We denote by $\left|\mathscr{O}^{\times}\right|$the order of the unit group of $\mathscr{O}$, equal to 24 if $D_{A}=2$, to 12 if $D_{A}=3$, or else to 2,4 or 6 . See for instance [25]. As usual, by $\prod_{p \mid D_{A}}$, we mean a product where $p$ ranges over the prime positive numbers dividing $D_{A}$.

For every chain $C$ in $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{2}$, let $L_{C}$ be the quaternionic projective line in $\mathbb{P}_{\mathrm{r}}^{2}(\mathbb{H})$ such that $C=L_{C} \cap \partial_{\infty} \mathbf{H}_{\mathbb{H}}^{2}$, and let $D_{C}=L_{C} \cap \mathbf{H}_{\mathbb{H}}^{2}$ be the associated quaternionic geodesic line. For every finite index subgroup $G$ of the arithmetic lattice $\mathrm{PU}_{q}(\mathscr{O})$, we denote by $G_{C}$ the stabiliser of $C$ in $G$, by $G_{\infty}$ the stabiliser of $\infty$ in $G$, and by $\operatorname{Covol}_{G}(C)$ the volume of the orbifold $G_{C} \backslash D_{C}$ for the Riemannian metric of constant sectional curvature -1 on the real hyperbolic 4 -space $D_{C}$. Recall that a chain $C$ is arithmetic over $\mathscr{O}$ if and only if the stabiliser in $\mathrm{PU}_{q}(\mathscr{O})$ (or equivalently in $G$ ) of the quaternionic geodesic line $D_{C}$ has finite covolume on $D_{C}$.

Theorem 4.2. Let $C_{0}$ be an arithmetic chain over a maximal order $\mathscr{O}$ in a definite quaternion algebra over $\mathbb{Q}$. Let $G$ be a finite index subgroup of $\mathrm{PU}_{q}(\mathscr{O})$. Then there exists a constant $\kappa>0$ such that, as $\epsilon>0$ tends to 0 , the number $\psi_{C_{0}, G}(\epsilon)$ of chains modulo $G_{\infty}$ in the $G$-orbit of $C_{0}$ with $d_{\mathrm{Cyg}}$-diameter at least $\epsilon$ is equal to

$$
\frac{352^{23} 3^{6} D_{A}^{2} \operatorname{Covol}_{G}\left(C_{0}\right)\left[\mathrm{PU}_{q}(\mathscr{O})_{\infty}: G_{\infty}\right]}{\pi^{6} m_{C_{0}, G} m_{A}\left|\mathscr{O}^{\times}\right|^{2} \prod_{p \mid D_{A}}(p-1)\left(p^{2}+1\right)\left(p^{3}-1\right)\left[\mathrm{PU}_{q}(\mathscr{O}): G\right]} \epsilon^{-10}\left(1+\mathrm{O}\left(\epsilon^{\kappa}\right)\right),
$$

where $m_{C_{0}, G}$ is the order of the pointwise stabiliser of $D_{C_{0}}$ in $G$.
Recall that the center cen $(C)$ of a finite chain $C$ is the image of $\infty=[1: 0: 0]$ under the reflexion on $L_{C}$. The following result is an equidistribution result in the quaternionic Heisenberg group of the centers of the arithmetic chains in a given orbit under (a finite index subgroup of) $\mathrm{PU}_{q}(\mathscr{O})$.

Theorem 4.3. Let $C_{0}, G$ and $m_{C_{0}, G}$ be as in Theorem 4.2. As $\epsilon>0$ tends to 0 , we have

$$
\left.\frac{m_{C_{0}, G} m_{A} \pi^{6} \prod_{p \mid D_{A}}(p-1)\left(p^{2}+1\right)\left(p^{3}-1\right)\left[\mathrm{PU}_{q}(\mathscr{O}): G\right]}{352^{24} 3^{6} \operatorname{Covol}_{G}\left(C_{0}\right)} \epsilon^{10} \sum_{\substack{C \in G \cdot C_{0} \\ \operatorname{diam}_{d_{C \text { Cyg }}(C) \geq \epsilon} \\(C)}} \Delta_{\text {cen }(C)}\right)
$$

As in Theorem 4.1, there exist $\kappa>0$ and $\ell \in \mathbb{N}$ such that for every smooth function $\psi$ with compact support on $\mathbb{H e i s}_{7}$, there is an error term in this equidistribution result when the measures on both sides are evaluated on $\psi$, of the form $\mathrm{O}\left(s^{-\kappa}\|\psi\|_{\ell}\right)$ where $\|\psi\|_{\ell}$ is the Sobolev norm of $\psi$.

We begin by a technical result used in the proofs of the above theorems, which does not require the assumption $n=2$. Recall that $d_{\mathrm{Cyg}}^{\prime \prime}$ is the modified Cygan distance defined in Section 2.

Lemma 4.4. For every m-chain $C$ in $\mathbf{H}_{\mathbb{H}}^{n}$, we have $\operatorname{diam}_{d_{\text {cyg }}}(C)=\sqrt{2} \operatorname{diam}_{d_{\text {Cyg }}^{\prime \prime}}(C)$.
Proof. If $C$ is a vertical $m$-chain, then both diameters are $+\infty$. We hence assume that $C$ is finite. Since the Heisenberg translations and rotations preserve $d_{\mathrm{Cyg}}$ and $d_{\mathrm{Cyg}}^{\prime \prime}$, and by the transitivity properties of the Heisenberg translations and rotations on the set of $m$-chains (see Section 3.2), we may assume that $C$ is a Euclidean sphere centered at $(0,0)$ with dimension $4 m-1$, contained in the horizontal plane $\mathbb{H}^{n-1} \times\{0\}$ of $\mathbb{H e i s}_{4 n-1}$. Since the Heisenberg dilations $(\zeta, u) \mapsto\left(\lambda \zeta, \lambda^{2} u\right)$ with $\lambda>0$ are homotheties of ratio $\lambda$ for $d_{\text {Cyg }}$ and $d_{\text {Cyg }}^{\prime \prime}$, we may assume that the radius of $C$ is equal to 1 .

For every $(\zeta, 0) \in C$, we thus have $d_{\text {Cyg }}((\zeta, 0),(0,0))=1$ by Equation (2.7), hence $\operatorname{diam}_{d_{\mathrm{Cyg}}}(C) \leq 2$ by the triangle inequality. Since

$$
d_{\mathrm{Cyg}}((\zeta, 0),(-\zeta, 0))=d_{\mathrm{Cyg}}((\zeta, 0) \cdot(\zeta, 0),(0,0))=d_{\mathrm{Cyg}}((2 \zeta, 0),(0,0))=2,
$$

we have $\operatorname{diam}_{d_{\text {Cyg }}}(C)=2$.
Using the transitivity properties of $\operatorname{Sp}(n-1)$ on the unit sphere $C$ of the Euclidean space $\mathbb{H}^{n-1}$ in the same way as in the proof of [20, Lem. 8] in the complex hyperbolic case, we may assume that $n=3$, and that

$$
\operatorname{diam}_{d_{\mathrm{Cyg}}^{\prime \prime}}(C)=\sup _{u \in \mathbb{H}, \phi \in[0, \pi]: \mathrm{n}(u)=1} d_{\mathrm{Cyg}}^{\prime \prime}((1,0,0),(u \cos \phi, \sin \phi, 0)) .
$$

By a computation similar to the one in [20, Lem. 8], using Equation (2.8) and the fact that $4 \mathrm{n}(\operatorname{Im} u)=4-(\operatorname{tr} u)^{2}$ for any unit quaternion $u$, we have

$$
\begin{aligned}
d_{\text {Cyg }}^{\prime \prime}((1,0,0),(u \cos \phi & \sin \phi, 0))^{2} \\
& =d_{\mathrm{Cyg}}^{\prime \prime}((0,0,0),(-1,0,0) \cdot(u \cos \phi, \sin \phi, 0))^{2} \\
& =d_{\mathrm{Cyg}}^{\prime \prime}((0,0,0),(u \cos \phi-1, \sin \phi,-2 \cos \phi \operatorname{Im} u))^{2} \\
& =\frac{(2-\cos \phi \operatorname{tr} u)^{2}+4 \cos ^{2} \phi \mathrm{n}(\operatorname{Im} u)}{\left((2-\cos \phi \operatorname{tr} u)^{2}+4 \cos ^{2} \phi \mathrm{n}(\operatorname{Im} u)\right)^{\frac{1}{2}}+(2-\cos \phi \operatorname{tr} u)} \\
& =\frac{2}{\frac{1}{\left(1+\cos ^{2} \phi-\cos \phi \operatorname{tr} u\right)^{\frac{1}{2}}}+\frac{2-\operatorname{tr} u \cos \phi}{2\left(1+\cos ^{2} \phi-\cos \phi \operatorname{tr} u\right)}} .
\end{aligned}
$$

As $1+\cos ^{2} \phi-\cos \phi \operatorname{tr} u \leq 2-\cos \phi \operatorname{tr} u \leq 4$, we have

$$
d_{\mathrm{Cyg}}^{\prime \prime}((1,0,0),(u \cos \phi, \sin \phi, 0))^{2} \leq 2 .
$$

Furthermore, the equality holds when $u=1$ and $\phi=\pi$. This proves the result.
Proof of Theorem 4.2 and Theorem 4.3. The diameter of a chain for the Cygan distance is invariant under the stabiliser in $\mathrm{PU}_{q}$ of the horosphere $\partial \mathscr{H}_{1}$, hence is invariant under $G_{\infty}$. The counting function $\psi_{C_{0}, G}$ is thus well defined.

Note that $\mathscr{H}_{1}$ is a horoball centered at the fixed point of a parabolic element in $\mathrm{PU}_{q}(\mathscr{O})$ (take the vertical Heisenberg translation by $(0,2 u)$ for any nonzero $u \in \mathscr{O} \cap \operatorname{Im} \mathbb{H}$ ). We will apply Theorem 4.1 with $\Gamma=G$, with $D^{-}=\mathscr{H}_{1}$, which is hence a horoball centered at the fixed point of a parabolic element in $G$, and with $D^{+}=D_{C_{0}}$, which is the quaternionic geodesic line in $\mathbf{H}_{\mathbb{H}}^{2}$ with boundary at infinity equal to $C_{0}$. In particular $m^{+}=m_{C_{0}, G}$.

Let us compute the constant $c\left(D^{-}, D^{+}\right)$appearing in the statement of Theorem 4.1. We have $\operatorname{Vol}\left(G \backslash \mathbf{H}_{\mathbb{H}}^{2}\right)=\left[\mathrm{PU}_{q}(\mathscr{O}): G\right] \operatorname{Vol}\left(\mathrm{PU}_{q}(\mathscr{O}) \backslash \mathbf{H}_{\mathbb{H}}^{2}\right)$, where, by [22, Thm. 1.4],

$$
\operatorname{Vol}\left(\operatorname{PU}_{q}(\mathscr{O}) \backslash \mathbb{H}_{\mathbb{H}}^{2}\right)=\frac{\pi^{4} m_{A}}{1752^{13} 3^{5}} \prod_{p \mid D_{A}}(p-1)\left(p^{2}+1\right)\left(p^{3}-1\right),
$$

and by [22, Lem. 8.4],

$$
\begin{equation*}
\operatorname{Vol}\left(\Gamma_{D^{-}} \backslash D^{-}\right)=\left[\mathrm{PU}_{q}(\mathscr{O})_{\infty}: G_{\infty}\right] \operatorname{Vol}\left(\mathrm{PU}_{q}(\mathscr{O})_{\left.\mathscr{H}_{1} \backslash \mathscr{H}_{1}\right)}=\frac{D_{A}^{2}\left[\mathrm{PU}_{q}(\mathscr{O})_{\infty}: G_{\infty}\right]}{160\left|\mathscr{O}^{\times}\right|^{2}}\right. \tag{4.2}
\end{equation*}
$$

By definition, we have

$$
\operatorname{Vol}\left(\Gamma_{D^{+}} \backslash D^{+}\right)=16 \operatorname{Covol}_{G}\left(C_{0}\right)
$$

since the sectional curvature of $D^{+}$is constant -4 and $D^{+}$has real dimension 4 . We hence have

$$
\begin{equation*}
c\left(D^{-}, D^{+}\right)=\frac{352^{13} 3^{6} D_{A}^{2} \operatorname{Covol}_{G}\left(C_{0}\right)\left[\mathrm{PU}_{q}(\mathscr{O})_{\infty}: G_{\infty}\right]}{\pi^{6} m_{C_{0}, G} m_{A}\left|\mathscr{O}^{\times}\right|^{2} \prod_{p \mid D_{A}}(p-1)\left(p^{2}+1\right)\left(p^{3}-1\right)\left[\mathrm{PU}_{q}(\mathscr{O}): G\right]} \tag{4.3}
\end{equation*}
$$

Let $g \in G$ be such that the quaternionic geodesic line $g D^{+}$is disjoint from $\mathscr{H}_{1}$ (which is the case except for $g$ in finitely many double classes in $G_{\mathscr{H}_{1}} \backslash G / G_{D^{+}}$). Let $\delta_{g}$ be the common perpendicular from $\mathscr{H}_{1}$ to $g D^{+}$. Its length $\ell\left(\delta_{g}\right)$ is the minimum of the distances from $\mathscr{H}_{1}$ to a geodesic line between two points of $\partial_{\infty}\left(g D^{+}\right)=g C_{0}$. Hence, by Lemmas 2.1 and 4.4 , we have

$$
\begin{align*}
\ell\left(\delta_{g}\right) & =\min _{x, y \in g C_{0}, x \neq y} d\left(\mathscr{H}_{1},\right] x, y[)=-\max _{x, y \in g C_{0}, x \neq y} \ln \frac{d_{\mathrm{Cyg}}^{\prime \prime}(x, y)}{\sqrt{2}} \\
& =-\ln \frac{\operatorname{diam}_{d_{\mathrm{Cyg}}^{\prime \prime}}\left(g C_{0}\right)}{\sqrt{2}}=-\ln \frac{\operatorname{diam}_{d_{\mathrm{Cyg}}}\left(g C_{0}\right)}{2} . \tag{4.4}
\end{align*}
$$

Respectively by the definition of the counting function $\psi_{C_{0}, G}$ in the statement of Theorem 4.2, since the stabiliser of $C_{0}$ in $G$ is equal to $G_{D_{C_{0}}}=G_{D^{+}}$, by Equation (4.4), by Theorem 4.1, and by Equation (4.3), we have, as $\epsilon>0$ tends to 0 ,

$$
\begin{aligned}
& \psi_{C_{0}, G}(\epsilon) \\
& =\operatorname{card} G_{\infty} \backslash\left\{C \in G \cdot C_{0}: \operatorname{diam}_{d_{\mathrm{Cyg}}}(C) \geq \epsilon\right\} \\
& =\operatorname{card}\left\{[g] \in G_{\infty} \backslash G / G_{D_{C_{0}}}: \operatorname{diam}_{d_{\mathrm{Cyg}}}\left(g C_{0}\right) \geq \epsilon\right\} \\
& =\operatorname{card}\left\{[g] \in G_{\mathscr{H}} \backslash G / G_{D_{C_{0}}}: \ell\left(\delta_{g}\right) \leq-\ln \frac{\epsilon}{2}\right\}+\mathrm{O}(1) \\
& =\mathscr{N}_{D^{-}, D^{+}}\left(-\ln \frac{\epsilon}{2}\right)+\mathrm{O}(1)=c\left(D^{-}, D^{+}\right) e^{-10 \ln \frac{\epsilon}{2}}\left(1+\mathrm{O}\left(e^{\kappa \ln \frac{\epsilon}{2}}\right)\right) \\
& =\frac{352^{23} 3^{6} D_{A}^{2} \operatorname{Covol}_{G}\left(C_{0}\right)\left[\mathrm{PU}_{q}(\mathscr{O})_{\infty}: G_{\infty}\right]}{\pi^{6} m_{C_{0}, G} m_{A}\left|\mathscr{O}^{\times}\right|^{2} \prod_{p \mid D_{A}}(p-1)\left(p^{2}+1\right)\left(p^{3}-1\right)\left[\mathrm{PU}_{q}(\mathscr{O}): G\right]} \epsilon^{-10}\left(1+\mathrm{O}\left(\epsilon^{\kappa}\right)\right) .
\end{aligned}
$$

This proves Theorem 4.2. Let us now prove Theorem 4.3.
We apply the equidistribution result in Equation (4.1) of the origins or $\left(\delta_{g}\right)$ of the common perpendiculars $\delta_{g}$ from $D^{-}=\mathscr{H}_{1}$ to the images $g D^{+}$for $g \in G$. As $s \rightarrow+\infty$, we hence have, using Equations (4.3) and (4.2),

$$
\frac{m_{C_{0}, G} m_{A} \pi^{6} \prod_{p \mid D_{A}}(p-1)\left(p^{2}+1\right)\left(p^{3}-1\right)\left[\mathrm{PU}_{q}(\mathscr{O}): G\right]}{352^{17} 3^{6} \operatorname{Covol}_{G}\left(C_{0}\right)} e^{-10 s} \sum_{[g] \in G / G_{D^{+}}: \ell\left(\delta_{g}\right) \leq s} \Delta_{\operatorname{or}\left(\delta_{g}\right)} \stackrel{*}{\stackrel{*}{v}} \operatorname{vol}_{\partial \mathscr{H}_{1}}
$$

Let $f: \partial_{\infty} \mathbf{H}_{\mathbb{H}}^{2}-\{\infty\}=\mathbb{H e i s}_{7} \rightarrow \partial \mathscr{H}_{1}$ be the orthogonal projection map, which is the homeomorphism $\left(w_{0}, w\right) \mapsto\left(w_{0}+\frac{1}{2}, w\right)$. The pushforward of the Haar measure Haar $_{\text {Heis }}$ by $f$ is

$$
\begin{equation*}
f_{*} \text { Haar }_{\text {Heis } 7}=8 \operatorname{vol}_{\partial \mathscr{H}_{1}}, \tag{4.6}
\end{equation*}
$$

see for example the end of the proof of Theorem 8.3 in [22].
Note that, for every chain $C$, if $r_{C}$ is the reflexion on the quaternionic projective line containing $C$, then the geodesic line from $\infty$ to $\operatorname{cen}(C)=r_{C}(\infty)$, being invariant under $r_{C}$, is orthogonal to the quaternionic geodesic line with boundary at infinity $C$. Hence for every $g \in G$, we have

$$
f^{-1}\left(\operatorname{or}\left(\delta_{g}\right)\right)=\operatorname{cen}\left(g C_{0}\right)
$$

Let us use in Equation (4.5) the change of variables $s=-\ln \frac{\epsilon}{2}$ and the continuity of the pushforward of measures by $f^{-1}$. By Equations (4.4) and (4.6), as $\epsilon>0$ tends to 0 , we obtain that the measures

$$
\frac{m_{C_{0}, G} m_{A} \pi^{6} \prod_{p \mid D_{A}}(p-1)\left(p^{2}+1\right)\left(p^{3}-1\right)\left[\mathrm{PU}_{q}(\mathscr{O}): G\right]}{352^{24} 3^{6} \operatorname{Covol}_{G}\left(C_{0}\right)} \epsilon^{10} \sum_{\begin{array}{c}
{[g] \in G / G_{D^{+}}} \\
\operatorname{diam}_{d_{\mathrm{Cyg}}}\left(g C_{0}\right) \geq \epsilon
\end{array}} \Delta_{\operatorname{cen}\left(g C_{0}\right)}
$$

weak-star converge to the Haar measure Haar Heis $_{7}$. This proves Theorem 4.3.
Example 4.5. Let $C_{0}=\left\{\left[w_{0}: 0: 1\right] \in \mathbb{P}_{\mathbf{r}}^{2}(\mathbb{H}): \operatorname{tr} w_{0}=0\right\}$ be the standard vertical chain in $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{2}$, which is the intersection of $\partial_{\infty} \mathbf{H}_{\mathbb{H}}^{2}$ with the quaternionic projective line $L_{C_{0}}=\left\{\left[z_{0}: z_{1}: z_{2}\right] \in \mathbb{P}_{\mathrm{r}}^{2}(\mathbb{H}): z_{1}=0\right\}$.

An element

$$
\pm\left(\begin{array}{lll}
a & \gamma^{*} & b \\
\alpha & M & \beta \\
c & \delta^{*} & d
\end{array}\right)
$$

of $\mathrm{PU}_{q}$ preserving the quaternionic geodesic line $L_{C_{0}} \cap \mathbf{H}_{\mathbb{H}}^{2}$ satisfies $\alpha w_{0}+\beta=0$ for all $w_{0} \in \mathbb{H}$ with $\operatorname{tr} w_{0}>0$. Thus, $\alpha=\beta=0$, and Equations (2.5) (or rather the similar equations obtained by the formula $X X^{*}=I_{n+1}$ instead of $X^{*} X=I_{n+1}$ ) imply that $\gamma=\delta=0$. Using again Equations (2.5), we see that the stabiliser of $L_{C_{0}}$ consists of the elements

$$
\left(\begin{array}{ccc}
a & 0 & b \\
0 & M & 0 \\
c & 0 & d
\end{array}\right)
$$

such that $\operatorname{tr}(\bar{c} a)=\operatorname{tr}(\bar{d} b)=0, \bar{c} b+\bar{a} d=1$ and $M \in \mathscr{O}^{\times}$. Thus,

$$
\operatorname{Covol}_{\mathrm{PU}_{q}(\mathscr{O})}\left(C_{0}\right)=\frac{\pi^{2}}{1080} \prod_{p \mid D_{A}}(p-1)\left(p^{2}+1\right)
$$

by [5, Thm. 2.5].
The pointwise stabiliser of $C_{0}$ in $\mathrm{PU}_{q}(\mathscr{O})$ consists of the diagonal elements with $a=d= \pm 1$ and $M \in \mathscr{O}^{\times}$, giving $m_{C_{0}, \mathrm{PU}}^{q}(\mathscr{O})=\left|\mathscr{O}^{\times}\right|$.

Theorems 4.2 and 4.3 then give

$$
\psi_{C_{0}, \mathrm{PU}}^{q}(\mathscr{O})(\epsilon)=\frac{1892^{20} D_{A}^{2}}{\pi^{4} m_{A}\left|\mathscr{O}^{\times}\right|^{3} \prod_{p \mid D_{A}}\left(p^{3}-1\right)} \epsilon^{-10}\left(1+\mathrm{O}\left(\epsilon^{\kappa}\right)\right),
$$

and

$$
\frac{\pi^{4} m_{A}\left|\mathscr{O}^{\times}\right| \prod_{p \mid D_{A}}\left(p^{3}-1\right)}{1892^{21}} \epsilon^{10} \sum_{C \in \mathrm{PU}_{q}(\mathscr{O}) \cdot C_{0}: \operatorname{diam}_{d_{\mathrm{Cyg}}} C \geq \epsilon} \Delta_{\mathrm{cen}(C)} \stackrel{*}{\rightharpoonup} \operatorname{Haar}_{\mathrm{Heceis}}^{7} \text {. }
$$

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[^0]:    Keywords: counting, equidistribution, Cartan chain, quaternionic Heisenberg group, Cygan distance, subRiemannian geometry, quaternionic hyperbolic geometry.
    2020 Mathematics Subject Classification: 11E39, 11F06, 11N45, 20G20, 53C17, 53C55.

